



Laboratoire de Chimie Théorique

Development of a Relativistic One-Dimensional Model Including Quantum Electrodynamics Effects

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June 18, 2025

Thesis work with Julien Toulouse

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Introduction



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Why are we interested in relativistic quantum chemistry?



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• Important for fast-moving particles: average electron velocity in hydrogen-like ground-state $v = (\alpha Z)\,c, \ \, \text{with}\,\,\alpha \approx 1/137$



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- Explains some physical properties of heavy elements
- yellow color of gold
 Pyykkö, ACIE, (2004)

P. Romaniello and P. L. de Boeij, JCP, (2005)

- lead-acid battery electro-chemical potential
 R. Ahuia, A. Blomovist, P. Larsson, P. Pvykkö and P. Zaleski-Eigierd, PRL. (2011)
 - R. Ahuja, A. Blomqvist, P. Larsson, P. Pyykkö and P. Zaleski-Ejgierd, PRL, (2011)
- liquid state of mercury at room temperature

K. Steenbergen, E. Pahl and P. Schwerdtfeger, JPCL., (2017)



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- Includes the electron spin and spin-orbit coupling

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Free particle energy

Non relativistic:

$$E = \frac{p^2}{2m} \Rightarrow E \in [0, +\infty),$$

• Relativistic:

$$E^2 = m^2c^4 + c^2p^2 \Rightarrow E \in (-\infty, -mc^2] \cup [mc^2, +\infty),$$



Free particle energy

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Link between both

• Taylor expansion of the square root for the positive part

$$E = \underbrace{mc^{2}}_{\text{Rest mass}} + \underbrace{\frac{p^{2}}{2m}}_{\text{Kinetic energy}} - \underbrace{\frac{p^{4}}{8m^{3}c^{2}}}_{\text{1st order}} + \dots$$



From classical to quantum - Klein-Gordon equation

$$\begin{array}{rcl} E^2 & = & p^2c^2 + m^2c^4 \\ -\hbar^2\partial_t^2 & = & -\hbar^2c^2\nabla^2 + m^2c^4 \end{array}$$

- Negative energy solutions give negative probability...
- Does not include spin



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$$-\hbar^{2}\partial_{t}^{2} = -\hbar^{2}c^{2}\nabla^{2} + m^{2}c^{4}$$

- Negative energy solutions give negative probability...
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From classical to quantum - Dirac equation

$$\begin{array}{rcl} E & = & \pm c\sqrt{p^2+m^2c^2} \\ \mathrm{i}\hbar\partial_t & = & c\sqrt{p^2+m^2c^2} \end{array}$$

• Expand the term under the square root as a perfect square:

$$p^{2} + m^{2}c^{2} = (\vec{\boldsymbol{\alpha}} \cdot \vec{p} + \boldsymbol{\beta}mc)^{2}$$
$$i\hbar \partial_{t} = c \vec{\boldsymbol{\alpha}} \cdot \vec{p} + \boldsymbol{\beta}mc^{2}$$



Dirac operator

$$D_0(x) = c \left(\vec{\boldsymbol{a}} \cdot \vec{p} \right) + \boldsymbol{\beta} \ mc^2, \tag{1}$$

•
$$\vec{a} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}$$
 where $\vec{\sigma} = \{\sigma_x, \ \sigma_y, \ \sigma_z\}$ are the Pauli matrices

$$\bullet \ \pmb{\beta} = \begin{pmatrix} \mathbb{I}_2 & 0 \\ 0 & -\mathbb{I}_2 \end{pmatrix}$$



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Dirac eigenvalue equation

$$(\mathbf{D}_{0}(x) + V(x))\mathbf{\psi}(x) = \mathcal{E}\mathbf{\psi}(x)$$

$$(D_0(x) + V(x))\psi(x) = \mathcal{E}\psi(x)$$
 (2)

•
$$\boldsymbol{\psi} = \begin{pmatrix} \psi^{L} \\ \psi^{S} \end{pmatrix}$$
 where $\psi^{L/S} = \begin{pmatrix} \psi^{L/S}_{\alpha} \\ \psi^{L/S}_{\beta} \end{pmatrix}$

• \mathscr{E} is the energy of the state ψ



Non-relativistic spectrum

$$\mathcal{H}(x) = -\frac{\Delta}{2m} + V(x) \tag{3}$$

Spectrum:

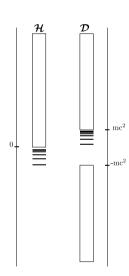
- If $\mathcal{E} > 0$: continuum
- If $\mathcal{E} < 0$: bound states

Relativistic spectrum

$$D(x) = c \left(\vec{\boldsymbol{\alpha}} \cdot \vec{p} \right) + \boldsymbol{\beta} \ mc^2 + V(x)$$
 (4)

Spectrum:

- If $\mathscr{E} \in (-\infty, -mc^2] \cup [mc^2, +\infty)$: continuum
- If $\mathcal{E} \in (-mc^2, mc^2)$: bound states
 - ⇒ Negative continuum spectrum

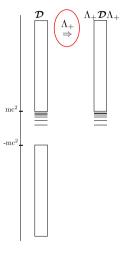


How to deal with the negative continuum

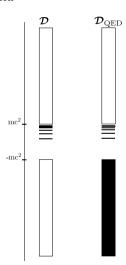


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• No-pair approximation



QED description





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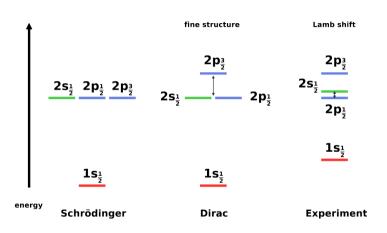


Figure 1: Lamb shift in hydrogen spectra¹

¹M. Salman, Ph.D. thesis (2022)

Motivations for the model



- We need accurate QED calculations for high precision spectroscopy
- So far, state-of-the-art relativistic quantum chemistry for many-electrons systems is based on no-pair approximation
- Next challenge in QC: can we go beyond the NPA?

Goal

Use a model system to understand how to develop new methods based on QED.

To do list:

- ☐ Study the 1D hydrogen and compute the QED effects.
- ☐ Study how to approximate these quantities in a finite basis set.
- \square Use our insight to solve the 3D problem.

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Free 1D Dirac operator

$$\mathbf{D}_{0}(x) = -\mathrm{i}c\sigma_{x}\frac{\mathrm{d}}{\mathrm{d}x} + \sigma_{z}mc^{2} \tag{5}$$

Lapidus, AJP (1983): T. Audinet, J. Toulouse, JCP (2023)

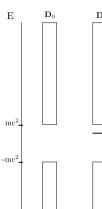
Dirac eigenvalue equations

$$\mathbf{D}_0(x)\psi_p^0(x) = \varepsilon_p \psi_p^0(x) \tag{6}$$

• Continuum: $\varepsilon_k = \pm \sqrt{m^2c^4 + k^2c^2}$

$$(\mathbf{D}_0(x) - Z\delta(x))\psi_p^Z(x) = \varepsilon_p \psi_p^Z(x) \tag{7}$$

- Bound State: $\varepsilon_b < mc^2$
- Continuum: $\varepsilon_k = \pm \sqrt{m^2c^4 + k^2c^2}$



 \mathbf{E}



- Spontaneous creation of electron positron pairs due to the external potential
- Creates a charge density that will interact with the bound states through the two-electron interaction

Vacuum Polarization Density

$$n^{\text{vp}}(x) = \sum_{e_n < 0} \psi_p^{Z\dagger}(x) \psi_p^{Z}(x) - \sum_{e_n < 0} \psi_p^{0\dagger}(x) \psi_p^{0}(x)$$

- In 3D this quantity diverges, needs to be renormalized
- We want to have a better understanding of this quantity and its influence on the energy spectrum



 \mathbf{E}

 mc^2

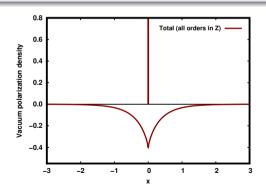
 $-mc^2$



Vacuum Polarization Density

$$n^{\mathrm{vp}}(x) = \sum_{\varepsilon_p < 0} \psi_p^{\mathrm{Z}\dagger}(x) \psi_p^{\mathrm{Z}}(x) - \sum_{\varepsilon_p < 0} \psi_p^{0\dagger}(x) \psi_p^{0}(x) = \mathcal{N}_0^{\mathrm{vp}} \delta(x) + n_{\mathrm{reg}}^{\mathrm{vp}}(x)$$
(8)

Audinet, Morellini, Levitt, Toulouse, JPA, (2025)



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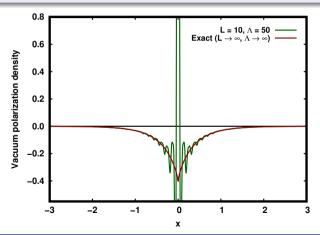
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Plane waves basis

$$\forall x \in \mathbb{R}, k \in \frac{2\mathbb{Z}\pi}{L}, |k| \le \Lambda, \quad \chi_k(x) = \frac{1}{\sqrt{L}} e^{ikx}$$
(9)

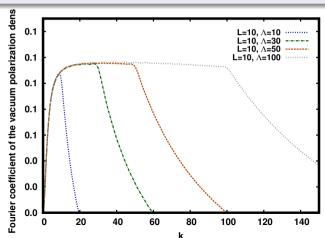




Regularization with Fourier Transform

$$n_{\text{reg}}^{\text{vp}}(x) = \mathscr{F}^{-1} \Big[\mathscr{F}[n^{\text{vp}}](k) - \mathscr{F}[n^{\text{vp}}](k_{\text{max}}) \Big] \theta(k_{\text{max}} - k)$$

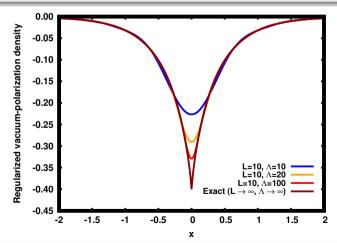
$$\tag{10}$$





Regularization with Fourier Transform

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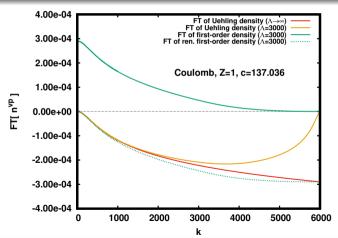
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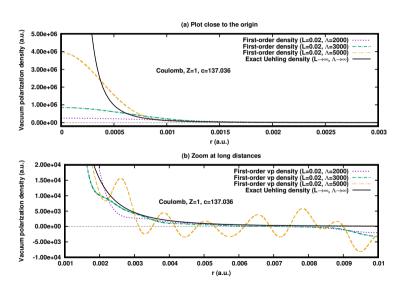
The truncated first-order vacuum-polarization density

$$\hat{n}_{\Lambda}^{\text{vp,(1)}}(\mathbf{k}) = \hat{n}_{\Lambda}^{\text{vp,U}}(\mathbf{k}) + \hat{n}_{\Lambda}^{\text{vp,renorm}}(\mathbf{k})$$
(12)

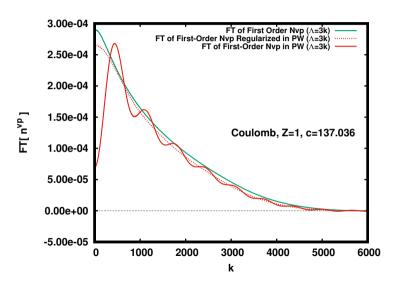




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- ☐ Use our insight to solve the 3D problem (in progress).

- We developed an effective QED theory including electron-positron pairs.
- Some singularities appears in a finite basis.
- This 1D model help us to understand this problem.
- Promising results for the 3D but limited by the size of the basis.