

## Problem Sheet 4

### Problem 9: A simple bremsstrahlung cascade

We want to simulate a simple bremsstrahlung cascade – a Markov process with a variable  $q$ , which we associate to a typical energy scale at which an emission occurs, and a variable  $z$  which is the momentum fraction left to the emitter after the emission. In the general language of Markov processes, our space of states thus consists of the variable  $z$ , and the time variable is  $q$ . We shall consider transition rates which read

$$P(z'|z, q') = \alpha \frac{1}{q'} \frac{1}{z - z'} \theta(z - z' - z q') \theta(z') \quad \alpha > 0. \quad (1)$$

Simulate the resulting Markov process using the methods derived in the lecture, starting from  $z = 1$  at  $q = 1$  and evolve down to small values  $\mu \sim 0.01$ : Generate subsequent emission steps, and terminate the cascade if the time  $\mu$  has been reached, *i.e.* if the process has performed the last emission at time  $q$  with momentum fraction  $z$  (or has started at these values), determine the next time  $q'$  and state  $z'$  from

$$\begin{aligned} \frac{dP}{dq' dz'} = & \delta(\mu - q') \delta(z - z') \exp \left( - \int_{\mu}^q dk \int d\xi P(\xi|z, k) \right) + \\ & P(z'|z, q') \exp \left( - \int_{q'}^q dk \int d\xi P(\xi|z, k) \right) \theta(q - q') \theta(q' - \mu). \end{aligned} \quad (2)$$

until a time  $q' = \mu$  has been generated. Interesting questions to be asked are:

- What is the average momentum fraction at the end of the cascade, what is its variance?
- What is the distribution of multiplicities of radiated particles?
- How do these quantities depend on the choice of  $\mu$ ?

### Problem 10: The Ising model

We want to simulate the Ising model in two dimensions, *i.e.* we consider a lattice of 'spins' which can take values  $\phi_i = \pm 1$  and which are governed by a Hamiltonian

$$H = -\kappa \sum_{\text{nearest neighbours } i,j} \phi_i \phi_j - h \sum_i \phi_i, \quad (3)$$

in the sense of a thermal average for which we consider distributions of configurations according to (we set  $k_B = 1$ )

$$\frac{\exp \left( -\frac{1}{T} H(\vec{\phi}) \right)}{\int \prod_i d\phi_i \exp \left( -\frac{1}{T} H(\vec{\phi}) \right)}. \quad (4)$$

Here,  $i$  indexes the lattice sites – for a lattice of  $n \times n$  sites, you might choose this as

$$i = n_x + n_y n \quad (5)$$

where  $(n_x, n_y)$  labels the individual lattice sites ranging from 0 to  $n - 1$ , and  $i$  ranges from 0 to  $n^2 - 1$ . We also want to impose periodic boundary conditions, such that lattice site indices are always given by  $n_x \bmod n$  and  $n_y \bmod n$ . You will need a data structure to store the configuration, which now can be just a simple vector, and a way to address the nearest neighbours (often referred to as a 'hopping' function or table).

- Provide all of the data structures.
- Use the Metropolis algorithm to implement local updates of a single lattice site.
- Provide updates of the entire lattice (a so-called 'sweep') by looping over local updates.
- Monitor the energy, *i.e.* measure  $H$  itself as a thermal average.