

ON LANDIS' CONJECTURE IN THE PLANE FOR REAL-VALUED POTENTIALS WITH DECAY

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ABSTRACT. We investigate the quantitative unique continuation properties of real-valued solutions to planar Schrödinger equations with potential functions that exhibit pointwise decay at infinity. That is, for equations of the form $-\Delta u + Vu = 0$ in \mathbb{R}^2 , where $|V(z)| \lesssim \langle z \rangle^{-N}$ for some $N > 0$, we prove that real-valued solutions satisfy exponential decay estimates with a rate that depends explicitly on N . Examples show that the estimates established here are essentially sharp. The case of $N = 0$ corresponds to the Landis conjecture, which was proved for real-valued solutions in the plane in [18], while the case of $N < 0$ was previously investigated by the author in [8]. Here, the proof techniques rely on the ideas presented in [18] combined with conformal transformations and an iteration scheme.

1. INTRODUCTION

This article is concerned with the quantitative unique continuation of solutions to elliptic Schrödinger equations in the plane. We consider equations with potentials that exhibits pointwise decay at infinity. The main result is a nearly sharp estimate for the optimal rate of decay at infinity for real-valued solutions. This theorem may be interpreted as a quantitative Landis-type theorem.

In the late 1960s, E. M. Landis [16] conjectured that if u and V are bounded functions that satisfy

$$(1.1) \quad -\Delta u + Vu = 0 \text{ in } \mathbb{R}^n,$$

and u decays faster than exponentially, i.e., $|u(x)| \lesssim \exp(-c|x|^{1+})$, then it follows that $u \equiv 0$. This conjecture was later disproved by Meshkov [19] who constructed non-trivial complex-valued functions u and V that solve (1.1) in \mathbb{R}^2 , where V is bounded and $|u(x)| \lesssim \exp(-c|x|^{4/3})$. Using Carleman estimate techniques, Meshkov also proved a *qualitative unique continuation* result: If (1.1) holds, where V is bounded and u satisfies a decay estimate of the form $|u(x)| \lesssim \exp(-c|x|^{4/3+})$, then necessarily $u \equiv 0$.

In their work on Anderson localization [2], Bourgain and Kenig established a quantitative version of Meshkov's result. As a first step in their proof, they used three-ball inequalities derived from Carleman estimates to establish *order of vanishing* estimates for local solutions to Schrödinger equations of the form (1.1). Then, through a scaling argument, they proved a *quantitative unique continuation* result. More specifically, they showed that if u and V are bounded and satisfy (1.1), and u is normalized so that $|u(0)| \geq 1$, then for sufficiently large values of R ,

$$(1.2) \quad \inf_{|x_0|=R} \|u\|_{L^\infty(B_1(x_0))} \geq \exp(-CR^\beta \log R),$$

where $\beta = \frac{4}{3}$. Since $\frac{4}{3} > 1$, the constructions of Meshkov, in combination with the qualitative and quantitative unique continuation theorems just described, indicate that Landis' conjecture cannot be true for complex-valued solutions in \mathbb{R}^2 . However, at the time, Landis' conjecture still remained open in the real-valued and higher-dimensional settings.

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The first significant step towards resolving Landis' conjecture in the real-valued planar setting was made by Kenig, Silvestre and Wang in [14] where they proved a quantitative form of Landis' conjecture under the assumption that the zeroth-order term satisfies $V \geq 0$ a.e. The techniques in [14] exploit the relationship between real-valued solutions to second-order elliptic PDEs in the plane and solutions to complex-valued Beltrami equations. Using similar ideas, analogous results were established in the settings with drift terms [15], variable coefficients [10], singular lower-order terms [12], and when V_- exhibits decay at infinity [4, 7]. Then, in [18], Logunov, Malinnikova, Nadirashvili, and Nazarov proved Landis' conjecture in the real-valued planar setting. Their proof uses the nodal structure of the domain along with a domain reduction technique to eliminate any sign condition on the zeroth-order term. Many of the ideas from [18] are used in this article.

In [5], we studied the quantitative unique continuation properties of solutions to elliptic equations of the form

$$\Delta u + W \cdot \nabla u + Vu = \lambda u \text{ in } \mathbb{R}^n,$$

where V and W exhibit pointwise decay at infinity, and $\lambda \in \mathbb{C}$. With $\langle x \rangle = \sqrt{1 + |x|^2}$, it was shown that if $|V(x)| \lesssim \langle x \rangle^{-N}$ and $|W(x)| \lesssim \langle x \rangle^{-P}$ for $N, P \geq 0$, then the quantitative estimate (1.2) holds with $\beta = \max\{1, \frac{4-2N}{3}, 2-2P\}$ and $\log R$ replaced by a different slowly-decaying function. These quantitative estimates were generalized in [17], where Lin and Wang proved analogous estimates for solutions to the corresponding equations with variable-coefficient leading terms. Qualitative estimates for similar equations are given in [3]. The constructions presented in [5, 6] show that the estimates described in this paragraph are sharp.

In [8], we studied quantitative unique continuation at infinity for real-valued solutions to (1.1) when $n = 2$ and the potential exhibits growth at infinity, i.e., $|V(z)| \lesssim |z|^N$ for $N > 0$. The techniques in that article rely heavily on the ideas in [18] and careful scaling arguments. Here, we address the more difficult setting where the potential exhibits decay at infinity. The precise statement of the main theorem is as follows.

Theorem 1 (Main Theorem). *For some $a_0 \geq 1$ and $N \in (0, 2)$, let $V : \mathbb{R}^2 \rightarrow \mathbb{R}$ satisfy*

$$(1.3) \quad |V(w)| \leq a_0^2 \langle w \rangle^{-N}.$$

Assume that $u : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a solution to

$$(1.4) \quad -\Delta u + Vu = 0 \text{ in } \mathbb{R}^2$$

with the properties that

$$(1.5) \quad |u(0)| = 1$$

and for each $w \in \mathbb{R}^2$,

$$(1.6) \quad |u(w)| \leq \exp\left(c_0 |w|^{1-\frac{N}{2}}\right).$$

For every $\varepsilon \in (0, \frac{N}{2})$, there exists $R_0(N, a_0, c_0, \varepsilon) > 0$ so that whenever $R \geq R_0$, it holds that

$$(1.7) \quad \inf_{|w_0|=R} \|u\|_{L^\infty(B_1(w_0))} \geq \exp\left(-R^{1-\frac{N}{2}+\varepsilon}\right).$$

The results of [5] prove an estimate of the form (1.2) with $\beta = \frac{4-2N}{3} = \frac{4}{3}(1 - \frac{N}{2}) > 1 - \frac{N}{2} + \varepsilon$. In that article, the assumptions are the same as those in Theorem 1, except that u may be complex-valued. Thus, as in the case of bounded V , Theorem 1 illustrates that better estimates hold in the real-valued planar setting.

As illustrated by the following example, Theorem 1 is sharp (up to ε) for all $N \in (0, 2)$. Fix $N \in (0, 2)$, then set $u(z) = \exp\left(-|z|^{1-\frac{N}{2}}\right)$. A computation shows that u satisfies (1.4) where

$$V(z) := \left(1 - \frac{N}{2}\right)^2 \left(|z|^{1-\frac{N}{2}} - 1\right) |z|^{-1-\frac{N}{2}}$$

satisfies $|V(z)| \lesssim |z|^{-N}$. On the other hand, for any $\beta > 0$, with $u(z) = |z|^{-\beta}$ on $|z| > 1$, we see that $\Delta u = \beta^2 |z|^{-\beta-2}$ and therefore u satisfies (1.4) on an exterior domain with $V(z) := \beta^2 |z|^{-2}$. In particular, we may not have exponential behavior when V decays fast enough, which explains why we restrict ourselves to $N < 2$.

To prove Theorem 1, we use an iterative argument that is reminiscent of the one in [5], see also [4, 7, 9, 17]. To initialize the iteration, we apply a quantitative estimate of the form (1.2) with $\beta = 1$. This result, which verifies Landis' conjecture in the real-valued planar setting, was originally proved by Logunov, Malinnikova, Nadirashvili, and Nazarov in [18], and we formulate it in Theorem 3 below. The iteration argument then relies on repeated applications of Proposition 2 which is proved using the ideas from [18]. Roughly speaking, Proposition 2 shows that if an estimate like (1.2) holds with $\beta = \beta_0$, then for some x_1 with $|x_1| \gg |x_0|$, another estimate like (1.2) holds with x_0 replaced by x_1 and $\beta = \beta_1 \in [1, \beta_0]$. When $\beta_0 = 1$, Proposition 2 isn't useful, but when $\beta_0 > 1$, we can decrease the exponent, i.e., make $\beta_1 < \beta_0$. Therefore, to benefit from the iteration argument, we need to transform to a situation where $\beta_0 > 1$. We observe that if u is composed with the real-variable version of the conformal transformation $z \mapsto z^\alpha$, then the new function also satisfies a Schrödinger equation. By choosing $\alpha > 1$ appropriately, we can ensure that the new potential function is bounded and that the new solution function satisfies a version of (1.2) with $\beta > 1$. By repeatedly applying Proposition 2 to the transformed equation, we can make β arbitrarily close to 1. Finally, to reach the conclusion, we undo the change of variables.

We use the notation $B_r(z)$ to denote a ball of radius $r > 0$ centered at the point z , abbreviated by B_r when the center is clear. Generic constants are denoted by c, C and may change from line to line without comment. Specific constants will be indicated by subscripts.

The article is organized as follows. In Section 2, we present a unique continuation theorem for harmonic functions in punctured domains. The content of this section is very similar to [18, Section 5] and [8, Section 2]. The iterative result described by Proposition 2 is the content of Section 3. Proposition 2 is a three-ball inequality for solutions to Schrödinger equations and its proof relies on the results from Section 2. In Section 4, we introduce the real-valued versions of $z \mapsto z^\alpha$ and record some of their properties. In particular, we show how solutions behave when they are composed with these transformations. Finally, the proof of Theorem 1 is presented in Section 5.

2. DECAY PROPERTIES OF HARMONIC FUNCTIONS IN PUNCTURED DOMAINS

In this section, we present and prove quantitative unique continuation results (in the form of three-ball inequalities) for harmonic functions in punctured domains. We recall the following application of the Harnack inequality which appears in [18, Claim 5.2] and is repeated in [8, Lemma 2.1].

Lemma 1 (Harnack application). *Let $\{D_j\}$ be a finite collection of 100-separated unit disks in the plane. Assume that h is real-valued and harmonic in $\mathbb{R}^2 \setminus \bigcup D_j$ and that for each index j , h doesn't change sign in $5D_j \setminus D_j$. There exists an absolute constant $C_H \geq 10$ for which*

- (1) $\max_{\partial(3D_j)} |h| \leq C_H \min_{\partial(3D_j)} |h|$
- (2) $\max_{\partial(3D_j)} |\nabla h| \leq C_H \min_{\partial(3D_j)} |h|$.

Proof. An application of the Harnack inequality shows that there exists $C_H > 0$ so that for every j

$$\max_{\partial(3D_j)} |h| \leq \sup_{4D_j \setminus 2D_j} |h| \leq C_H \inf_{4D_j \setminus 2D_j} |h| \leq C_H \min_{\partial(3D_j)} |h|.$$

For each $z \in \partial(3D_j)$, since h doesn't change signs in $B_2(z)$, an application of Cauchy's inequality as in [13, Lemma 1.11] shows that $|\nabla h(z)| \leq |h(z)|$ and the conclusion follows. \square

Now we state and prove the main result of this section. The following is a slight modification of [8, Lemma 2.1], which resembles the result [18, Theorem 5.3].

Proposition 1 (Three-ball inequality for harmonic function in punctured domain). *Let $\{D_j\}$ be a finite collection of 100-separated unit disks in the plane for which $0 \notin \bigcup 3D_j$. For some $R \geq 2^{10}$, let h be a harmonic function in $B_R \setminus \bigcup D_j$ with the property that for each index j , h doesn't change sign in*

$(5D_j \setminus D_j) \cap B_R$. Assume that for some $S \in (2^8, \frac{R}{2})$ and some $M > \log(16R)$, with $T := R - \frac{S}{32}$, it holds that

$$(2.1) \quad \sup_{B_T \setminus \bigcup 3D_j} |h| \geq e^{-M} \sup_{B_R \setminus \bigcup 3D_j} |h|.$$

Then for every $r \in (0, \frac{R}{2^{10}})$, we have

$$(2.2) \quad \sup_{B_r \setminus \bigcup 3D_j} |h| \geq \left(\frac{16r}{R} \right)^\kappa \sup_{B_T \setminus \bigcup 3D_j} |h|,$$

where $\kappa(R, S, M) = \max \{6C_H R, 2^{13} MRS^{-1}\}$ and $C_H \geq 10$ is from Lemma 1.

Remark 2. In contrast to [8, Lemma 2.1], we no longer assume that S is bounded from below with respect to R and we impose a lower bound on M . This flexibility is useful for the iterative argument.

Proof. We may assume without loss of generality that

$$(2.3) \quad \sup_{B_T \setminus \bigcup 3D_j} |h| = 1.$$

Set $k = \max \{2C_H R, 2^{10} MRS^{-1}\}$. For the sake of contradiction, assume that

$$(2.4) \quad \sup_{B_r \setminus \bigcup 3D_j} |h| \leq \left(\frac{16r}{R} \right)^{3k}.$$

Define the punctured annular region

$$\Omega := \left\{ \frac{r}{2} < |z| < R - 1 \right\} \setminus \bigcup 3D_j$$

and the function

$$f(z) = \frac{h_x - ih_y}{z^k}.$$

Observe that f is analytic in Ω and $|f(z)| = |\nabla h(z)| |z|^{-k}$. We'll analyze the behavior of f over Ω . We begin with bounding h and ∇h over the innermost and outermost parts of the boundary of Ω .

Let W_1 be the connected component of $\partial\Omega$ that intersects the inner circle $\{|z| = \frac{r}{2}\}$. If $z \in W_1$, then there are three cases to consider:

- (a) $|z| \neq \frac{r}{2}$.
- (b) $|z| = \frac{r}{2}$ and there exists j for which $z \in 4D_j \setminus 3D_j$.
- (c) $|z| = \frac{r}{2}$ and $z \cap 4D_j$ is empty for all j .

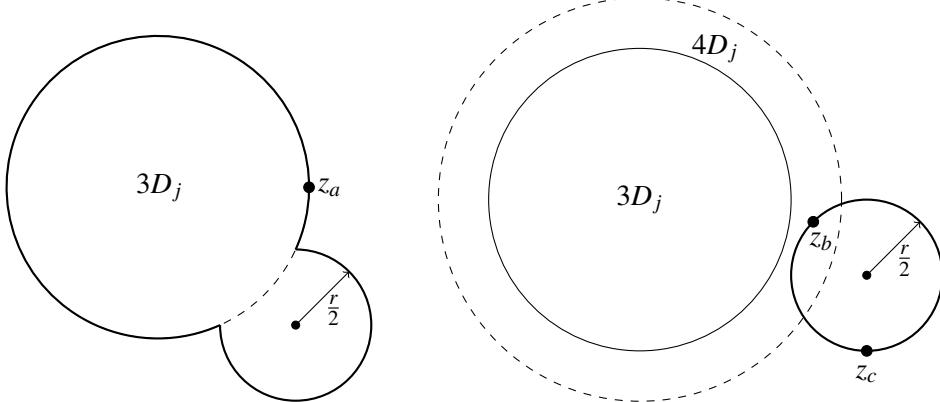


FIGURE 1. Possible images of W_1 with cases (a), (b) and (c) illustrated by the points z_a , z_b , and z_c , respectively.

Case (a): There exists j for which $z \in \partial(3D_j)$ and $3D_j \cap \{|z| = \frac{r}{2}\}$ is non-empty. An application of Lemma 1 combined with the fact that $\partial(3D_j) \cap B_R$ is non-empty shows that

$$|h(z)|, |\nabla h(z)| \leq C_H \min_{\partial(3D_j)} |h| \leq C_H \sup_{B_r \setminus \cup 3D_j} |h| \leq C_H \left(\frac{16r}{R}\right)^{3k},$$

where the last inequality follows from (2.4).

Case (b): Since h doesn't change signs in $B_1(z)$, then an application of [13, Lemma 1.11] shows that

$$|\nabla h(z)| \leq 2|h(z)| \leq 2 \sup_{B_R \setminus \cup 3D_j} |h| \leq 2 \left(\frac{16r}{R}\right)^{3k},$$

where the second inequality uses that $z \notin \cup 3D_j$ and we have again applied (2.4).

Case (c): Let $d = \min\{1, \frac{r}{2}\}$ and observe that $B_d(z) \subset B_R \setminus \cup 3D_j$, so an application of Cauchy's inequality, [13, Lemma 1.10], shows that

$$|\nabla h(z)| \leq \frac{2}{d} \sup_{B_d(z)} |h| \leq \frac{2}{d} \sup_{B_R \setminus \cup 3D_j} |h| \leq \frac{2}{d} \left(\frac{16r}{R}\right)^{3k}.$$

If $d = 1$, since $k > 1$, then $\frac{2}{d} \left(\frac{16r}{R}\right)^k = 2 \left(\frac{16r}{R}\right)^k < \frac{32r}{R} < \frac{1}{2^5} < \frac{1}{2}$. On the other hand, if $d = \frac{r}{2}$, then

$$\frac{2}{d} \left(\frac{16r}{R}\right)^k = \frac{4}{r} \left(\frac{16r}{R}\right)^k = \frac{64}{R} \left(\frac{16r}{R}\right)^{k-1} < \frac{1}{2^4} < \frac{1}{2}.$$

Since $k \geq 2C_H R \geq 2^{11}C_H$, then $2^{10k} \geq \max\{C_H, 2\} = C_H$ and $\left(\frac{r}{R}\right)^k \leq 2^{-10k} \leq \frac{1}{\max\{C_H, 2\}}$. Therefore, by combining all three cases, we see that

$$(2.5) \quad \sup_{W_1} |h|, \sup_{W_1} |\nabla h| \leq \left(\frac{16r}{R}\right)^{2k}.$$

Let W_2 be the connected component of $\partial\Omega$ that intersects the outer circle $\{|z| = R - 1\}$ and note that $W_2 \subset \overline{B_{R-1}} \setminus B_{R-7}$. Now if $z \in W_2$, there are two cases to consider:

- (a) there exists j for which $z \in 4D_j$.
- (b) $|z| = R - 1$ and $z \cap 4D_j$ is empty for all j .

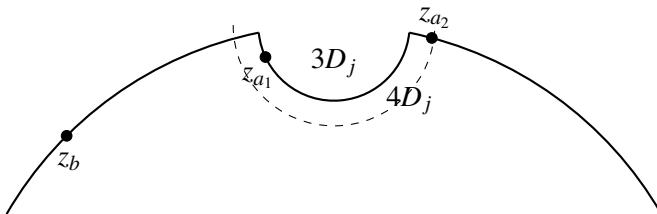


FIGURE 2. A possible image of W_2 with case (a) illustrated by the points z_{a_1} and z_{a_2} , and case (b) illustrated by z_b .

Case (a): Since h doesn't change sign in $B_1(z)$, then an application of [13, Lemma 1.10] shows that

$$|\nabla h(z)| \leq 2|h(z)| \leq 2 \sup_{B_R \setminus \cup 3D_j} |h| \leq 2e^M,$$

where we have applied (2.1).

Case (b): Since $B_1(z) \subset B_R \setminus \cup 3D_j$, then

$$|\nabla h(z)| \leq 2 \sup_{B_1(z)} |h| \leq 2 \sup_{B_R \setminus \cup 3D_j} |h| \leq 2e^M.$$

By combining both cases, we see that

$$(2.6) \quad \sup_{W_2} |h| \leq e^M, \sup_{W_2} |\nabla h| \leq 2e^M.$$

We now use these estimates on h from above to understand the behavior of the analytic function $f = \frac{h_x - ih_y}{z^k}$. Define the set $\Omega_1 \subset \Omega$ as

$$\Omega_1 := \left\{ \frac{r}{2} < |z| < T \right\} \setminus \bigcup 3D_j.$$

Using containment, assumption (2.4), and our normalization (2.3), we see that

$$\sup_{B_{\frac{r}{2}} \setminus \bigcup 3D_j} |h| \leq \sup_{B_r \setminus \bigcup 3D_j} |h| \leq \left(\frac{16r}{R} \right)^{3k} < 1 = \sup_{B_T \setminus \bigcup 3D_j} |h|.$$

Therefore, there exists $z_0 \in \Omega_1$ for which $|h(z_0)| = 1$. By (2.5), we have

$$\sup_{W_1} |h| \leq \left(\frac{16r}{R} \right)^{2k} < \frac{1}{2},$$

so there exists $z_1 \in W_1$ for which $|h(z_1)| = \alpha < \frac{1}{2}$. Let Γ be a path in Ω_1 from z_0 to z_1 for which $\ell(\Gamma) \leq 4R$. If we assume that $|\nabla h(z)| < \frac{1}{8R}$ for all $z \in \Gamma$, then

$$\frac{1}{2} < |h(z_1) - h(z_0)| = \left| \int_{\Gamma} \nabla h(w) \cdot dw \right| \leq \int_{\Gamma} |\nabla h(w)| |dw| < \frac{1}{8R} \ell(\Gamma) < \frac{1}{2},$$

which is impossible so it follows that

$$\sup_{\Omega} |\nabla h| \geq \sup_{\Omega_1} |\nabla h| \geq \frac{1}{8R}.$$

Therefore,

$$(2.7) \quad \sup_{\Omega} |f| \geq \sup_{\Omega_1} |f| \geq \sup_{\Omega_1} |\nabla h| T^{-k} \geq \frac{1}{8R} T^{-k}.$$

An application of (2.5) shows that

$$(2.8) \quad \max_{W_1} |f| \leq \max_{W_1} |\nabla h| \left(\frac{2}{r} \right)^k \leq \left(\frac{16r}{R} \right)^{2k} \left(\frac{2}{r} \right)^k = \left(\frac{2^9 r}{R} \right)^k R^{-k} < 2^{-k} R^{-k},$$

where we have used that $\frac{R}{r} > 2^{10}$. Since $k \geq 2C_H R \geq 20R$, then $2^{-k} < \frac{1}{8R}$. In particular, by combining (2.8) with (2.7), we deduce that

$$\max_{W_1} |f| < \frac{1}{8R} T^{-k} \leq \sup_{\Omega} |f|.$$

An application of (2.6) shows that

$$(2.9) \quad \max_{W_2} |f| \leq \max_{W_2} |\nabla h| (R-7)^{-k} \leq 2e^M \left(\frac{R - \frac{S}{32}}{R-7} \right)^k T^{-k}.$$

Now

$$(2.10) \quad 2e^M \left(\frac{R - \frac{S}{32}}{R-7} \right)^k < \frac{1}{8R} \iff k \log \left(\frac{R-7}{R - \frac{S}{32}} \right) > M + \log(16R).$$

Since $S \geq 2^8$, then $\frac{S}{32} - 7 \geq \frac{S}{2^8}$ and we see that

$$\log \left(\frac{R-7}{R - \frac{S}{32}} \right) = \log \left(1 + \frac{\frac{S}{32} - 7}{R - \frac{S}{32}} \right) \geq \log \left(1 + \frac{S}{2^8 R} \right) \geq \frac{S}{2^8 R} \left(1 - \frac{S}{2^9 R} \right).$$

As $S \leq \frac{R}{2}$, then $1 - \frac{S}{2^9 R} \geq 1 - 2^{-10} > \frac{1}{2}$ and we deduce that

$$\log \left(\frac{R-7}{R - \frac{S}{32}} \right) > \frac{S}{2^9 R}.$$

Since $k \geq 2^{10}MRS^{-1}$, then

$$\frac{k}{2} \log \left(\frac{R-7}{R-\frac{S}{32}} \right) > M > \log(16R),$$

where the second inequality is by assumption. Therefore, (2.10) holds. Combining (2.10) with (2.9) and (2.7) shows that

$$\max_{W_2} |f| < \sup_{\Omega} |f|.$$

Since f is a holomorphic function in Ω , then the maximum principle guarantees that $\sup_{\Omega} |f| = \sup_{\partial\Omega} |f|$. As shown above, the maximum doesn't occur on W_1 or W_2 , so there must exist a disk $3D_j \subset \{\frac{r}{2} < |z| < R-1\}$ for which $\sup_{\Omega} |f| = \sup_{\partial(3D_j)} |f|$.

Considering only the disks D_j for which $3D_j \subset \{\frac{r}{2} < |z| < R-1\}$, define $z_j \in \partial(3D_j)$ to be the point that is closest to the origin, i.e. has the smallest modulus. Then set $m_j = \min_{\partial(3D_j)} |h|$. Define j_0 to be the index for which

$$(2.11) \quad m_{j_0} |z_{j_0}|^{-k} = \max_j m_j |z_j|^{-k}$$

and let j_1 be the index for which

$$\sup_{\Omega} |f| = \sup_{\partial(3D_{j_1})} |f|.$$

For any $z \in \Omega$, an application of Lemma 1 shows that

$$(2.12) \quad |\nabla h(z)| |z|^{-k} \leq \sup_{\Omega} |f| = \sup_{\partial(3D_{j_1})} |\nabla h(z)| |z_{j_1}|^{-k} \leq C_H m_{j_1} |z_{j_1}|^{-k} \leq C_H m_{j_0} |z_{j_0}|^{-k},$$

so we see that

$$|\nabla h(z)| \leq C_H m_{j_0} \left(\frac{|z|}{|z_{j_0}|} \right)^k.$$

Since $\sup_{\Omega} |\nabla h| \geq \frac{1}{8R}$ and this estimate holds for every $z \in \Omega$, then

$$\frac{1}{8R} \leq C_H m_{j_0} \left(\frac{|z|}{|z_{j_0}|} \right)^k \leq C_H m_{j_0} \left(\frac{2R}{r} \right)^k.$$

In particular,

$$(2.13) \quad |h(z_{j_0})| \geq m_{j_0} \geq \frac{1}{8C_H R} \left(\frac{r}{2R} \right)^k.$$

Define $s = \inf \{ \tau \leq 1 : tz_{j_0} \in \Omega \text{ for all } t \in (\tau, 1) \}$ so that the straight line path given by $\gamma(t) = tz_{j_0}$ for $s < t < 1$ is contained in Ω while $sz_{j_0} \in \partial\Omega$. Integrating ∇h along γ shows that

$$h(z_{j_0}) - h(sz_{j_0}) = \int_{\gamma} \nabla h(z) \cdot dz = \int_s^1 \nabla h(tz_{j_0}) \cdot z_{j_0} dt.$$

Applications of (2.12) and (2.13) show that

$$(2.14) \quad \begin{aligned} |h(sz_{j_0})| &\geq |h(z_{j_0})| - \left| \int_s^1 \nabla h(tz_{j_0}) \cdot z_{j_0} dt \right| \geq |h(z_{j_0})| - |z_{j_0}| \left| \int_s^1 C_H m_{j_0} t^k dt \right| \\ &\geq m_{j_0} - \frac{m_{j_0} C_H R}{k+1} = m_{j_0} \left(1 - \frac{C_H R}{k+1} \right) > \frac{m_{j_0}}{2}, \end{aligned}$$

where the last inequality uses that $k+1 > 2C_H R$. Since $k \geq 2C_H R$, then $2^k > 16C_H R$ and it follows that

$$\left(\frac{R}{r} \right)^k \geq 2^{10k} > 16C_H R \cdot 2^{9k}.$$

Combining (2.14) with (2.13) shows that

$$|h(sz_{j_0})| > \frac{m_{j_0}}{2} \geq \frac{1}{16C_H R} \left(\frac{r}{2R}\right)^k > 2^{9k} \left(\frac{r}{R}\right)^k \left(\frac{r}{2R}\right)^k = \left(\frac{16r}{R}\right)^{2k}.$$

By comparing this bound with (2.5), we conclude that $sz_{j_0} \notin W_1$ so it must hold that $sz_{j_0} \in \partial(3D_{j_2})$ for some $3D_{j_2} \subset \{\frac{r}{2} < |z| < R - 1\}$. Then Lemma 1, that $|z_{j_2}| \leq |sz_{j_0}|$, and (2.14) show that

$$(2.15) \quad C_H m_{j_2} |z_{j_2}|^{-k} \geq \sup_{\partial(3D_{j_2})} |h| |z_{j_2}|^{-k} \geq |h(sz_{j_0})| |sz_{j_0}|^{-k} \geq \frac{m_{j_0}}{2} s^{-k} |z_{j_0}|^{-k}.$$

Since $z_{j_0} \in \partial(3D_{j_0})$ and $sz_{j_0} \in \partial(3D_{j_2})$ where $j_0 \neq j_2$, and the balls $\{D_j\}$ are of unit radius and 100-separated, then $|z_{j_0} - sz_{j_0}| \geq 96$. After rearrangement, we see that $s^{-k} \geq (1 - \frac{96}{R})^{-k}$. Since $10 \leq C_H$, $2C_H R \leq k$, and $\frac{96}{R} < -\log(1 - \frac{96}{R})$, then

$$\log(2C_H) < 96 \cdot 2C_H \leq \frac{96}{R} k < -k \log\left(1 - \frac{96}{R}\right) \leq -k \log s,$$

from which it follows that $s^{-k} > 2C_H$. We then conclude from (2.15) that $m_{j_2} |z_{j_2}|^{-k} > m_{j_0} |z_{j_0}|^{-k}$ which contradicts (2.11) and gives the desired contradiction. In other words, (2.4) fails to hold and we see that

$$\sup_{B_r \setminus \bigcup 3D_j} |h| > \left(\frac{16r}{R}\right)^{3k} = \left(\frac{16r}{R}\right)^{3k} \sup_{B_T \setminus \bigcup 3D_j} |h|,$$

which implies (2.2) by our choice of k . \square

3. THE ITERATIVE PROPOSITION

Here we present the proposition that is used repeatedly in the iteration argument. This unique continuation result takes the form of a three-ball inequality for solutions to Schrödinger equations. The techniques used to prove this theorem are very similar to those that appear in the proofs of [18, Theorem 2.2] and [8, Theorem 1.2].

Proposition 2 (Iterative Proposition). *Given $a > 0$ and $R \geq \hat{R}_0(a)$, let $W : B_R \rightarrow \mathbb{R}$ satisfy $\|W\|_{L^\infty(B_R)} \leq a^2$ and let $v : B_R \rightarrow \mathbb{R}$ be a solution to*

$$-\Delta v + Wv = 0 \text{ in } B_R$$

with

$$(3.1) \quad \|v\|_{L^\infty(B_R)} \leq e^{c_1 R}$$

for some $c_1 \geq 1$. Assume that for some $S \in (2^8, \frac{R}{2})$, there exists $z_0 \in \bar{B}_{R-S}$ and $L \geq 0$ such that

$$(3.2) \quad |v(z_0)| \geq e^{-L}.$$

Then there exists $r_0(a, R) > 1$ so that whenever $r \in (0, r_0)$, it holds that

$$(3.3) \quad \|v\|_{L^\infty(B_r)} \geq \left(\frac{r}{R}\right)^\tau,$$

with

$$\tau(R, S, L, a, c_1) = \max \left\{ 3aC_H \sqrt{\frac{C_K}{32 \ln 2}} R \sqrt{\log R}, 2^{14} \left(c_1 R + L + \frac{c_d}{\log R} \right) R S^{-1} \right\} + \frac{1}{5} \left(L + \frac{c_d}{\log R} \right),$$

where C_K and c_d are universal constants, and C_H is the Harnack constant from Lemma 1.

Proof. Let $v : B_R \rightarrow \mathbb{R}$, $W : B_R \rightarrow \mathbb{R}$ and $z_0 \in \bar{B}_{R-S}$ be as given in the statement. Let F_0 denote the nodal set of v , i.e.

$$F_0 = \{z \in \mathbb{R}^2 : v(z) = 0\}.$$

Set $c_s = 2^{10} 10^2$. For $\rho > 0$ to be specified below, there exists a set $F_1 \subset B_R$ that consists of a collection of $c_s \rho$ -separated closed disks of radius ρ that are also $c_s \rho$ -separated from 0 , z_0 , F_0 , and ∂B_R . Moreover, the set $F_0 \cup F_1 \cup \partial B_R$ is a $10c_s \rho$ -net in B_R . A more detailed description of this process is given in [18, §2, Act I].

Define $\Omega = B_R \setminus (F_0 \cup F_1)$ and $\Omega_1 = B_R \setminus F_1$. As shown in [18, §3.1], there exists a universal constant c_P (that depends on c_s) so that Ω has Poincaré constant bounded above by $c_P \rho^2$. In particular, since $c_P \rho^2 \|W\|_{L^\infty(B_R)} \leq c_P \rho^2 a^2$, then by choosing $\rho \leq \rho_0$, a universal constant, we can apply [18, Lemma 3.2]. For $\varepsilon \ll 1$ to be defined later on, let

$$(3.4) \quad \rho = \varepsilon a^{-1}.$$

An application of the arguments in [18, §3.2] then shows that there exists $\phi : \Omega \rightarrow \mathbb{R}$ with the properties that

$$(3.5) \quad \begin{aligned} \Delta \phi - W\phi &= 0 \text{ in } \Omega \\ \phi - 1 &\in W_0^{1,2}(\Omega) \\ \|\phi - 1\|_\infty &\leq c_b (\rho a)^2 = c_b \varepsilon^2, \end{aligned}$$

where c_b is a universal constant that depends on c_P , and we have used (3.4). By extending ϕ to equal 1 across $F_0 \cup F_1$, it is then shown in [18, Lemma 4.1] that $f := \frac{\phi}{\phi - 1} \in W_{\text{loc}}^{1,2}(B_R)$ is a weak solution to the divergence-form equation

$$\operatorname{div}(\phi^2 \nabla f) = 0 \text{ in } \Omega_1.$$

Moreover, the bound in (3.5) implies that for any $z \in B_R$,

$$(3.6) \quad (1 - c_b \varepsilon^2) |f(z)| \leq |v(z)| \leq (1 + c_b \varepsilon^2) |f(z)|.$$

We introduce the Beltrami coefficient μ , defined as follows:

$$\mu = \begin{cases} \frac{1-\phi^2}{1+\phi^2} \frac{f_x+if_y}{f_x-if_y} & \text{in } \Omega_1 \text{ when } \nabla f \neq 0 \\ 0 & \text{otherwise} \end{cases}.$$

Since $|\mu| \lesssim \varepsilon^2$, then as shown in [1], there exists a K -quasiconformal homeomorphism of the complex plane where $K \leq 1 + C_K \varepsilon^2$ and C_K depends on c_b . That is, there exists some $w \in W_{\text{loc}}^{1,2}$ that satisfies the Beltrami equation $\frac{\partial w}{\partial \bar{z}} = \mu \frac{\partial w}{\partial z}$. In fact, an application of the Riemann uniformization theorem shows that there exists a K -quasiconformal homeomorphism $g : B_R \rightarrow B_R$ that is onto with $g(0) = 0$. Moreover, the function $h := f \circ g^{-1}$ is harmonic in $g(\Omega_1)$.

Mori's Theorem implies that

$$(3.7) \quad \frac{1}{16} \left| \frac{z_1 - z_2}{R} \right|^K \leq \frac{|g(z_1) - g(z_2)|}{R} \leq 16 \left| \frac{z_1 - z_2}{R} \right|^{\frac{1}{K}}.$$

Thus, if we set

$$(3.8) \quad \varepsilon = \sqrt{\frac{\ln 2}{2C_K \log R}},$$

then $K \in \left[1, 1 + \frac{\ln 2}{2 \log R}\right]$ and $R \simeq R^K \simeq R^{\frac{1}{K}}$. Since the arguments above hold when $\rho = \frac{\varepsilon}{a} \leq \rho_0$, then we ensure that $R \geq R_1 := \exp\left[\frac{\ln 2}{2C_K(a\rho_0)^2}\right]$. If $|z_1 - z_2| = \rho$, (3.7) shows that

$$|g(z_1) - g(z_2)| \leq 16R \left(\frac{\rho}{R}\right)^{\frac{1}{K}} = 16\rho \left(\frac{aR}{\varepsilon}\right)^{1-\frac{1}{K}}.$$

If $R \geq R_2 := \inf\left\{R > 0 : \frac{R^2}{\log R} \geq \frac{2C_K a^2}{\ln 2}\right\}$, then $\frac{a}{\varepsilon} \leq R$ and

$$|g(z_1) - g(z_2)| \leq 16\rho \exp\left[2 \log R \left(1 - \frac{1}{K}\right)\right] \leq 32\rho.$$

On the other hand, if $|z_1 - z_2| = c_s \rho$, (3.7) implies that

$$|g(z_1) - g(z_2)| \geq \frac{R}{16} \left(\frac{c_s \rho}{R}\right)^K \geq c_s \frac{\rho}{16} \left(\frac{\varepsilon}{aR}\right)^{K-1} \geq c_s \frac{\rho}{16} \exp[-2 \log R (K-1)] \geq c_s \frac{\rho}{32}.$$

Recalling that $c_s = 2^{10}10^2$, we may conclude that

$$|g(z_1) - g(z_2)| \geq 3200\rho.$$

Therefore, h is harmonic in $B_R \setminus \bigcup D_j$, where each D_j is a disk of radius 32ρ . Moreover, the disks are 3200ρ -separated from each other, 0, and $g(z_0)$, while h doesn't change sign in any of the annuli $100D_j \setminus D_j$. Since $z_0 \in \bar{B}_{R-S} \cap \Omega$, then for every $z \in \partial B_R$, $|z_0 - z| \geq S \geq 2^8$ and the distortion estimate described by (3.7) shows that

$$|g(z_0) - g(z)| \geq \frac{R}{16} \left| \frac{z_0 - z}{R} \right|^K \geq \frac{R}{16} \left(\frac{S}{R} \right)^K = \frac{S}{16} \left(\frac{2^8}{R} \right)^{K-1} \geq \frac{S}{16} R^{1-K} \geq \frac{S}{16} \exp\left(-\frac{\ln 2}{2}\right) > \frac{S}{32}.$$

It follows that $g(z_0) \in B_T \setminus \bigcup 3D_j$, where $T := R - \frac{S}{32}$.

Since $g : B_R \rightarrow B_R$, then we may rescale the map to get

$$\tilde{g} := \frac{g}{32\rho} : B_R \rightarrow B_{\frac{R}{32\rho}}$$

which is onto with $\tilde{g}(0) = 0$. Using (3.4) and (3.8), set $C_1 = \frac{a}{32} \sqrt{\frac{2C_K}{\ln 2}}$ and define

$$\tilde{R} = \frac{R}{32\rho} = C_1 R \sqrt{\log R}.$$

From here, we see that $\tilde{h} := f \circ \tilde{g}^{-1}$ is harmonic in $\tilde{g}(\Omega_1)$. In particular, \tilde{h} is harmonic in $B_{\tilde{R}} \setminus \bigcup \tilde{D}_j$, where the \tilde{D}_j are unit disks that are 100-separated from each other, from 0, and from $\tilde{g}(z_0)$. Moreover, \tilde{h} doesn't change signs on any annuli $5\tilde{D}_j \setminus \tilde{D}_j$. With $\tilde{S} := C_1 S \sqrt{\log R}$ and $\tilde{T} := \tilde{R} - \frac{\tilde{S}}{32} = C_1 T \sqrt{\log R}$, $\tilde{g}(z_0) \in B_{\tilde{T}} \setminus \bigcup 3\tilde{D}_j$. Set $\hat{R}_0 = \max\{R_1, R_2, 2^{10}, \exp(C_1^{-2})\}$ so the above arguments hold. Since $R \geq \hat{R}_0$ implies that $C_1 \sqrt{\log R} \geq 1$ and $R \geq 2^{10}$, then $\tilde{R} = C_1 R \sqrt{\log R} \geq 2^{10}$ as well. Because $S \in (2^8, \frac{R}{2})$, then $\tilde{S} \in (2^8, \frac{\tilde{R}}{2})$.

As z_0 satisfies the bound in (3.2), then (3.6) shows that

$$(1 + c_b \varepsilon^2) |f(z_0)| \geq |v(z_0)| \geq \exp(-L).$$

Since $\tilde{g}(z_0) \in B_{\tilde{T}} \setminus \bigcup 3\tilde{D}_j$, then

$$(3.9) \quad \sup_{B_{\tilde{T}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \geq |\tilde{h} \circ \tilde{g}(z_0)| = |f(z_0)| \geq \frac{e^{-L}}{1 + c_b \varepsilon^2}.$$

An application of (3.1) shows that

$$(3.10) \quad \sup_{B_{\tilde{R}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \leq \sup_{B_{\tilde{R}}} |\tilde{h}| = \sup_{\tilde{g}(B_R)} |f \circ \tilde{g}^{-1}| = \sup_{B_R} |f| \leq \frac{1}{1 - c_b \varepsilon^2} \sup_{B_R} |v| \leq \frac{e^{c_1 R}}{1 - c_b \varepsilon^2}.$$

Combining (3.9) and (3.10) then shows that

$$\sup_{B_{\tilde{T}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \geq \frac{1 - c_b \varepsilon^2}{1 + c_b \varepsilon^2} e^{-(c_1 R + L)} \frac{e^{c_1 R}}{1 - c_b \varepsilon^2} \geq e^{-M} \sup_{B_{\tilde{R}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}|,$$

where we introduce $M := c_1 R + L + \frac{c_d}{\log R}$ and c_d depends on c_b and C_K , see (3.8). Because $R \geq R_2$ implies that $16\tilde{R} = 16C_1 R \sqrt{\log R} < R^2$, then because $c_1 \geq 1$, $M > R \geq 2^{10}$, and we have $M > \log(16\tilde{R})$. Therefore, all of the hypotheses of Proposition 1 hold with h , $\{D_j\}$, R , S , and T replaced by \tilde{h} , $\{\tilde{D}_j\}$, \tilde{R} , \tilde{S} , and \tilde{T} , respectively.

For $r \ll R$, an application of (3.7) shows that $g(B_r)$ contains a disk of radius r_1 , where

$$r_1 \geq \frac{R}{16} \left(\frac{r}{R} \right)^K \geq \frac{R}{16} \left(\frac{r}{R} \right)^2.$$

It follows that $\tilde{g}(B_r) \supset B_{\tilde{r}}$, where

$$(3.11) \quad \frac{16\tilde{r}}{\tilde{R}} = \left(\frac{r}{R} \right)^2.$$

If $r \leq \sqrt{\frac{98 \cdot 16R}{C_1 \sqrt{\log R}}}$, then $\tilde{r} \leq 98$, so that $B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j = B_{\tilde{r}}$, and then

$$(3.12) \quad \sup_{B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| = \sup_{B_{\tilde{r}}} |\tilde{h}| \leq \sup_{\tilde{g}(B_r)} |f \circ \tilde{g}^{-1}| = \sup_{B_r} |f|.$$

Since $r \in (0, \frac{R}{2^3})$ implies that $\tilde{r} \in \left(0, \frac{\tilde{R}}{2^{10}}\right)$, then in this case we may apply Proposition 1 with r replaced by \tilde{r} . In particular, the conclusion (2.2) from Proposition 1 shows that

$$(3.13) \quad \sup_{B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \geq \left(\frac{16\tilde{r}}{\tilde{R}}\right)^{\tilde{\kappa}} \sup_{B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}|,$$

where

$$\tilde{\kappa} := \kappa(\tilde{R}, \tilde{S}, M) = \max \left\{ 6C_H C_1 R \sqrt{\log R}, 2^{13} \left(c_1 R + L + \frac{c_d}{\log R} \right) RS^{-1} \right\}.$$

Since $R \geq 2^{10}$ implies $\frac{R}{8} > 1$, while $R \geq R_2$ implies $\sqrt{\frac{98 \cdot 16R}{C_1 \sqrt{\log R}}} > 1$, then $r_0 := \min \left\{ \frac{R}{8}, \sqrt{\frac{98 \cdot 16R}{C_1 \sqrt{\log R}}} \right\} > 1$. Assume that $R \geq \hat{R}_0$ and $r < r_0$. Combining (3.6), (3.12), (3.13), (3.11), and (3.9) shows that

$$\frac{1}{1 - c_b \epsilon^2} \sup_{B_r} |v| \geq \sup_{B_r} |f| \geq \sup_{B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \geq \left(\frac{16\tilde{r}}{\tilde{R}}\right)^{\tilde{\kappa}} \sup_{B_{\tilde{r}} \setminus \bigcup 3\tilde{D}_j} |\tilde{h}| \geq \left(\frac{r}{R}\right)^{2\tilde{\kappa}} \frac{e^{-L}}{1 + c_b \epsilon^2}.$$

Since $C_1 = \frac{a}{32} \sqrt{\frac{2C_K}{\ln 2}}$, $\frac{1 - c_b \epsilon^2}{1 + c_b \epsilon^2} \geq e^{-\frac{c_d}{\log R}}$, and $e < 4 \leq (2^{10})^{\frac{1}{5}} \leq (\frac{R}{r})^{\frac{1}{5}}$, then (3.3) follows, as required. \square

4. TRANSFORMATION MAPS

In this section, we introduce the transformation maps that serve as real-valued versions of the conformal maps $z \mapsto z^\alpha$. Once the maps are defined, we show how they transform solutions to elliptic PDEs and how they transform balls. These results will be used at the beginning and at the end of the proof of Theorem 1.

Given $z = (x, y) \in \mathbb{R}^2$ in Cartesian coordinates, the polar coordinates for z are $(r, \theta) \in \mathbb{R}_+ \times (-\pi, \pi]$, where $r = \sqrt{x^2 + y^2}$ and $\theta = \operatorname{sgn}(y) \arccos(\frac{x}{r})$. Let $\mathbb{R}_+^2 = \{z = (r, \theta) : r > 0, \theta \in (-\frac{\pi}{2}, \frac{\pi}{2})\}$ denote the right half-plane. For $\alpha \in (1, \infty)$, we define $T_\alpha : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ to be the map associated to the conformal transformation on \mathbb{C} given by $z \mapsto z^\alpha$. In polar coordinates, T_α is described as

$$(4.1) \quad T_\alpha(r, \theta) = (r^\alpha, \alpha\theta).$$

To avoid continuity issues at $\theta = \pi$, we restrict the domain and only consider $T_\alpha : \mathbb{R}_+^2 \rightarrow \mathbb{R}^2$.

Lemma 2 (Transformation of PDEs). *If $u : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a solution to (1.4), then with $v : \mathbb{R}_+^2 \rightarrow \mathbb{R}$ and $W : \mathbb{R}_+^2 \rightarrow \mathbb{R}$ defined by $v(z) = u(T_\alpha(z))$ and $W(z) = \alpha^2 |z|^{2\alpha-2} V(T_\alpha(z))$, it holds that*

$$-\Delta v + Wv = 0 \text{ in } \mathbb{R}_+^2.$$

Proof. Let $w := T_\alpha(z)$ be given in polar coordinates by (ρ, φ) . Since u is defined on \mathbb{R}^2 , then in polar coordinates, $u(\rho, \varphi)$ can be defined for all $\rho \in \mathbb{R}^+$ and all $\varphi \in \mathbb{R}$ using periodicity. We then see that $v(z) = v(r, \theta) = u(r^\alpha, \alpha\theta) = u(T_\alpha(z))$ is well-defined on \mathbb{R}_+^2 . Since

$$\begin{aligned} \partial_r v(r, \theta) &= \alpha r^{\alpha-1} \partial_\rho u(\rho, \varphi) \\ \partial_r^2 v(r, \theta) &= \alpha^2 r^{2\alpha-2} \partial_\rho^2 u(\rho, \varphi) + \alpha(\alpha-1) r^{\alpha-2} \partial_\rho u(\rho, \varphi) \\ \partial_\theta v(r, \theta) &= \alpha^2 \partial_\varphi^2 u(\rho, \varphi), \end{aligned}$$

then

$$\begin{aligned}
\Delta v(z) &= \partial_r^2 v(r, \theta) + \frac{1}{r} \partial_r v(r, \theta) + \frac{1}{r^2} \partial_\theta^2 v(r, \theta) \\
&= \alpha^2 r^{2\alpha-2} \partial_\rho^2 u(\rho, \varphi) + \alpha(\alpha-1) r^{\alpha-2} \partial_\rho u(\rho, \varphi) + \alpha r^{\alpha-2} \partial_\rho u(\rho, \varphi) + r^{-2} \alpha^2 \partial_\varphi^2 u(\rho, \varphi) \\
&= \alpha^2 r^{2\alpha-2} \left[\partial_\rho^2 u(\rho, \varphi) + \rho^{-1} \partial_\rho u(\rho, \varphi) + \rho^{-2} \partial_\varphi^2 u(\rho, \varphi) \right] \\
&= \alpha^2 r^{2\alpha-2} \Delta u(w) = \alpha^2 r^{2\alpha-2} V(w) u(w),
\end{aligned}$$

where we have used (1.4). With $W(z) = \alpha^2 |z|^{2\alpha-2} V(T_\alpha(z)) = \alpha^2 r^{2\alpha-2} V(w)$ as given, the conclusion follows. \square

Lemma 3 (Ball containment). *Let $z_0 = (r_0, 0)$ and set $\tilde{r} = \frac{r_0^{1-\alpha}}{2\sqrt{3}\alpha}$. There exists $r_\alpha > 0$ so that whenever with $r_0 \geq r_\alpha$, it holds that $T_\alpha(B_{\tilde{r}}(z_0)) \subset B_1(T_\alpha(z_0))$ and $B_1(T_\alpha(z_0)) \subset T_\alpha(B_1(z_0))$.*

Proof. With \tilde{r} as given, it can be shown that

$$\begin{aligned}
(4.2) \quad B_{\tilde{r}}(z_0) &\subset \left\{ r \in [r_0 - \tilde{r}, r_0 + \tilde{r}], |\theta| \leq \frac{2\tilde{r}}{\sqrt{3}r_0} \right\} \\
&= \left\{ r \in \left[r_0 \left(1 - \frac{r_0^{-\alpha}}{2\sqrt{3}\alpha} \right), r_0 \left(1 + \frac{r_0^{-\alpha}}{2\sqrt{3}\alpha} \right) \right], |\theta| \leq \frac{1}{3\alpha r_0^\alpha} \right\}.
\end{aligned}$$

With $w_0 = T_\alpha(z_0) = (r_0^\alpha, 0) = (\rho_0, 0)$, it follows from (4.2) that

$$\begin{aligned}
T_\alpha(B_{\tilde{r}}(z_0)) &\subset \left\{ r^\alpha \in \left[r_0^\alpha \left(1 - \frac{r_0^{-\alpha}}{2\sqrt{3}\alpha} \right)^\alpha, r_0^\alpha \left(1 + \frac{r_0^{-\alpha}}{2\sqrt{3}\alpha} \right)^\alpha \right], |\alpha\theta| \leq \frac{\alpha}{3\alpha r_0^\alpha} \right\} \\
&= \left\{ \rho \in \left[\rho_0 \left(1 - \frac{\rho_0^{-1}}{2\sqrt{3}\alpha} \right)^\alpha, \rho_0 \left(1 + \frac{\rho_0^{-1}}{2\sqrt{3}\alpha} \right)^\alpha \right], |\varphi| \leq \frac{1}{3\rho_0} \right\}.
\end{aligned}$$

Since

$$\rho_0 \left(1 \pm \frac{\rho_0^{-1}}{2\sqrt{3}\alpha} \right)^\alpha = \rho_0 \pm \frac{1}{2\sqrt{3}} + \frac{(1 - \frac{1}{\alpha})}{24} r_0^{-\alpha} + \dots$$

then there exists $r_1(\alpha) \gg 1$ so that whenever $r_0 \geq r_1$, we have $\rho_0 - \frac{1}{2} \leq \rho_0 \left(1 - \frac{\rho_0^{-1}}{2\sqrt{3}\alpha} \right)^\alpha$ and $\rho_0 \left(1 + \frac{\rho_0^{-1}}{2\sqrt{3}\alpha} \right)^\alpha \leq \rho_0 + \frac{1}{2}$. Because,

$$(4.3) \quad B_1(w_0) \supset \left\{ \rho \in \left[\rho_0 - \frac{1}{2}, \rho_0 + \frac{1}{2} \right], |\varphi| \leq \frac{1}{3\rho_0} \right\},$$

then it follows that $T_\alpha(B_{\tilde{r}}(z_0)) \subset B_1(w_0)$.

On the other hand, with $T_\alpha^{-1}(B_1(w_0)) = \{z : T_\alpha(z) \in B_1(w_0)\}$, (4.2) (with \tilde{r} replaced by 1) implies that

$$\begin{aligned}
T_\alpha^{-1}(B_1(w_0)) &\subset \left\{ r^\alpha \in [\rho_0 - 1, \rho_0 + 1], |\alpha\theta| \leq \frac{2}{\sqrt{3}\rho_0} \right\} \\
&= \left\{ r \in \left[(r_0^\alpha - 1)^{\frac{1}{\alpha}}, (r_0^\alpha + 1)^{\frac{1}{\alpha}} \right], |\theta| \leq \frac{2}{\sqrt{3}\alpha r_0^\alpha} \right\}.
\end{aligned}$$

Since

$$(r_0^\alpha \pm 1)^{\frac{1}{\alpha}} = r_0 \left(1 \pm \frac{1}{r_0^\alpha} \right)^{\frac{1}{\alpha}} = r_0 \pm \frac{1}{\alpha r_0^{\alpha-1}} - \frac{1 - \frac{1}{\alpha}}{2\alpha r_0^{2\alpha-1}} + \dots$$

then there exist $r_2(\alpha) \gg 1$ so that whenever $r_0 \geq r_2$, we have $(r_0^\alpha + 1)^{\frac{1}{\alpha}} \leq r_0 + \frac{1}{2}$ and $(r_0^\alpha - 1)^{\frac{1}{\alpha}} \geq r_0 - \frac{1}{2}$. If $r_0 \geq r_3 := \left(\frac{2\sqrt{3}}{\alpha}\right)^{\frac{1}{\alpha-1}}$, then

$$T_\alpha^{-1}(B_1(w_0)) \subset \left\{ r \in \left[r_0 - \frac{1}{2}, r_0 + \frac{1}{2}\right], |\theta| \leq \frac{1}{3r_0} \right\} \subset B_1(z_0),$$

where we have used (4.3). It follows that $B_1(w_0) \subset T^\alpha(B_1(z_0))$. To complete the proof, choose $r_\alpha = \max\{r_1, r_2, r_3\}$. \square

5. PROOF OF THEOREM 1

In this section, we prove the main result, Theorem 1. Before proceeding to the rigorous details, we describe the main steps:

- Initialize the iterative argument by applying Theorem 3 to the solution function u .
- Compose the solution with the transformation map given in (4.1) to get a solution v to a different Schrödinger equation, see Lemma 2.
- Apply the iterative result, Proposition 2, to v many times to reduce the exponent.
- Undo the change of variables to get the desired estimate for u .

We recall the following result originally proved in [18], see also [8, Theorem 1.1] in the case where $N = 0$ for this formulation. This estimate serves as the initialization step in our iteration argument.

Theorem 3 (Initialization step). *For some $a_0 > 0$, let $V : \mathbb{R}^2 \rightarrow \mathbb{R}$ satisfy $\|V\|_{L^\infty} \leq a_0^2$. Assume that $u : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a solution to (1.4) that satisfies (1.5) and for each $w \in \mathbb{R}^2$,*

$$|u(w)| \leq \exp(c_0 |w|).$$

Then there exists constants $\bar{C}_0 = \bar{C}_0(a_0, c_0) > 0$ and $\bar{R}_0 > 0$ so that whenever $R \geq \bar{R}_0$, it holds that

$$(5.1) \quad \inf_{|w_0|=R} \|u\|_{L^\infty(B_1(w_0))} \geq \exp\left(-\bar{C}_0 R \log^{\frac{3}{2}} R\right).$$

Observe that if u satisfies the hypothesis of Theorem 1, then it also satisfies the hypotheses of Theorem 3. We now have all of the tools we need to prove Theorem 1.

Proof of Theorem 1. With $N \in (0, 2)$ as given, let $\alpha = (1 - \frac{N}{2})^{-1} = \frac{2}{2-N} \in (1, \infty)$. Set $\delta = \min\{\frac{\varepsilon}{2-N}, 1\}$. With r_α from Lemma 3, $\bar{C}_0(a_0, c_0)$ and \bar{R}_0 from Theorem 3, and $\hat{R}_0(a)$ from Proposition 2, define $\bar{r}_1 \geq \max\{r_\alpha, \bar{R}_0^{\frac{1}{\alpha}}, \hat{R}_0(a_0\alpha)\}$ as small as possible so that each of the following conditions hold:

$$(5.2) \quad r^\delta \geq \bar{C}_0 \alpha^{\frac{3}{2}} \log^{\frac{3}{2}} r$$

$$(5.3) \quad r \log r \geq c_d$$

$$(5.4) \quad r^{\frac{\delta}{1+\delta}} \geq \frac{3a_0\alpha C_H}{2^{14}(5c_0+4)} \sqrt{\frac{C_K}{32 \ln 2}} \sqrt{\log r}$$

$$(5.5) \quad r^{\frac{\delta^2}{2}} \geq \left[2^{14}(5c_0+4) + \frac{2}{5}\right] \log r$$

$$(5.6) \quad r^{\frac{5\delta^2}{6}} \geq \left[2^{14}(5c_0+4) + \frac{2}{5}\right] \log \left[\frac{4\alpha}{\sqrt{3}} \left(\frac{3}{2}r\right)^\alpha\right].$$

Here C_K and c_d are universal constants from Proposition 2 and C_H is the universal Harnack constant from Lemma 1.

For each $k \in \mathbb{N}$, define

$$\alpha_k = \frac{2 + (k-1)N}{2 + (k-2)N} > 1.$$

Observe that $\alpha_1 = \alpha$, $\alpha_{k+1} = 2 - \frac{1}{\alpha_k} < \alpha_k$, and $\lim_{k \rightarrow \infty} \alpha_k = 1$. Choose $\ell \in \mathbb{N}$ so that $\frac{N}{2 + (\ell-2)N} \leq \delta < \frac{N}{2 + (\ell-3)N}$

or $\alpha_\ell \leq 1 + \delta < \alpha_{\ell-1}$. For each $k \in \{1, 2, \dots, \ell-1\}$, recursively define $\bar{r}_{k+1} = \bar{r}_k^{\alpha_k + \delta} + \frac{1}{2}\bar{r}_k - 1$, where

\bar{r}_1 is from above. Take $w_0 \in \mathbb{R}^2$ with $|w_0| \geq R_0 := \bar{r}_\ell^\alpha$. After a rotation, we may assume that w_0 belongs to the positive x -axis. Choose $r_1 \geq \bar{r}_1$ so that with $r_{k+1} = r_k^{\alpha_{k-1}+\delta} + \frac{1}{2}r_k - 1$ for each $k \in \{1, \dots, \ell-1\}$, it holds that $|w_0| = r_\ell^\alpha$. Set $z_k = (r_k, 0) \in \mathbb{R}_+^2$ for each $k \in \{1, \dots, \ell-1\}$ so that

$$(5.7) \quad |z_k| = |z_{k-1}|^{\alpha_{k-1}+\delta} + \frac{1}{2}|z_{k-1}| - 1.$$

With $\alpha = \alpha_1 > 1$, let $T_\alpha : \mathbb{R}_+^2 \rightarrow \mathbb{R}^2$ be as given in (4.1), see Section 4. Define $v : \mathbb{R}_+^2 \rightarrow \mathbb{R}$ by $v(z) = u(T_\alpha(z))$.

Claim. For every $k \in \{1, \dots, \ell-1\}$,

$$(5.8) \quad \|v\|_{L^\infty(B_1(z_k))} \geq \exp(-|z_k|^{\alpha_{k-1}+\delta}).$$

To prove the claim, we proceed by induction. Define $w_1 = T_\alpha(z_1) = (r_1^\alpha, 0)$. Since $r_1 \geq r_\alpha$, Lemma 3 implies that $B_1(w_1) \subset T_\alpha(B_1(z_1))$ and then

$$\|v\|_{L^\infty(B_1(z_1))} = \|u \circ T_\alpha\|_{L^\infty(B_1(z_1))} = \|u\|_{L^\infty(T_\alpha(B_1(z_1)))} \geq \|u\|_{L^\infty(B_1(w_1))}.$$

As $|w_1| = r_1^\alpha \geq \bar{R}_0$, an application of (5.1) from Theorem 3 shows that

$$\|u\|_{L^\infty(B_1(w_1))} \geq \exp(-\bar{C}_0|w_1|\log^{\frac{3}{2}}|w_1|).$$

Because $|w_1| = |z_1|^\alpha$, then $\bar{C}_0|w_1|\log^{\frac{3}{2}}|w_1| = \bar{C}_0|z_1|^\alpha\log^{\frac{3}{2}}|z_1|^\alpha = \bar{C}_0\alpha^{\frac{3}{2}}|z_1|^\alpha\log^{\frac{3}{2}}|z_1| \leq |z_1|^{\alpha+\delta}$, where we have used (5.2) to reach the last bound. Combining these observations shows that

$$\|v\|_{L^\infty(B_1(z_1))} \geq \|u\|_{L^\infty(B_1(w_1))} \geq \exp(-\bar{C}_0|w_1|\log^{\frac{3}{2}}|w_1|) \geq \exp(-|z_1|^{\alpha+\delta}).$$

In other words, (5.8) holds for $k = 1$.

Now assume that (5.8) holds for $k-1 \in \mathbb{N}$. Define $S_k = \frac{1}{2}|z_{k-1}|$ and $R_k = |z_{k-1}|^{\alpha_{k-1}+\delta}$. Comparing with (5.7), we see that $|z_k| = R_k + S_k - 1$. Since $B_1(z_{k-1}) \subset B_{R_k-S_k}(z_k)$, then the inductive hypothesis shows that there exists $z_0 \in \bar{B}_{R_k-S_k}(z_k)$ so that

$$|v(z_0)| \geq \exp(-|z_{k-1}|^{\alpha_{k-1}+\delta}) = \exp(-R_k).$$

The bound in (1.6) shows that $|v(z)| = |u(T_\alpha(z))| \leq \exp(c_0|z|^{\alpha(1-\frac{N}{2})}) = \exp(c_0|z|)$ from which it follows that

$$\|v\|_{L^\infty(B_{R_k}(z_k))} \leq \exp(c_0(|z_k|+R_k)) \leq \exp\left(\frac{5c_0}{2}R_k\right).$$

Let $W : \mathbb{R}_+^2 \rightarrow \mathbb{R}$ be given by $W(z) = \alpha^2|z|^{2\alpha-2}V(T_\alpha(z)) = \alpha^2|w|^{2-\frac{2}{\alpha}}V(w)$ so that by (1.3),

$$|W(z)| = \alpha^2|w|^{2-\frac{2}{\alpha}}|V(w)| \leq a_0^2\alpha^2|w|^{2-N-\frac{2}{\alpha}} = (a_0\alpha)^2.$$

An application of Lemma 2 with assumption (1.4) shows that

$$-\Delta v + Wv = 0 \text{ in } \mathbb{R}_+^2.$$

As $R_k \geq \bar{r}_1$ implies that $R_k \geq \hat{R}_0(a_0\alpha)$ and $S_k \in \left(2^8, \frac{R_k}{2}\right)$, Proposition 2 is applicable with $a = a_0\alpha$, $R = R_k$, $S = S_k$, $L = R_k$, and $c_1 = \frac{5c_0}{2}$. With $r = 1$, Proposition 2 shows that

$$(5.9) \quad \|v\|_{L^\infty(B_1(z_k))} \geq \exp(-\tau_k(R_k)\log R_k),$$

where

$$(5.10) \quad \begin{aligned} \tau_k(R_k) &= \max \left\{ 3a_0\alpha C_H \sqrt{\frac{C_K}{32\ln 2}} R_k \sqrt{\log R_k}, 2^{15} \left(\frac{5c_0}{2} + 1 + \frac{c_d}{R_k \log R_k} \right) R_k^{2-\frac{1}{\alpha_{k-1}+\delta}} \right\} \\ &\quad + \frac{1}{5} \left(R_k + \frac{c_d}{\log R_k} \right). \end{aligned}$$

Because $R_k \geq \bar{r}_1$ and $\alpha_{k-1} > 1$, conditions (5.3) and (5.4) imply that

$$(5.11) \quad \tau_k(R_k) \leq \left[2^{14}(5c_0 + 4) + \frac{2}{5} \right] R_k^{2 - \frac{1}{\alpha_{k-1} + \delta}}$$

while condition (5.5) implies that

$$\left[2^{14}(5c_0 + 4) + \frac{2}{5} \right] \log R_k \leq R_k^{\frac{\delta^2}{2}}.$$

Since $\alpha > 1$ and $\delta \leq 1$ implies that $2 - \frac{1}{\alpha+\delta} + \frac{\delta^2}{2} < 2 - \frac{1}{\alpha} + \delta$, then combining these bounds shows that

$$\tau_k(R_k) \log R_k \leq R_k^{2 - \frac{1}{\alpha_{k-1}} + \delta} = R_k^{\alpha_k + \delta}.$$

Returning to (5.9), we conclude that

$$\|v\|_{L^\infty(B_1(z_k))} \geq \exp(-R_k^{\alpha_k + \delta}) \geq \exp(-|z_k|^{\alpha_k + \delta}),$$

establishing the claim given by (5.8).

Next we consider the behavior of v on a ball centered at z_ℓ . The arguments in the previous paragraph show that Proposition 2 is applicable with $a = a_0\alpha$, $R = R_\ell$, $S = S_\ell$, $L = R_\ell$, and $c_1 = \frac{5c_0}{2}$. In particular, with $r = \tilde{r} := \frac{|z_\ell|^{1-\alpha}}{2\sqrt{3}\alpha}$, Proposition 2 shows that

$$\|v\|_{L^\infty(B_{\tilde{r}}(z_\ell))} \geq \left(\frac{\tilde{r}}{R_\ell} \right)^{\tau_\ell(R_\ell)},$$

where $\tau_\ell(R_\ell)$ is defined by (5.10) and the bound in (5.11) holds. Because $R_\ell < |z_\ell| < \frac{3}{2}R_\ell$, then $\frac{R_\ell}{\tilde{r}} < \frac{4\alpha}{\sqrt{3}} \left(\frac{3}{2}R_\ell \right)^\alpha$ and condition (5.6) implies that

$$\left[2^{14}(5c_0 + 4) + \frac{2}{5} \right] \log \left(\frac{R_\ell}{\tilde{r}} \right) \leq R_\ell^{\frac{5\delta^2}{6}}.$$

Since $\alpha_{\ell-1} > 1 + \delta$ and $\delta \leq 1$ implies that $2 - \frac{1}{\alpha_{\ell-1} + \delta} + \frac{5\delta^2}{6} < \alpha_\ell + \delta \leq 1 + 2\delta$, then $\tau_\ell(R_\ell) \log \left(\frac{R_\ell}{\tilde{r}} \right) \leq R_\ell^{1+2\delta}$ and we conclude that

$$\|v\|_{L^\infty(B_{\tilde{r}}(z_\ell))} \geq \exp(-R_\ell^{1+2\delta}) \geq \exp(-|z_\ell|^{1+2\delta}).$$

As $|z_\ell| \geq r_\alpha$, an application of Lemma 3 shows that $T_\alpha(B_{\tilde{r}}(z_\ell)) \subset B_1(T_\alpha(z_\ell))$ and then

$$\|u\|_{L^\infty(B_1(T_\alpha(z_\ell)))} \geq \|u\|_{L^\infty(T_\alpha(B_{\tilde{r}}(z_\ell)))} = \|v\|_{L^\infty(B_{\tilde{r}}(z_\ell))} \geq \exp(-|z_\ell|^{1+2\delta}) = \exp(-r_\ell^{1+2\delta}).$$

Because we assumed that w_0 belongs to the positive x -axis and has $|w_0| = r_\ell^\alpha$, then $T_\alpha(z_\ell) = (r_\ell^\alpha, 0) = w_0$. Since $\delta = \frac{\epsilon}{2-N}$, then $\frac{1+2\delta}{\alpha} = (1 - \frac{N}{2})(1 + 2\delta) = 1 - \frac{N}{2} + \epsilon$ and $r_\ell^{1+2\delta} = |w_0|^{1 - \frac{N}{2} + \epsilon}$. Combining these observations shows that $\|u\|_{L^\infty(B_1(w_0))} \geq \exp(-|w_0|^{1 - \frac{N}{2} + \epsilon})$. Since $w_0 \in \mathbb{R}^2$ with $|w_0| \geq R_0$ was arbitrary, then (1.7) follows, as required. \square

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