## Linear algebra

## **Linear vector spaces**

## Number field (K)

- The elements are scalar
- Addition and Multiplication are consistent with arithmetic law (交換律、結合律、分配律)
- Closed under addition and multiplication
- Identity element with addition and multiplication (addition: 0, multiplication: 1)
- Has inverse elements (addition: negative numbers, multiplication: reciprocal numbers)

## Vector space (V)

- The elements are vectors
- Closed under addition and multiplication
  - $\forall x, y \in V, z = x + y \text{ where } z \in V$
  - $\forall a \in K \text{ and } x \in V, z = ax \text{ where } z \in V$
- Identity element with addition and multiplication (addition: 0)

V is a vector space on K if V satisfy:

- A1 加法結合律: $\forall u,v,w\in V, (u+v)+w=u+(v+w)$
- A2 交換律: $\forall u, v \in V, u+v=v+u$
- A3 零向量(加法單位元素): $\forall u \in V, \exists 0 \in u + 0 = u$

• A4 逆向量: $\forall u \in V, \exists -u \in V, u + (-u) = 0.$ 

• A5 分配律 $1: orall a \in K \ \& \ u,v \in V, a(u+v) = au + av$ 

• A6 分配律2: $orall a,b\in K\ \&\ u\in V, (a+b)u=au+bu$ 

• A7 乘法結合律 $orall a,b\in K\ \&\ u\in V, (ab)u=a(bu)$ 

• A8 單位純量: $orall u \in V, 1u = u ext{ where } 1 \in K$ 

For example: n次多項式的集合

$$a_0 + a_1 t + a_2 t^2 + \ldots + a_n t^n, \, a_i \in K$$

## **Linear combinations and Linear spans**

vector x is the linear combination of  $v_1,\ldots,v_m\in V$ . Hence, all possible x could form a subset S (also vector space), which is the linear spans of V.

## Linear independence

For

$$a_1v_1 + a_2v_2 + \ldots + a_nv_n = 0, a_n \in K$$

If their only solution is  $a_i=0,\ \forall i$ , then  $v_i,\dots,v_n$  is a set of linearly independent vectors

## Dimension, basis and components/coordinates

n-dimentional vector space V  $(\dim V=n)$  can be generated from a set of linearly independent vectors  $e_1,\ldots,e_n\Longrightarrow \{e_1,\ldots,e_n\}$  is the basis of V

That is,  $\forall v \in V, \ v = x_1e_1 + \cdots + x_ne_n \implies \{x_1, \ldots, x_n\}$  is called the components or coordinates of  $\boldsymbol{v}$  relative to the basis

For example: Fourier series

$$f(x) = \sum_{0}^{\infty} f_n e^{inx}$$

its basis is  $e^{inx}$  , component is  $f_n$  and  $\dim=\infty$ 

## **Linear maps**

V and U are vector spaces,  $x \in V$  is the domain, and  $y \in U$  is the codomain

Definition of mapping:  $x \rightarrow y$ , or Ax = y

Linear mapping properties:

$$ullet A(x_1+x_2) = Ax_1 + Ax_2, ext{where } x_1, x_2 \in V$$

• 
$$A(ax) = aAx$$
, where  $x \in V, a \in K$ 

### **Matrices**

For an  $n \times n$  matrix

$$A=(a_{ij})=egin{pmatrix} a_{11} & a_{12} & \cdots & a_{1n} \ a_{21} & a_{22} & \cdots & a_{2n} \ dots & dots & \ddots & dots \ a_{n1} & a_{n2} & \cdots & a_{nn} \end{pmatrix}$$

$$Det(A) = |A| = \sum_{i,j,k...} \epsilon_{i,j,k...} a_{1i} \ a_{2j} \ a_{3k} \dots$$

•  $\epsilon$  is Levi-Civita tensor (recall 胡德邦)

 If any two columns or rows of A are swapped, then Det(A) changes sign

• 
$$\operatorname{Det}(A) = \operatorname{Det}(A^T)$$

• 
$$\operatorname{Det}(AB) = \operatorname{Det}(A)\operatorname{Det}(B)$$

Equality:  $A=B \Leftrightarrow a_{ij}=b_{ij}$ 

Addition: 
$$A+B=C \Leftrightarrow a_{ij}+b_{ij}=c_{ij}$$

Multiplication: trivival, we skip

#### **Direct product:**

$$C_{(mn imes mn)}=A_{(m imes m)}\otimes B_{(n imes n)}=egin{pmatrix} a_{11}B & a_{12}B & \cdots & a_{1n}B\ a_{21}B & a_{22}B & \cdots & a_{2n}B\ dots & dots & dots & dots\ a_nB & a_{n2}B & \cdots & a_{nn}B \end{pmatrix}$$

column vector: 
$$|x
angle = egin{pmatrix} x_1 \ x_2 \ dots \ x_n \end{pmatrix}_{n imes 1}$$
 row vector  $\langle x| = (x_1 \quad x_2 \quad \cdots)_{1 imes n}$ 

#### unit matrix:

$$I = egin{pmatrix} 1 & 0 & \cdots & 0 \ 0 & 1 & \cdots & 0 \ dots & dots & \ddots & dots \ 0 & 0 & \cdots & 1 \end{pmatrix}$$

#### **Null matrix:**

$$O = egin{pmatrix} 0 & 0 & \cdots & 0 \ 0 & 0 & \cdots & 0 \ \vdots & \vdots & \ddots & \vdots \ 0 & 0 & \cdots & 0 \end{pmatrix}$$

#### **Diagonal matrices:**

$$I=egin{pmatrix} a_{11} & 0 & \cdots & 0 \ 0 & a_{22} & \cdots & 0 \ dots & dots & \ddots & dots \ 0 & 0 & \cdots & a_{nn} \end{pmatrix}$$

if A, B are diagonal  $\Rightarrow AB = BA$ 

Trace: 矩陣對角元素的和  $Trace(A) = \sum_i a_{ii}$ 

- Tr(AB) = Tr(BA)
- Cyclic properties:

$$Tr(ABCD) = Tr(DABC) = Tr(CDAB) = Tr(BCDA)$$

- ullet If AB=BA,  $A^2=B^2=I$  then Tr(A)=Tr(B)=0
- $Tr(A \otimes B) = Tr(A)Tr(B)$

**Minor of**  $a_{ij}$ :  $|M_{ij}| =$  delete  $i^{th}$  row and  $j^{th}$  column, compute determinant

Cofactor of  $a_{ij}: c_{ij}=(-1)^{i+j}|M_{ij}|$ 

### Classical adjoint:

$$\mathrm{adj} \ A = egin{pmatrix} c_{11} & c_{21} & \cdots & c_{n1} \ c_{12} & c_{22} & \cdots & c_{n2} \ dots & dots & \ddots & dots \ c_{1n} & c_{2n} & \cdots & c_{nn} \end{pmatrix}$$

Inversion:  $A^{-1} \Rightarrow AA^{-1} = A^{-1}A = I$ 

Transpose:  $A=a_{ij}\Rightarrow A^T=\widetilde{A}=a_{ji}$ 

A is symmetric:  $A = A^T$ 

A is anti-symmetric:  $A=-A^T$ 

# **Orthogonal matrices**

定義一座標轉換

$$egin{cases} \hat{e}_1' = \ket{e_1'} = a_{11}\hat{e}_1 + a_{12}\hat{e}_2 + a_{13}\hat{e}_3 \ \hat{e}_2' = \ket{e_2'} = a_{21}\hat{e}_1 + a_{22}\hat{e}_2 + a_{23}\hat{e}_3 \ \hat{e}_3' = \ket{e_3'} = a_{31}\hat{e}_1 + a_{32}\hat{e}_2 + a_{33}\hat{e}_3 \end{cases}$$

可以進一步定義一個 M

$$M = egin{pmatrix} a_{11} & a_{12} & a_{13} \ a_{21} & a_{22} & a_{23} \ a_{31} & a_{32} & a_{33} \end{pmatrix}$$

假設一向量 v ,在原座標中的表示為

$$v = v_1 e_1 + v_2 e_2 + v_3 e_3$$

則其在新座標系下的表示即為

$$v_1'(a_{11}\hat{e}_1+a_{12}\hat{e}_2+a_{13}\hat{e}_3)+\ v=v_1'e_1'+v_2'e_2'+v_3'e_3'=v_2'(a_{21}\hat{e}_1+a_{22}\hat{e}_2+a_{23}\hat{e}_3)+\ v_3'(a_{31}\hat{e}_1+a_{32}\hat{e}_2+a_{33}\hat{e}_3)$$

注意到

$$v_1'(a_{11}\hat{e}_1) + v_2'(a_{21}\hat{e}_1) + v_3'(a_{31}\hat{e}_1) = v_1\hat{e}_1$$

寫成以下關係式

$$v = egin{pmatrix} v_1 \ v_2 \ v_3 \end{pmatrix} = egin{pmatrix} a_{11}v_1' + a_{12}v_2' + a_{13}v_3' \ a_{21}v_1' + a_{22}v_2' + a_{23}v_3' \ a_{31}v_1' + a_{32}v_2' + a_{33}v_3' \end{pmatrix} = M^Tv' \Leftrightarrow v = M^Tv'$$

同時我們也知道 v'=Mv ,與上式作比較後得 orthogonal matrix 的條件為  $M^T=M^{-1}$  ,在所有類型矩陣之中,只有旋轉矩陣與鏡射矩陣為 orthogonal matrices 。

• 旋轉矩陣:Det(A) = +1

• 鏡射矩陣:Det(A) = -1

# 旋轉矩陣 (Rotations)

向量旋轉太簡單,跳過

坐標系旋轉:在 3D space 當中,只能分別依 x, y, z 軸旋轉,因其 number of degrees of freedom (dof) =3,所以只需要 3 個 parameter 來 完成座標系的旋轉

- 1. Rotation about z by lpha
- 2. Rotation about y by  $\beta$
- 3. Rotation about z by  $\gamma$

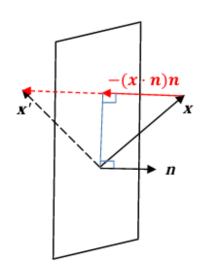
$$R_z(lpha) = egin{pmatrix} \coslpha & \sinlpha & 0 \ -\sinlpha & \coslpha & 0 \ 0 & 0 & 1 \end{pmatrix} R_y(eta) = egin{pmatrix} \coseta & 0 & -\sineta \ 0 & 1 & 0 \ \sineta & 0 & \coseta \end{pmatrix} R_z(\gamma) =$$

且必定能寫成以下形式

$$R = R_z(\gamma)R_y(\beta)R_z(\alpha)$$

這就是標準的  $O_3^+$  (SO(3)) group,稱為正交純旋轉群

# 鏡射矩陣 (Reflections)



注意到  $x' = x - 2n(x \cdot n)$  ,我們可以將這操作化為矩陣形式

$$x' = Ox, \ O = I - 2nn^T$$

再次注意到  $O^T=O$  ,且 OO=I 。由兩條件推得  $O^{-1}=O^T$  為 orthogonal matrix

$$|O|=-1$$

## **Similarity Transformation**

考慮一矩陣  $A\Rightarrow v_1=Av$ , A 為旋轉或鏡射的任一種,另考慮一正 交座標旋轉矩陣  $M\Rightarrow v'=Mv$ ,則有以下關係:

$$v_1' = Mv_1 = MAv = (MAM^{-1})(Mv) = (MAM^{-1})v' = A'v'$$

就酷酷的

## Linear independence

#### Wronskian 判別法:

If  $\{u_i(x)\}\equiv\{u_1(x),u_2(x),\ldots,u_n(x)\}$  are  $n^{th}-{
m order}$  differentiable function set, consider

$$W(\{u_i\}) = egin{bmatrix} u_1 & u_2 & \dots & u_n \ u_1' & u_2' & \dots & u_n' \ dots & & & & \ dots & & & & \ u_1^{(n-1)} & u_2^{(n-1)} & \dots & u_n^{(n-1)} \ \end{pmatrix}$$

If it's doesn't equal to 0, then  $\{u_i(x)\}$  is linearly independent. Notice that if it's equal to 0, we can not determine whether it's linear dependent or independent.

## **Hermitian matrices & Unitary matrices**

### 定義一些新性質:

Complex conjugates:  $A^* = (a_{ij})$ 

Hermitian conjugate matrix:  $A^+ \Rightarrow a^+_{ij} = a^*_{ji} \;,\;\; |x
angle^+ = \langle x|$ 

•  $(y^+ABCx)^+ = x^+C^+B^+A^+y$ 

Hermiatian matrix:  $A^+=A$ 

Anti-Hermitian matrix:  $A^+=-A$ 

Unitary matirx:  $U^+ = U^{-1}$ 

M is normal  $\Leftrightarrow$  M commute with  $M^+$ :  $[M,M^+]=0$ 

- [A,B] = AB BA
- $[A,B]^+ = [B^+,A^+]$
- Tr([A, B]) = Tr(AB) Tr(BA) = 0
- Hermitian, anti-Hermitian, and unitary matrices are all normal

If [A,B]=C , A and B are Hermitian, then C is anti-Hermitian, that is,  $[A,B]^+=-[A,B]$ 

## Inner product (scalar product)

General Definition of inner product

First we define basis  $e_i$  in order to evaluate  $\langle x|y\rangle$ . Thus, we have  $|x\rangle\equiv x_ie_i,\ |y\rangle\equiv y_ie_i.$ 

Let  $\langle x|y\rangle=G_{ij}$ , where G is a matrix tensor and  $G_{ij}$  is metric coefficients. Its value will depend on i and j, then

$$\langle x|y
angle = \langle x_ie_i|y_je_j
angle = x_i^*\langle e_i|e_j
angle y_j = x_i^*G_{ij}y_j = \mathbf{x}^\dagger\mathbf{G}\mathbf{y}$$

We formally define the inner product. Notice that when  $e_i$  is orthogonal and normalized, then

$$\langle e_i | e_j 
angle = G_{ij} \, \Longrightarrow \, \langle x | y 
angle = x_i^* y_i$$

• *G* is Hermitian

In addition, for a non-orthogonal basis, we could turn it into orthogonal one for easy calculation

Consider  $e_i$  is orthogonal basis,  $e_i' \equiv M e_i$  , where M is arbitrary matrix, so that

$$|x
angle = (M^\dagger)^{-1} |x
angle$$

then

$$egin{aligned} \langle x'|y'
angle &= x'^\dagger G y' = x^\dagger (M^{-1}) G (M^{-1})^\dagger y \equiv \langle x|y
angle = x^\dagger y \ \implies M^{-1} G (M^\dagger)^{-1} = I \implies G = M M^\dagger \end{aligned}$$

Recall orthogonal matrices, but we only consider real number there

#### Some useful skill

- $\langle x|y
  angle = \langle x,y
  angle = (x,y)$  is **bilinear** (Be linear on each variable)
- $ullet \langle x|y
  angle = \langle y|x
  angle^*$
- $\langle x|x
  angle \geq 0$
- ullet  $\langle x,y
  angle^{1/2}=|x|=||x||$  is the length or norm of the |x
  angle
- $\langle x|y\rangle = |x||y|\cos\theta$
- ullet If  $\langle x|y
  angle=0$  (in any vector space) the vectors are said to be orthogonal

The standard (Euclidean) inner product on  $\mathbb{R}^n$  is called **dot product** 

$$\langle x,y
angle = \sum x_i y_i$$

which is generalized to  $\mathbb{C}^n$  as

$$\langle x,y
angle = \sum x_i^*y_i$$

Cauchy-Schwarz inequality

$$|\langle x|x
angle\langle y|y
angle\geq |\langle x|y
angle|^2$$

### **Adjoint operator**

The adjoint of a linear operator A satisfy with

$$egin{aligned} \langle A^\dagger x | y 
angle \stackrel{ ext{def}}{=} \langle x | A y 
angle & orall x, y \ & \Leftrightarrow (A_{ji})^{\dagger *} x_j^* y_j \stackrel{ ext{def}}{=} x_i^* A_{ij} y_j \ & \Leftrightarrow (A_{ji})^{\dagger *} = A_{ij} \Leftrightarrow (A_{ji})^\dagger = A_{ji}^* \end{aligned}$$

That is, A is Hermitian

## **Unitary Transformation**

在 Similarity transformation 中,如果進行座標轉換的 matrix 是 unitary ,則這個轉換叫做 unitary transformation,在內積空間中保持向量長度和角度不變,其有以下關係:

$$A' = UAU^{-1} = UAU^{\dagger}$$

若  $U_1$  和  $U_2$  皆為 unitary  $\Rightarrow$   $U_3 = U_1 U_2$  也是 unitary

### **Pauli matrices and Dirac matrices**

Pauli matrices

To describe a particle of spin 1/2 (in a non-relativistic theory):

$$\sigma_x = egin{pmatrix} 0 & 1 \ 1 & 0 \end{pmatrix}, \quad \sigma_y = egin{pmatrix} 0 & -i \ i & 0 \end{pmatrix}, \quad \sigma_z = egin{pmatrix} 1 & 0 \ 0 & -1 \end{pmatrix}$$

- $(\sigma_i)^2 = I$
- ullet Anti-commutation:  $(\sigma_i\sigma_j+\sigma_j\sigma_i)=2\delta_{ij}I$
- Cyclic permutation of indices:  $\sigma_i \sigma_j = i \sigma_k$

**Dirac matrices** 

首先就是 Dirac 覺得 Pauli 的矩陣不夠用,在他的 relativistic electron theory 中,需要 4 個 anti-commuting matrices,因此他來了

出發點,任一2×2的矩陣必可被表示為

$$M = C_0 I + C_1 \sigma_1 + C_2 \sigma_2 + C_3 \sigma_3$$

故 I,  $\sigma_1$ ,  $\sigma_2$ ,  $\sigma_3$  形成一組 basis,但其中只有  $\sigma_i$  是 anticommutative,所以我們要以這四個為基礎來建構一組  $4\times 4$  的矩 陣,使用以下算式

$$\left\{egin{aligned} \sigma_{i, ext{Dirac}} &= I \otimes \sigma_{i, ext{Pauli}} \ 
ho_{i, ext{Dirac}} &= \sigma_{i, ext{Pauli}} \otimes I \end{aligned}
ight.$$

這樣我們可以造出6個 4×4 的矩陣,滿足:

- ullet Anti-commutation:  $\sigma_i\sigma_j+\sigma_j\sigma_i=2\delta_{ij}I,\quad 
  ho_i
  ho_j+
  ho_j
  ho_i=2\delta_{ij}I$
- commutation:  $[\sigma_i, \rho_j] = 0$
- ullet Cyclic permutation of indices:  $\sigma_i\sigma_j=i\sigma_k,\quad 
  ho_i
  ho_j=i
  ho_k$

### 接著建構一個乘法表,撰我們要的東東

$E_{ij}$	j=	$I_0$	$\sigma_1$	$\sigma_2$	$\sigma_3$
i=					
$I_0$		I	$\sigma_1$	$\sigma_2$	$\sigma_3$
$ ho_1$		$ ho_1, -Y_5$	$lpha_1$	$lpha_2$	$lpha_3$
$ ho_2$		$\rho_2,\alpha_5$	$Y_1$	$Y_2$	$Y_3$
$ ho_3$		$\rho_3,\alpha_4,Y_4,\beta$	$\delta_1$	$\delta_2$	$\delta_3$

狄拉克選了  $\alpha_i$  (i=1,2,3,4) ,作為他 relativistic electron theory 中的一組anti-commutative matrices.

最後我們來討論一下  $E_{ij}$  的性質

- $Det(E_{ij}) = 1$
- $E_{ij}^2 = I$
- ullet  $E_{ij}$  is Hemitian and Unitary
- $Tr(E_{ij}) = 0$  除了  $Tr(E_{00}) = 4$
- 任二個  $E_{ij}$  的乘積必等於另一個  $E_{ij}$  乘上  $\{-1,\pm i\}$  三者其中之一
- 彼此線性獨立
- 任一  $4 \times 4$  的矩陣必可寫成這 16 個  $E_{ij}$  的線性和: $M = \sum_{i,j=0}^3 c_{ij} E_{ij}$

## Eigenvalues, Eigenvectors & Diagonalisation

### Introduction

牽涉到矩陣運算的物理問題中,常藉由將矩陣對角化來簡化計算或 釐清其中的物理觀念,像是剛體轉動問題,在數學上這個步驟是將 一個給定的 matrix 做 orthogonal similarity transformation (if real) 或 unitary transformation 來轉換成一個 diagonal matrix。

範例: Inertial matrix

## **Eigenvalues and eigenvectors**

If  $A|v
angle=\lambda|v
angle$  where A is n imes n,  $\lambda\in C^1$ , |v
angle
eq 0, then

 $\lambda$ : eigenvalue

|v
angle : eigenvector

That is, 
$$(A-\lambda I)|v
angle=0$$

且注意到  $A-\lambda I$  中的每一列都和  $|v\rangle$  正交,因此這 n 個向量不可能全部彼此正交,以致該矩陣的 Det 必為零:

$$\operatorname{Det}(A-\lambda I) = egin{array}{c|cccc} a_{11} - \lambda & a_{12} & \cdots & a_{1n} \ a_{21} & a_{22} - \lambda & \cdots & a_{2n} \ dots & dots & \ddots & dots \ a_{n1} & a_{n2} & \cdots & a_{nn} - \lambda \end{array} = P_A(\lambda) = \prod_{i=1}^n (\lambda_i - \lambda) =$$

- $P_A(\lambda)$  為  $\lambda$  的 n 次多項式,他的所有解即為所有 eigenvalue,其集合  $\{\lambda_i\}$  稱為 spectrum of A
- 對於對稱矩陣,由於其解得的 eigenvectors 必互相正交,因此我們一般選取  $\langle v_i|v_i \rangle=1$ ,這樣  $\{|v_i \rangle\}$  形成 orthonormal basis 也叫做 eigenvector basis

### 一些解題重要性質

- $\operatorname{Det}(A) = \prod_{i=1}^n \lambda_i$
- If  $\operatorname{Det}(A) = 0 \implies \exists \lambda_i = 0 \implies (A \lambda I) |v\rangle = A |v\rangle = 0$
- $Tr(A) = \sum_{i=1}^{n} \lambda_i$
- For a real  $\pmb{A}$ , its complex  $\lambda_i$  (if any) must appear in complex-conjugate pairs.
- If  $\lambda i$  are the eigenvalues of A, then:
  - (a) the eigenvalues of  $A^{\dagger}$  are  $\lambda_i^*$
  - ullet (b) the eigenvalues of  $A^{-1}$  are  $\lambda_i^{-1}$
  - (c)  $A, A^{\dagger}, A^{-1}$  have the same eigenvectors
- Eigenvalues remain the same under orthogonal/unitary transformation of coordinate system
- If A is Hermitian, then: (skip proof)
  - $\lambda_i$  is real number

- Each eigenvectors must be orthogonal
- 若 $P_A(\lambda)=|A-\lambda I|=0$ 有 m 重根,則對應到該重根  $\lambda$  的 eigenspace 可能有 1 至m 個線性獨立的eigenvectors
- 若 A 為 normal, 則 m 重根 λ 剛好有 m 個線性獨立的 eigenvectors。
- If A is normal, each eigenvector corresponded to different  $\lambda_i$  are orthogonal
- If A is anti-Hermitian,  $\lambda_i$  will be pure imaginary number or zero
- ullet If A is unitary,  $|\lambda_i|=1$
- The eigenvalues of a lower or upper triangular matrix are the diagonal entries of the matrix

## Rank, nullity and eigenvalues

Triangular factorization  $A_{n imes m} = L_{n imes n} U_{n imes m}$ 

- ullet U: upper triangular matrix,可由對 A 做 Gaussian elimination 求得
- rank: U 中非零列的數量
- pivot: 每一非零列的第一個數
- L: Lower triangular matrix,規則看下面範例,這很酷

E.g.1

$$A = egin{pmatrix} 1 & 0 & 0 \ 1 & 3/2 & 1/2 \ 0 & 1/2 & 3/2 \end{pmatrix}$$

執行Gaussian elimination:

Step 1: 將第2列扣除第1列的1倍: $L_{21}$ =1

Step 2: 將第3列扣除第1列的0倍: $L_{31}$ =0

Step 3: 將第3列扣除第2列的1/3倍: $L_{32}$ =1/3

$$L = egin{pmatrix} 1 & 0 & 0 \ 1 & 1 & 0 \ 0 & 1/3 & 1 \end{pmatrix}\!, \quad U = egin{pmatrix} 1 & 0 & 0 \ 0 & 3/2 & 1/2 \ 0 & 0 & 4/3 \end{pmatrix}$$

Rank=3, Pivots=1, 3/2, 4/3

藉此方式,我們可以快速得到 A 的 Det,因為  $Det(A) = \prod_{i=0}^n U_{ii}$ 

然而,有些時候 A 不能直接做高斯消去,因此我們需要引入一個 P 矩陣,其作用在改變 A 的行順序,式子變為 PA=LU

E.g.2

$$A = egin{pmatrix} 1 & 0 & 0 \ 0 & 1/2 & 3/2 \ 1 & 3/2 & 1/2 \end{pmatrix}, \quad P = egin{pmatrix} 1 & 0 & 0 \ 0 & 0 & 1 \ 0 & 1 & 0 \end{pmatrix} \implies PA = egin{pmatrix} 1 & 0 & 0 \ 1 & 3/2 & 1/2 \ 0 & 1/2 & 3/2 \end{pmatrix}.$$

這時我們發現,在寫程式的時候 U 這逼會有點麻煩,因為他長得有點醜。所以我們引入一個  $D=\mathrm{diag}(\mathrm{pivots})$  ,然後讓 U 變成 U' (pivots 變成1),這樣記憶體運用效率比較高,因此 e.g.2 變成

$$PA = LDU' \implies egin{pmatrix} 1 & 0 & 0 \ 1 & 3/2 & 1/2 \ 0 & 1/2 & 3/2 \end{pmatrix} = egin{pmatrix} 1 & 0 & 0 \ 1 & 1 & 0 \ 0 & 1/3 & 1 \end{pmatrix} egin{pmatrix} 1 & 0 & 0 \ 1 & 3/2 & 0 \ 0 & 0 & 4/3 \end{pmatrix} egin{pmatrix} 1 \ 0 \ 0 \end{bmatrix}$$

最後是 Triangular factorization 的存在性及唯一性:

存在性1:對於可逆矩陣,若且唯若它的所有(包含原矩陣和反矩陣) leading principal minors 都非零,則其 LU 分解存在。

存在性2:對於不可逆矩陣,rank 為 k,若其前 k 個首要 leading principal minors 都非零,則其 LU 分解存在。

存在性3:對於  $n \times n$  方形 (square) 矩陣,若其前 n-1 行為線性獨立,但至少有一個 leading principal minors 為零,則其 LU 分解不存在。

唯一性1:對於方形的可逆矩陣,若其 LU 分解存在,則為唯一。

唯一性2:對於  $n \times n$  方形矩陣,若其前 n-1 行非線性獨立,則其有無窮多個 LU 分解。

### 解方程式

n imes n 矩陣 A 的 eigenvalues 若有 m 個 0 解,則定義 rank=n-m。若 rank < n,則  $A^{-1}$  不存在,也就是  $\mathrm{Det}(A)=0$ 

rank 和m元一次聯立方程式之解的關係:解帶有 (m - rank) 個獨立變數。以下我們將以求解一方程式做舉例

$$egin{cases} x+y+z &=1 \ 2x+2y+3z &=3 \implies Aw=b \implies egin{cases} x+y+z=1 \ z=1 \ 0=0 \end{cases}$$

我們發現在第一行式子只會得到 x = -y 的結果,因此此解帶有 1 個獨立變數,與關係式 m - rank 相等,選擇 x 為獨立變數,其解為:

$$egin{cases} x = x \ y = -x \implies w = w_{ ext{particular}} + w_{ ext{nullspace}} \ z = 1 \end{cases}$$

此處即為

$$w_p = egin{pmatrix} 0 \ 0 \ 1 \end{pmatrix}, \quad w_n = x egin{pmatrix} 1 \ -1 \ 0 \end{pmatrix} \implies w = egin{pmatrix} x \ -x \ 1 \end{pmatrix}$$

其中需要注意的是  $w_n$  為 Aw=0 的解,所有 w 的集合構成 nullspace

最後是 rank-nullity theorem: n = rank + nullity

• kernel(A): 所有滿足  $A\vec{x} = \vec{0}$  的向量

• nullspace: kernel 的同義詞

• nullity: kernel 的維度(有幾個自由變數)

## Diagonalization of a matrix

Given a n imes n normal matrix A (e.g. n=3), where the orthonormal eigenvectors are  $|v_1\rangle, |v_2\rangle, |v_3\rangle$ . Let  $U=(|v_1\rangle, |v_2\rangle, |v_3\rangle)$ , which is unitary. Hence, U can be used to perform unitary transformation on A

$$\Lambda = U^\dagger A U = egin{pmatrix} \lambda_1 & 0 & 0 \ 0 & \lambda_2 & 0 \ 0 & 0 & \lambda_3 \end{pmatrix}$$

Important: A and  $\Lambda$  are similar

 ${\cal A}$  is diagonalized by  ${\cal U}.$  The eigenvectors of  ${\cal A}$  become the basis of new coordinates

Something you should notice:

- Normal matrices are always diagonalizable
- The eigenvectors of  $\Lambda$  are canonical basis

$$(1,0\ldots 0),(0,1\ldots 0),\ldots (0,0\ldots 1)$$

- ullet The eigenvalues of  $\Lambda$  is the same as A
- $\operatorname{Det}(A) = \operatorname{Det}(\Lambda)$
- $\operatorname{Tr}(A) = \operatorname{Tr}(\Lambda)$
- ullet  $A^m = U \Lambda^m U^\dagger$
- $ullet ext{Tr}(A^m) = ext{Tr}(\Lambda^m) = \sum_{i=1}^n \lambda_i^m$
- 求反矩陣

$$A^{-1} = U egin{pmatrix} \lambda_1^{-1} & 0 & 0 \ 0 & \lambda_2^{-1} & 0 \ 0 & 0 & \lambda_3^{-1} \end{pmatrix} U^\dagger$$

### **Quadratic and Hermitian forms**

Homogeneous function

$$f(sx) = s^k f(x)$$
, k: degree

When k=2, f(x) is a homogeneous quadratic function, which is always the quadratic form of a symmetric matrix.

Quadratic form of a real symmetric A is

$$f(x) = Q(x) = x^TAx = (Ux')^T(U\Lambda U^T)(Sx') = x^{'T}\Lambda x' = \lambda_i x_i^{'2}$$

E.g.1: Diagonalize the quadratic form  $Q=2x^2+4xy+5y^2$ 

$$Q = (x,y) egin{pmatrix} a & b/2 \ b/2 & c \end{pmatrix} egin{pmatrix} x \ y \end{pmatrix} = (x,y) egin{pmatrix} 2 & 2 \ 2 & 5 \end{pmatrix} egin{pmatrix} x \ y \end{pmatrix} = x^T A x$$

$$\implies A = S\Lambda S^T = egin{pmatrix} rac{2}{\sqrt{5}} & rac{1}{\sqrt{5}} \ -rac{1}{\sqrt{5}} & rac{2}{\sqrt{5}} \end{pmatrix} egin{pmatrix} 1 & 0 \ 0 & 6 \end{pmatrix} egin{pmatrix} rac{2}{\sqrt{5}} & -rac{1}{\sqrt{5}} \ rac{1}{\sqrt{5}} & rac{2}{\sqrt{5}} \end{pmatrix}$$

$$\lambda = 1, 6 > 0 \Rightarrow A$$
 is positive definite

$$\implies Q = x'^2 + 6y'^2 = \left[rac{1}{\sqrt{5}}(2x - y)
ight]^2 + 6\left[rac{1}{\sqrt{5}}(x + 2y)
ight]^2$$

接著我們可以來複習酷酷的二次曲面

$$Q(x) = k \Leftrightarrow \lambda_i^2 x_i^2 = k$$

沒了,剩下隨便

然後我們來討論一下什麼時候 symmetric A 會是 positive definite ,只要滿足以下之一:

- $\lambda_i > 0$
- $Q(x) > 0 \ \forall x$
- All pivots > 0
- $A = R^T R$

### 求R的方法:

法一 Eigenvalues

$$A = U \Lambda U^\dagger \implies R = U \sqrt{\Lambda} U^\dagger$$

法二 Eigen decomposition

$$A = U \Lambda U^\dagger \implies R = \sqrt{\Lambda} U^\dagger$$

法三 Triangular factorization 複習一下

$$A=LU=LDU'$$
,若 $A$ 為 symmetric,則

$$A = LDL^T \implies R = \sqrt{D}L^T$$
 (Cholesky decomposition)

我們也可以用這個方法來快速的找到 A 的反矩陣

$$A^{-1} = R^{-1}R^{-1^T} = P^TP$$
 where  $P = R^{-1^T} = \sqrt{D^{-1}}L^{-1}$