Bottomonium suppression from the 3-loop QCD potential

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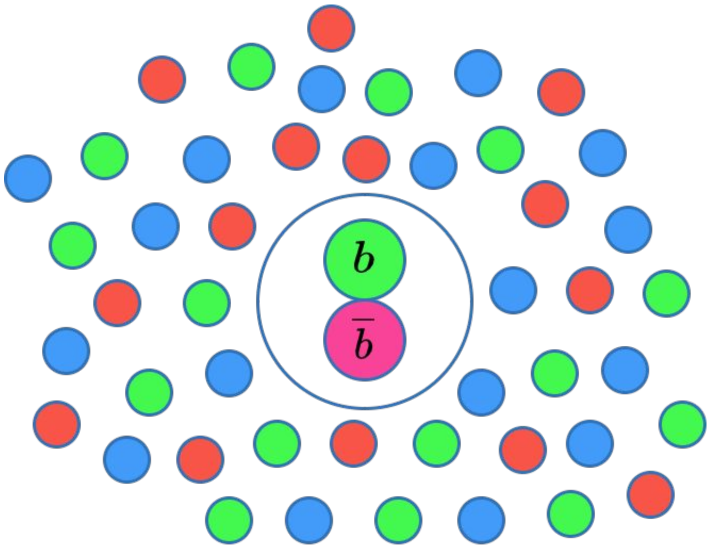
Quarkonium Suppression from first principles

- Quarkonium can be used to probe the QGP
- Matsui & Satz proposed Quarkonium suppression as a signal for the QGP
- Quarkonium dissolves in the QGP
- Measured Quarkonium yields are lower in HIC compared to pp collisions

T. Matsui, H. Satz, Phys. Lett. B 178 (1986) 416

Propagation through QGP

$T \approx O(100\text{MeV})$

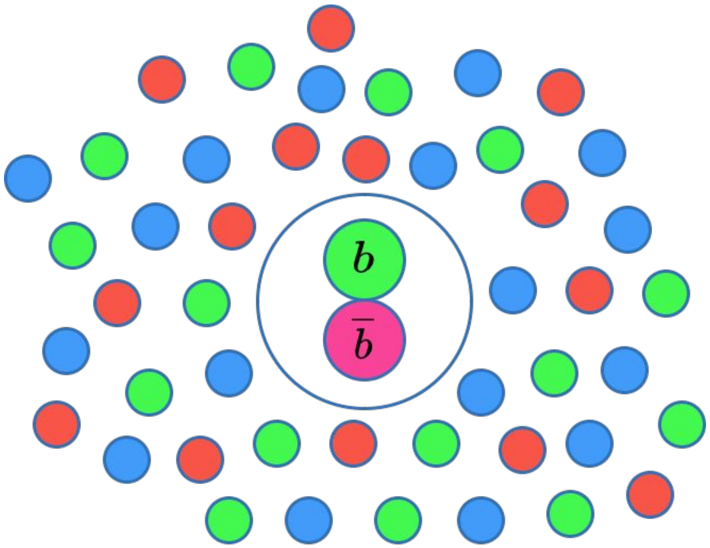


Quarkonium Suppression from first principles

- We aim to describe this phenomenon from **first principles**
- Provide predictions for experiments
- We focus on **bottomonium** since the high mass allows for simplifications

Propagation through QGP

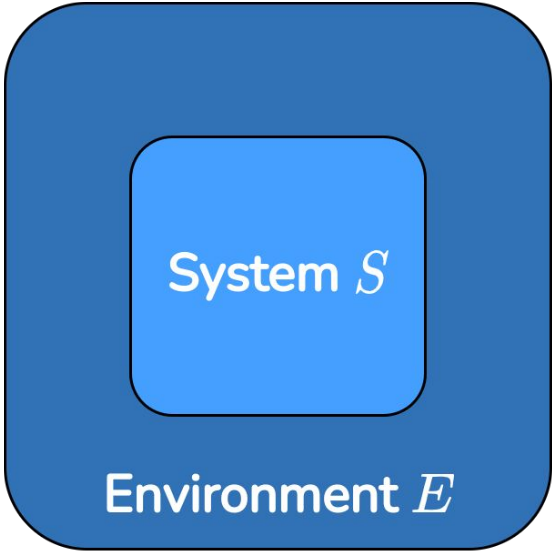
$T \approx O(100\mathrm{MeV})$



Open Quantum Systems

- Quantum system not isolated
- Split into System S and Environment E

$$H = H_S \otimes I_E + I_S \otimes H_E + H_{\mathrm{int}}$$



- Time evolution by Von-Neumann Equation

$$\frac{d}{dt}\rho = -i[H,\rho]$$

- Not interested in environmental d.o.f.: **Trace out**!

$$\rho_S = \mathrm{Tr}_E[\rho]$$

Open Quantum Systems

- Time evolution by Von-Neumann Equation

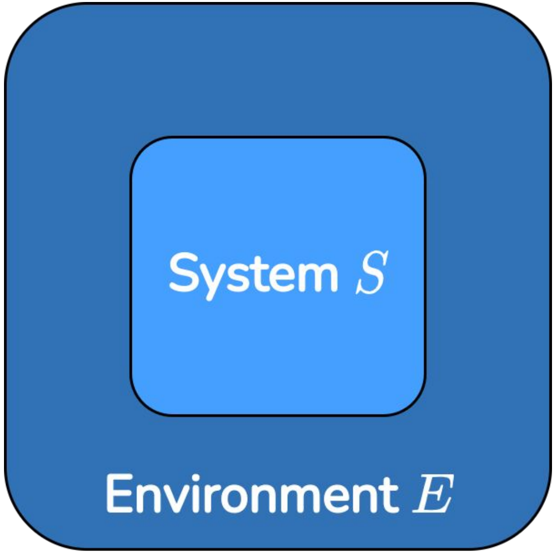
$$\frac{d}{dt}\rho = -i[H, \rho]$$

- Environmental d.o.f. not needed Trace out!

$$\rho_S = \mathrm{Tr}_E[\rho]$$

- "Master equation" for the System: **Lindblad Equation**

$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2}\left\{ C_n^\dagger C_n, \rho_S \right\} \right)$$



Open Quantum Systems

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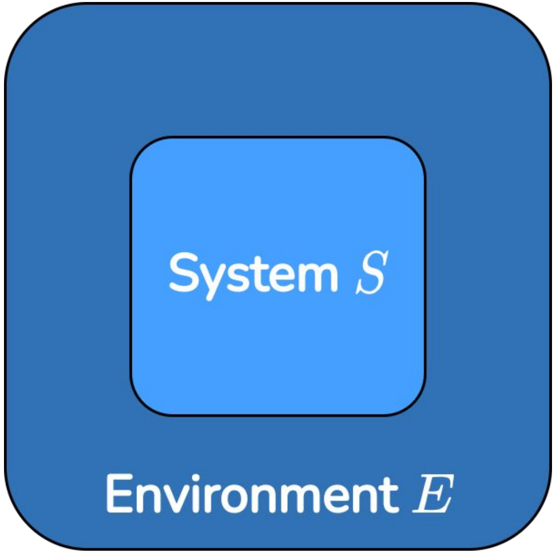
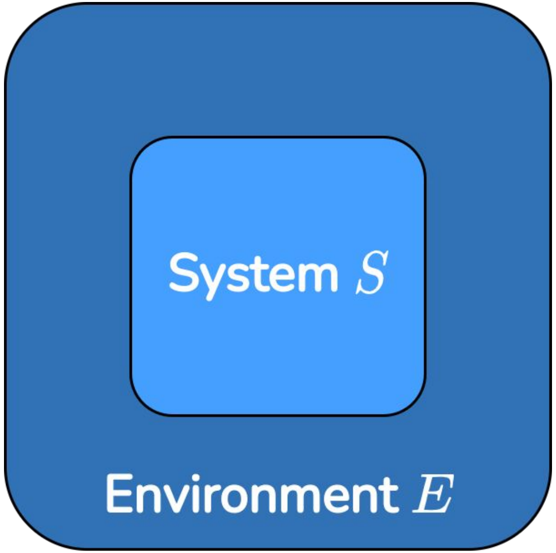


Open Quantum Systems



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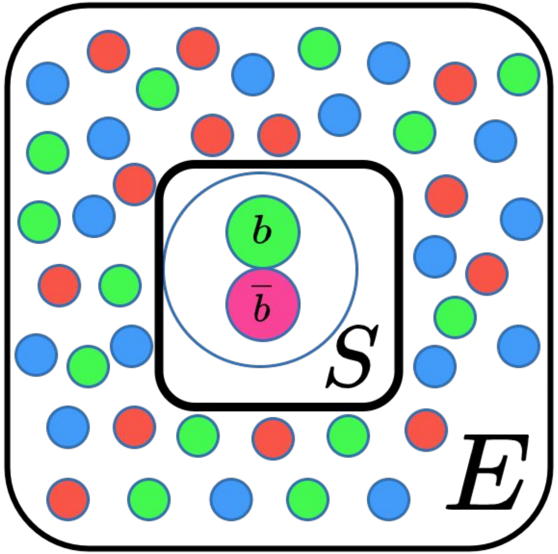
- “Master equation” for the System: **Lindblad Equation**

non-unitary

$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2}\left\{C_n^\dagger C_n, \rho_S\right\} \right)$$

OQS for quarkonium

- Quarkonium: System S
- QGP: Environment E



Aim to describe Quarkonium Suppression by a master equation for encoding the interaction with the QGP

$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2}\{C_n^\dagger C_n, \rho_S\} \right)$$

EFTs for Quarkonium Suppression

- Use NREFTs to exploit hierarchy of scales

$$M \gg 1/a_0 \gg \pi T \gg E$$

- Inverse radius:

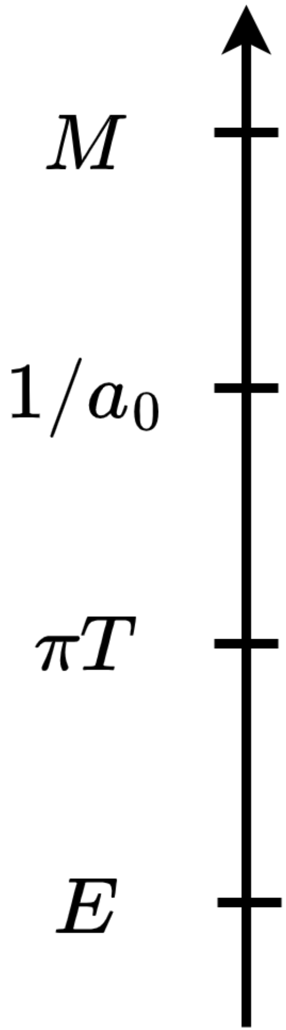
$$1/a_0 \approx 1.2\mathrm{GeV}$$

- Temperature regime:

$$250\mathrm{MeV} < T < 425\mathrm{MeV}$$

- Binding Energy:

$$E \sim 0.4\mathrm{GeV}$$

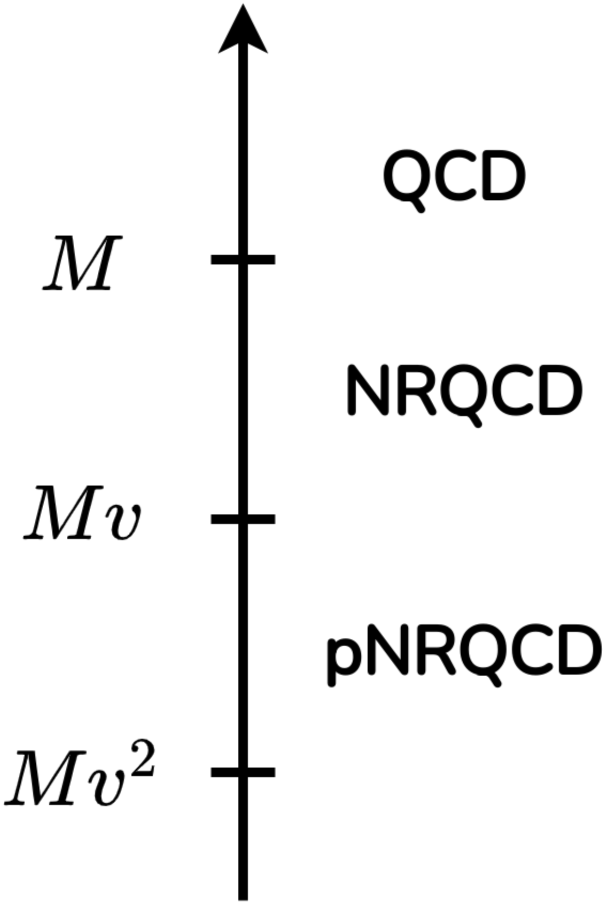


pNRQCD

N. Brambilla, A. Pineda, J. Soto, and A. Vairo, Nuclear Physics B 566, 275 (2000)

- We use pNRQCD, an EFT from full QCD
- pNRQCD is obtained by integrating out the hard scale M and soft scale Mv
- Degrees of freedom: Singlet and octet bound states
- Using pNRQCD one can derive a master equation for the quarkonium density matrix

Brambilla, Escobedo, Soto, Vairo: Phys. Rev. D 97 (2018) 7, 074009



$v \ll 1$: Relative Quark-Antiquark velocity

pNRQCD master equation

Brambilla, Escobedo, Soto, Vairo: Phys. Rev. D 97 (2018) 7, 074009

- In general the master Eq. is not of Lindblad form
- Simplify using hierarchy of scales $T \gg E$

$$A_i^{uv} = \frac{g^2}{6N_c} \int_0^\infty \mathrm{d}s e^{-ih_u s} r_i e^{ih_v s} \left\langle \tilde{E}_j^a(0, \vec{0}) \tilde{E}_j^a(s, \vec{0}) \right\rangle$$

- Expand exponentials in E/T
- At LO in E/T we get

$$A_i^{uv} = \frac{g^2}{6N_c} \int_0^\infty \mathrm{d}s r_i \left\langle \tilde{E}_j^a(0, \vec{0}) \tilde{E}_j^a(s, \vec{0}) \right\rangle$$

$$= \frac{r_i}{2}(\kappa - i\gamma)$$

Transport coefficients

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2}\left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

$$H = \begin{pmatrix} h_s + \frac{r^2}{2}\gamma & 0 \\ 0 & h_o + \frac{N_c^2-2}{2(N_c^2-1)}\frac{r^2}{2}\gamma \end{pmatrix}$$

$$h_{s,o} = \vec{p}^{\,2}/M + V_{s,o}$$

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2}\left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

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Quarkonium Potential

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H,\rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2}\left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

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$$C_i^0 = \sqrt{\frac{\kappa}{N_c^2-1}} r_i \begin{pmatrix} 0 & 1 \\ \sqrt{N_c^2-1} & 0 \end{pmatrix},$$

$$C_i^1 = \sqrt{\frac{\kappa(N_c^2-4)}{2(N_c^2-1)}} r_i \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

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pNRQCD master equation

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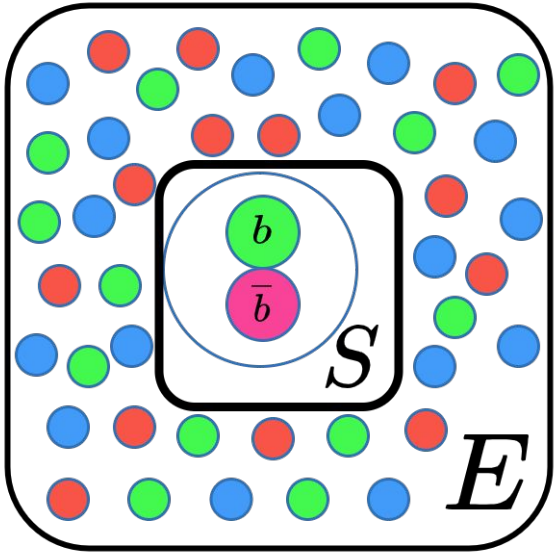
Transport coefficients

pNRQCD master equation



- Hilbert Space:
  - Singlet and octet states

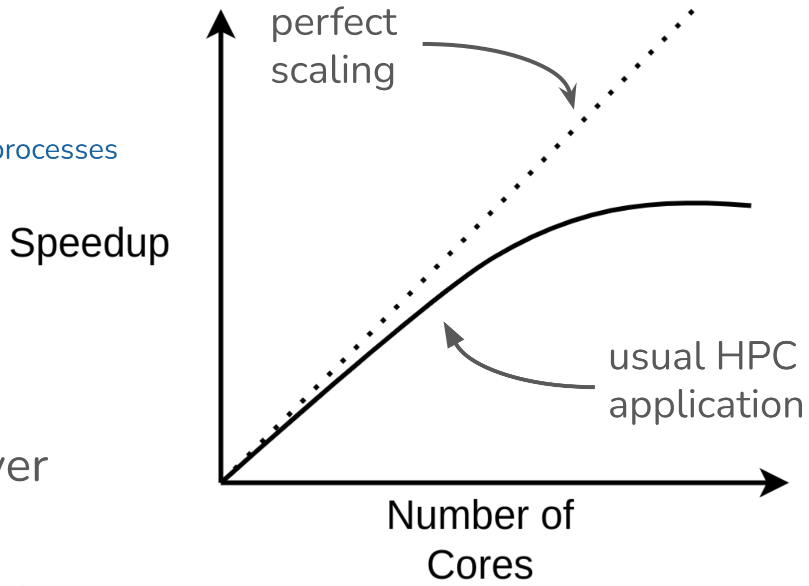
$$\rho = \begin{pmatrix} \rho_s & 0 \\ 0 & \rho_o \end{pmatrix}$$

  - Discretizing radial part of the wavefunction (e.g. 2048 lattice)
  - Angular momentum quantum numbers
- **Very large Hilbert space**

$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2} \left\{ C_n^\dagger C_n, \rho_S \right\} \right)$$

Quantum trajectory algorithm

J. Dalibard, Y. Castin, and K. Mølmer, Wave-function approach to dissipative processes in quantum optics, Phys. Rev. Lett. 68 (1992), pp. 580–583.



- Idea:
  1. Evolve individual trajectories $|\phi(t)\rangle$ stochastically
  2. Calculate observables by averaging over trajectories $\overline{\langle\phi(t)|A|\phi(t)\rangle}$

can evolve to arbitrary l

- Averaging over the density matrix $\sigma(t) = |\phi(t)\rangle\langle\phi(t)|$ restores the Lindblad equation

Advantages:

- Evolve vector of size N_H instead N_H^2 density matrix
- Simulation of individual trajectories is **embarrassingly parallel**

Previous work

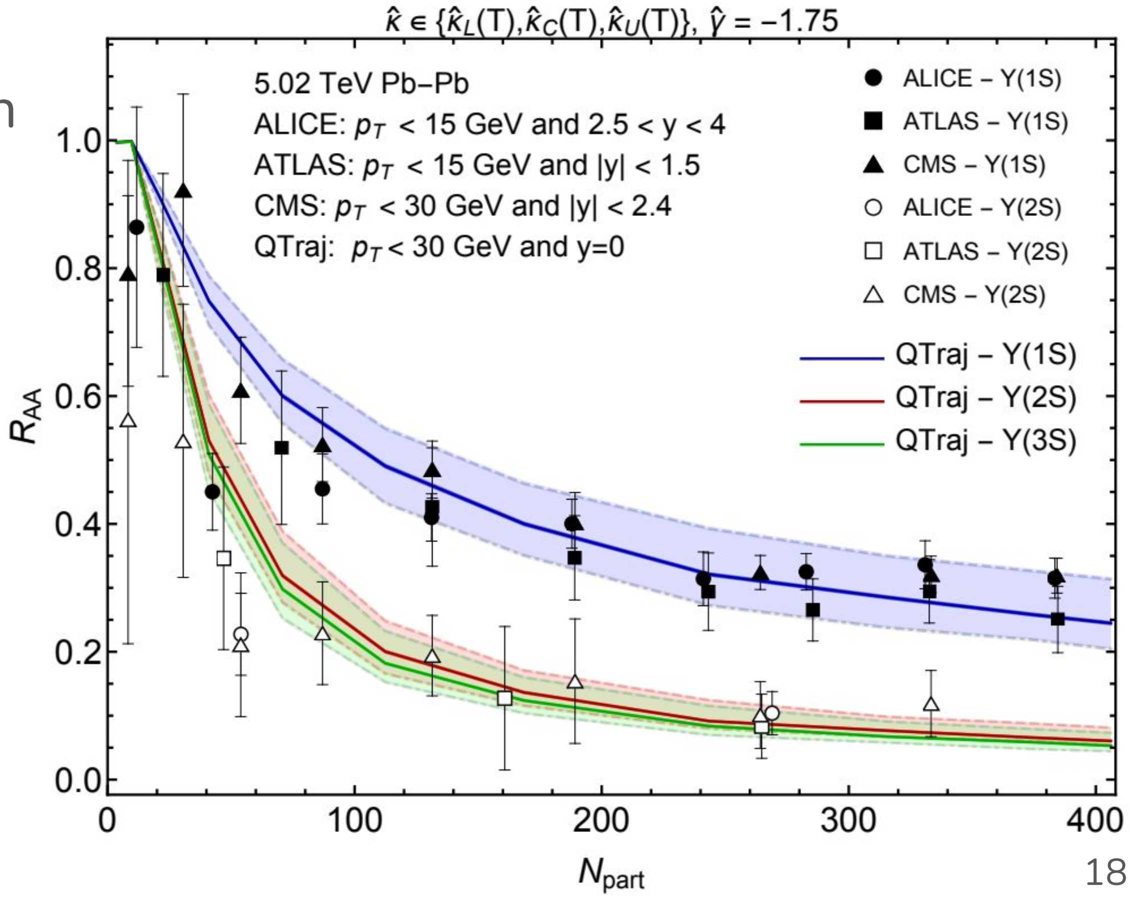
N. Brambilla, M. A. Escobedo, M. Strickland, A. Vairo, P. Vander Griend and J. H. Weber, Phys. Rev. D 104 (2021) 094049

- temperature evolution from hydrodynamics simulation

M. Alqahtani and M. Strickland, The European Physical Journal C 81 (2021)

$$\text{Survival Probability} = \frac{\langle \psi(t) | 1S \rangle}{\langle \psi(0) | 1S \rangle}$$

- Including Feed-down from PDG data



Previous work

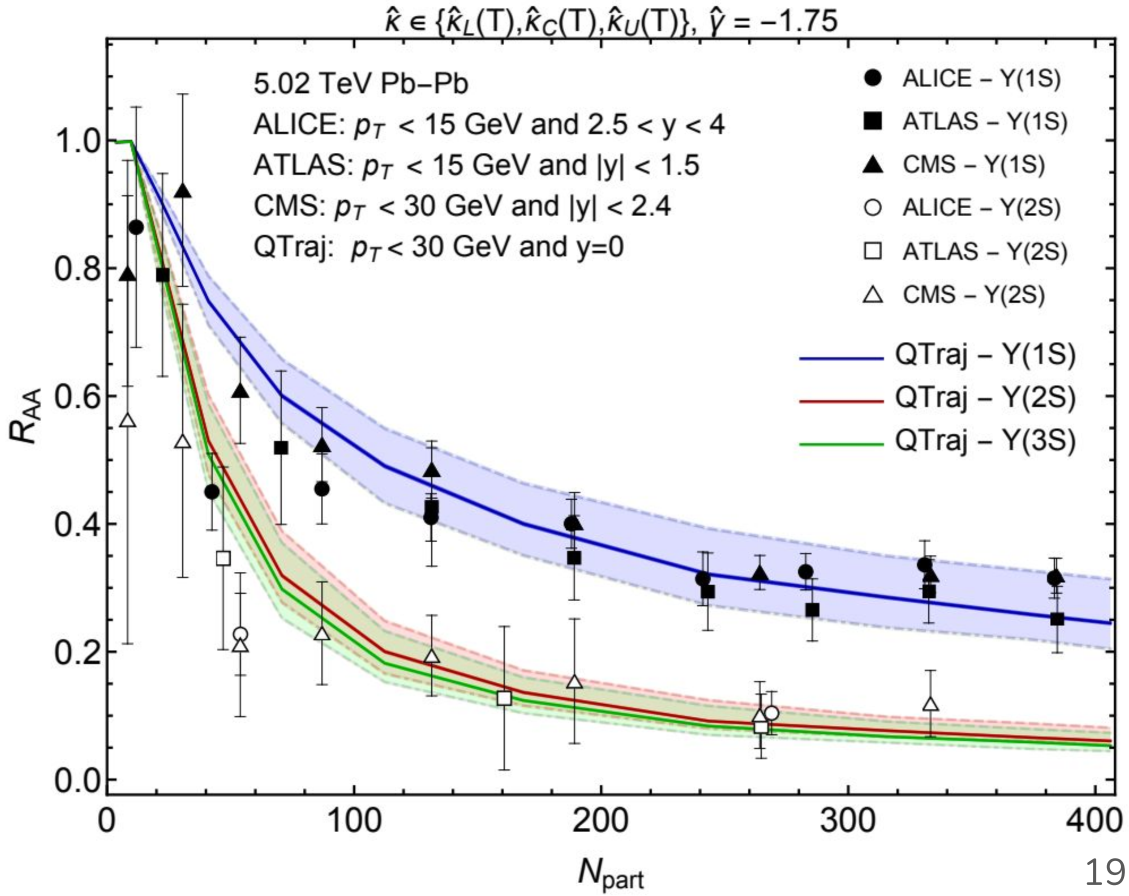
N. Brambilla, M. A. Escobedo, M. Strickland, A. Vairo, P. Vander Griend and J. H. Weber, Phys. Rev. D 104 (2021) 094049

- Coulomb potential

$$V_s = -C_f \alpha_s / r$$
$$V_o = \alpha_s / (2N_c r)$$

- Temperature dependent $\hat{\kappa}$

$$\hat{\gamma} = -1.75$$



New Potential

J. Segovia, S. Steinbeißer, and A. Vairo, Physical Review D 99 (2019)

- Motivation: Implement a higher order potential with a more realistic spectrum

$$V_s^{3L}(r) = V_s^{\text{pert}}(r) + V_s^{\text{non-pert}}(r)$$

$$V_s^{\text{pert}}(\nu, \nu_r, r) = \begin{cases} \sum_{k=0}^{3} V_{s,\text{RS}'}^{(k)} \alpha_s^{k+1}(1/r) & \text{if } r < \nu_r^{-1} \\ \sum_{k=0}^{3} V_{s,\text{RS}'}^{(k)} \alpha_s^{k+1}(\nu) & \text{if } r > \nu_r^{-1} \end{cases}$$

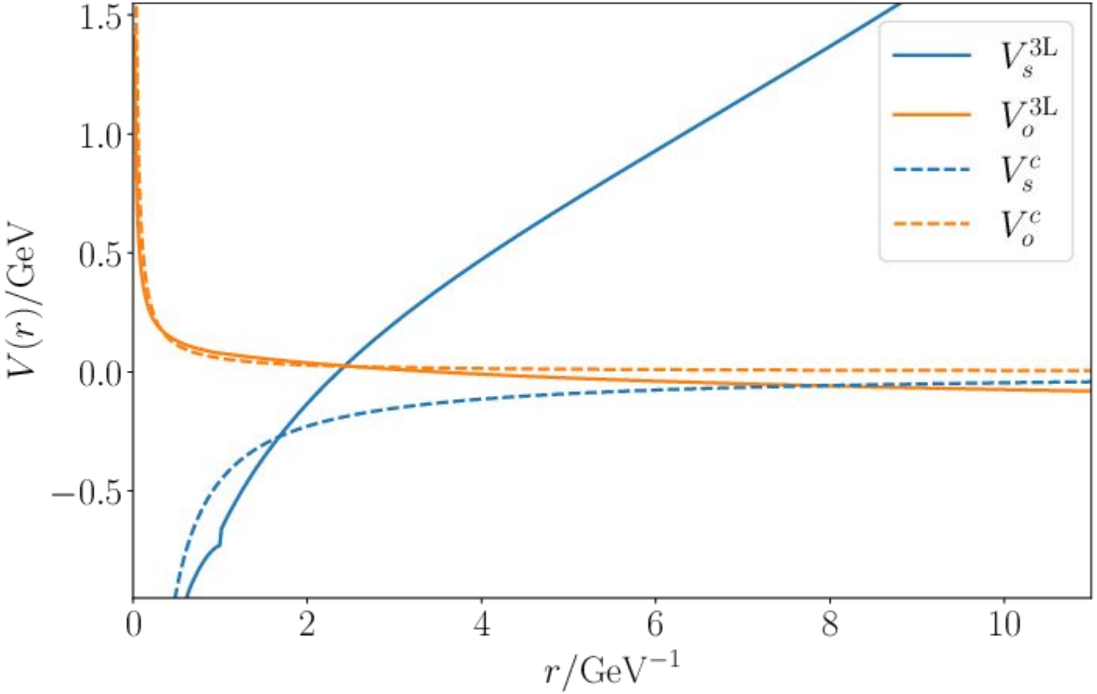
three loop pNRQCD

$$\text{Re}\left(V_s^{\text{non-pert}}(r)\right) = \frac{\gamma}{2} r^2$$

leading non-perturbative correction

New Potential

J. Segovia, S. Steinbeißer, and A. Vairo, Physical Review D 99 (2019)



Spectrum:

	PDG	V_s^c	V_s^{3L}
$m(1S)/\,\mathrm{GeV}$	9.445	9.446	9.445
$m(2S)/\,\mathrm{GeV}$	10.017	9.637	10.066
$m(3S)/\,\mathrm{GeV}$	10.355	9.672	10.451
$m(1P)/\,\mathrm{GeV}$	9.888	9.636	9.892
$m(2P)/\,\mathrm{GeV}$	10.251	9.672	10.320

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm

$$\kappa = \hat{\kappa} T^3$$

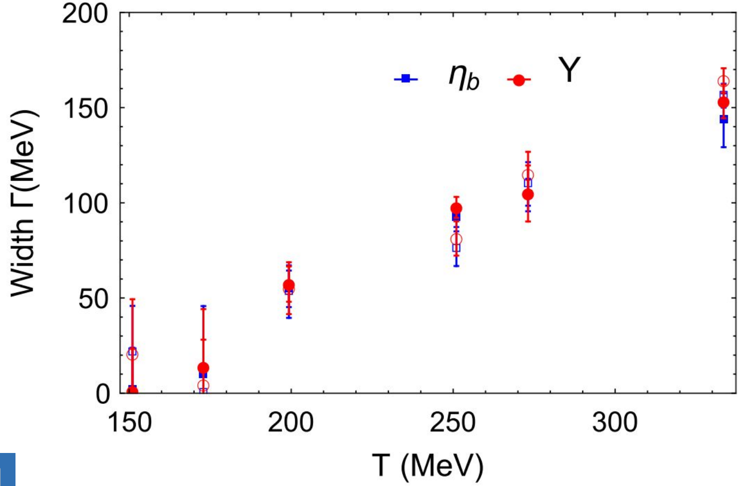
no vacuum part

$$\gamma = \gamma(T=0) + \hat{\gamma} T^3$$

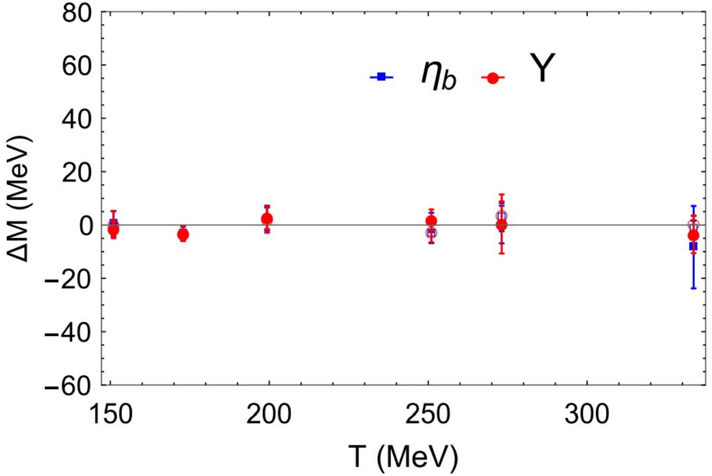
Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm

$$\Gamma \propto \hat{\kappa}$$



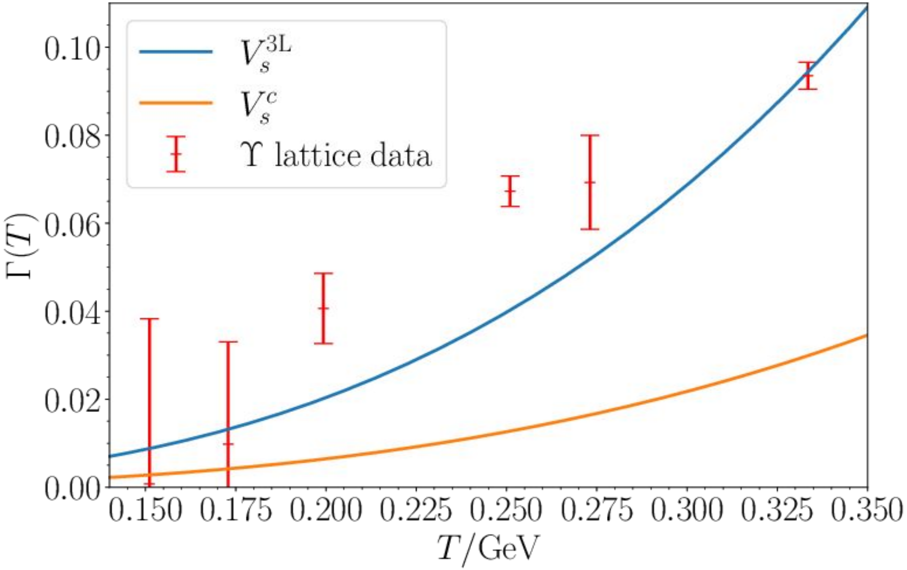
$$\delta m \propto \hat{\gamma} \Rightarrow \hat{\gamma} = 0$$

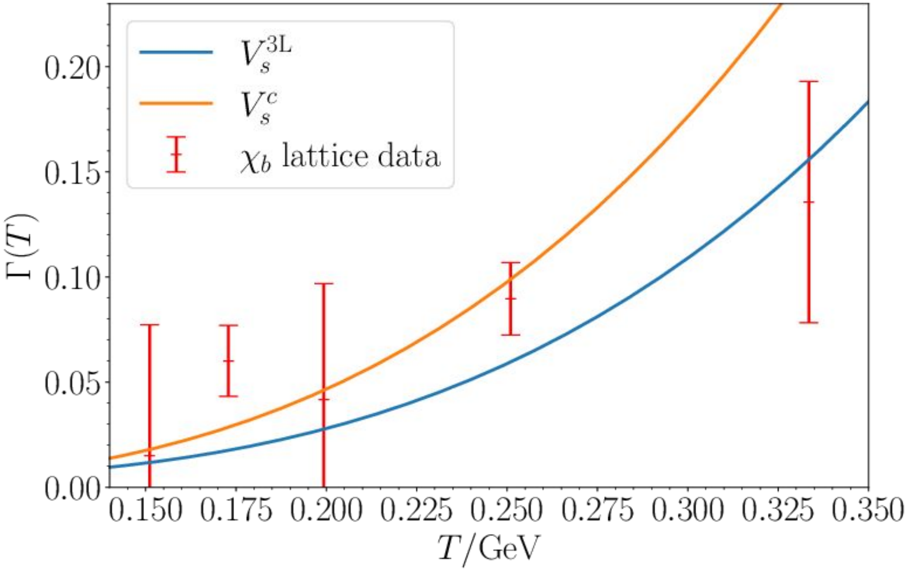


R. Larsen, S. Meinel, S. Mukherjee, and P. Petreczky, Physical Review D 100 (2019)

Determination of transport coefficients

- Obtain $\hat{\kappa}$ from fits to $1S$ and $1P$ data and average





Coulomb: $\hat{\kappa} = 0.33 \pm 0.04$

New potential: $\hat{\kappa} = 1.93 \pm 0.16$

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm

$$\kappa = \hat{\kappa} T^3$$

no vacuum part

$$\gamma = \gamma(T=0) + \hat{\gamma} T^3$$

- Assume simple model for the vacuum part $\gamma(T=0)$

$$\langle E^a(t)\Omega(t,0)^{ab}E^b(0)\rangle = \langle E^2(0)\rangle \mathrm{e}^{-i\Lambda_E t}$$

G. S. Bali and A. Pineda, Physical Review D 69 (2004)

$$\langle g^2 E^2(0)\rangle = -0.2\mathrm{GeV}^4 \qquad \Lambda_E = 1.25\mathrm{GeV}$$

S. Narison, Qcd parameters and sm-high precisions from $e_+e_-\to$hadrons :Summary(2023),2309.05342

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm

$$\kappa = \hat{\kappa} T^3$$

no vacuum part

$$\gamma = \gamma(T=0) + \hat{\gamma} T^3$$

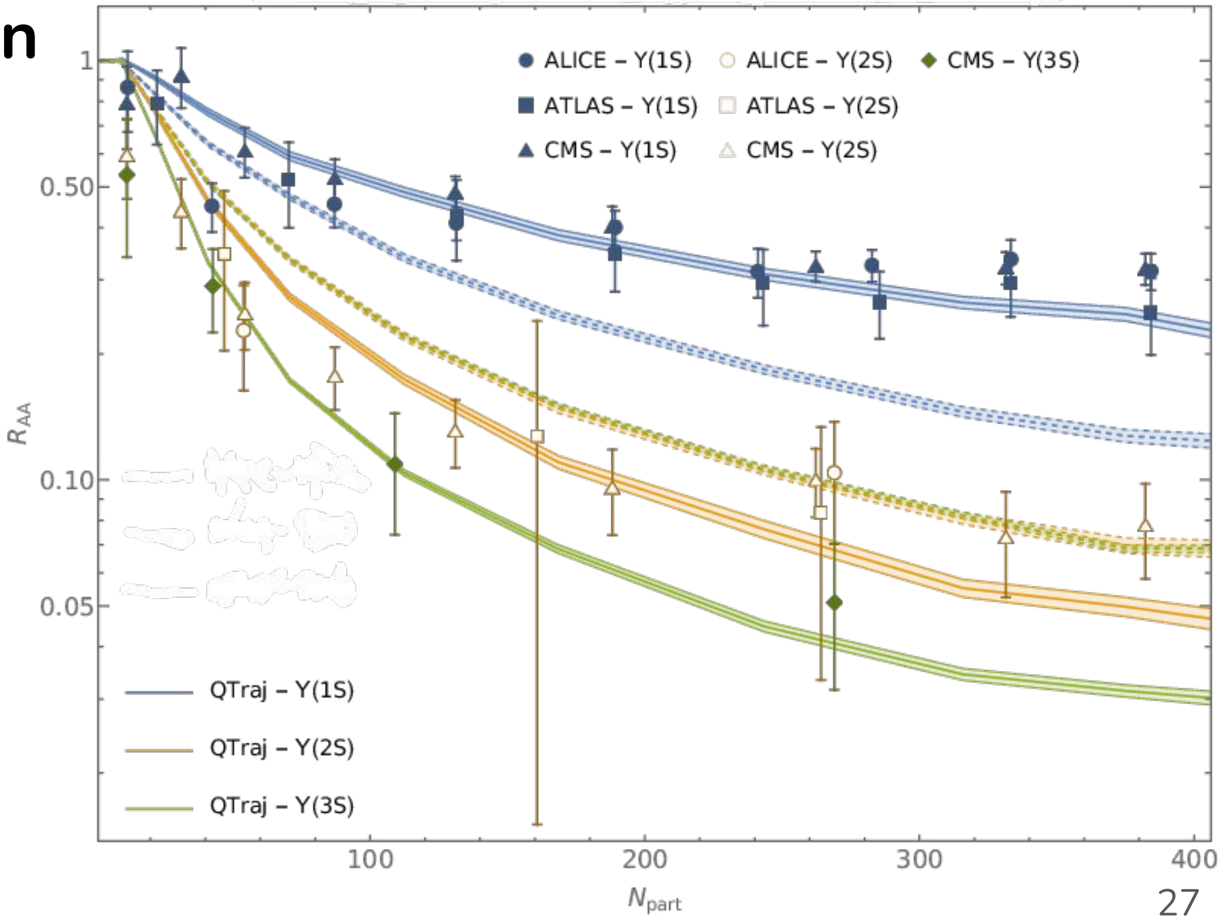
- Assume simple model for the vacuum part $\gamma(T=0)$ leading to

$$\gamma(T=0) = 0.017 \mathrm{GeV}^3$$

Nuclear modification factor results

- New potential can describe the experimental data
- Coulomb potential with $\hat{\kappa} = 0.33 \pm 0.04$ can not describe the data
- Comparison with previous LO results

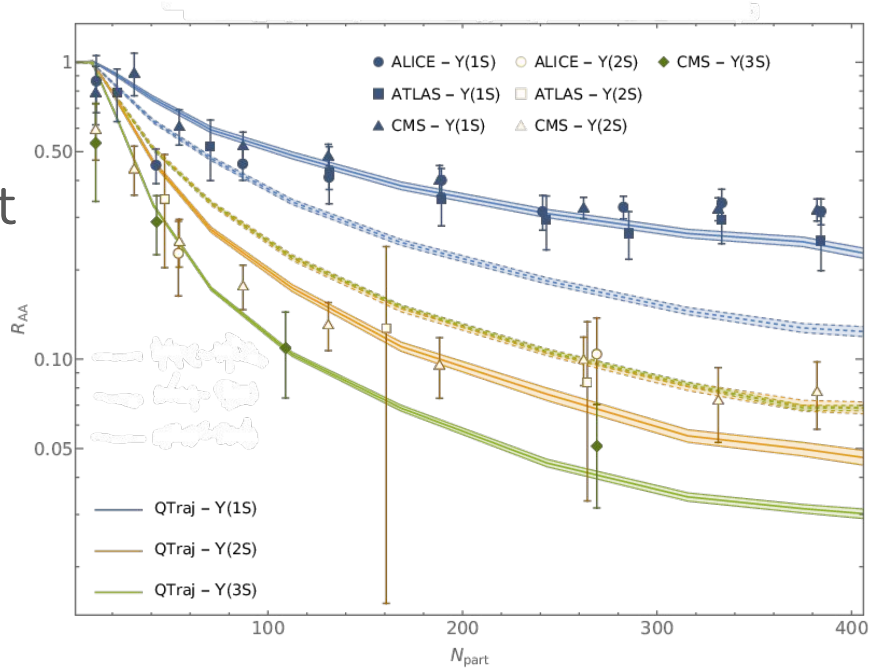
dashed: Coulomb solid: New potential



Summary and outlook

- We implemented a new potential which gives a realistic spectrum
- We extracted transport coefficient values from lattice data
- Our results agree well with the experimental data

- Future: Extend analysis to NLO description in E/T expansion



Thanks!

Backup slides

Quantum Trajectories

$$U(\Theta) = 1 - iH_{\text{eff}}\delta t$$

$$\psi_0$$

1. Evolve state $|\psi(t)\rangle$ with $U(\Theta)$

$$|\psi(t+\delta t)\rangle = U(\Theta)|\psi(t)\rangle$$

2. Compute norm

$$\langle\psi(t+\delta t)|\psi(t+\delta t)\rangle = 1 - \delta p(\Theta) < 1$$

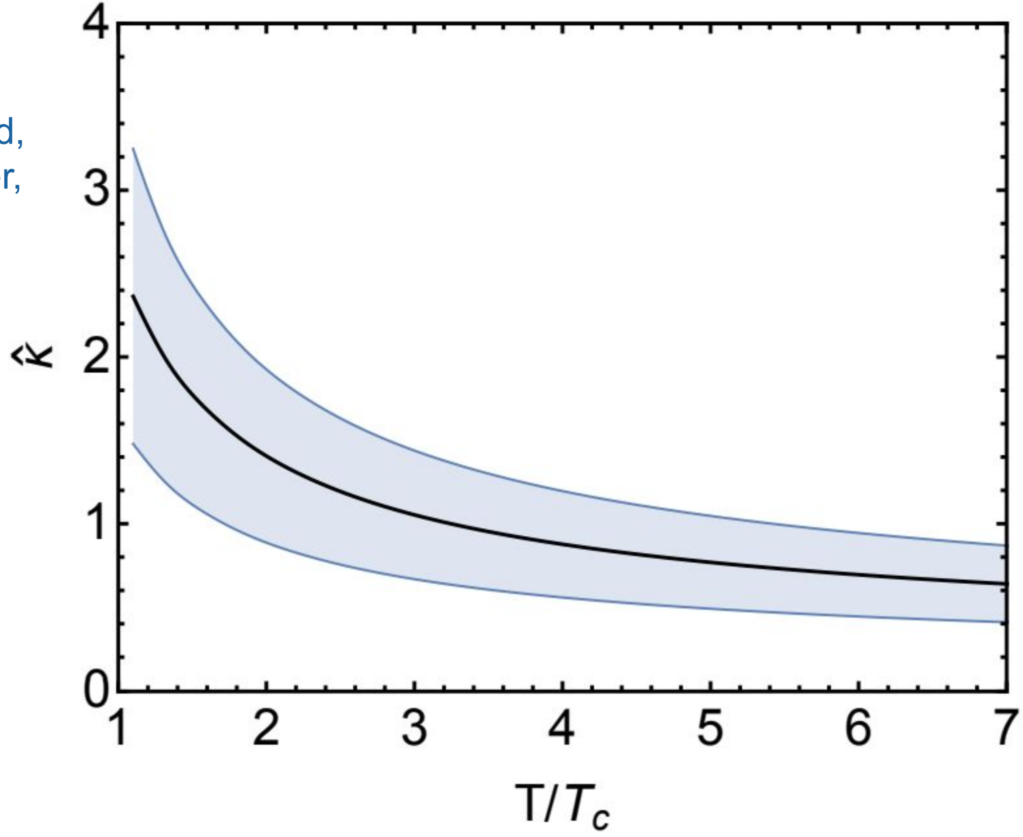
3. Apply jump operator $C(\Theta)$ with probability $\delta p(\Theta)$

$$|\psi(t+\delta t)\rangle = C(\Theta)|\psi(t)\rangle$$

4. Normalize $|\psi(t+\delta t)\rangle$

Non perturbative correction

N. Brambilla, M. A. Escobedo, M. Strickland, A. Vairo, P. Vander Griend, and J. H. Weber, JHEP 05, 136 (2021), 2012.01240



Heavy quark diffusion coefficient

$$V_s^{\text{non-pert}}(r) = -i\frac{g^2 T_F}{3N_c} r^2 \int_0^\infty dt \langle E^a(t)\Omega(t,0)^{ab}E^b(0)\rangle$$

$$\gamma = \frac{g^2}{3N_c} \mathrm{Im} \int_0^\infty dt \, \langle E^a(t)\Omega(t,0)^{ab}E^b(0)\rangle$$

In medium width

- Width given by collapse operators

$$\Gamma = \sum_n C_n^\dagger C_n$$

- At LO in E/T

$$\Gamma = \hat{\kappa} T^3 r^2$$