PDE HOMEWORK 4

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Exercise 1.

Proof.

(1): Since U is bounded and u continuous so we can use DCT to pass limit, thus

$$e'(t) \stackrel{DCT}{=} \int_{U} \partial_{t} u^{2}(x, t) dx = 2 \int_{U} u \partial_{t} u = 2 \int_{U} u \nabla \cdot (a(x) \nabla u)$$

$$\stackrel{IBP}{=} 2 \int_{\partial U} u(a(x) \nabla u) - 2 \int_{U} a(x) |\nabla u|^{2} = -2 \int_{U} a |\nabla u|^{2}$$

since *u* vanishes on the boundary. One more derivative is similar:

$$\begin{split} e''(t) &\stackrel{above}{=} \partial_t \left[2 \int_U u \nabla \cdot (a(x) \nabla u) \right] \stackrel{DCT}{=} 2 \int_U \partial_t u \nabla \cdot (a(x) \nabla u) + u \nabla \cdot (a(x) \nabla u_t) dx \\ &= 2 \int_U (\nabla \cdot a \nabla u)^2 dx + 2 \int_U u \nabla \cdot (a \nabla (\nabla \cdot a \nabla u)) dx \\ \stackrel{IBP}{=} 2 \int_U (\nabla \cdot a \nabla u)^2 dx - 2 \int_U a \nabla u \cdot [\nabla (\nabla \cdot a \nabla u)] dx + 0 \\ \stackrel{IBP}{=} 2 \int_U (\nabla \cdot a \nabla u)^2 dx + 2 \int_U (\nabla \cdot a \nabla u)^2 dx - \sum_{i=1}^n 2 \int_{\partial U} a \partial_i u (\nabla \cdot a \nabla u) \\ &= 4 \int_U (\nabla \cdot a \nabla u)^2 dx - \sum_{i=1}^n 2 \int_{\partial U} a \partial_i u \partial_t u dx = 4 \int_U (\nabla \cdot a \nabla u)^2 dx \end{split}$$

because u(x, t) = 0 on $\mathbb{R}^n \times [0, T]$ implies $\partial_t u = 0$ on ∂U .

(2):

$$(e'(t))^{2} = 4\left(\int_{U} u \nabla \cdot (a(x)\nabla u)\right)^{2} = 4\left(\int_{U} u dx \int_{U} \nabla \cdot (a(x)\nabla u) dx\right)^{2}$$

$$\stackrel{C.S.}{\leq} 4\left(\int_{U} u dx\right)^{2} \left(\int_{U} \nabla \cdot (a(x)\nabla u) dx\right)^{2} \leq 4\left(\int_{U} \int_{U} u^{2} dx\right) \cdot \left(\int_{U} (\nabla \cdot a\nabla u)^{2} dx\right)$$

$$= e(t)e''(t)$$

(3):

Just take derivative and by (2)

$$f''(t) = \frac{e(t)e''(t) - [e'(t)]^2}{\rho^2} \ge 0$$

so f is convex.

Now we deduce the formula. Convexity gives (since everywhere defined)

$$f((1-\tau)t_1 + \tau t_2) \le (1-\tau)f(t_1) + \tau f(t_2)$$

now take exponential on both sides we have

$$e\left((1-\tau)t_1+\tau t_2\right) \leq \left[e(t_1)\right]^{1-\tau}\cdot \left[e(t_2)\right]^{\tau}$$

Let $\tau = \frac{t}{T}$ and $t_1 = 0, t_2 = T$, then this gives

$$e(t) \le e(T)^{1-t/T} e(0)^{t/T} = M^{1-t/T} \varepsilon^{t/T}$$

(4):

 $\varepsilon > 0$:

Either by what is the energy or the fact that there is a square so that it is non-negative. So at the place that it is not strictly positive, it is 0. But then since $e(T) = \varepsilon > 0$, by smoothness we know that at some point e' > 0, which is impossible by our computation in (1). Thus Even if we do not assume e(t) > 0 on the whole domain, it is implicitly implied, so proof in (3) works.

 $\varepsilon = 0$:

This is a little bit trickier but fine. Since $e' \le 0$, if e reaches 0 at some time t, then it stays 0 after that point. Thus, we can define (due to completeness of \mathbb{R})

$$s := \sup\{t | e(t) > 0\}.$$

We know by continuity that e(s) = 0 and $e(s - \delta) > 0$ for any δ . Now the only problem with our above argument is that f is not defined on [s, T]. But this does not matter since we can use approximation, i.e. let's say

$$\zeta = e(s - \delta) > 0$$

then we have that f is defined everywhere in $[0, s - \delta]$, so the above conclusion holds due to convexity:

$$e(t) \le e(T)^{1-t/(s-\delta)} e(0)^{t/(s-\delta)} = M^{1-t/(s-\delta)} \zeta^{t/(s-\delta)} \le C \zeta^{2t/s}$$

for δ small enough. Now we take $\delta \to 0$, then $\zeta \to 0$ by continuity, so we know for all $t \in (0, s)$, we have the property

$$e(t) \le C\zeta^{2t/s} \to 0$$

thus e(t) = 0 everywhere on (0, T]. But e is continuous so e(0) = M = 0. Thus

$$0 = e(t) \le e(T)^{1-t/T} e(0)^{t/T} = M^{1-t/T} \varepsilon^{t/T} = 0.$$

Exercise 2.

Proof.

(1): Using chain rule we have

$$\dot{e}(t) \stackrel{DCT}{=} \frac{1}{2} \int_{B(x_0), \alpha(t_0 - t)} \partial_t \left(\frac{1}{c^2(x)} (\partial_t u)^2 + |\nabla u|^2 \right) dx$$
$$- \frac{\alpha}{2} \int_{\partial B(x_0), \alpha(t_0 - t)} \left(\frac{1}{c^2(x)} (\partial_t u)^2 + |\nabla u|^2 \right) d\sigma$$

and the first term can be simplified by

$$\begin{split} &\frac{1}{2} \int_{B(x_0),\alpha(t_0-t)} \partial_t \left(\frac{1}{c^2(x)} (\partial_t u)^2 + |\nabla u|^2 \right) dx \\ &= \frac{1}{2} \int_{B(x_0),\alpha(t_0-t)} \frac{1}{c^2(x)} (2\partial_t u \cdot \partial_t^2 u) + 2\nabla u_t \cdot \nabla u dx \\ &= \int_{B(x_0),\alpha(t_0-t)} \frac{1}{c^2(x)} (\partial_t u \cdot c^2(x) \Delta u) + \nabla u_t \cdot \nabla u dx \\ &= \int_{B(x_0),\alpha(t_0-t)} \partial_t u \nabla \cdot \nabla u + \nabla u_t \cdot \nabla u dx \\ &= \int_{B(x_0),\alpha(t_0-t)} \nabla u_t \cdot \nabla u dx + \int_{\partial B(x_0),\alpha(t_0-t)} \partial_t u \nabla u d\sigma - \int_{B(x_0),\alpha(t_0-t)} \nabla u_t \cdot \nabla u dx \\ &= \int_{\partial B(x_0),\alpha(t_0-t)} u_t \frac{\partial u}{\partial v} d\sigma \end{split}$$

so putting things together we have

$$\dot{e}(t) = \int_{\partial B(x_0), \alpha(t_0 - t)} u_t \frac{\partial u}{\partial \nu} - \frac{\alpha}{2} \left(\frac{1}{c^2(x)} (\partial_t u)^2 + |\nabla u|^2 \right) d\sigma.$$

(2):

We use Cauchy to get

$$|u_t \partial_n u| \stackrel{C.S.}{\leq} \frac{1}{2} \left(u_t^2 + |\partial_n u|^2 \right) \leq \frac{1}{2} \left(u_t^2 + |\nabla u|^2 \right)$$

SO

$$\dot{e}(t) = \int_{\partial B(x_0)} \frac{c^2(x) - \alpha}{2c^2(x)} (\partial_t u)^2 + \frac{1 - \alpha}{2} |\nabla u|^2 d\sigma$$

and so as long as $\alpha \ge 1$ and $\alpha \ge c^2(x) \ge c_0^2$ then we have the desired result, thus a reasonable range is

$$\alpha \ge \max\{1, c_0^2\}.$$

(3):

When t = 0 we know that $\partial_t = 0$ and since u = 0 is constant we have $\nabla u = 0$. Thus e(0) = 0.

If we have e(t) = 0 for $0 \le t \le t_0$ then we know $\partial_t u = 0$ and $\nabla u = 0$ in the cone so there is no direction in the spacetime along which u will change in the cone, so u is constant there and hence is 0.

Thus this requires exactly that $\dot{e}(t) = 0$ for every t in the domain. If there is a point $z \in C$ the cone such that $c^2(z) - \alpha > 0$ then we can choose c(x) to be the constant function $c(x) = c_0$ then we see that if we just choose to go only along the t direction and not at all the t direction locally at t, then Cauchy Schwartz is attained at t and we get a strict positive derivative, which by continuity result in a strict change of t. Thus, we will need

$$c^2(x) - \alpha \le 0$$

everywhere in $\partial B(x_0)$, $\alpha(t_0 - t)$ for all t (even at endpoints due to continuity). So if we let

$$s := \inf\{c(x) : x \in C\}$$

then $\alpha = \max\{1, s^2\}$ is minimal.