

# Overview

Intro to  
Comp. Phys.

James  
Farrell &  
Jure  
Dobnikar

Numerical  
Differenti-  
ation

Numerical  
Integration

## 1 Numerical Differentiation

## 2 Numerical Integration

The **Taylor expansion** of continuous and differentiable function of single variable,  $f(x)$ , around the point  $x + h$ , is given by,

$$\begin{aligned} f(x+h) &= f(x) + hf'(x) + \frac{1}{2!}h^2f''(x) + \frac{1}{3!}h^3f'''(x) + \dots \\ &= \sum_{n=0}^{\infty} \frac{h^n}{n!} f^{(n)}(x) \end{aligned}$$

We can **truncate** a Taylor expansion to arrive at **approximations of different orders**,

$$f(x+h) = f(x) + \mathcal{O}(h)$$

$$f(x+h) = f(x) + hf'(x) + \mathcal{O}(h^2)$$

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \frac{1}{6}h^3f'''(x) + \mathcal{O}(h^4)$$

The approximations converge to the correct answer as  $h \rightarrow 0$ .

Rearranging the **first-order** approximation yields a **first-order** approximation to the derivative,

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$$\boxed{f'(x) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)} \quad (1)$$

FIRST FORWARD DIFFERENCES

# Difference Formulae

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Expanding  $f(x - h)$  instead, we get a different approximation with the same error characteristics,

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$$f(x-h) = f(x) - hf'(x) + \mathcal{O}(h^2)$$

$$\boxed{f'(x) = \frac{f(x) - f(x-h)}{h} + \mathcal{O}(h)} \quad (2)$$

FIRST BACKWARD DIFFERENCES

We can eliminate higher-order derivatives and get higher-order approximations by taking **linear combinations** of approximations at **different points**, giving us a **second-order approximation** to the derivative,

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$\begin{aligned} f(x+h) - f(x-h) &= +f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \\ &\quad - f(x) + hf'(x) - \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3) \end{aligned}$$

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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$f(x+h) - f(x-h) = 2hf'(x) + \mathcal{O}(h^3)$$

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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$\boxed{f'(x) = \frac{f(x+h) - f(x-h)}{2h} + \mathcal{O}(h^2)} \quad (3)$$

FIRST CENTRAL DIFFERENCES

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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$\boxed{f'(x) = \frac{f(x+h) - f(x-h)}{2h} + \mathcal{O}(h^2)} \quad (3)$$

## FIRST CENTRAL DIFFERENCES

Not only do the zeroth and second order terms of the Taylor series cancel, **ALL other even order terms** also cancel.

First central differences converges to the true value of  $f'(x)$  faster than first forward or first backward differences.

# Comparison of Forward, Backward, and Central Differences

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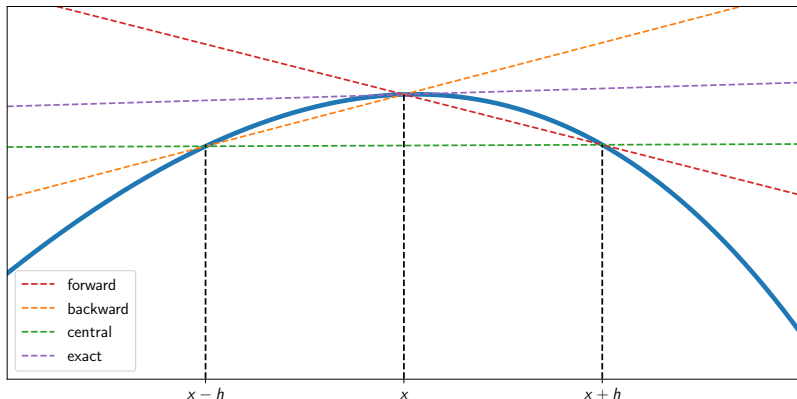


Figure 1: Comparison of the first forward, backward, and central differences methods for  $f(x) = x - x^2 + x^3 - x^4$  centred at the point  $x = 0.6$  with step size  $h = 0.1$ .

	forward	backward	central	exact
$f'(0.6)$	-0.135	0.139	0.002	0.016

Table 1:  $f'(0.6)$  at several levels of approximation.

Higher order derivatives can also be obtained in this way,

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$



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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$\begin{aligned} f(x+h) + f(x-h) &= +f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \\ &\quad + f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3) \end{aligned}$$

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$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

$$f''(x) = \frac{f(x+h) - 2f(x) + f(x-h)}{h^2} + \mathcal{O}(h^2) \quad (4)$$

SECOND CENTRAL DIFFERENCES

We can construct approximations to **arbitrary order** by adding more and more combinations, e.g., this fourth-order approximation to the first derivative,

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \frac{1}{6}h^3f'''(x) + \frac{1}{24}h^4f^{(4)}(x) + \mathcal{O}(h^5)$$

$$f(x-h) = f(x) - hf'(x) + \frac{1}{2}h^2f''(x) - \frac{1}{6}h^3f'''(x) + \frac{1}{24}h^4f^{(4)}(x) + \mathcal{O}(h^5)$$

$$f(x+2h) = f(x) + 2hf'(x) + 2h^2f''(x) + \frac{4}{3}h^3f'''(x) + \frac{2}{3}h^4f^{(4)}(x) + \mathcal{O}(h^5)$$

$$f(x-2h) = f(x) - 2hf'(x) + 2h^2f''(x) - \frac{4}{3}h^3f'''(x) + \frac{2}{3}h^4f^{(4)}(x) + \mathcal{O}(h^5)$$

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$$f(x+h) - f(x-h) = 2hf'(x) + \frac{1}{3}h^3f'''(x) + \mathcal{O}(h^5)$$

$$f(x+2h) - f(x-2h) = 4hf'(x) + \frac{8}{3}h^3f'''(x) + \mathcal{O}(h^5)$$

We can construct approximations to **arbitrary order** by adding more and more combinations, e.g., this fourth-order approximation to the first derivative,

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$$8(f(x+h) - f(x-h)) - (f(x+2h) - f(x-2h)) = 12hf'(x) + \mathcal{O}(h^5)$$

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$$f'(x) = \frac{f(x-2h) - 8f(x-h) + 8f(x+h) - f(x+2h)}{12h} + \mathcal{O}(h^4) \quad (5)$$

FIVE POINT FORMULA

A solution for **boundary points** with the same error order as first central differences can be obtained by taking differences of  $f(x+h)$ ,  $f(x+2h)$ ,

$$f(x+h) = f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3)$$

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$$f(x+2h) = f(x) + 2hf'(x) + 2h^2f''(x) + \mathcal{O}(h^3)$$

$$f(x+2h) - 4f(x+h) = -3f(x) - 2hf'(x) + \mathcal{O}(h^3)$$

A solution for **boundary points** with the same error order as first central differences can be obtained by taking differences of  $f(x+h)$ ,  $f(x+2h)$ ,

$$\begin{aligned}f(x+h) &= f(x) + hf'(x) + \frac{1}{2}h^2f''(x) + \mathcal{O}(h^3) \\f(x+2h) &= f(x) + 2hf'(x) + 2h^2f''(x) + \mathcal{O}(h^3)\end{aligned}$$

$$f'(x) = \frac{4f(x+h) - 3f(x) - f(x+2h)}{2h} + \mathcal{O}(h^2) \quad (6)$$

which expression approximates the derivative at  $x$  **using only  $x$  and points to the right of  $x$ .**

# Non-uniform grids

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It is easy enough to modify the second central difference equations to incorporate **non-uniform spacings**  $a, b$ ,

$$b^2 f(x+a) = b^2 f(x) + ab^2 f'(x) + \frac{1}{2} a^2 b^2 f''(x) + a^3 b^2 f'''(x) + \mathcal{O}(a^4)$$

$$a^2 f(x-b) = a^2 f(x) - ba^2 f'(x) + \frac{1}{2} b^2 a^2 f''(x) - b^3 a^2 f'''(x) + \mathcal{O}(b^4),$$

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rearranging,

$$f'(x) = \frac{b^2 f(x+a) - a^2 f(x-b) + (b^2 - a^2) f(x)}{ab(b+a)} + \mathcal{O}(h^2) \quad (7)$$

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rearranging,

$$f'(x) = \frac{b^2 f(x+a) - a^2 f(x-b) + (b^2 - a^2) f(x)}{ab(b+a)} + \mathcal{O}(h^2) \quad (7)$$

Notice that, with uneven spacings, **only the second order terms** in the Taylor series cancel—other, **higher-order even terms remain**.

# Richardson Extrapolation

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Say we can express a problem in the form,

$$A = A(h) + Kh^k + K'h^{k+1} + K''h^{k+2} + \dots$$

where  $h, k, A(h)$  are known, and the constants  $K^n$  are **in general not known**.

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where  $h, k, A(h)$  are known, and the constants  $K^n$  are **in general not known**.

Truncating at  $k$ th order in  $h$ , we get,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1}).$$

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Truncating at  $k$ th order in  $h$ , we get,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1}).$$

How can we **improve the error characteristics** of this expression?



# Richardson Extrapolation

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Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$A = A\left(\frac{h}{2}\right) + K\left(\frac{h}{2}\right)^k + \mathcal{O}(h^{k+1})$$

Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$A = A\left(\frac{h}{2}\right) + \frac{1}{2^k}K(h)^k + \mathcal{O}(h^{k+1})$$

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Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$2^k A = 2^k A\left(\frac{h}{2}\right) + K(h)^k + \mathcal{O}(h^{k+1})$$

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Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$(2^k - 1) A = 2^k A\left(\frac{h}{2}\right) - A(h) + K(h)^k - K(h)^k + \mathcal{O}(h^{k+1})$$

Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

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Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$A = \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} + \mathcal{O}(h^{k+1})$$

Once again, we take **linear combinations** to eliminate the lowest-order unknown term,

$$A = A(h) + Kh^k + \mathcal{O}(h^{k+1})$$

$$A = B(h) + \mathcal{O}(h^{k+1}) \quad (8)$$

$$B(h) = \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} \quad (9)$$

RICHARDSON EXTRAPOLATION

# Richardson Extrapolation: Example

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$$\begin{aligned} A &= B(h) + \mathcal{O}(h^{k+1}) \\ B(h) &= \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} \end{aligned}$$

let

$$A(h) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$



# Richardson Extrapolation: Example

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let

$$A(h) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$

then  $k = 1$  and

$$B(h) = 2A\left(\frac{h}{2}\right) - A(h)$$

# Richardson Extrapolation: Example

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$$\begin{aligned} A &= B(h) + \mathcal{O}(h^{k+1}) \\ B(h) &= \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} \end{aligned}$$

let

$$A(h) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$

substituting,

$$B(h) = 2 \frac{f(x + h/2) - f(x)}{\frac{h}{2}} - \frac{f(x+h) - f(x)}{h}$$

# Richardson Extrapolation: Example

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$$\begin{aligned} A &= B(h) + \mathcal{O}(h^{k+1}) \\ B(h) &= \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} \end{aligned}$$

let

$$A(h) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$

rearranging,

$$B(h) = \frac{4f(x+h/2) - 3f(x) - f(x+h)}{h}$$

# Richardson Extrapolation: Example

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$$\begin{aligned} A &= B(h) + \mathcal{O}(h^{k+1}) \\ B(h) &= \frac{2^k A\left(\frac{h}{2}\right) - A(h)}{2^k - 1} \end{aligned}$$

let

$$A(h) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$

setting  $h \rightarrow 2h$ ,

$$B(2h) = \frac{4f(x+h) - 3f(x) - f(x+2h)}{2h}$$

which you should remember from equation 6.

# Catastrophic Cancellation

Intro to  
Comp. Phys.

James  
Farrell &  
Jure  
Dobnikar

Numerical  
Differenti-  
ation

Numerical  
Integration

All of the difference methods suffer from catastrophic cancellation.

$$f'(x) = \frac{f(x+h) - f(x)}{h} + \mathcal{O}(h)$$

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$$x - (x - h) \neq h$$

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Remember, float addition is **NOT associative**, so the parentheses are meaningful.



We can avoid this problem by taking an **imaginary step**,

$$f(x + ih) = f(x) + ihf'(x) - \frac{h^2}{2}f''(x) - i\frac{h^3}{6}f'''(x) + \mathcal{O}(h^4)$$

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FIRST COMPLEX DIFFERENCES

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## FIRST COMPLEX DIFFERENCES

Addition of a real and imaginary float **does not incur precision loss!**

# Overview

Intro to  
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Numerical  
Differenti-  
ation

Numerical  
Integration

## 1 Numerical Differentiation

## 2 Numerical Integration

# Numerical Integration

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Differenti-  
ation

Numerical  
Integration

A wealth of methods exist for approximating integrals with no closed form (or any other integral, for that matter).

$$\int_a^b \sqrt{1-x^4} dx$$

$$\int_a^b \frac{1}{\log x} dx$$

$$\int_a^b \exp(-x^2) dx$$

$$\int_a^b \frac{\sin x}{x} dx$$

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$$\int_a^b \frac{\sin x}{x} dx$$

We will look at **two** classes of approximations,

**NEWTON-COTES** and **GAUSSIAN QUADRATURE**

which both work by cutting up the integral into **slices**,

$$\int_a^b f(x) dx = \int_a^{a+h} f(x) dx + \dots + \int_{b-h}^b f(x) dx$$

small intervals on which  $f$  varies **smoothly**.



# From Taylor series

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ation

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Integration

Idea: integrate a Taylor expansion, then truncate the integral.

$$\int_a^b f(x) dx = \int_{-h}^h f(x_0 + y) dy$$

where  $h = (a - b)/2$ ,  $x_0 = (a + b)/2$ .

# From Taylor series

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$$\int_{-h}^h f(x_0 + y) dy = \int_{-h}^h \sum_{n=0}^{\infty} \frac{f^n(x_0)}{n!} y^n$$

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where  $h = (a - b)/2$ ,  $x_0 = (a + b)/2$ .

$$\int_{-h}^h f(x_0 + y) dy = \sum_{n=0}^{\infty} \frac{f^n(x_0)}{n!} \int_{-h}^h y^n \quad (11)$$

$$= \sum_{n=0}^{\infty} \frac{f^n(x_0)}{n!} \left[ \frac{y^{n+1}}{n+1} \right]_{-h}^h \quad (12)$$

$$= \sum_{n=0}^{\infty} \frac{f^n(x_0)}{(n+1)!} (h^{n+1} - (-h)^{n+1}) \quad (13)$$

$$= \sum_{n=0}^{\infty} \frac{f^{2n}(x_0)}{(2n+1)!} 2h^{n+1} \quad (14)$$

# First order

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$$\int_{-h}^h f(x_0 + y) dy = \sum_{n=0}^{\infty} \frac{f^{(2n)}(x_0)}{(2n+1)!} 2h^{n+1}$$

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$$\int_{-h}^h f(x_0 + y) dy = \sum_{n=0}^{\infty} \frac{f^{(2n)}(x_0)}{(2n+1)!} 2h^{n+1}$$

Expanding to first order,

$$\int_{-h}^h f(x_0 + y) dy = 2hf(x_0) + \mathcal{O}(h^3)$$

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$$\boxed{\int_a^b f(x) dx \approx 2hf\left(\frac{a+b}{2}\right)}$$
$$h = (b - a)/2$$

(15)

MIDPOINT RULE

## Second order

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Numerical  
Differenti-  
ation

Numerical  
Integration

$$\int_{-h}^h f(x_0 + y) dy = \sum_{n=0}^{\infty} \frac{f^{2n}(x_0)}{(2n+1)!} 2h^{n+1}$$

## Second order

$$\int_{-h}^h f(x_0 + y) dy = \sum_{n=0}^{\infty} \frac{f^{(2n)}(x_0)}{(2n+1)!} 2h^{n+1}$$

Expanding to second order,

$$\int_{-h}^h f(x_0 + y) dy = 2hf(x_0) + \frac{f''(x_0)}{3} h^3 + \mathcal{O}(h^5)$$



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But!

$$f''(x_0) = \frac{f(x_0 - h) - 2f(x_0) + f(x_0 + h)}{h^2} + \mathcal{O}(h^2)$$

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$$\int_a^b f(x) dx \approx \frac{h}{3} \left[ f(a) + 4f\left(\frac{a+b}{2}\right) + f(b) \right]$$

$$h = (a - b)/2$$

(16)

SIMPSON'S RULE

# Summing up: composite rules

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- The midpoint rule and Simpson's rule allow us to approximate the integral over a **single interval**.

# Summing up: composite rules

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- The approximation **improves** as the interval becomes **smaller** (proportional to  $h^3$  for the midpoint rule,  $h^5$  for Simpson's rule).

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- We can take **higher-order expansions** to get formulae with better **asymptotic errors**...

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- We can take **higher-order expansions** to get formulae with better **asymptotic errors**...
- ...or **sum up** the approximations for many intervals to get smaller errors for each small interval.



# Summing up: composite rules

Take the midpoint rule as an example. If the approximation is made as a sum over many intervals,

$$\int_a^b f(x) dx = \sum_{i=0}^{n-1} [2hf(a + ih) + \mathcal{O}(h^3)] ; \quad h = \frac{b-a}{2n}$$

# Summing up: composite rules

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COMPOSITE MIDPOINT RULE

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### COMPOSITE MIDPOINT RULE

Notice the leading error term **goes down** one order of  $h$ .

- When we derived equation 16, Simpson's rule for the integral of a function on an interval, we **truncated** a Taylor expansion of an integral and substituted an **approximation for a second derivative**.

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- Idea: approximate the integral as a **weighted sum** of function values at  $m$  **evenly-spaced** points on the interval,

$$A = \int_a^b f(x) dx \approx \sum_{i=1}^m w_i f(x_i) = \sum_{i=1}^m w_i f_i$$

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- Q. How do we choose the weights?
- A. Choose the weights so that **integrals of polynomials** up to a given order are **exact**.
- Q. Why is that appealing?
- A. The Taylor expansion of a function about a point is an infinite order polynomial; close to the point, the function and its integral are approximated well by **a few low-order monomials**.

# Deriving Newton-Cotes Weights

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Differenti-  
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Numerical  
Integration

For  $m$  evenly-spaced points, we get exact results for order- $(n-1)$  polynomials by solving  $m$  equations:

$$f(x) = 1 \quad \Rightarrow \quad \int_{x_1}^{x_m} 1 dx = x_m - x_1 = \sum_{i=1}^m w_i \cdot 1;$$

$$f(x) = x \quad \Rightarrow \quad \int_{x_1}^{x_m} x dx = \frac{1}{2} (x_m^2 - x_1^2) = \sum_{i=1}^m w_i x_i;$$

$$\vdots$$
$$\vdots$$

$$f(x) = x^{m-1} \quad \Rightarrow \quad \int_{x_1}^{x_m} x^{m-1} dx = \frac{1}{m} (x_m^m - x_1^m) = \sum_{i=1}^m w_i x_i^{m-1}.$$

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Solving for a single weight,

$$x_2 - x_1 = \sum_{i=1}^{m+1} w_i f_i$$

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Solving for a single weight,

$$x_2 - x_1 = w_1 \cdot 1$$

$$\implies w_1 = x_2 - x_1$$

# Deriving Newton-Cotes Weights

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Numerical  
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Numerical  
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Solving for a single weight,

$$x_2 - x_1 = w_1 \cdot 1$$

$$\implies w_1 = x_2 - x_1$$

$$\boxed{\int_{x_1}^{x_2} f(x) dx \approx hf_1} \quad (18)$$
$$h = x_2 - x_1$$

RECTANGLE RULE

# Deriving Newton-Cotes Weights

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Numerical  
Differenti-  
ation

Numerical  
Integration

Solving for two weights,

$$x_2 - x_1 = w_1 \cdot 1 + w_2 \cdot 1$$

$$\frac{1}{2} (x_2^2 - x_1^2) = w_1 x_1 + w_2 x_2$$

by symmetry,

# Deriving Newton-Cotes Weights

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Jure  
Dobnikar

Numerical  
Differenti-  
ation

Numerical  
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$$w_1 = w_2$$

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by symmetry,

$$w_1 = w_2 = \frac{1}{2} (x_2 - x_1)$$

$$\boxed{\int_{x_1}^{x_2} f(x) dx \approx \frac{h}{2} (f_1 + f_2)}$$
$$h = x_2 - x_1$$

(19)

TRAPEZIUM RULE

# Deriving Newton-Cotes Weights

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Integration

Solving for three weights,

$$x_3 - x_1 = w_1 \cdot 1 + w_2 \cdot 1 + w_3 \cdot 1$$

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$$\frac{1}{3} (x_3^3 - x_1^3) = w_1 x_1^2 + w_2 x_2^2 + w_3 x_3^2$$

By symmetry,  $w_1 = w_3$ ; we also require the  $x_i$  to be **evenly-spaced**, so can write  $x_1 = x_2 - h$ ;  $x_3 = x_2 + h$ ;  
Substituting and rearranging,

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Substituting and rearranging,

$$2h = 2w_1 + w_2$$

$$2hx_2 = 2w_1 x_2 + w_2 x_2$$

$$2hx_2^2 + \frac{2}{3}h^3 = 2w_1 (x_2^2 + h^2) + w_2 x_2^2$$



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Substituting and rearranging,

$$\implies w_1 = \frac{1}{3}h = w_3; \quad w_2 = \frac{4}{3}h.$$

# Deriving Newton-Cotes Weights

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Substituting and rearranging,

$$\boxed{\int_{x_1}^{x_3} f(x) dx \approx \frac{h}{3} (f_1 + 4f_2 + f_3)} \quad (20)$$
$$2h = x_3 - x_1$$

SIMPSON'S RULE

# Error in the Trapezium Rule

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Integration

Compare the trapezium rule,

$$\int_{-h}^h f(x+y) dy \approx h[f(x-h) + f(x+h)] = A_{\text{est}}$$

with the exact integral,

$$\int_{-h}^h f(x+y) dy = 2hf(x) + \frac{f''(x)}{3}h^3 + \mathcal{O}(h^5) = A_{\text{exact}}$$

# Error in the Trapezium Rule

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Q. What is the error in the trapezium rule approximation?

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Q. What is the error in the trapezium rule approximation?

$$A_{\text{exact}} - A_{\text{est}} = 2hf(x) + \frac{f''(x)}{3}h^3 - h[f(x-h) + f(x+h)] + \mathcal{O}(h^5)$$

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but, from the central differences formula,

$$-h^3 f''(x) = 2hf(x) - h[f(x-h) + f(x+h)] + \mathcal{O}(h^5)$$



# Error in the Trapezium Rule

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so,

$$A_{\text{exact}} - A_{\text{est}} = -\frac{2f''(x)}{3}h^3$$

# Romberg's method

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- applying **Richardson extrapolation** to the trapezium rule yields Simpson's rule...

# Romberg's method

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- applying Richardson extrapolation to Simpson's rule yields **Boole's rule**, a Newton–Cotes formula that is exact to fifth order.

# Romberg's method

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- ...doesn't lead to a Newton–Cotes formula.

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- at high orders, the Newton–Cotes formulae containing large weights of different signs, leading to **loss of precision**

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- applying Richardson extrapolation to Boole's rule...
- ...doesn't lead to a Newton–Cotes formula.
- at high orders, the Newton–Cotes formulae containing large weights of different signs, leading to **loss of precision**
- the formulae that come from Richardson extrapolation, **Romberg's methods**, are relatively stable

In deriving the weights for the Newton-Cotes formulae, our goal was to make an  $m$ -point approximation to the integral

$$\int_{x_1}^{x_m} f(x) dx$$

that was exact for **order- $(m - 1)$  polynomials**.



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By **fixing the integration limits to  $(-1, 1)$** , we can make the approximation **exact for all odd monomials**,

$$\int_{-1}^1 x^n = 0 \quad \text{for all odd } n$$

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By **fixing the integration limits to  $(-1, 1)$** , we can make the approximation **exact for all odd monomials**,

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If a change of coordinates that gives these limits is possible, then

## GAUSSIAN QUADRATURE

gives an  $m$ -point approximation exact for **order- $(2m-1)$  polynomials**.

# Deriving Gaussian Quadrature nodes and weights

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Integration

For two points, in this context called **nodes**,  $x_1, x_2$ ;  $x_1 < x_2$  somewhere on the interval  $(-1, 1)$ , we have **four equations with four unknowns**,

$$f(x) = 1 \rightarrow A = \int_{-1}^1 dx = 2 = w_1 + w_2$$

$$f(x) = x \rightarrow A = \int_{-1}^1 x dx = 0 = w_1 x_1 + w_2 x_2$$

$$f(x) = x^2 \rightarrow A = \int_{-1}^1 x^2 dx = \frac{2}{3} = w_1 x_1^2 + w_2 x_2^2$$

$$f(x) = x^3 \rightarrow A = \int_{-1}^1 x^3 dx = 0 = w_1 x_1^3 + w_2 x_2^3$$

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$$f(x) = x^3 \rightarrow A = \int_{-1}^1 x^3 dx = 0 = w_1 x_1^3 + w_2 x_2^3$$

Solving, we find,

$$x_2 = -x_1 = \frac{1}{\sqrt{3}}; \quad w_1 = w_2 = 1.$$

# Deriving Gaussian Quadrature nodes and weights

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$$\boxed{\boxed{\int_{-1}^1 f(x) dx \approx f\left(-\frac{1}{\sqrt{3}}\right) + f\left(\frac{1}{\sqrt{3}}\right)}} \quad (21)$$

GAUSSIAN QUADRATURE ON TWO NODES

In the general case of limits  $[a, b]$ , we must first apply the coordinate transformation,

$$\begin{aligned} t &= \frac{b-a}{2}x + \frac{b+a}{2} \\ \int_a^b f(t) dt &= \int_{-1}^1 f\left(\frac{b-a}{2}x + \frac{b+a}{2}\right) \frac{b-a}{2} dx \\ &= \int_{-1}^1 g(x) dx \end{aligned} \tag{22}$$

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The approximate expression for the integral then becomes,

$$\int_a^b f(t) dt \approx \frac{b-a}{2} \sum_{i=1}^n w_i f\left(\frac{b-a}{2}x_i + \frac{b+a}{2}\right) \quad (23)$$

# Newton–Cotes vs. Gaussian quadrature

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- Gaussian quadrature is more accurate with fewer points, incurring **less computational effort**...



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- ...but requires the function to be sampled at **specific points**, which is generally **not possible** with **experimental** data.

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- Gaussian quadrature is more accurate with fewer points, incurring **less computational effort**...
- ...but requires the function to be sampled at **specific points**, which is generally **not possible** with **experimental** data.
- Newton–Cotes methods can be generalised to work with **non-uniform data points**, making Simpson's method a good choice from integrating experimental data.