

Bottomonium suppression from the 3-loop QCD potential

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in collaboration with Nora Brambilla, Michael Strickland,
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QWG

28.02.2024



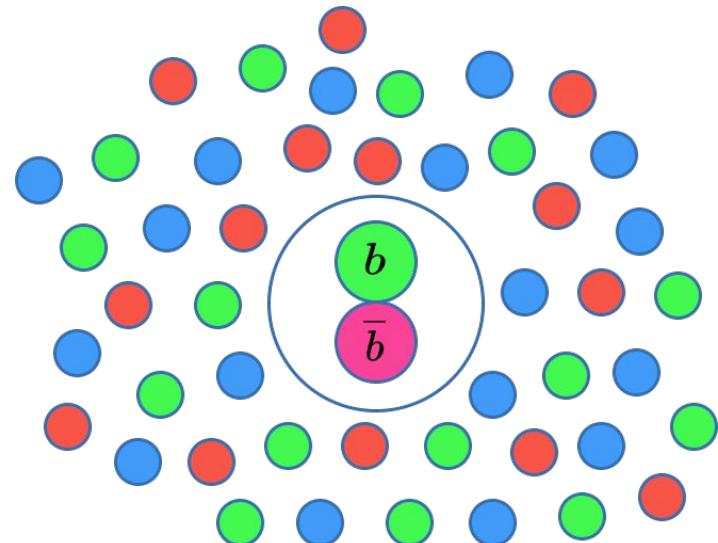
Quarkonium Suppression from first principles

T. Matsui, H. Satz, Phys. Lett. B 178 (1986) 416

- Matsui & Satz proposed Quarkonium suppression as a signal for the QGP
- Quarkonium dissolves in the QGP
- Measured Quarkonium yields are lower in HIC compared to pp collisions
- We aim to describe this phenomenon from **first principles**

Propagation through QGP

$$T \approx O(100\text{MeV})$$



Open Quantum Systems

- Quantum system not isolated
- Split into System S and Environment E

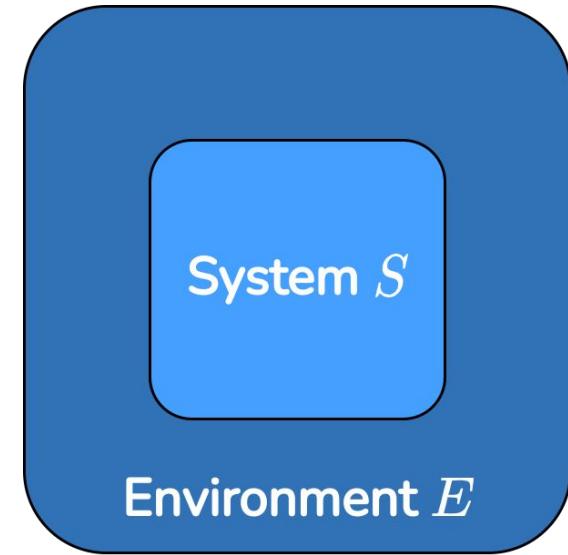
$$H = H_S \otimes I_E + I_S \otimes H_E + H_{\text{int}}$$

- Time evolution by Von-Neumann Equation

$$\frac{d}{dt} \rho = -i[H, \rho]$$

- Not interested in environmental d.o.f.: **Trace out!**

$$\rho_S = \text{Tr}_E[\rho]$$



Open Quantum Systems

- Time evolution by Von-Neumann Equation

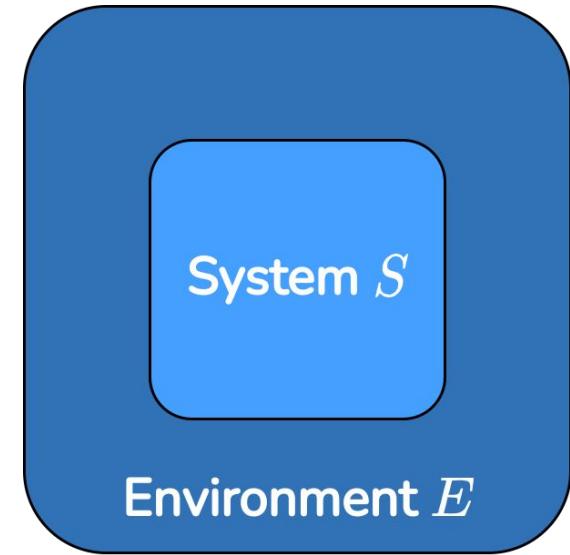
$$\frac{d}{dt}\rho = -i[H, \rho]$$

- Environmental d.o.f. not needed Trace out!

$$\rho_S = \text{Tr}_E[\rho]$$

- “Master equation” for the System: **Lindblad Equation**

$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2} \{C_n^\dagger C_n, \rho_S\} \right)$$



Open Quantum Systems

- Time evolution by Von-Neumann Equation

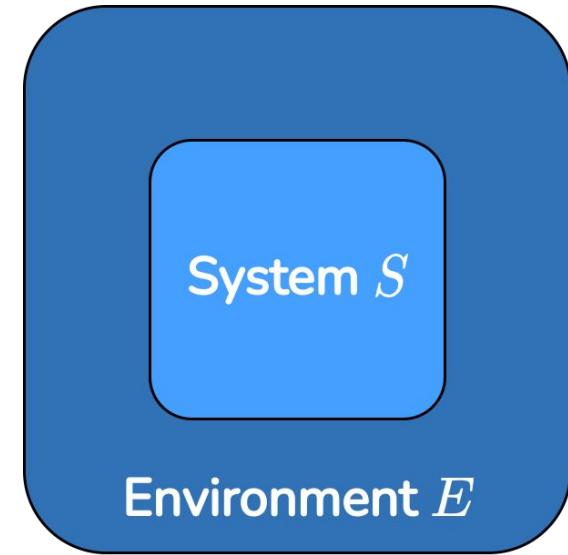
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Open Quantum Systems

- Time evolution by Von-Neumann Equation

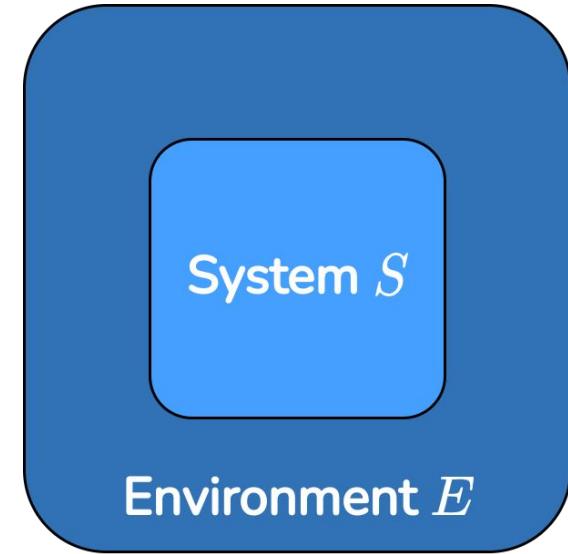
$$\frac{d}{dt}\rho = -i[H, \rho]$$

- Environmental d.o.f. not needed Trace out!

$$\rho_S = \text{Tr}_E[\rho]$$

- “Master equation” for the System: Lindblad Equation *non-unitary*

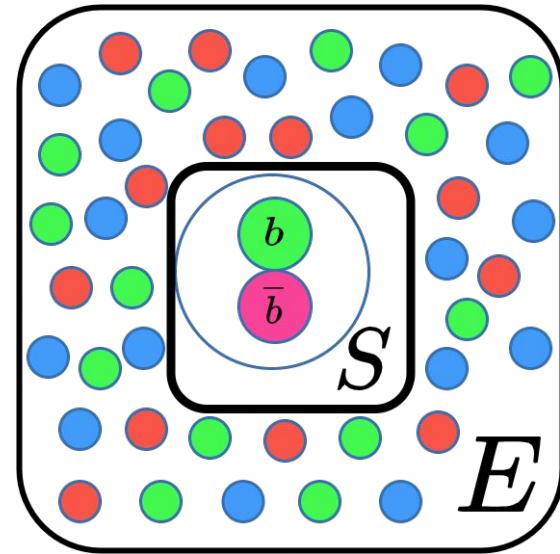
$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2} \{ C_n^\dagger C_n, \rho_S \} \right)$$



OQS for quarkonium

- Quarkonium: System S
- QGP: Environment E

Aim to describe Quarkonium Suppression by
a master equation for encoding the
interaction with the QGP



$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2} \{C_n^\dagger C_n, \rho_S\} \right)$$

EFTs for Quarkonium Suppression

- Use NREFTs to exploit hierarchy of scales

$$M \gg 1/a_0 \gg \pi T \gg E$$

- Inverse radius:

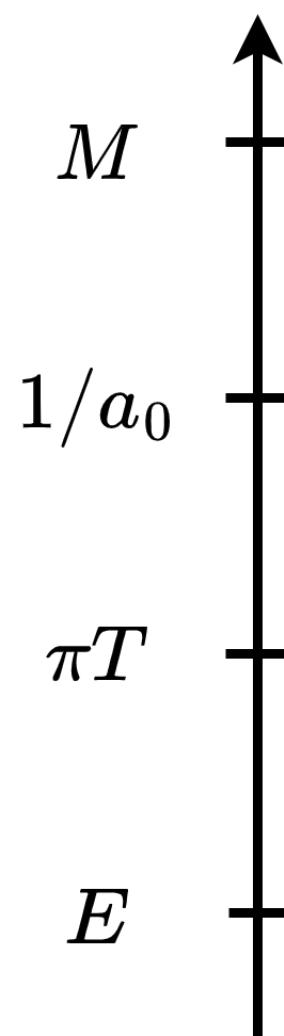
$$1/a_0 \approx 1.2 \text{ GeV}$$

- Temperature regime:

$$250 \text{ MeV} < T < 425 \text{ MeV}$$

- Binding Energy:

$$E \sim 0.4 \text{ GeV}$$

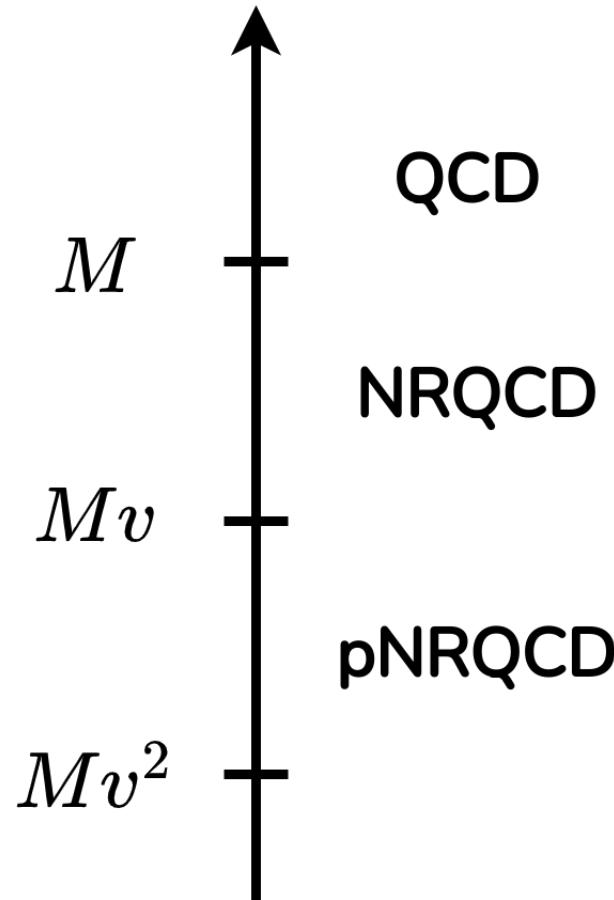


pNRQCD

N. Brambilla, A. Pineda, J. Soto, and A. Vairo,
Nuclear Physics B 566, 275 (2000)

- We use pNRQCD, an EFT from full QCD
- pNRQCD is obtained by integrating out the hard scale M and soft scale Mv
- Degrees of freedom: Singlet and octet bound states
- Using pNRQCD one can derive a master equation for the quarkonium density matrix

Brambilla, Escobedo, Soto, Vairo: Phys. Rev.
D 97 (2018) 7, 074009



pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2} \left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

$$H = \begin{pmatrix} h_s + \frac{r^2}{2}\gamma & 0 \\ 0 & h_o + \frac{N_c^2-2}{2(N_c^2-1)} \frac{r^2}{2}\gamma \end{pmatrix}$$

$$h_{s,o} = \vec{p}^2/M + V_{s,o}$$

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2} \left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

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$$h_{s,o} = \vec{p}^2/M + V_{s,o}$$

Quarkonium Potential

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2} \left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

$$H = \begin{pmatrix} h_s + \frac{r^2}{2}\gamma & 0 \\ 0 & h_o + \frac{N_c^2 - 2}{2(N_c^2 - 1)} \frac{r^2}{2}\gamma \end{pmatrix} \quad C_i^0 = \sqrt{\frac{\kappa}{N_c^2 - 1}} r_i \begin{pmatrix} 0 & 1 \\ \sqrt{N_c^2 - 1} & 0 \end{pmatrix},$$

$$C_i^1 = \sqrt{\frac{\kappa(N_c^2 - 4)}{2(N_c^2 - 1)}} r_i \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

$$h_{s,o} = \vec{p}^2/M + V_{s,o}$$

pNRQCD master equation

$$\frac{d\rho(t)}{dt} = -i[H, \rho(t)] + \sum_n \left[C_i^n \rho(t) C_i^{n\dagger} - \frac{1}{2} \left\{ C_i^{n\dagger} C_i^n, \rho(t) \right\} \right],$$

$$H = \begin{pmatrix} h_s + \frac{r^2}{2}\gamma & 0 \\ 0 & h_o + \frac{N_c^2 - 2}{2(N_c^2 - 1)} \frac{r^2}{2}\gamma \end{pmatrix}$$

↑
↑
Transport
coefficients

$$C_i^0 = \sqrt{\frac{\kappa}{N_c^2 - 1}} r_i \begin{pmatrix} 0 & 1 \\ \sqrt{N_c^2 - 1} & 0 \end{pmatrix},$$

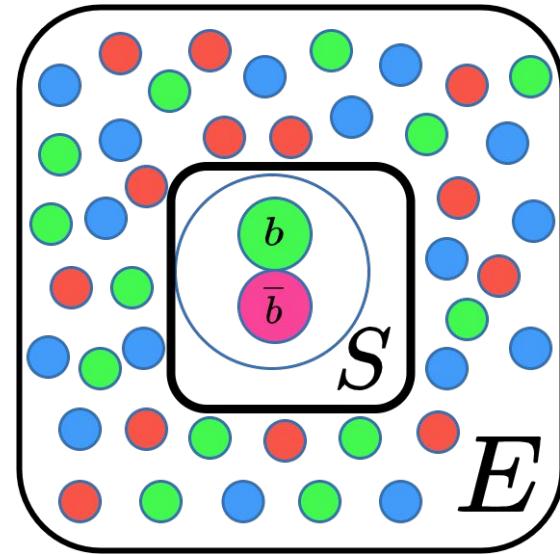
$$C_i^1 = \sqrt{\frac{\kappa(N_c^2 - 4)}{2(N_c^2 - 1)}} r_i \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

pNRQCD master equation

- Hilbert Space:
 - Singlet and octet states

$$\rho = \begin{pmatrix} \rho_s & 0 \\ 0 & \rho_o \end{pmatrix}$$

- Discretizing radial part of the wavefunction (e.g. 2048 lattice)
- Angular momentum quantum numbers
- Very large Hilbert space



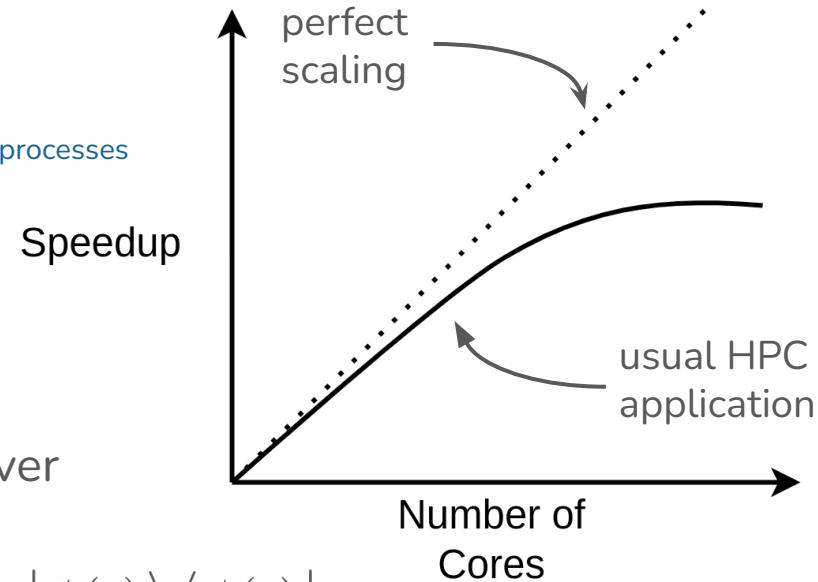
$$\frac{d\rho_S}{dt} = -i[H_S, \rho_S] + \sum_n \left(C_n \rho_S C_n^\dagger - \frac{1}{2} \{ C_n^\dagger C_n, \rho_S \} \right)$$

Quantum trajectory algorithm

J. Dalibard, Y. Castin, and K. Mølmer, Wave-function approach to dissipative processes in quantum optics, Phys. Rev. Lett. 68 (1992), pp. 580–583.

- Idea:
 - Evolve individual trajectories $|\phi(t)\rangle$ stochastically
 - Calculate observables by averaging over trajectories $\overline{\langle\phi(t)|A|\phi(t)\rangle}$
- Averaging over the density matrix $\sigma(t) = |\phi(t)\rangle\langle\phi(t)|$ restores the Lindblad equation

can evolve
to arbitrary l



Advantages:

- Evolve vector of size N_H instead N_H^2 density matrix
- Simulation of individual trajectories is **embarrassingly parallel**

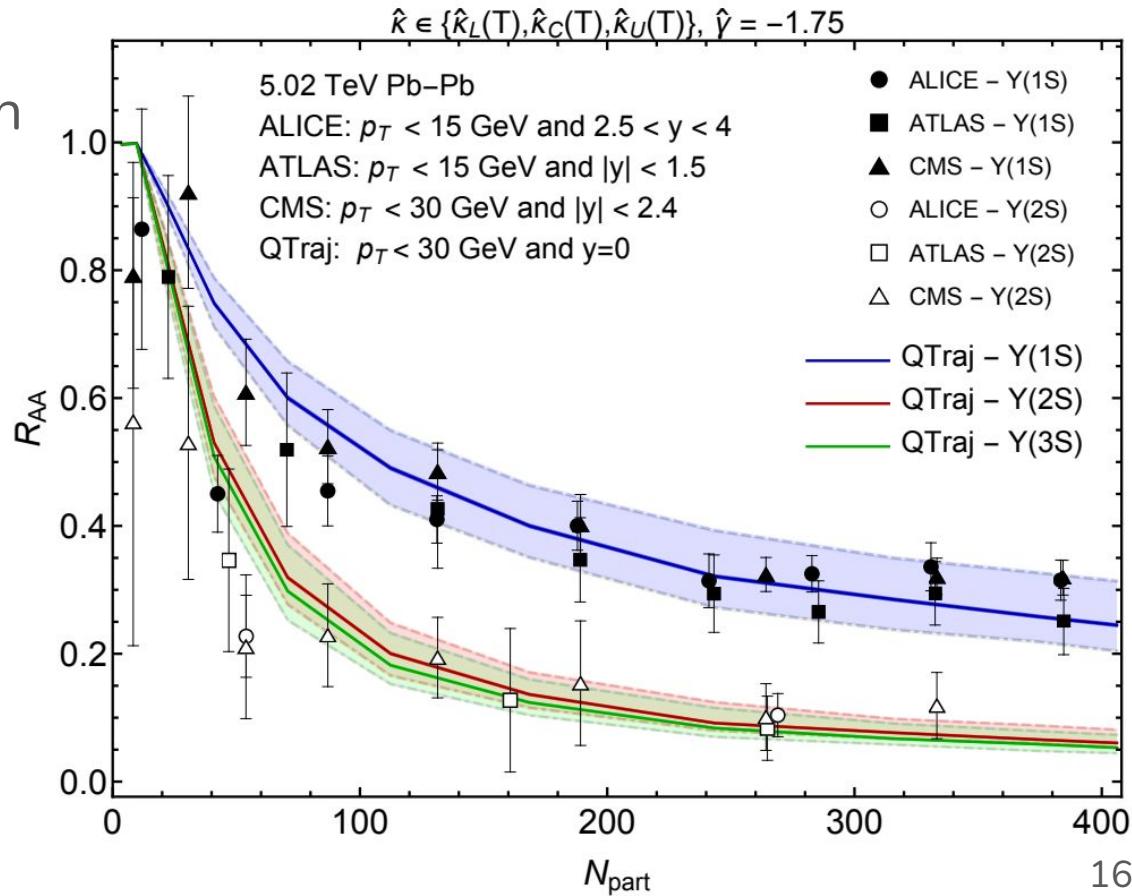
Previous work

- temperature evolution from hydrodynamics simulation

M. Alqahtani and M. Strickland, The European Physical Journal C 81 (2021)

$$\text{Survival Probability} = \frac{\langle \psi(t) | 1S \rangle}{\langle \psi(0) | 1S \rangle}$$

- Including Feed-down from PDG data



Previous work

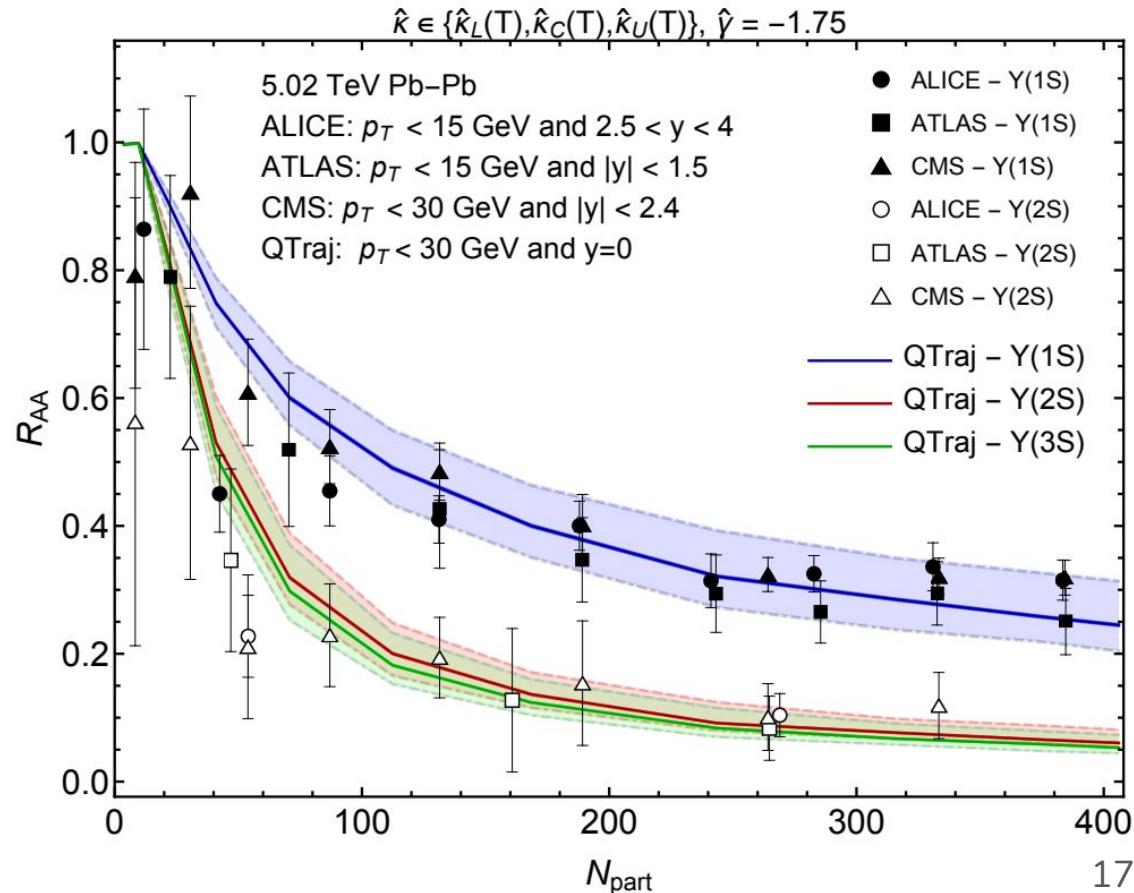
- Coulomb potential

$$V_s = -C_f \alpha_s / r$$

$$V_o = \alpha_s / (2N_c r)$$

- Temperature dependent $\hat{\kappa}$

$$\hat{\gamma} = -1.75$$



New Potential

J. Segovia, S. Steinbeißer, and A. Vairo, Physical Review D 99 (2019)

- Motivation: Implement a higher order potential with a more realistic spectrum

$$V_s^{3L+np}(r) = V_s^{3L}(r) + V_s^{np}(r)$$

$$V_s^{3L}(r) = \begin{cases} \sum_{k=0}^3 V_{s,RS'}^{(k)} \alpha_s^{k+1}(1/r) & \text{if } r < \nu_r^{-1} \\ \sum_{k=0}^3 V_{s,RS'}^{(k)} \alpha_s^{k+1}(\nu) & \text{if } r > \nu_r^{-1} \end{cases}$$

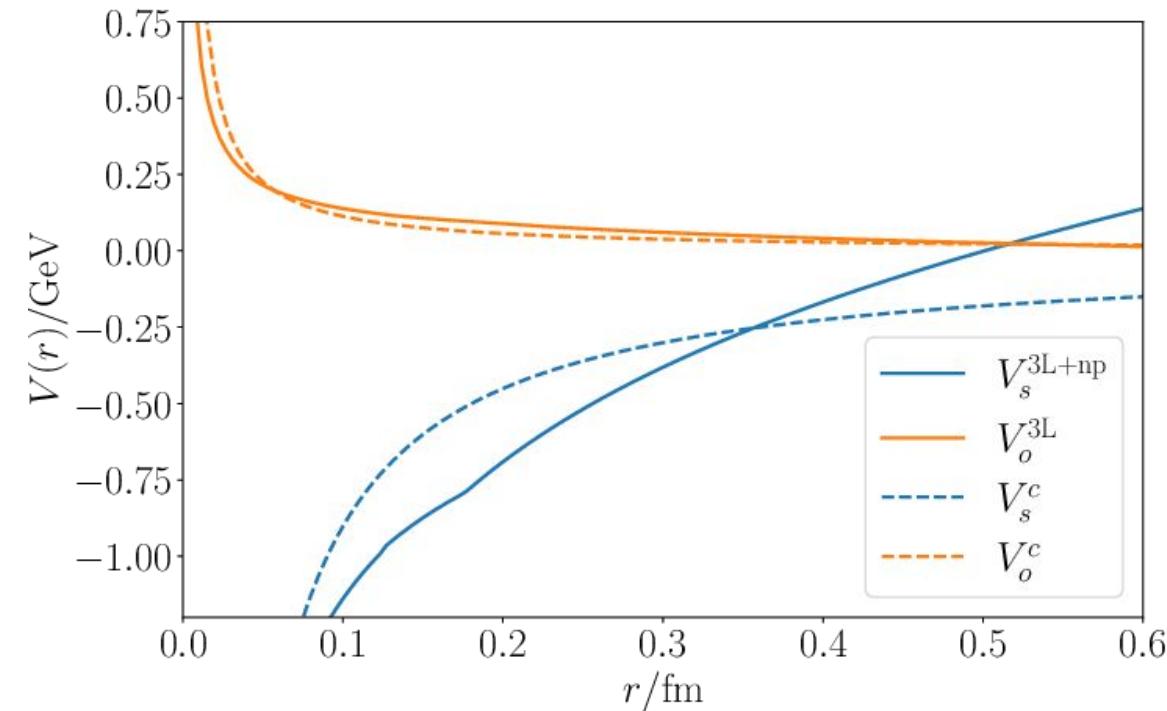
three loop pNRQCD

$$\text{Re}[V_s^{np}(r)] = \frac{\gamma}{2} r^2$$

leading non-perturbative correction

New Potential

J. Segovia, S. Steinbeißer, and A. Vairo, Physical Review D 99 (2019)



Spectrum:

	PDG	V_s^c	$V_s^{3\text{L+np}}$
$M(1S)/\text{GeV}$	9.445	9.446	9.446
$M(2S)/\text{GeV}$	10.017	9.637	10.044
$M(3S)/\text{GeV}$	10.355	9.672	10.397
$M(1P)/\text{GeV}$	9.888	9.636	9.888
$M(2P)/\text{GeV}$	10.251	9.672	10.281

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm

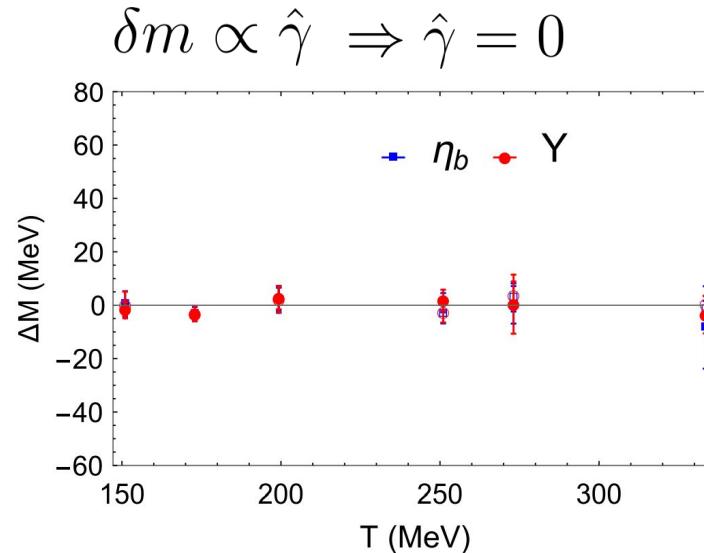
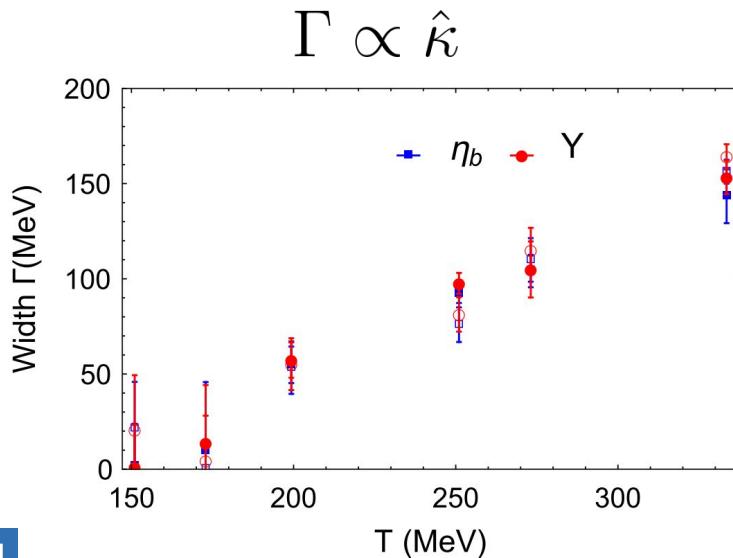

$$\kappa = \hat{\kappa} T^3$$

no vacuum
part

$$\gamma = \gamma(T = 0) + \hat{\gamma} T^3$$

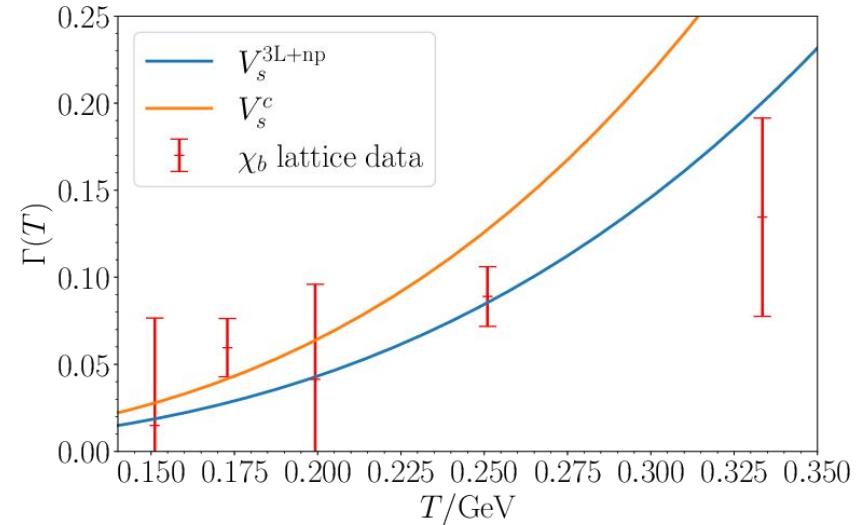
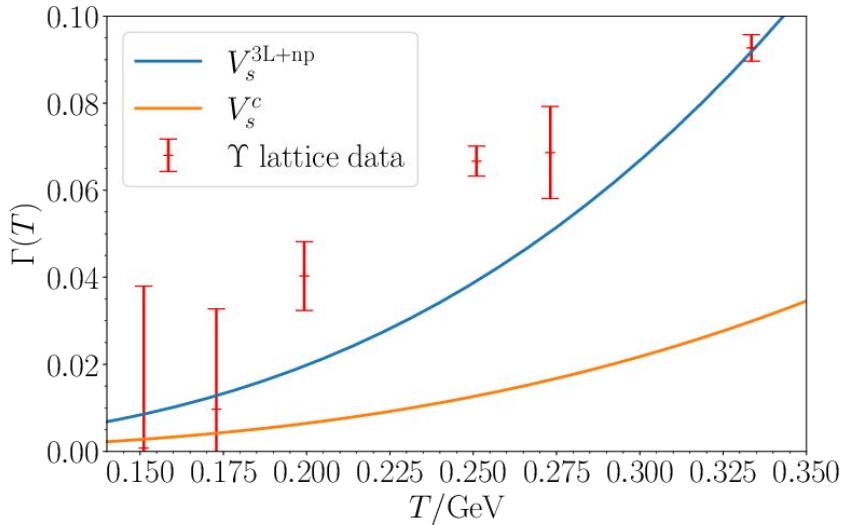
Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm



Determination of transport coefficients

- Obtain $\hat{\kappa}$ from fits to $1S$ and $1P$ data and average

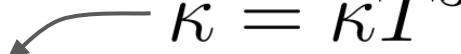


Coulomb: $\hat{\kappa} = 0.33 \pm 0.04$

New potential: $\hat{\kappa} = 1.88 \pm 0.16$

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm


no vacuum part $\kappa = \hat{\kappa}T^3$
 $\gamma = \gamma(T=0) + \hat{\gamma}T^3$

- Assume simple model for the vacuum part $\gamma(T=0)$

$$\langle E^a(t)\Omega(t,0)^{ab}E^b(0)\rangle = \langle E^2(0)\rangle e^{-i\Lambda_E t}$$

$\langle g^2 E^2(0)\rangle = -0.2 \text{ GeV}^4$

G. S. Bali and A. Pineda, Physical Review D 69 (2004)

$\Lambda_E = 1.25 \text{ GeV}$

Determination of transport coefficients

- Indirectly determine $\hat{\kappa}$ and $\hat{\gamma}$ from lattice measurements of the **in medium width** Γ and **mass shift** δm


$$\kappa = \hat{\kappa}T^3$$

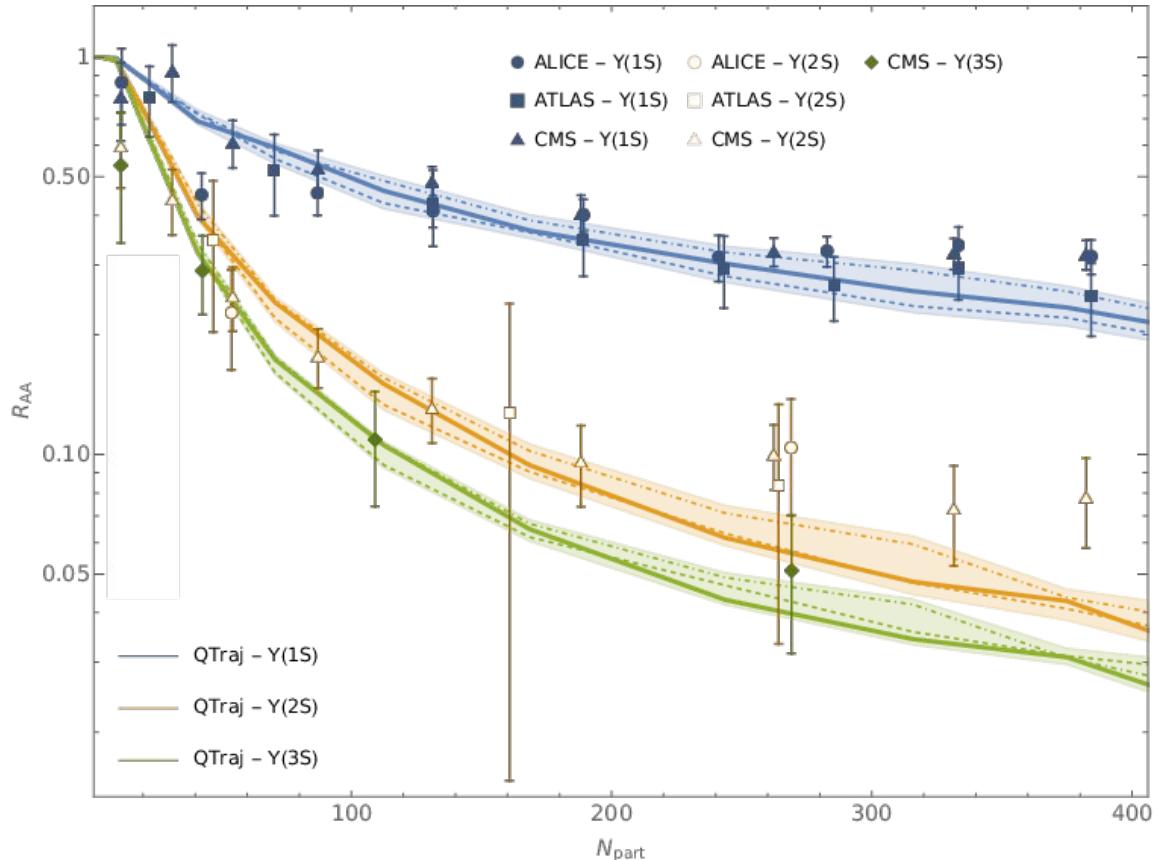
no vacuum part $\gamma = \gamma(T=0) + \hat{\gamma}T^3$

- Assume simple model for the vacuum part $\gamma(T=0)$ leading to

$$\gamma(T=0) = 0.017\text{GeV}^3$$

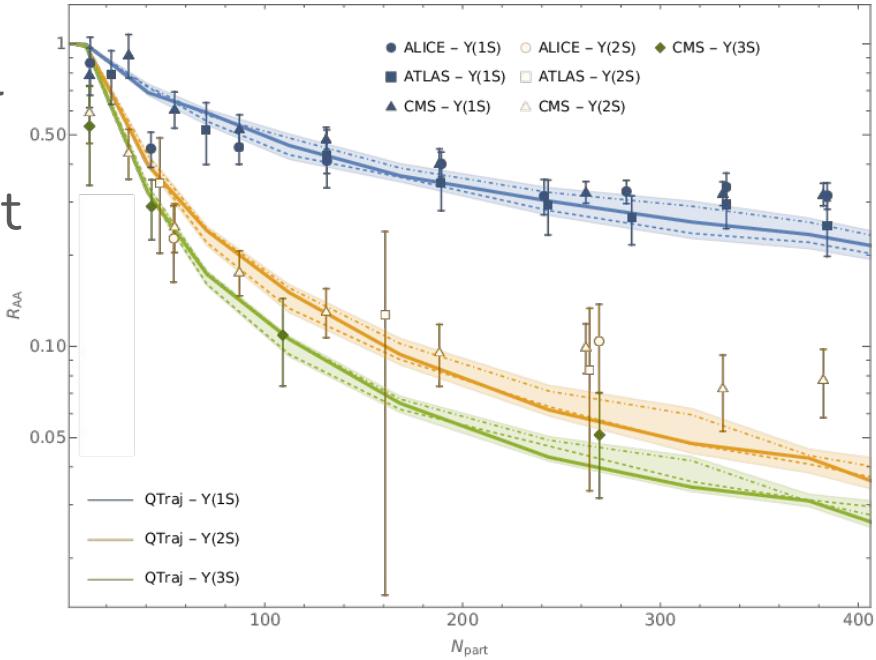
Nuclear modification factor results

- New potential can describe the experimental data
- Coulomb potential with $\hat{\kappa} = 0.33 \pm 0.04$ can not describe the data



Summary and outlook

- We implemented a new potential which gives a realistic spectrum
- We extracted transport coefficient values from lattice data
- Our results agree well with the experimental data
- Future: Extend analysis to NLO description in $E/(\pi T)$ expansion



Thanks!

Backup slides

Quantum Trajectories

$$U(\Theta) = 1 - iH_{\text{eff}}\delta t$$

ψ_0
↓

1. Evolve state $|\psi(t)\rangle$ with $U(\Theta)$

$$|\psi(t + \delta t)\rangle = U(\Theta)|\psi(t)\rangle$$

2. Compute norm

$$\langle\psi(t + \delta t)|\psi(t + \delta t)\rangle = 1 - \delta p(\Theta) < 1$$

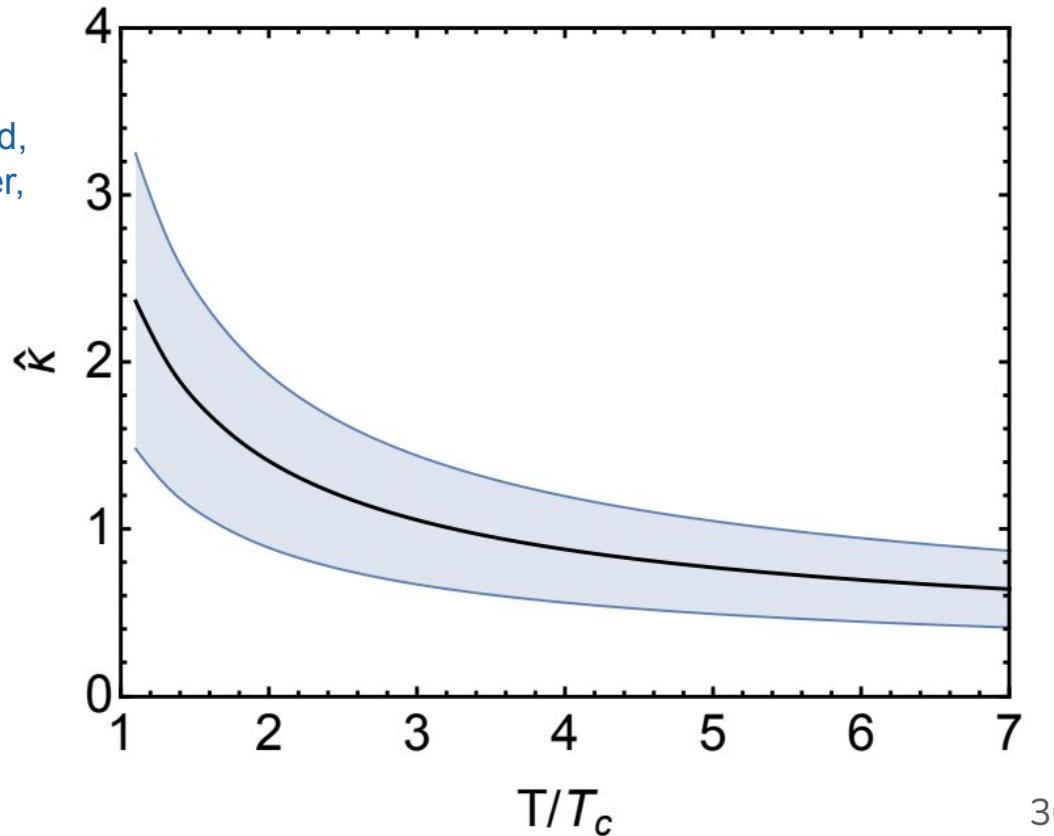
3. Apply jump operator $C(\Theta)$ with probability $\delta p(\Theta)$

$$|\psi(t + \delta t)\rangle = C(\Theta)|\psi(t)\rangle$$

4. Normalize $|\psi(t + \delta t)\rangle$

Non perturbative correction

N. Brambilla, M. A. Escobedo, M. Strickland,
A. Vairo, P. Vander Griend, and J. H. Weber,
JHEP 05, 136 (2021), 2012.01240



Heavy quark diffusion coefficient

$$V_s^{\text{non-pert}}(r) = -i \frac{g^2 T_F}{3N_c} r^2 \int_0^\infty dt \langle E^a(t) \Omega(t, 0)^{ab} E^b(0) \rangle$$

$$\gamma = \frac{g^2}{3N_c} \text{Im} \int_0^\infty dt \langle E^a(t) \Omega(t, 0)^{ab} E^b(0) \rangle$$

In medium width

- Width given by collapse operators

$$\Gamma = \sum_n C_n^\dagger C_n$$

- At LO in E/T

$$\Gamma = \hat{\kappa} T^3 r^2$$

pNRQCD master equation

Brambilla, Escobedo, Soto, Vairo: Phys. Rev. D 97 (2018) 7, 074009

- In general the master Eq. is not of Lindblad form
- Simplify using hierarchy of scales $T \gg E$

$$A_i^{uv} = \frac{g^2}{6N_c} \int_0^\infty ds e^{-ih_u s} r_i e^{ih_v s} \left\langle \tilde{E}_j^a(0, \vec{0}) \tilde{E}_j^a(s, \vec{0}) \right\rangle$$

- Expand exponentials in E/T
- At LO in E/T we get

$$\begin{aligned} A_i^{uv} &= \frac{g^2}{6N_c} \int_0^\infty ds r_i \left\langle \tilde{E}_j^a(0, \vec{0}) \tilde{E}_j^a(s, \vec{0}) \right\rangle \\ &= \frac{r_i}{2} (\kappa - i\gamma) \end{aligned}$$

Transport
coefficients