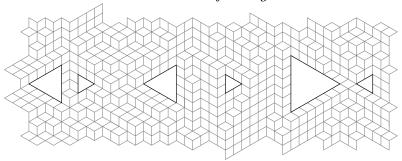
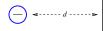
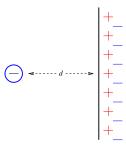
$Symmetric\ rhombus\ tilings\ of\ holey\ hexagons\ and$ $the\ method\ of\ images.$

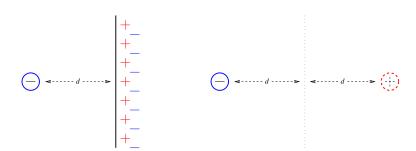


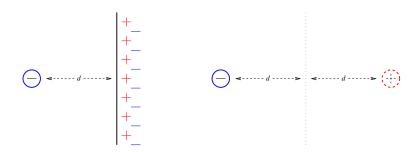
Tomack Gilmore Universität Wien







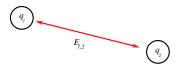


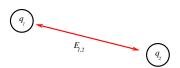


$$E_{\ominus,|} = \frac{1}{2} E_{\ominus,\Theta}$$







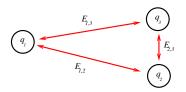


Coulomb's Law

The electrostatic energy of the system consisting of two point charges with signed magnitudes q_1 and q_2 is

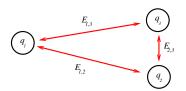
$$E_{1,2} = k_e \frac{q_1 q_2}{d(q_1, q_2)},$$

where k_e denotes Coulomb's constant and $d(q_1, q_2)$ is the Euclidean distance between the two charges.



The energy of a system of three point charges is given by

$$E_{1,2} + E_{1,3} + E_{2,3} = k_e \left(\frac{q_1 q_2}{d(q_1, q_2)} + \frac{q_1 q_3}{d(q_1, q_3)} + \frac{q_2 q_3}{d(q_2, q_3)} \right).$$

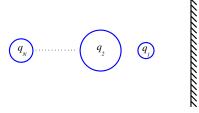


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More generally, the electrostatic energy of a system of point charges q_1, \ldots, q_n is given by

$$k_e \sum_{1 \le i \le j \le n} \frac{q_i q_j}{d(q_i, q_j)}.$$



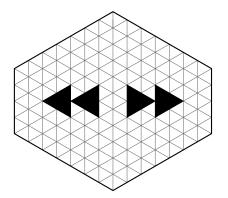
$$E_{\{q_1,\dots,q_N\},|}$$

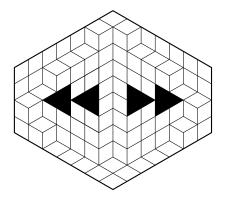


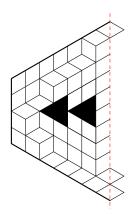
$$E_{\{q_1,\dots,q_N\},|} = \frac{1}{2} E_{\{q_1,\dots,q_N\},\{q'_1,\dots,q'_N\}}$$

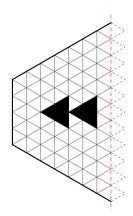


$$\begin{split} E_{\{q_1,\dots,q_N\},|} &= \frac{1}{2} E_{\{q_1,\dots,q_N\},\{q'_1,\dots,q'_N\}} \\ &= \frac{k_e}{2} \left(\sum_{1 \le i < j \le N} \frac{q_i q_j}{d(q_i,q_j)} + \sum_{1 \le i < j \le N} \frac{q'_i q'_j}{d(q'_i,q'_j)} \right) \\ &- \sum_{1 \le i,j \le N} \frac{|q_i q'_j|}{d(q_i,q'_j)} \right). \end{split}$$

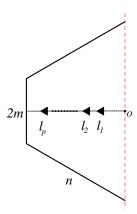






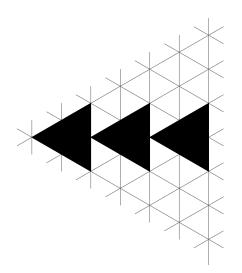


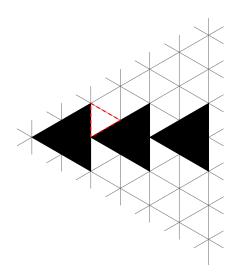
$$V_{7.4}^{\{-1,-3\}}$$

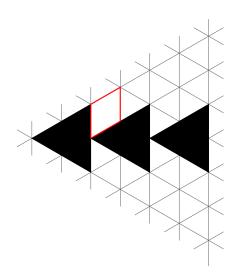


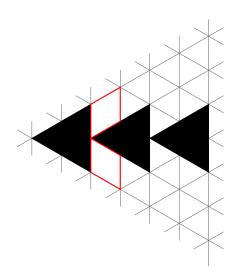
 $V_{n,2m}^L$

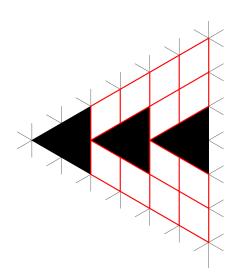
(here $L = \{l_1, \dots, l_p\}$ indexes the left pointing triangular holes by their vertical lattice distance from the origin).

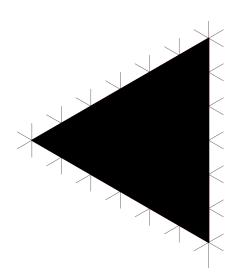












Theorem (TG 2016)

$$M(V_{n,2m}^{L}) = \left(\prod_{i=1}^{n} \frac{2i+2m-1}{2i-1} \prod_{1 \le i < j \le n} \frac{i+j+2m-1}{i+j-1}\right) \cdot \det \widehat{E}_{R,L}$$

where $R = \{-l_1, \ldots, -l_{|L|}\}$, $\widehat{E}_{R,L} = (\hat{e}_{i,j})_{1 \leq i,j \leq |L|}$ with (i,j)-entries given by

$$\hat{e}_{i,j} = \frac{\Gamma(m + \frac{1}{2})\Gamma(\frac{n}{2} - \frac{l_i}{2} + \frac{3}{2})\Gamma(m + n + 1)\Gamma(\frac{n}{2} + \frac{r_j}{2} + \frac{3}{2})}{2^{l_i - r_j - 2}\pi\Gamma(m)\Gamma(\frac{n}{2} - \frac{l_i}{2} + 2)\Gamma(m + n + \frac{1}{2})\Gamma(\frac{n}{2} + \frac{r_j}{2} + 2)} \times \sum_{s=0}^{\infty} \frac{(2 + \frac{r_j}{2} - \frac{l_i}{2})_s(\frac{1}{2})_s(m + n + 1)_s(1 - m)_s}{(\frac{n}{2} - \frac{l_i}{2} + 2)_s(\frac{n}{2} + \frac{r_j}{2} + 2)_s(\frac{3}{2})_s(s!)}$$

(here
$$\Gamma(x) = (x-1)!$$
 and $(x)_y = \frac{\Gamma(x+y)}{\Gamma(x)}$ for $x \in \mathbb{Q}, y \in \mathbb{Z}$).

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As the distances between the sets of holes grows large,

$$\omega(\xi; R, L) \sim \det \left(\frac{(\xi(\xi+2))^{-1/2}}{\pi(r_j - l_i)} \left(\frac{2}{\xi+1} \right)^{r_j - l_i + 2} \right)_{1 \le i, j \le |L|}$$

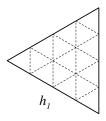
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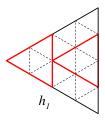
and so for $\xi = 1$,

$$\omega(1; R, L) \sim \left(\frac{1}{2\pi}\right)^{|L|} \frac{\prod_{i=2}^{|L|} \prod_{j=1}^{i-1} d(r_i, r_j) d(l_i, l_j)}{\prod_{1 \le i, j \le |L|} d(r_i, l_j)}.$$



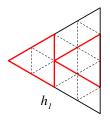


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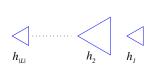




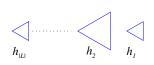
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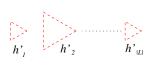
$$\omega(1; R, L) \sim \prod_{h \in \mathcal{H}} C_h \prod_{1 \le j < i \le |\mathcal{H}|} d(h_i, h_j)^{\frac{1}{4}q(h_i)q(h_j)},$$

where \mathcal{H} is the set of holes induced by the holes of side length two indexed by L and R.

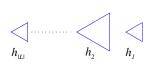


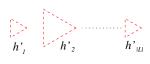
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