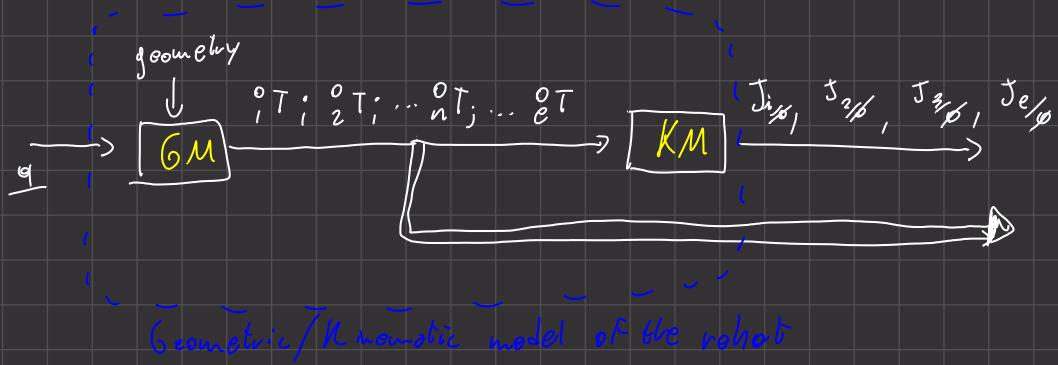


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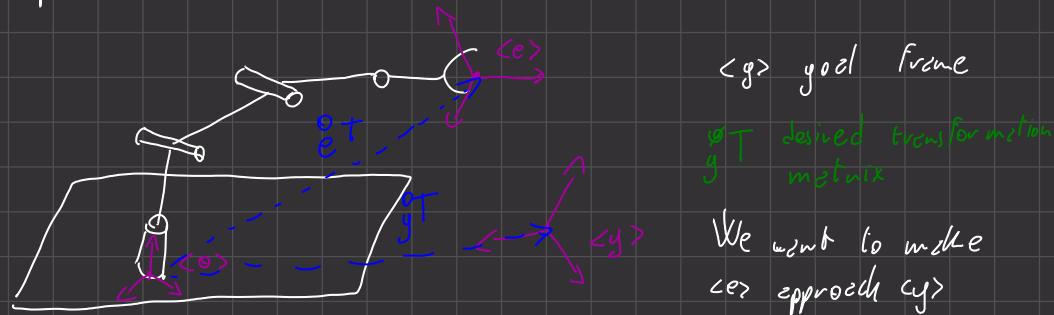
$$J_{1/6}(q) = J_{1/6}(\overset{0}{T}(q); \overset{1}{T}(q); \overset{2}{T}(q); \dots \overset{n}{T}(q))$$

The info that we need to describe 2 Jacobian is not only the angle of every joint but also the distances and all geometrical distances and all the transf.
matrix contain this information. Inside J there are all these geometric distances. To contain all the parameters of the geometry. This connects to the geometric
model.

We can now show a new block that can take the info from the 6M.



We have the means for building two blocks. One to know where the frames start and on the basis of that we can work on the kinematic model. The Jacobians provide the link from the cartesian to joint space?

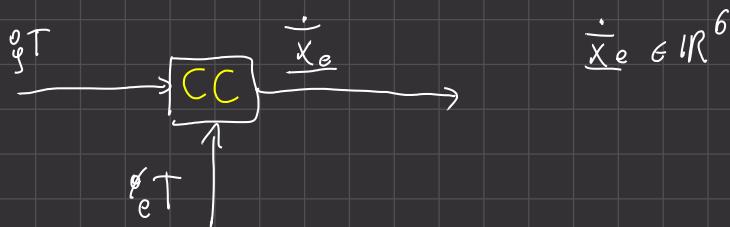


ce depends on how the robot moves and is constrained by that. Let's think for a moment that ce was not constrained. First we want to understand how a frame not constrained should move to reach cg .

If I want to go to position of cg with ce how should I move? For this we should have a vector in the direction from ce to cg . But keep in mind that this is not the only solution because any vector orthogonal to the distance of ce would turn but get closer if there is a component towards cg . In addition to this examples there are infinite vectors that work. We will also see solutions for rotation.

The block outputs the desired velocity for the frame

We have introduced the Cartesian Controller



This outputs the desired cartesian velocity (\dot{x}), but we need the desired joint velocity vector (\dot{q}).
We have seen that Jacobian maps from joint to cartesian space.

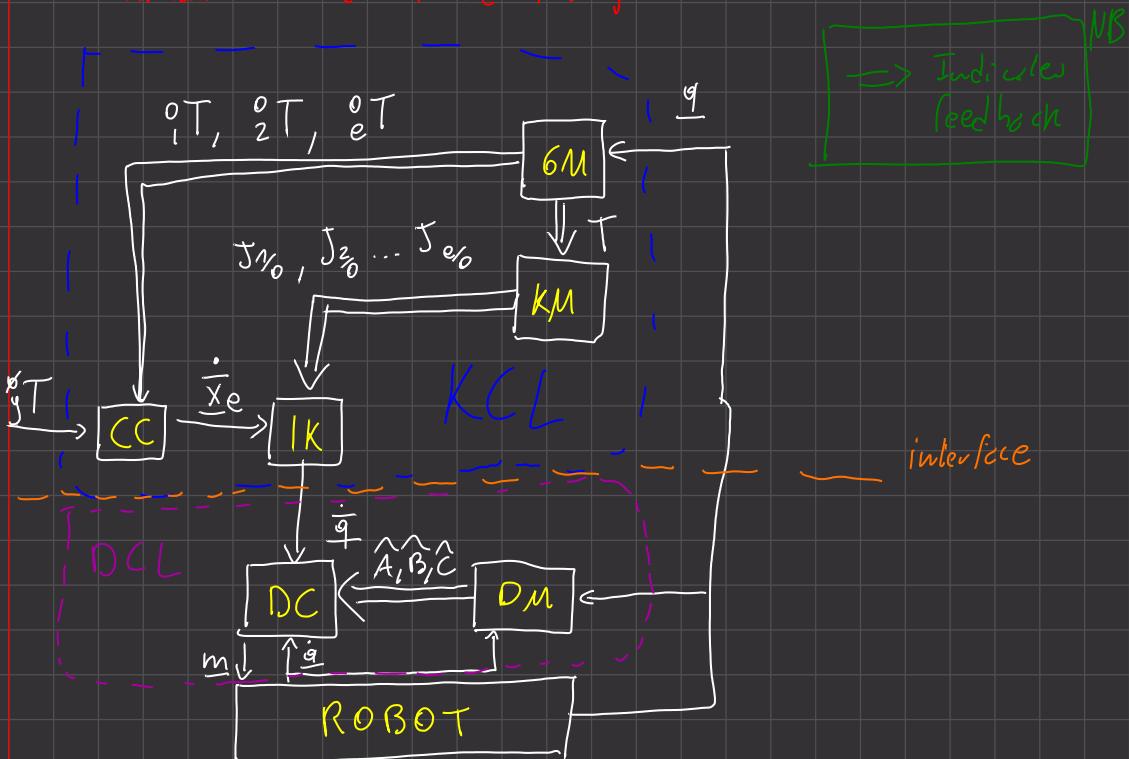
We will see that the inverse of Jacobian will offer in some way the inverse, even if the Jacobian won't be square. This block will output the velocity that best approximate the desired velocity.

If we have 2 q-degree of freedom what we can go in a subset of the 6-dof space.

In the other hand if we have 2 more than 6-dof what sometimes we cannot guarantee the velocity required. For example think when your arm is stretched and you want to extend even more. You can't.

We are still missing the inverse kinematics block. That provides the best approximation to this problem. Remember the workspace of the robot.

The overall control scheme is the following:



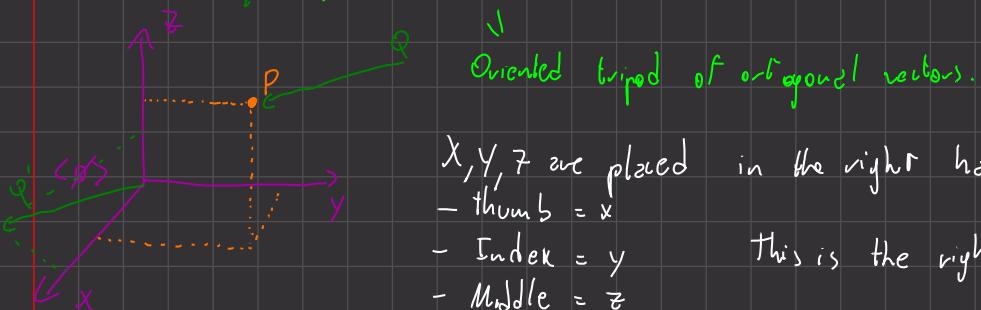
The robot exerts torque (m). From the robot we measure q and \dot{q} . The first thing we do in the DCL is building a controller that builds the desired velocity - the objective of the DCL is to follow as close as possible the reference velocity. Once you have done this you need to drive the ee to the goal. Now I need to prevent \dot{q} . From yearlong we work

position of the frames and with KM their movement with friction. Feeding this to the zidhun we see how the system should move in space. We get the derived velocity that gets mapped to the derived velocity and the inverse kinematic.

TODO

CH. 7 Geometric Fundamentals

Coordinate Systems (Frame)



With the concept of frame we can describe vectors and points. For example

$$\phi P = \begin{bmatrix} x_p \\ y_p \\ z_p \end{bmatrix}$$

The components are the length of the vectors that allow us to reach P in the reference of this particular frame.

We also have vectors.

If we imagine Q in ϕ of the frame it's easy to identify the components. This is called a projected vector.

$$\phi(P-Q) = \phi P - \phi Q$$

TODO ?

The point on which I apply the force is very important, but when I want to know the distance the point on which is applied is not very important.

Often we can also say:

Projected vector ???

$$\phi(P-Q) = \phi P - \phi Q \triangleq \phi \underbrace{v}_{\substack{\uparrow \\ \text{the vector that joins } Q \text{ to } P}}_{P,Q}$$

the vector that joins Q to P

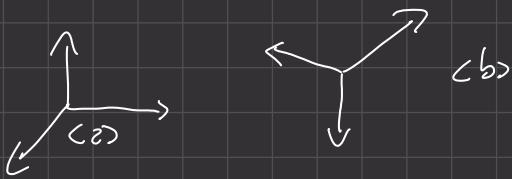
We can now recall the gressman rule to join point and vectors

$$\left| \begin{array}{c} \text{Gressman Rule} \\ \hline \phi P = \phi Q + \underbrace{v}_{P,Q} \end{array} \right|$$

In the frame unit vectors are immediately representable:

$$\begin{array}{ll} x & \begin{bmatrix} 1 & 0 & 0 \end{bmatrix}^T \\ y & \begin{bmatrix} 0 & 1 & 0 \end{bmatrix}^T \\ z & \begin{bmatrix} 0 & 0 & 1 \end{bmatrix}^T \end{array}$$

Let's now consider two frames



Set of information needed to describe how c_b is with respect to c_2 is

$$\left\{ \begin{bmatrix} i_b \\ j_b \\ k_b \end{bmatrix}, \begin{bmatrix} i_{c_2} \\ j_{c_2} \\ k_{c_2} \end{bmatrix} \right\}$$

unit vectors of b projected on c_2 to code the rotation of frame c_b in respect to c_2 .

This information is encoded into the rotation or orientation matrix.

Orientation matrices (Rotation Matrices)

Let's consider a vector v and represent it as a function of c_b .

$$v = i_b x_b + j_b y_b + k_b z_b \quad i, j, k \text{ form a base.}$$

If now I perform projection of v on c_2

$$v = i_{c_2} x_{c_2} + j_{c_2} y_{c_2} + k_{c_2} z_{c_2}$$

Now representing as vector-matrix multiplication we get:

$$v = \begin{bmatrix} i_{c_2} & j_{c_2} & k_{c_2} \end{bmatrix} \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix}$$

↗ c_2
 ↗ b
 ↗ R = ROTATION MATRIX
 ↗ v
 ↗ c_b
 ↗ v

retained frame
referring frame

$$v = R b v \quad \text{This is another way of seeing the multiplication.}$$

In Matlab this will be represented like $v - R - b$.

It's important that in the vector we include the name of the frame on which the vector should be projected on. For example b -distance.

So a multiplication could be:

$$z\text{-distance} = {}^aR_b \quad b\text{-distance}.$$

R is a ortho-normal matrix
columns are unit length
orthogonal to each other

$${}^aR {}^bR^T = I_{n \times n}$$

Given this the matrix has some properties:

Reading list:

$$\det(AB) = \det(A) \det(B)$$

$$\det(A) = \det(A^T)$$

Being ortho-normal means:

$$\begin{cases} {}^aR {}^bR^T = I_{3 \times 3} \\ {}^aR^T {}^bR = I_{3 \times 3} \end{cases} \Rightarrow {}^bR^T = {}^aR^{-1} = {}^bR$$

↑ this because of $\det(R) = 1$

This is important because:

$$\forall M: MM^T = I \Rightarrow \det(M) = \pm 1$$

For Rot. Mat the determinant equals 1 $\rightarrow \det({}^aR) = 1$

↓
this to preserve the direction meaning
the right handed frame.

For example

$$\begin{bmatrix} x & y & z \text{ on } b \\ -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \rightarrow \det = -1$$

We can see that y and z are the same while x is the opposite, but this is impossible if we consider 2 right handed frames. So this cannot happen. $\det({}^aR) = -1$ does not respect right-hand rule.

From these properties we can also conclude that the transpose is the inverse.

It also makes sense if you think about it geometrically.

$$SO(n) = \left\{ R \in \mathbb{R}^{n \times n} : R^T R = I_{n \times n}, \underbrace{\det(R) = +1} \right\}$$

↑
Special orthogonal group

Special for this limitation

This group has some properties:

N.B

Topo on what? MULT DIV

- **CLOSURE:** given any two elements R_1, R_2 of $SO(n)$ then also $R_1 \cdot R_2 \in SO(n)$
- **IDENTITY:** there exists an element I of $SO(n)$ such that $IR \in SO(n)$ for any R
- **INVERSE:** for any $R \in SO(n)$ there exists R^{-1} such that $RR^{-1} = I$
- **ASSOCIATIVITY:** given three elements R_1, R_2, R_3 in $SO(n)$ it follows that $(R_1 R_2) R_3 = R_1 (R_2 R_3)$

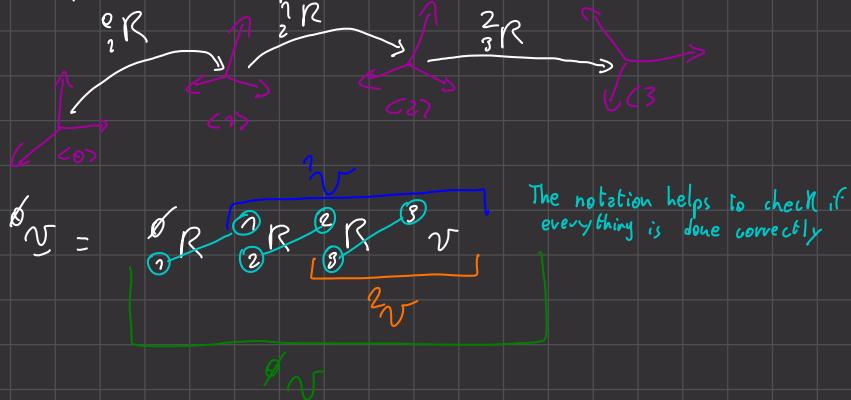
A member of $SO(3)$ is a rotation Matrix

A consideration about $\begin{smallmatrix} {}^2b \\ {}^2b \end{smallmatrix} R^T = {}^b_2 R$ is that this is also equivalent to:

$${}^b_2 R = \begin{bmatrix} {}^2\mathbf{i}_b & {}^2\mathbf{j}_b & {}^2\mathbf{k}_b \end{bmatrix} = \begin{bmatrix} {}^b\mathbf{i}_2^T \\ {}^b\mathbf{j}_2^T \\ {}^b\mathbf{k}_2^T \end{bmatrix}$$

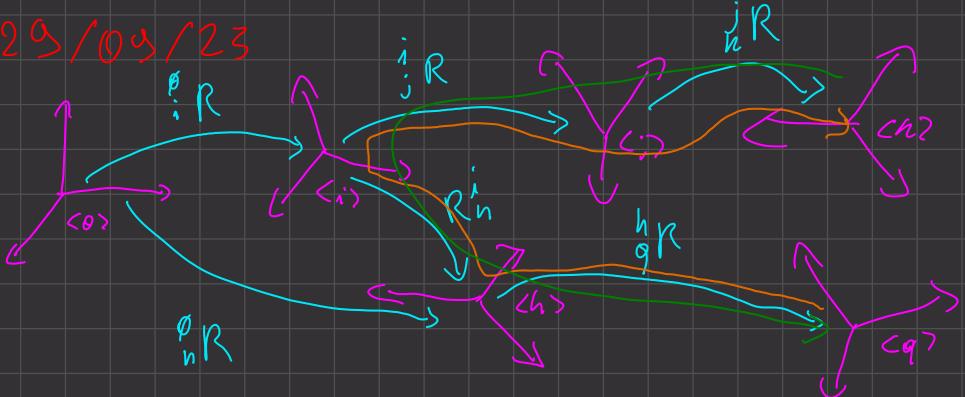
On the transpose you will see what you see on columns now on the rows.

Our manipulator is composed by a tree of frames:



N.B. This holds for every \mathcal{V} so for definition ${}^0\mathcal{R} {}^1\mathcal{R} {}^2\mathcal{R} = {}^0\mathcal{R}$

29/03/23



Here we have some reference frames and suppose that we know some rotation matrices.

postmultiplication

$$\text{How do you compute: } {}^q_h R = {}^h_h R^+ {}^i_h R^T {}^j_h R = \underbrace{{}^q_h R}_\text{postmultiplication} \underbrace{{}^i_h R}_\text{postmultiplication} \underbrace{{}^j_h R}_\text{postmultiplication}$$

We can also do the same thing following this path:

premultiplication

$${}^q_h R = \underbrace{{}^q_h R^+}_\text{premultiplication} \underbrace{{}^i_h R^T}_\text{premultiplication} \underbrace{{}^j_h R}_\text{premultiplication}$$



To get from P to \bar{P} the position of all the frames is an essential parameter. I could for example use another vector P and now the description of P becomes $P + \bar{P}$. This can be done iteratively going back towards $\langle 0 \rangle$.

$${}^K P = (P - O_n)$$

$${}^0 P = {}^0 O_n + {}^0 (P - O_n)$$

$${}^0 P = {}^0 O_n + {}^0 R {}^n P$$

Homogeneous Coordinates

$${}^K \bar{P} = \begin{bmatrix} {}^n P \\ 1 \end{bmatrix} \in \mathbb{R}^4$$

$${}^0 \bar{P} = {}^0 T {}^n \bar{P}$$

$${}^0 T \triangleq \begin{bmatrix} {}^0 R & | & {}^0 O_n \\ \hline 0_{3 \times 1} & | & 1 \end{bmatrix} \in \mathbb{R}^{4 \times 4}$$

$$\text{So if: } d_0 \begin{bmatrix} {}^0 R & | & {}^0 O_n \\ \hline 0_{3 \times 1} & | & 1 \end{bmatrix} \begin{bmatrix} {}^n P \\ 1 \end{bmatrix} = \begin{bmatrix} {}^0 R {}^n P + {}^0 O_n \\ 1 \end{bmatrix}$$

$${}^0 T = {}^0 \bar{P}$$

Inverse transformation matrix

The increase in dimension is a gimmick to obtain the previous formula with a matrix by vector multiplication.

We know that:

$$\begin{matrix} {}^z b \\ {}^z b \end{matrix}^T = \begin{matrix} {}^b b \\ {}^b b \end{matrix}^T$$

N.B.

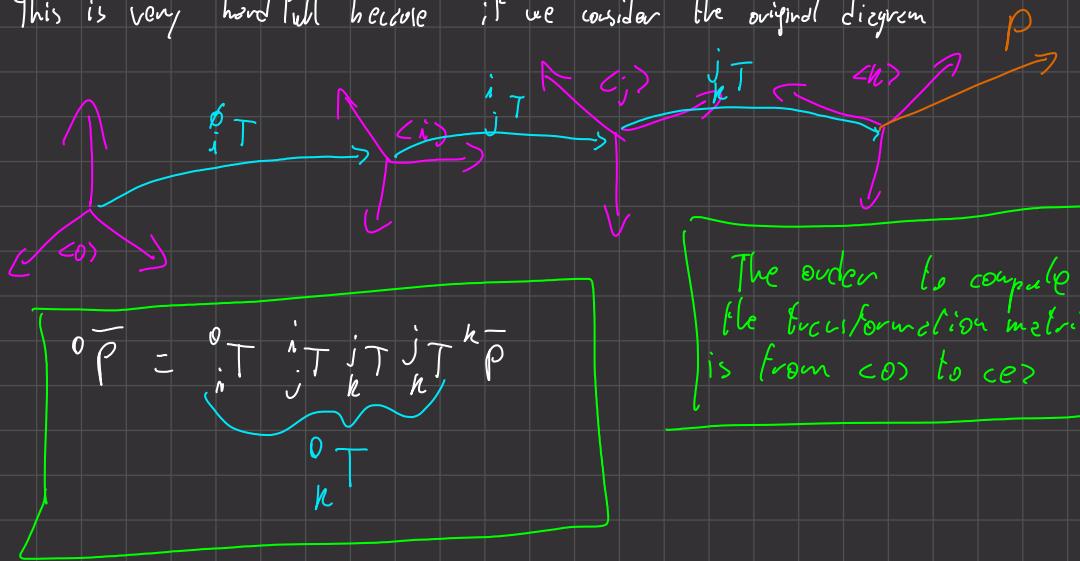
BUT

$$\begin{matrix} {}^z b \\ {}^b b \end{matrix}^T \neq \begin{matrix} {}^z b \\ {}^z b \end{matrix}^T$$

$$\boxed{\begin{matrix} {}^z b \\ {}^b b \end{matrix}^T = \begin{bmatrix} {}^z b^T & {}^z R^T & {}^z O_b \\ 0_{3 \times 1} & I \end{bmatrix}}$$

The origin of ${}^z b$ on ${}^b b$ would be the opposite vector and to put it in the right direction I need a minus in front.

This is very hard full because if we consider the original diagram



The order to compute the transformation matrix is from c_0 to c_n

It's also possible to describe a vector using homogeneous coordinates.
It's done by just adding ϕ to the vector.

$$P = \begin{bmatrix} {}^n P \\ \phi \end{bmatrix}_{4 \times 1}$$

This because the origin in this case does not come into play.

The inertial frame it's not different from any other frame.

Three parameter representation of rotation matrices

$\overset{3}{\underset{b}{R}}$

- Rotation matrices are orthonormal. vectors are orthogonal to each other
modulus = 1 for every vector.

So I have:

$$\left\{ \begin{array}{l} \overset{2}{\lambda_b^T} \overset{2}{\lambda_b} = 1 \\ \overset{2}{j_b^T} \overset{2}{j_b} = 1 \\ \overset{2}{k_b^T} \overset{2}{k_b} = 1 \end{array} \right. \quad \left\{ \begin{array}{l} \downarrow \text{dot product} \\ \overset{2}{\lambda_b^T} \overset{2}{j_b} = 0 \\ \overset{2}{j_b^T} \overset{2}{k_b} = 0 \\ \overset{2}{k_b^T} \overset{2}{\lambda_b} = 0 \end{array} \right.$$

$\overset{2}{\lambda_b}$ satisfy all these equations

We have 6 constraints and 9 numbers so we only have to pick 3 that are unknown.

For this reason we can talk about 3 number representation of rotation matrices.

The easiest of these representation is called

→ Pronounced Euler

Euler angles

They express that any rotation can be represented by three rotations between 3 different axes.

We can choose between:

- extrinsic rotation
- intrinsic rotation

- I describe the final orientation with rotation around a fixed coordinate system. This unfortunately is not so easy to implement because when we multiply R. matrices we are not doing intrinsic rotation. In this case we first rotate around the original x, then around the original y and finally around the original z.
- this is easier to implement. We first rotate around x, then the rotated y and then the rotated z. When we multiply with Rot. Matrices we are doing intrinsic rotations.

Intrinsic angles can be divided in:

- proper Euler angles

$z-x-z$; $y-x-y$; $x-z-x$; ...

These represent the sequence of axes around which you do the rotation

- Tait-Bryan angles

$x-y-z$; $z-y-x$; ...

$$z - y^1 - x^1$$

This is the most used one and we will use this one

$${}^3 R = R_z(\psi) \ R_y(\theta) \ R_x(\phi)$$

YAW

PITCH

ROLL

The open stresses that we are talking about intrinsic rotation

It's possible to demonstrate that an extrinsic rotation x-y-z is the same as an intrinsic rotation z-y-x.

02/10/23

$${}^3 R = R_z(\psi) \ R_y(\theta) \ R_x(\phi) =$$

$$= \begin{bmatrix} C\psi C\theta & -s\psi C\phi + c\psi s\theta s\phi & s\psi s\phi + c\psi c\phi s\theta \\ S\psi C\theta & C\psi C\phi + s\phi s\theta s\psi & -c\psi s\phi + s\theta s\psi c\phi \\ -s\theta & c\theta s\phi & c\theta c\phi \end{bmatrix}$$

Not needed to know by heart

It's useful to know the structure to know the inverse.

$$\theta = \text{atan2}(-r_{31}, \sqrt{r_{11}^2 + r_{21}^2}) \in \left[-\frac{\pi}{2}, \frac{\pi}{2}\right]$$

software function to compute atan2 on a quadrant instead of two.

$$\text{If } \cos\theta \neq 0 : \quad \psi = \text{atan2}(r_{21}, r_{11}) \\ \phi = \text{atan2}(r_{32}, r_{31})$$

If $\cos\theta = 0$ it's an interesting configuration because leads to a singularity.

In the sequence of rotation if we rotate in pitch of $\pm \frac{\pi}{2}$ the x axis becomes the same as the z axis of the original frame. We have infinite combination of first and third axis that end up representing the same output. We get a singularity when the pitch is $\pm \frac{\pi}{2}$.

Suppose now that I want to control roll to a certain angle, when i get $\frac{\pi}{2}$ in pitch yaw and roll become the same. It's an ill defined situation.

The use of Euler angles is fine but we must have regard for pitch being close to $\frac{\pi}{2}$

Vector operations

- Linear combination

$$\underline{v} = c_1 \underline{v}_1 + c_2 \underline{v}_2 + \dots + c_n \underline{v}_n$$

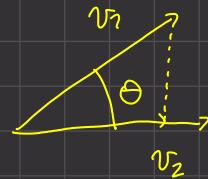
If they are linear independent they can form the base of your n-based dimension space.

If I take into consideration the projection of a frame I also need to project on the same frame all the components.

$${}^2 \underline{v} = c_1 {}^2 \underline{v}_1 + c_2 {}^2 \underline{v}_2 + \dots + c_n {}^2 \underline{v}_n$$

- Scalar product

$$(\underline{v}_1 \cdot \underline{v}_2) = (\underline{v}_2 \cdot \underline{v}_1) = |\underline{v}_1| |\underline{v}_2| \cos \theta$$



$$(\underline{v}_1 \cdot \underline{v}_2) = {}^2 \underline{v}_1^\top {}^2 \underline{v}_2$$

projected on
c2)

$$\left[\begin{array}{l} \text{I can consider } (\underline{v}, \cdot) = (\underline{v}^\top)_{1 \times 3} \\ \text{linear operation} \end{array} \right]$$

- Vector product

$$(\underline{v}_1 \wedge \underline{v}_2) = (\underline{v}_1 \times \underline{v}_2) = \begin{cases} \text{modulus} & |\underline{v}_1 \times \underline{v}_2| = |\underline{v}_1| |\underline{v}_2| \sin(\theta) \xrightarrow{\text{minimum angle}} \\ \underline{v}_1 \times \underline{v}_2 \perp (\underline{v}_1, \underline{v}_2) & \end{cases}$$

The direction is orthogonal to both \underline{v}_1 and \underline{v}_2 according to the right hand rule

- Properties:

$$\bullet \text{ Anti-commutative: } (\underline{v}_1 \times \underline{v}_2) = -(\underline{v}_2 \times \underline{v}_1)$$

$$\bullet {}^2(\underline{v}_1 \times \underline{v}_2) = {}^2[\underline{v}_1 \times] {}^2 \underline{v}_2 = \begin{bmatrix} 0 & -z & y \\ z & 0 & -x \\ -y & x & 0 \end{bmatrix} {}^2 \underline{v}_2$$

Components of ${}^2 \underline{v}_1$

For this version

↳ This matrix is skew-symmetric

$$\bullet {}^2 [\underline{v}_1 \times]^T = - {}^2 [\underline{v}_1 \times]$$

axial vector

sometimes this can also be represented like: $[\underline{v}_1 \times] = S(\underline{v}_1)$

$$\bullet \text{ Inverse mapping: } \text{vex operation} \rightarrow \text{vex}(S(\underline{v}_1)) = \underline{v}_1$$

$$\boxed{\forall M: M \in \mathbb{R}^{3 \times 3} / M = -M^T}$$
$$M \underline{y} = \text{vex}(M) \times \underline{y}$$

Any skew-symmetric has an axial vector. Vex is needed to get x, y and z

Power of skew-symmetric 3×3 real matrices associated with a unit vector

$$h \in \mathbb{R}^{3 \times 1}$$

$$\underline{[h \ x]} = \underline{h \ x}$$

$$\underline{[h \ x]}^2 = (\underline{h} \ \underline{h}^\top - I)$$

After the q^{th} power they repeat from the first one.

$$\underline{[h \ x]}^3 = - \underline{[h \ x]}$$

$$\underline{[h \ x]}^4 = - \underline{[h \ x]}^2$$

$$\underline{[h \ x]}^5 = \underline{[h \ x]}$$

$$\underline{[h \ x]}^6 = \underline{[h \ x]}^2$$

⋮

- In general we state that:

$$[\underline{h \ x}]^{2i+1} = (-1)^i [\underline{h \ x}]$$

Rule for odd numbers
(skew symmetric)

$$[\underline{h \ x}]^{2i+2} = (-1)^i (\underline{h} \ \underline{h}^\top - I)$$

Rule for even numbers
(symmetric)

$(I - \underline{h} \underline{h}^\top)$ projects a vector on the plane orthogonal to \underline{h} (TODOLICHT)

- $(I - \underline{h} \underline{h}^\top)$ is the projection operator defined by \underline{h}

$$(I - \underline{h} \underline{h}^\top) \underline{h} = \emptyset \text{ because } (\underline{h} - \underline{h} \underbrace{\underline{h}^\top \underline{h}}_{\sim}) = \emptyset$$

- true $([\underline{h \ x}]^{2i+2}) = 0$ for \rightarrow sum of all the elements on the main diagonal in a matrix

$$t_r([\underline{h \ x}]^{2i+2}) = (-1)^{(i+1)} \cdot 2$$

- $A \in \mathbb{R}^{n \times n}$

$$A = A^\top \text{ symmetrical}$$

$$A = -A^\top \text{ skew symmetrical}$$

$$\forall A \in \mathbb{R}^{n \times n} \quad A = \underbrace{\frac{A+A^\top}{2}}_{\text{symmetric}} + \underbrace{\frac{A-A^\top}{2}}_{\text{skew-symmetric}}$$

no symm
skew-symm

Linear Operators between Vectors

Consider: frame \mathbf{b}

- any projected vector ${}^b\mathbf{v}$
- ${}^b\mathbf{w}$ subject to ${}^b\mathbf{w} = {}^b\mathbf{L} {}^b\mathbf{v}$

$${}^2\mathbf{w} = {}^2\mathbf{R} {}^b\mathbf{w} = {}^2\mathbf{R} {}^b\mathbf{L} {}^b\mathbf{v} = \boxed{{}^2\mathbf{R} {}^b\mathbf{L} {}^b\mathbf{R}^T} {}^2\mathbf{v} = {}^2\mathbf{L} {}^2\mathbf{v}$$

$$\boxed{{}^2\mathbf{L} = {}^2\mathbf{R} {}^b\mathbf{L} {}^b\mathbf{R}^T}$$

This is the operation to change the operator from one frame to another.

For example:

vector product operator

$${}^b(\mathbf{w} \times \mathbf{v}) = {}^b[\mathbf{w} \times] \mathbf{v}$$

$${}^2(\mathbf{w} \times \mathbf{v}) = {}^2\mathbf{R} {}^b[\mathbf{w} \times] {}^b\mathbf{R}^T {}^2\mathbf{v} = {}^2[\mathbf{w} \times] {}^2\mathbf{v}$$

06/10/23

Matrix exponential operator

$${}^b\mathbf{L} = e^A \quad A \in \mathbb{C}^{3 \times 3}$$

$${}^b\mathbf{L} = \sum_{k=0}^{\infty} \frac{({}^bA)^k}{k!}$$

As any other linear operator we can write:

$${}^b\mathbf{w} = {}^b\mathbf{L} {}^b\mathbf{v} = e^A {}^b\mathbf{v}$$

Now we want to know ${}^2\mathbf{L}$:

$${}^2\mathbf{L} = {}^2\mathbf{R} {}^b\mathbf{L} {}^b\mathbf{R}^T = {}^2\mathbf{R} e^A {}^b\mathbf{R}^T$$

Now if we substitute the definition of the exponential:

$$= {}^2\mathbf{R} \left(\sum_{k=0}^{\infty} \frac{({}^bA)^k}{k!} \right) {}^b\mathbf{R}^T = \sum_{k=0}^{\infty} \frac{{}^2\mathbf{R} ({}^bA)^k {}^b\mathbf{R}^T}{k!}$$

$$= \sum_{k=0}^{\infty} \frac{({}^2\mathbf{R} {}^bA {}^b\mathbf{R}^T)^k}{k!} \quad \text{this step is possible because :}$$

$$\underbrace{({}^2\mathbf{R} {}^bA {}^b\mathbf{R}^T)}_{\text{Identity}} \underbrace{({}^2\mathbf{R} {}^bA {}^b\mathbf{R}^T)}_{\text{Identity}} \dots \underbrace{({}^2\mathbf{R} {}^bA {}^b\mathbf{R}^T)}_{\text{Identity}} \underbrace{({}^2\mathbf{R} {}^bA {}^b\mathbf{R}^T)}_{\text{Identity}}$$

k times

To conclude

$$L = e^{\theta R} A R^T = e^{\theta A}$$

Properties of the matrix exponential

- $M e^{\theta M} = e^{\theta M} M \quad \theta \in \mathbb{R}$
- $(e^{\theta M})^T = e^{\theta M^T}$
- $(e^{\theta M})^{-1} = e^{-\theta M}$

Exponential representation of rotation matrices.

Every $e^{\theta[\underline{h}\times]} \in SO(3)$, so this is a rotation matrix

To prove this:

$$(e^{\theta[\underline{h}\times]})^T e^{\theta[\underline{h}\times]} = e^{\theta[\underline{h}\times]^T} e^{\theta[\underline{h}\times]} = \\ = e^{-\theta[\underline{h}\times]} e^{\theta[\underline{h}\times]} = I_{3 \times 3}$$

This determinant is either +1 or -1. If $\theta=0$ the determinant is +1. There is a continuity rule also for the determinant. It cannot jump from +1 to -1. For this continuity rule we can state that the determinant is always +1.

So we have $\det = 1$ and the product for the transpose is $I_{3 \times 3}$ so it's a Rotation Matrix.

Angle axis parameters.

$\theta \in \mathbb{R}$ angle

$\underline{h} \in \mathbb{R}^{3 \times 1}$, $|\underline{h}|=1$ axis

together they are the angle-axis representation of the rotation matrix $e^{\theta[\underline{h}\times]}$

\underline{h} gives you the axis around which to rotate and θ gives you the quantity.

Rotation Vector

$\underline{g} = \theta \underline{h} \in \mathbb{R}^{3 \times 1}$ is called the rotation vector associated to the rotation matrix $e^{\theta[\underline{h}\times]}$

Angle axes vector

$$\underline{v} = \begin{bmatrix} h \\ \theta \end{bmatrix} \in \mathbb{R}^{3+1}$$

This is also called 4D-vector associated to the matrix $e^{\underline{\theta}[\underline{h}x]}$

Let's now consider how the exponential is computed:

$$e^{\underline{\theta}[\underline{h}x]} = \sum_{k=0}^{\infty} \frac{(\underline{\theta}\underline{A})^k}{k!} =$$

We split even and odd powers of this infinite summation:

$$= \mathbb{I}_{3 \times 3} + \left(\underline{\theta} - \frac{\underline{\theta}^3}{3!} + \frac{\underline{\theta}^5}{5!} \dots \right) [\underline{h}x] + \left(\frac{\underline{\theta}^2}{2!} - \frac{\underline{\theta}^4}{4!} + \dots \right) [\underline{h}x]^2$$

We can do this for the properties of skew-symmetrical matrices

Rodrigues Formula

$$= \mathbb{I}_{3 \times 3} + \underbrace{\sin \underline{\theta}}_{\text{the first sum converges to } \sin} [\underline{h}x] + \underbrace{(1-\cos \underline{\theta})}_{\text{the second one converges to } (1-\cos \underline{\theta})} [\underline{h}x]^2$$

N. B.

\underline{h} is already a unit vector

From the Rodrigues formula we can find the mere mapping from a rotation matrix to the angle-axes mapping.

On the diagonal you get the square of the components.

Equivalent Angle-axes inverse problem

$$R \xrightarrow{?} \underline{\theta}, \underline{h}$$

$$(\underline{h}\underline{h}^T - \mathbb{I})$$

$$\text{Let's look at the } \text{tr}(R) = \text{tr} (\mathbb{I} + \sin \underline{\theta} [\underline{h}x] + (1-\cos \underline{\theta}) [\underline{h}x]^2) =$$

$$= 3 + 0 + (1-\cos \underline{\theta})(1-3) = 3 - 2(1-\cos \underline{\theta}) = 2 + 2\cos \underline{\theta}$$

$$\text{So: } \underline{\theta} = \cos^{-1} \left(\frac{\text{tr}(R)-2}{2} \right)$$

Like this we can find more than one parameter

For example if:

$\underline{\theta} = \emptyset$ then \underline{h} is arbitrary

$\underline{\theta} = \pi$ then

Let's try it:

$$e^{\hat{\theta}[\underline{h}x]} = I_{3 \times 3} + 2(\underline{h}\underline{h}^T - I) = 2\underline{h}\underline{h}^T - I_{3 \times 3} =$$

$$= \begin{bmatrix} 2h_1^2 - 1 & 2h_1h_2 & 2h_1h_3 \\ 2h_1h_2 & 2h_2^2 - 1 & 2h_2h_3 \\ 2h_1h_3 & 2h_2h_3 & 2h_3^2 - 1 \end{bmatrix} = R$$

From this we can conclude that the modulus of h =

$$|h_i| = \sqrt{\frac{r_{ii} + 1}{2}} \quad \forall i = 1, 2, 3$$

To find the sign of this vector we can select a random component $i \in \{1, 2, 3\}$.

$$h_i = \pm \sqrt{\frac{r_{ii} + 1}{2}}$$

$$\text{The generic } h_j = \text{sign}(h_i) \text{sign}(r_{ij}) \sqrt{\frac{r_{jj} + 1}{2}} \quad \forall j \in \{1, 2, 3\} \setminus i$$

If you pick $i = 1$
then you try for 2 and 3

What about all the other cases?

Let's consider $R - R^T$

First we expand R and R^T with the Rodriguez formula:
we are left with

$$R - R^T = \sin \theta [\underline{h}x] - \sin \theta [\underline{h}x]^T$$

$$\begin{aligned} R - R^T &= \cancel{I_{3 \times 3}} + \sin \theta [\underline{h}x] + (1 - \cos \theta)[\underline{h}x]^T \\ &\quad - \cancel{-I_{3 \times 3}} - \sin \theta [\underline{h}x] - (1 - \cos \theta)[\underline{h}x]^T = \\ &= 2 \sin \theta [\underline{h}x] \end{aligned}$$

$$\frac{R - R^T}{2} = \sin \theta [\underline{h}x] \quad \text{so} \quad \sin \theta \underline{h} = \text{vex}\left(\frac{R - R^T}{2}\right)$$

This extracts the
axis vector from the
skew symmetric Matrix

The goal at the end will be to make the rotation matrix converge to I and the norm of the rotation vector to 0.

Both yaw-pitch-roll and rotation matrices show that the mapping from R to these 2/3 parameters is not unique

Euler Theorem

It states that any 3d frame can be aligned with any other arbitrary fixed frame by one rotation θ around one axis.

08/10/23

Unit quaternion

These are defined by $q \in \mathbb{R}^{4 \times 1}$

$$q = \begin{bmatrix} u \\ \varepsilon \end{bmatrix} \text{ where } \begin{cases} u \in \mathbb{R}^3 \\ \varepsilon \in \mathbb{R} \end{cases}$$

$$\|u\|^2 + \|\varepsilon\|^2 = 1 \leftarrow \text{this is why it's called unit quaternion}$$

The relation between this and the rotation vector are the following:

Trigonometric Properties

First let's recall:

$$\begin{array}{l} \text{I} \quad \sin(\varphi + \beta) = \sin \varphi \cos \beta + \sin \beta \cos \varphi \\ \text{II} \quad \cos(\varphi + \beta) = \cos \varphi \cos \beta - \sin \varphi \sin \beta \end{array} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{Ptolemaic Formulas}$$

Let's now consider:

$$\varphi = \beta = \frac{\theta}{2}$$

Now we get what appears in the Rodriguez formula:

$$\text{I} \quad \sin \theta = 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2}$$

$$\text{II} \quad \cos \theta = \cos^2 \frac{\theta}{2} - \sin^2 \frac{\theta}{2}$$

Now if we do

$$\cos \theta = \cos^2 \frac{\theta}{2} - \sin^2 \frac{\theta}{2} + \cos^2 \left(\frac{\theta}{2} \right) - \cos^2 \left(\frac{\theta}{2} \right)$$

$$\Rightarrow \cos \theta = 2 \cos^2 \left(\frac{\theta}{2} \right)$$

Now if we add and subtract $\sin^2 \left(\frac{\theta}{2} \right)$ here we get

$$\Rightarrow \cos \theta = 2 \sin^2 \left(\frac{\theta}{2} \right)$$

$$R = I + \sin\theta [\underline{h}x] + (1 - \cos\theta) [\underline{h}x]^2 = \\ = I + 2 \sin\left(\frac{\theta}{2}\right) \cos\left(\frac{\theta}{2}\right) [\underline{h}x] + 2 \sin^2\left(\frac{\theta}{2}\right) [\underline{h}x]^2$$

Now if we define $\begin{cases} \mu \triangleq \cos\frac{\theta}{2} \\ \varepsilon \triangleq \sin\frac{\theta}{2} h \end{cases}$ we can write: ???

$\sin\theta$ goes inside $[\underline{h}x]$

$R = I + 2\mu[\varepsilon x] + 2[\varepsilon x]^2$

→ this is a not singular representation of the rotation matrix.

If $\theta=0$ then $\varepsilon=0$. There is no undefined of the axis vector \underline{h} because it's always well defined.

Inverse mapping: from $R \rightarrow$ quaternions

To compute the inverse mapping we do:

$$\text{tr}(R) = 3 + 0 + 2 \text{tr}([\varepsilon x]^2)$$

$$\text{tr}(R) = 3 + 2 \text{tr}(\underline{\varepsilon} \underline{\varepsilon}^T - \underline{\varepsilon}^T \underline{\varepsilon} I)$$

$$\text{tr}(R) = 3 + 2 (\|\underline{\varepsilon}\|^2 - 3\|\underline{\varepsilon}\|^2)$$

$$\text{tr}(R) = 3 - 4\|\underline{\varepsilon}\|^2$$

$$\text{tr}(R) = 4\mu^2 - 1$$

$$\mu = \frac{1}{2} \sqrt{\text{tr}(R) + 1}$$

let's now work on $\underline{\varepsilon}$

$$R - R^T = \cancel{I} + 2\mu[\varepsilon x] + 2[\varepsilon x]^2 + \\ - \cancel{I} + 2\mu[\varepsilon x] - 2[\varepsilon x]^2$$

$$R - R^T = 4\mu[\varepsilon x]$$

$$\frac{R - R^T}{2} = 2\mu[\varepsilon x]$$

$$\underline{\varepsilon} = \frac{1}{2\mu} \text{vex}\left(\frac{R - R^T}{2}\right) \quad \text{if } \mu \neq 0$$

ϵ eigenvector of the matrix $\frac{R \cdot R^T}{2}$ and then pick the one with eigenvalue = 1 if $\mu=0$
 \hookrightarrow eigenvectors

Flipping the signs will give you the other eigenvalues?

Recalling euler angles, there is a problem that if we have the sequence $x-y-x$ if $y=0$ then we have two times the same x

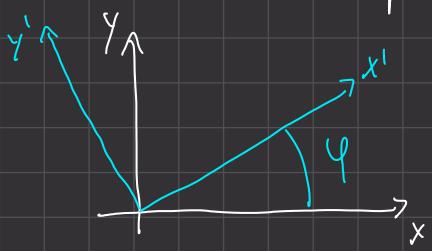
Let's also recall how an elementary rotation is written. So for example a rotation around the z axis of φ is:

$$R_z(\varphi) = \begin{bmatrix} \cos \varphi & -\sin \varphi & 0 \\ \sin \varphi & \cos \varphi & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$R_y(\theta) = \begin{bmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{bmatrix}$$

$$R_x[\phi] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \phi & -\sin \phi \\ 0 & \sin \phi & \cos \phi \end{bmatrix}$$

To find them for example:



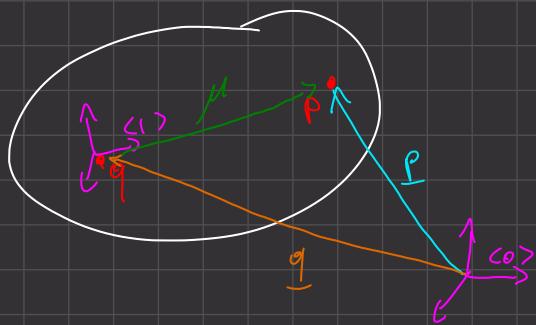
Right hand rule. the rotation is out of the page.

So there are the x' coordinates on the x axis

These are the y' on the y axis z' on the z .

Angular velocity vector

Let's consider a generic object with specific points where we have some frames



$$p = q + \mu$$

If we now start projecting:

$${}^0\dot{p} = {}^0\dot{q} + {}^0\dot{\mu} = {}^0\dot{q} + {}^1R{}^0\dot{\mu}$$

Now we want to consider:

$${}^0\dot{p} \triangleq \underbrace{\frac{d}{dt}({}^0R{}^1p)}_{\text{II}} \neq {}^1R \underbrace{\frac{d}{dt}{}^1p}_{\text{I}}$$

this is \downarrow the correct one because we are switching from frame ϕ ?
In the second one we are calculating from the frame c_1

This ambiguity will be solved with ${}^0Vp/{}_0$ (this means seen from ϕ)

Using this notation the previous becomes:

$${}^0\dot{p} \triangleq \underbrace{\frac{d}{dt}({}^0R{}^1p)}_{\text{II}} \neq \underbrace{{}^1R \frac{d}{dt}{}^1p}_{\text{I}}$$

$$\quad \quad \quad {}^0Vp/{}_0 \quad \quad \quad {}^1Vp/{}_1$$

Let's continue:

$$\begin{aligned} {}^0\dot{p} &= {}^0\dot{q} + \frac{d}{dt}({}^0R{}^1\mu) \\ &= {}^0\dot{q} + {}^0R{}^1\dot{\mu} + {}^0R{}^1\ddot{\mu} \\ &= {}^0\dot{q} + \underbrace{{}^0R \underbrace{{}^0R^T {}^0R}_{{}^0I_{3x3}}}^{{}^0R} \dot{\mu} + {}^0R{}^1\ddot{\mu} \end{aligned}$$

0q is the position of frame
 \hookrightarrow to c_0

μ position from p to c_1

This tells us c_0 is changing in this term times respect to c_0 into account that c_1 is changing orientation to c_0

here μ changes in } this is ϕ if in the respect to c_1 than } rigid space of projected on c_0 } frame c_1

Now R is an element of $SO(3)$ so:

$$RR^T = I_{3 \times 3}$$

So if now I do:

$$\frac{d}{dt} ({}^0 R {}^0 R^T) = 0 = {}^0 \dot{R} {}^0 R^T + {}^0 R {}^0 \dot{R}^T$$
$$\underbrace{{}^0 \dot{R} {}^0 R^T}_{\text{skew-symmetric matrix}} = - {}^0 R {}^0 \dot{R}^T = - \underbrace{({}^0 \dot{R} {}^0 R^T)^T}_{\text{skew-symmetric matrix}}$$

We can always associate an axial vector to a skew symmetrical matrix.

$$({}^0 \dot{R} {}^0 R^T) = {}^0 \underline{\omega}_{1/0} \times \underline{\mu} \quad ???$$

Definition of angular velocity vector

$\forall R \in SO(3)$ there exists a unique $\omega \in \mathbb{R}^{3 \times 1}$ such

that

$$\boxed{{}^0 \dot{R} {}^0 R^T = [\underline{\omega}_{1/0}]} \quad \text{Steady equation}$$

So if we have:

$${}^0 \dot{R} {}^0 R^T = \left[\begin{array}{c} \underline{\omega}_{1/0} \\ \underline{\mu} \end{array} \right]$$

↔

inverse = transpose

$${}^0 \dot{R} = \left[\begin{array}{c} \underline{\omega}_{1/0} \\ \underline{\mu} \end{array} \right] {}^0 R$$

$${}^0 \dot{R} = {}^0 R \left[\begin{array}{c} \underline{\omega}_{1/0} \\ \underline{\mu} \end{array} \right]$$

$$\left[\begin{array}{c} \underline{\omega}_{1/0} \\ \underline{\mu} \end{array} \right] = {}^0 R \left[\begin{array}{c} \underline{\omega}_{1/0} \\ \underline{\mu} \end{array} \right] {}^0 R^T$$

Transformation of the operator

Example:

$$R_z(\alpha) = \begin{bmatrix} \cos(\alpha) & -\sin(\alpha) & 0 \\ \sin(\alpha) & \cos(\alpha) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$\alpha \neq 0$

$$R_z(\alpha) = \begin{bmatrix} -\sin(\alpha) & -\cos(\alpha) & 0 \\ \cos(\alpha) & -\sin(\alpha) & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

If we directly calculate:

$$\dot{R}_2(\alpha) R_2^T(\alpha) = \dot{\alpha} \cdot \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} \omega x \\ 0 \\ 0 \end{bmatrix}$$

$$\text{So } \omega = \begin{bmatrix} 0 \\ 0 \\ \dot{\alpha} \end{bmatrix}$$

When you have the angular velocity (ω) you cannot directly integrate it but use something like the stepdown matrix.

Except when you are constantly rotating around one axis, meaning that I don't change axes in respect to which I rotate.

To better say in general you have to apply the stepdown for every infinitesimal rotation and compute ω .

Angular velocity properties and composition rules

$${}^0\dot{R} {}^0R^T = \begin{bmatrix} {}^0\omega_{0/1} \times \end{bmatrix} \xrightarrow{\text{N.B.}}$$

$${}^0\dot{R} \left[{}^1\dot{R} {}^0R^T \right] {}^0R^T \xrightarrow{\text{N.B.}} \begin{bmatrix} {}^0\omega_{0/1} \end{bmatrix}$$

$$= {}^0\dot{R} {}^1\dot{R} \left({}^0R^T {}^0R^T \right) = {}^0\dot{R} {}^1\dot{R} \xrightarrow{\text{N.B.}} A^T B = (B^T A)^T$$

$$= \left({}^0\dot{R} {}^1\dot{R} \right)^T = \underbrace{\left({}^0\dot{R} {}^0R^T \right)^T}_{\text{skew symmetric}} = - \underbrace{\left({}^0\dot{R} {}^0R^T \right)}_{\text{stepdown}} =$$

$$= - \left[{}^0\omega_{0/1} \times \right]$$

N.B.

$$\phi \omega_{0/1} = - \omega_{0/0}$$

We can conclude

The angular velocity is a relative concept

Let's see the composition:

$${}^0R = {}^1R {}^2R$$

$$\dot{{}^0R} {}^0R^T = \left[\frac{d}{dt} \left({}^1R {}^2R \right) \right] ({}^0R {}^1R)^T$$

$$= \left({}^0\dot{R} {}_2^1 R + {}^0R {}_2^1 \dot{R} \right) {}_2^1 R^T {}_1^0 R^T$$

$$= \underbrace{{}^0\dot{R} {}^0 R^T}_{\text{stepdown}} + {}^0R \underbrace{\left[{}_2^1 \dot{R} {}_2^1 R^T \right]}_{\text{stepdown}} {}_1^0 R^T$$

this would be:

$$\phi \omega_{z0} + \phi \begin{bmatrix} {}^1\omega_{x1} \\ {}^1\omega_{y1} \end{bmatrix}$$

Rotated to be on ϕ

From this follows:

$$\underline{\omega}_z = \underline{\omega}_{z0} + \underline{\omega}_{x1}$$

Galilean velocity composition

Angular velocity and the rotation vector

$$\underline{\omega} = \dot{\theta} \underline{h} + \sin \theta \underline{i} + (1 - \cos \theta) \underline{h} \times \underline{h}$$

Base of a 3-dimensional space. \Rightarrow

$\begin{cases} \underline{h} \text{ and } \underline{i} \text{ are orthogonal} \\ \text{because } \underline{h} \text{ is unitary in length and can only rotate, not change the length.} \end{cases}$

If $\underline{h} = 0$, means that the axis around which one frame is rotating respect to the other is not changing, and we can integrate.

Otherwise if $\underline{h} \neq 0$ integrating ω it does not give θ , and we are rotating around moving axes.

Mapping between ω and the 3-D axis vector

$$\underline{v} = \begin{bmatrix} \underline{h} \\ \theta \end{bmatrix}$$

This works only if

$$\cos \theta \neq 1$$

$$\underline{v} = \begin{bmatrix} -\frac{\sin \theta}{2(1 - \cos \theta)} [\underline{h} \times]^2 - \frac{1}{2} [\underline{h} \times] \\ \underline{h}^T \end{bmatrix} \underline{\omega}$$

You cannot integrate R from the stepdown because otherwise you won't be in $SO(3)$. One solution could be to normalize the matrix. If you want to know the were orientation it could be convenient to find the mapping of your pitch-roll angles.

If we know how your pitch-roll relate in relation to ω then we are good.

When you know your pitch-roll you can now do a forward Euler integration. What you will get is another $so(3)$ element.

Angular velocity vector and time derivative of minimal parametrization vectors

$${}^2\omega_{b/2} = {}^bL_{b/2} \dot{\bar{\varphi}} \quad \bar{\varphi} = \begin{bmatrix} \psi \\ \theta \\ \phi \end{bmatrix} \quad N.B.,$$

The order is
yaw-pitch-roll ($z-y-x$)

$$\begin{aligned} {}^2\omega_{b/2} &= \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \dot{\psi} + R_z(\psi) \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} \dot{\theta} + R_z(\psi) R_y(\theta) \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \dot{\phi} \\ &= \begin{bmatrix} 0 & -s\psi & c\psi c\theta \\ 0 & c\psi & s\psi c\theta \\ 1 & 0 & -s\theta \end{bmatrix} \begin{bmatrix} \dot{\psi} \\ \dot{\theta} \\ \dot{\phi} \end{bmatrix} \end{aligned}$$

$$\det(L_{b/2}) = -\cos(\theta) \quad \leftarrow \text{This mapping becomes undefined whenever } \theta = \frac{\pi}{2} \text{ or better I cannot find the inverse mapping but from where I don't know the derivatives of } \psi, \theta, \phi$$

26/20

Unit quaternions and the angular velocity vector

$$\begin{cases} \dot{\mu} = -\frac{1}{2} \underline{\varepsilon}^T \underline{\omega} \\ \dot{\underline{\varepsilon}} = \frac{1}{2} (\mu \mathbb{I}_{3 \times 3} - [\underline{\varepsilon} \times]) \underline{\omega} \end{cases}$$

Inverse mapping

$$\underline{\omega} = 2(\mu \underline{\varepsilon} - \underline{\varepsilon} \mu) + 2 \underline{\varepsilon} \times \dot{\underline{\varepsilon}}$$

Keep in mind that doing a forward Euler integration you could spoil the relation:

$$\mu^2 + \|\underline{\varepsilon}\|^2 = 1$$

If you then use this not unit quaternion you would end up with a non rotation matrix so you need to normalize the unit quaternion before calculating the rotation matrix.

Euler angles do not have these problems, but they suffer from singularities.

For example to describe the rotation of an object to another with a unit quaternion is not too easy, while using yaw-pitch-roll it's easier.

23/20 - KINEMATICS

Revise rotation

Vectors can be described as pure geometrics so an abstract thing.

The geometric idea is useful to work models.

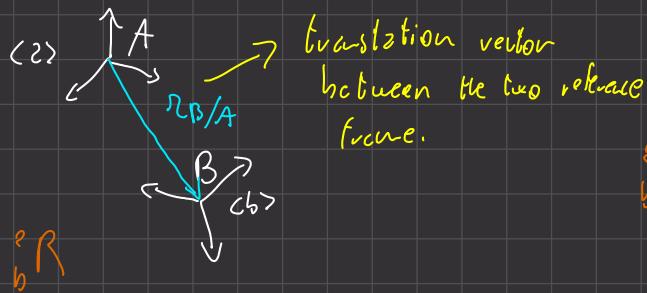
Then we have the algebraic representation

We are going to learn how to handle moving vectors in space.

We can observe these moving vectors from a frame called observer.

Each one of the body comprising a robot is home of one reference frame.

A, B, C are points
 $\{c_2, c_b\}$ are ref. frame



R_b this represent the axis of frame c_b as seen from frame c_2

The axis of frame c_2 will be described in two ways:

$$\left\{ \begin{array}{l} (\underline{i}_2, \underline{j}_2, \underline{k}_2) \\ (\underline{e}_1^2, \underline{e}_2^2, \underline{e}_3^2) \end{array} \right. \begin{array}{l} \text{in frame } 7 \\ \text{first vector} \end{array}$$

For example in frame B this have we:

$$\left\{ \begin{array}{l} (\underline{i}_1, \underline{j}_1, \underline{k}_1) \\ (\underline{e}_1^2, \underline{e}_2^2, \underline{e}_3^2) \end{array} \right.$$

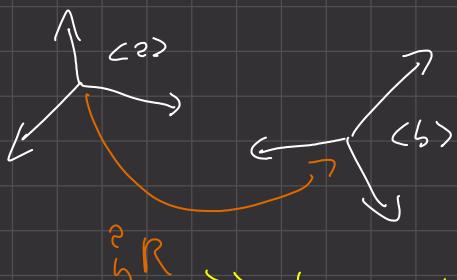
So \mathbb{R} can be written as:

$${}^2\mathbb{R} = \left[\left(e_i^1 \cdot e_j^2 \right) \right] = \begin{bmatrix} e_1^1 & e_1^2 \\ e_2^1 & e_2^2 \\ e_3^1 & e_3^2 \end{bmatrix} \begin{bmatrix} e_1^1 & e_2^2 \\ * & * \\ * & * \end{bmatrix} \begin{bmatrix} e_1^1 & e_2^2 \\ e_1^2 & e_3^2 \\ e_2^2 & e_3^2 \end{bmatrix}$$

${}^2e_1 = {}^1e_1$

${}^2e_2 = {}^1e_2$

${}^2e_3 = {}^1e_3$



→ This is also a way to indicate how to rotate vectors.

We can also define an homogeneous transformation:

$${}^2\mathbb{T} = \begin{bmatrix} {}^2\mathbb{R} & \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \\ \begin{pmatrix} 0 & 0 & 0 & 1 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \end{bmatrix}$$

sometimes this is called \mathbb{H}

Let's suppose that we get a vector $\underline{b_m}$.

We might be interested to transform this vector in frame ${}^2\mathbb{R}$

To do that first we have to write it in homogeneous coordinates:

$$\underline{\underline{b_m}} = \begin{bmatrix} \underline{b_m} \\ 1 \end{bmatrix}$$

$${}^2\mathbb{T} \underline{\underline{b_m}} = {}^2\underline{\underline{m}} = \begin{bmatrix} {}^2\mathbb{R} \underline{b_m} + {}^2\mathbb{r} \\ 1 \end{bmatrix} \rightarrow \underline{\underline{m}}$$

Assumptions:

- All measurements are consistent
- Non consider relativistic effects

Let's consider a vector that is changing its length in time.

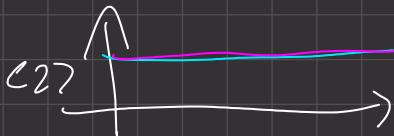
$\mu(t)$

Closed example

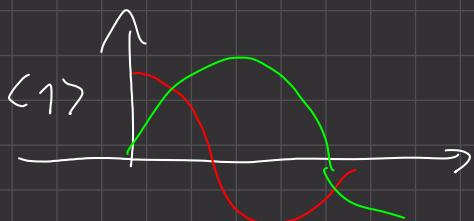
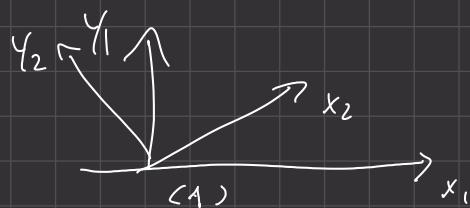


If we consider the frame of the axis the orientation

of the vectors is constant so if we map them



Now if I take another observer:



Now if I take the vectors taken from frame

$$\text{I get } \frac{d_2}{dt} x_2 = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

the first

If I take the derivatives from frame 1:

$$\frac{d}{dt} \underline{x}_2 = \begin{bmatrix} -\sin(\alpha) \\ \cos(\alpha) \end{bmatrix}$$

The derivatives are not absolute scalar functions

$$\mu(t) \rightarrow \frac{d}{dt} \underline{\mu}(t) \text{ in frame 2} = \frac{d}{dt} \sum_{i=1}^3 \mu_i(t) e_i^2 \quad \left\{ x, i_1 + y_1 j_1 + z_1 k_1 \right\}$$

$$= \sum_{i=1}^3 \left(\frac{d}{dt} \mu_i(t) e_i^2 \right) + \sum_{i=1}^3 \mu_i(t) \cdot \frac{d^2}{dt^2} e_i^2$$

these are the same
common for all the observer. These are p
because represent how the frame
changes with respect
with himself.

equivalent

$$= \frac{d}{dt} \sum_{i=1}^3 {}^b \mu_i(t) e_i^b = \sum_{i=1}^3 \left(\frac{d}{dt} {}^b \mu_i(t) \right) e_i^b + \sum_{i=1}^3 {}^b \mu_i(t) \frac{d}{dt} e_i^b$$

We can write that:

$${}^b \left(\frac{d}{dt} \underline{\mu} \right) = \begin{pmatrix} \frac{d}{dt} {}^b \mu_1 \\ \frac{d}{dt} {}^b \mu_2 \\ \frac{d}{dt} {}^b \mu_3 \end{pmatrix}$$

We know we want to compute the time derivative in frame 2 of a quantity expressed in frame b.

$$\frac{d^2}{dt^2} \underline{\mu}(t) = \frac{d^2}{dt^2} \sum_{i=1}^3 {}^b \mu_i(t) e_i^b =$$

vector in b composed by derivatives of scalar functions that
are invariant to the reference frame.

$$= \sum_{i=1}^3 \left(\frac{d^2}{dt^2} {}^b \mu_i(t) \right) e_i^b + \sum_{i=1}^3 {}^b \mu_i(t) \left(\frac{d^2}{dt^2} e_i^b \right)$$

So now we can write:

$$\left(\begin{array}{c} \frac{d}{dt} \underline{\mu} \\ \hline \end{array} \right) = \left(\begin{array}{c} \frac{d}{dt} \underline{\mu}_1 \\ \frac{d}{dt} \underline{\mu}_2 \\ \frac{d}{dt} \underline{\mu}_3 \end{array} \right) = \frac{d^2}{dt^2} \left(\begin{array}{c} \underline{\mu}_1(t) \\ \underline{\mu}_2(t) \\ \underline{\mu}_3(t) \end{array} \right)$$

Now we know that:

$$\frac{d^2}{dt^2} \underline{\mu}(t) = \frac{d}{dt} \underline{\mu}(t) + \sum \underline{\mu}_i(t) \left(\frac{d}{dt} \underline{e}_i \right)$$

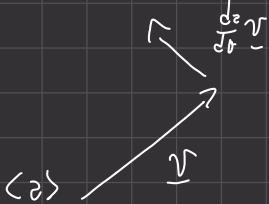
NOTE

Consider a vector \underline{v} s.t. it's $|\underline{v}| = \underline{v} \cdot \underline{v} = 1$

Let's now compute the time derivative:

$$\frac{d}{dt} (\underline{v} \cdot \underline{v}) = \phi \Rightarrow \underline{v} \cdot \frac{d}{dt} \underline{v} = 0$$

This tells us that the time derivative of a vector won't be orthogonal to the vector. The vector can only rotate, not change its length.



We can think of associating a cross product $\underline{\omega} \left(\frac{d}{dt} \underline{v} \right)$.

So we assume that I can find a vector $\underline{\omega}$ that:

$$\underline{v} \cdot \underbrace{\frac{d}{dt} \underline{v}}_{\underline{\omega} \times \underline{v}}$$

Assume $\omega_1 \neq \omega_2$ so $\omega_1 \times v = \omega_2 \times v$

$$(\omega_1 - \omega_2) \times v = 0 \quad \forall v \text{ so } \omega \text{ is unique}$$

ω represent the angular velocity of v to respect to b .

Now we have 2 way to describe:

$$\sum b_i \mu_i(t) \left(\underbrace{\frac{d}{dt} e_i^b}_{\omega_b \times e_i^b} \right)$$

$$\uparrow$$

$\omega_b \times e_i^b \rightarrow$ the angular velocity of frame b in respect to $?$

$$\begin{aligned} \frac{d}{dt} \mu(t) &= \frac{db}{dt} \mu(t) + \sum \omega_{b/2} \times (\mu_i \cdot e_i^b) = \\ &= \frac{d}{dt} \mu(t) + \omega_{b/2} \times \underbrace{\sum_{i=1}^3 (\mu_i \cdot e_i^b)}_{\mu(t)} \end{aligned}$$

We have just derived the law that associates a derivative from 2 frame to another.

the mother of all laws

$$\boxed{\frac{d}{dt} \mu(v) = \frac{db}{dt} \mu(v) + \omega_{b/2} \times \mu(v)}$$

value of vector
is the object value.

Sythetic representation this means scalar multiplying what's inside [] by e_1^b, e_2^b, e_3^b

$$\boxed{\left[\frac{d}{dt} \mu(t) \right] = {}^b_R \left[\frac{db}{dt} \mu(t) \right] + {}^b_R \left[\omega_{b/2} \times \mu(t) \right]}$$

this is what they

$$\frac{d}{dt} \begin{bmatrix} {}^2\mu_1(t) \\ {}^2\mu_2(t) \\ {}^2\mu_3(t) \end{bmatrix} = \frac{d}{dt} \begin{bmatrix} {}^b\mu_1(t) \\ {}^b\mu_2(t) \\ {}^b\mu_3(t) \end{bmatrix} \xrightarrow{3 \times 3} \begin{bmatrix} {}^2\omega_{b/2} x \\ {}^2\mu(t) \end{bmatrix}$$

skew symmetric

$$\frac{d}{dt} \mu = \sum \left(\frac{d}{dt} {}^e\mu_i \right) e_i^2 + \dots$$

this is always a standard derivative
because these quantities are scalar

$$\frac{d}{dt} {}^b\mu = \sum \left(\frac{d}{dt} {}^b\mu_i \right) e_i^b$$

$$\frac{d}{dt} {}^2\mu = \frac{d}{dt} \left(\sum {}^b\mu_i e_i^b \right) = \sum \left(\frac{d}{dt} {}^b\mu_i \right) {}^b e_i + \sum {}^b\mu_i \left(\frac{de}{dt} e_i^b \right)$$

Let's now consider an example.

$$\frac{d}{dt} {}^2 e_i^b = \frac{d}{dt} {}^b e_i^b + \left({}^2\omega_{b/2} x e_i^b \right) \quad \forall i = 1, 2, 3$$

\downarrow
this is ϕ

$${}^2 \left(\frac{d}{dt} {}^b e_i^b \right) = {}^2 \dot{R} = \begin{bmatrix} {}^2\omega_{b/2} x \end{bmatrix} {}^2 R$$

some trickery as before

Properties:

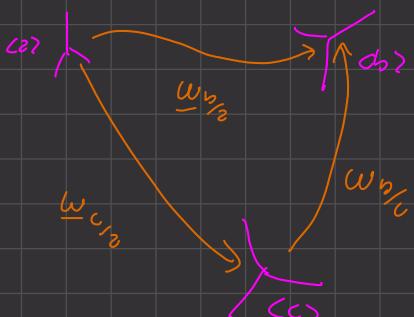
identity

$$) \quad {}^2 R {}^2 R = \mathbb{I} \rightarrow \frac{d}{dt} \left({}^2 R {}^2 R \right) = [\phi] \Rightarrow {}^2 R {}^2 R + {}^2 R {}^2 R = [\phi].$$

$${}^2 \dot{R} {}^2 R = - {}^2 R {}^2 \dot{R} = - {}^2 R \left[{}^2 \omega_{b/2} \right] {}^2 R$$

$${}^2 \dot{R} = - {}^2 R \left[{}^2 \omega_{b/2} \right]$$

2) Composition of angular velocities



$$\begin{cases} \frac{d\omega}{dt} \mu = \frac{d\omega_b}{dt} \mu + \omega_{b/2} \times \mu \\ - \frac{d\epsilon}{dt} \mu = \frac{d\omega_b}{dt} \mu + \omega_{b/2} \times \mu \\ + \frac{d\phi}{dt} \mu = \frac{d\omega_c}{dt} \mu + \omega_{c/2} \times \mu \end{cases}$$

$$\cancel{\frac{d\omega}{dt} \mu} + \cancel{\frac{d\epsilon}{dt} \mu} - \cancel{\frac{d\phi}{dt} \mu} = \cancel{\frac{d\omega_b}{dt} \mu} + \frac{d\omega_b}{dt} \mu + \left[(\omega_{c/2} + \omega_{b/2}) - \omega_{b/2} \right] \times \mu$$

$$[(\omega_{c/2} + \omega_{b/2}) - \omega_{b/2}] \times \mu = \emptyset$$

The assumed μ arbitrary so this implies that the square brackets is a null vector:

||

$$\omega_{b/2} = \omega_{c/2} + \omega_{b/2}$$

Angular velocities are additive quantities.

Very easy equality:

$$\omega_{2/2} = \omega_{b/2} + \omega_{2/2} \Rightarrow \omega_{2/2} = -\omega_{b/2}$$

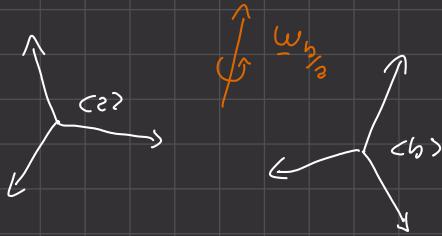
Knowing this we can immediately write the previous equation

$$\overset{b}{\omega} \overset{2}{R} \overset{2}{\dot{R}} = -\overset{b}{\omega} \overset{2}{R} \overset{2}{\dot{R}} = -\overset{b}{\omega} \overset{2}{R} [\overset{2}{\omega}_{b/2}] \overset{2}{\dot{R}} = \overset{b}{\omega} \overset{2}{R} [\overset{2}{\omega}_{b/2}]$$

Note

$$\overset{2}{A} = \overset{b}{\omega} \overset{2}{R} \overset{2}{A} \overset{b}{\omega}$$

30/20



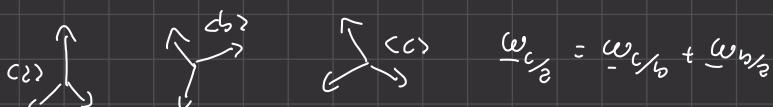
last lesson we proved that if we have a generic time varying vector then we can compute the derivative in both frames and we have relations between different frames

$$\frac{d^2}{dt^2} \mu = \frac{db}{dt} \mu(b) + \underbrace{\omega_{b/s} \times \mu(b)}_{\text{this takes into consideration the relative movement}}$$

$$\frac{d^2}{dt^2} \mu = {}^b R \frac{d^2}{dt^2} {}^b \mu + \left[{}^b \omega_{b/s} \times \right] \underbrace{{}^b \mu(t)}_{{}^b R \mu^b(t)}$$

the stepdown equation integrating?

$$\left. \begin{aligned} {}^b R &= \left[{}^2 \omega_{b/s} \times \right] {}^b R \\ {}^b \dot{R} &= \left[{}^b \omega_{b/s} \times \right] {}^b R \end{aligned} \right\} \text{we have exploited that } \omega_{c/b} = -\omega_{b/c}$$



ENVIRONMENT

RECAP



trajectory

$$\underline{\underline{r}}_{P/A} \triangleq (P-A)$$

$$\underline{\underline{r}}_{P/B} \triangleq (P-B)$$

$$\underline{\underline{r}}_{B/A} \triangleq (B-A)$$

frame b is both rotating and translating with respect to frame c2s.

$$\underline{\underline{r}}_{P/A} = \underline{\underline{r}}_{P/B} + \underline{\underline{r}}_{B/A}$$

position of p with respect to A

we also have:

$${}^2 R = \left[e_i^a \cdot e_j^b \right]$$

We want to compute velocity of P for observer in frame 2

$$\frac{d}{dt} {}_2 r_{P/A} = \frac{d}{dt} \left[{}_2 r_{P/B} + {}_2 r_{B/A} \right] = \frac{d}{dt} {}_2 r_{P/B} + \frac{d}{dt} {}_2 r_{B/A}$$

$\frac{d}{dt} {}_2 r_{P/B} + \omega_{b/2} \times {}_2 r_{P/B}$

$v_{P/b}$

This can be called
 $v_{B/2}$
velocity of point B computed
in frame 2.

NOTE: a is lowercase

Now we can say that:

$$\frac{d}{dt} {}_2 r_{P/A} = v_{P/b}$$

So now,

this are due to the effect of the motion of b

$$v_{P/b} = v_{P/b} + \omega_{b/2} \times {}_2 r_{P/B} + v_{B/b}$$

Galileo theorem for velocity distribution

$$v_{P/b} = v_{P/b} + \omega_{b/2} \times {}_2 r_{P/B} + v_{B/b}$$

$$\begin{aligned} \frac{d}{dt} v_{P/b} &= \frac{d}{dt} \left[v_{P/b} + \omega_{b/2} \times {}_2 r_{P/B} + v_{B/b} \right] = \\ &= \frac{d}{dt} v_{P/b} + \frac{d}{dt} \omega_{b/2} \times {}_2 r_{P/B} + \frac{d}{dt} v_{B/b} \\ &= \underbrace{\frac{d}{dt} v_{P/b}}_{\text{acceleration } \underline{\Omega} P/b} + \underbrace{\omega_{b/2} \times v_{P/b}}_{\text{angular acceleration of } b \text{ with respect to } 2} + \underbrace{\frac{d}{dt} v_{B/b}}_{\underline{\Omega} B/2} \end{aligned}$$

$$\underline{\Omega}_{P/B} = \underline{\Omega}_{P/B} + 2 \underbrace{\omega_{B/Z}}_{\text{Coriolis acc.}} \times \underline{v}_{P/B} + \left(\frac{d\omega}{dt} \right) \underline{\omega}_{B/Z} \times \underline{r}_{P/B} + \underbrace{\omega_{B/Z} \times [\omega_{B/Z} \times \underline{r}_{P/B}]}_{\text{centrifugal acc.}} + \underline{\Omega}_{B/Z}$$

Make some notes on the previous formula:

Note:

$$\frac{d\omega}{dt} \underline{\omega}_{B/Z} = \frac{d\omega}{dt} \underline{\omega}_{B/Z} + \cancel{\omega_{B/Z} \times \underline{\omega}_{B/Z}} = \phi \Rightarrow \frac{d\omega}{dt} \underline{\omega}_{B/Z} = \dot{\omega}_{B/Z}$$

\hookrightarrow let's work assume
this is a generic
vector $\omega(t)$

this allows us
to write this

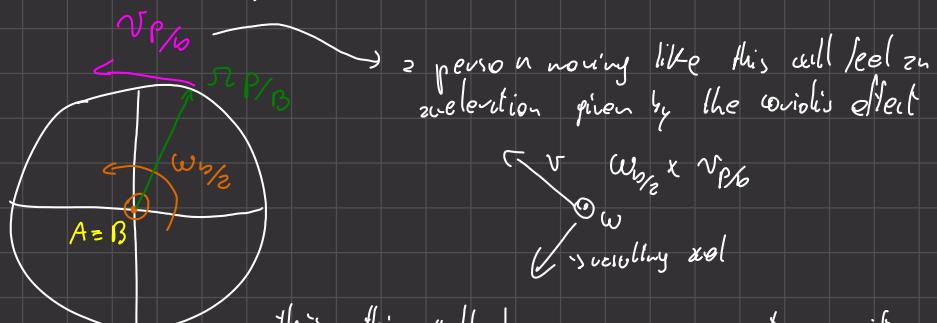
ω is in general the time derivative of a vector and ω is the time derivative of a vector. Note the rate of change of the angular velocity is the same for both observers.

$$\underline{\Omega}_{P/Z} = \underline{\Omega}_{P/B} + 2 \underbrace{\omega_{B/Z} \times \underline{v}_{P/B}}_{\text{Coriolis acc.}} + \cancel{\omega_{B/Z} \times \underline{r}_{P/B}} + \underbrace{\omega_{B/Z} \times (\omega_{B/Z} \times \underline{r}_{P/B})}_{\text{cent. acc.}} + \underline{\Omega}_{B/Z}$$

NOTE:

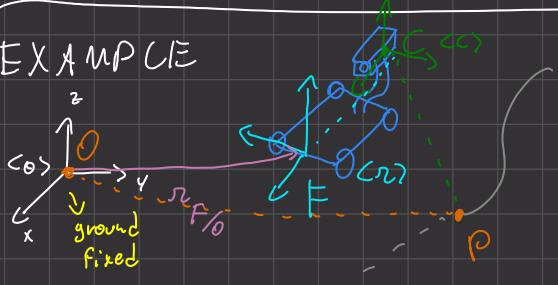
Sometime in science fiction movies you see elements in artificial gravity in centrifuge to generate centrifugal gravity.

This is done by using this quantity.



this thing will become even more strange if you start jumping.

EXAMPLE



Now this robot is moving along a certain trajectory. Now on top of this robot there is a camera.

The movement of the robot will be on the x-y plane.

Using the measurements of P we would like to give a quantitative measurement of how P is moving with respect to the ground station.

We assume that the camera will not move with respect to the platform.

We want to predict the motion of P in space given the camera movement.

We want to compute: $\underline{r}_{P/\theta}$, $\underline{v}_{P/\theta}$, $\underline{\alpha}_{P/\theta}$

- assume these quantities are known at $t = t_m$ \rightarrow time of the measurement.
- assume these is "smooth" time function
 $\underbrace{\text{differentiable and to third derivative continuity}}$

GOAL: "predict" $\underline{r}_{P/\theta}$ for $t > t_m$

Given that these functions are smooth we can use second order Taylor expansion:

$$\underline{r}_{P/\theta}(t) \approx \underline{r}_{P/\theta}(t_m) + \underline{v}_{P/\theta}(t_m)(t - t_m) + \frac{1}{2} \underline{\alpha}_{P/\theta}(t_m)(t - t_m)^2$$

In practice we are going to implement:

$$\underline{r}_{P/\theta} \approx \underline{r}_{C/F} + \underline{v}_{C/F}(t - t_m) + \frac{1}{2} \underline{\alpha}_{C/F}(t - t_m)^2$$

Note - that C and F are positioned arbitrary on the robot

- we assume that the camera is calibrated with respect to frame C/F (C_1, C_2, C_3) are calibrated w.r.t. (F, C_{22})

- assume $\underline{r}_{C/F}$ Known, $\underline{v}_{C/F}$ Known

$$\textcircled{1} \quad \underline{r}_{P/\theta} = \underline{r}_{P/C} + \underline{r}_{C/F} + \underline{r}_{F/\theta}$$

$$\underline{C/R} = \underline{r}_{C/F} + \underline{R}$$

\uparrow Known from calibration
assumed Known

\textcircled{2} Computing:

$$\underline{v}_{P/\theta} = \frac{d}{dt} \underline{r}_{P/\theta} = \frac{d}{dt} \left[\underline{r}_{P/C} + \underline{r}_{C/F} + \underline{r}_{F/\theta} \right] =$$

$$= \left[\frac{d}{dt} \underline{r}_{P/C} + \omega_{C/F} \times \underline{r}_{C/F} \right] + \left[\frac{d}{dt} \underline{r}_{C/F} + \omega_{F/\theta} \times \underline{r}_{F/\theta} \right] + \underline{v}_{F/\theta}$$

velocity of P is estimated by the camera

$$\omega_{C/F} + \omega_{F/\theta} \rightarrow \text{this is } 0 \text{ if the camera is not rotating.}$$

$$\underline{\nu}_{P/O} = \underline{\nu}_{P/C} + \underline{\omega}_{R/O} \times (\underline{\nu}_{P/C} + \underline{\nu}_{C/P}) + \cancel{\frac{d}{dt} \underline{\nu}_C} \quad \text{for us because the center is fixed}$$

$$^0\nu_{P/O} = {}^0R {}^C\nu_{P/C} + [{}^0\omega_{R/O} \times] [{}^0{}_C R {}^R \underline{\nu}_{P/C} + {}^0R {}^R \underline{\nu}_{C/P}] + {}^0\nu_{F/O}$$

try doing the anal. by yourself.

Final notes:

$$\frac{d\theta}{dt} \underline{\nu}_{F/O} = - \frac{d\theta}{dt} \underline{\nu}_{O/F} = \frac{d\underline{\nu}}{dt} \underline{\nu}_{O/F} + \underline{\omega}_{R/O} \times \underline{\nu}_{O/F}$$