Are we ready for Drone Racing? Optimal Control and Trajectory Estimation

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Abstract—In light of the recent interest in drone racing, both manned and unmanned, we present an approach to calculate the optimal control inputs and state trajectory of a quadcopter to finish a given race track in a minimum time. We assume the race track to consist of a set of arbitrarily positioned gates in 3D that the drone has to go through. We show how to solve this problem by using GPOPS-II, a matlab-based general purpose optimal control software. The hope of this work is to provide useful information for human pilots, and developers of autonomous drones, as would be provided by a tactician in an air race competition.

Index Terms—Optimal Control, Optimal Trajectory, Quadcopter, Drone Racing

SUPPLEMENTARY MATERIAL

Videos of the experiments: **TODO:** Add video url Open-source code: **TODO:** Add code url

I. Introduction

THE advent of drone racing competitions, both manned and unmanned, requires the estimation of the fastest possible trajectory through a given race track. Finding this optimal trajectory can be useful for human pilots, and their tacticians, to evaluate their performance and improve upon it. For example, the teams in the Red Bull Air Race [1] make use of such information to train and perfect their maneuvers. In this same way, drones can autonomously follow a dynamically feasible trajectory with high accuracy. This is especially interesting for autonomous drone races such as the Alpha Pilot race [2]. Therefore, being able to infer the optimal feasible trajectory could potentially make fully-autonomous drones beat human pilots. We are still not there, but great progress in this direction is being made [3].

Contributions.

- Formulation of the minimum-time problem for the specific problem of drone-racing through multiple gates.
- Presentation of a means to solve the problem using readily available software.

Paper Structure. Section II presents the mathematical formulation of our approach, and discusses the assumptions made. Section III reports and discusses the experimental results.

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II. APPROACH

We formulate our drone race optimization problem as a minimum time optimal control problem with nonlinear dynamics and constraints. Therefore, we must define both the dynamics of a drone and the constraints to pass through each gate. Section II-A, II-B and II-C present the nonlinear dynamics for a quadcopter. Section II-D casts the quadcopter dynamics using a state-space model. The state-space formulation allows us to present in a compact form the optimization problem that we detail in section II-E.

A. Mathematical Model of a Quadcopter

The absolute position ξ of the drone is given in an Euclidean inertial frame, referred to as *world frame* (G), by: x,y,z. This corresponds to the 3D coordinates of the center of mass of the drone. The attitude η of the drone with respect to the world frame is given using Euler angles: pitch θ , roll ϕ , and yaw ψ .

$$oldsymbol{\xi} = \left[egin{array}{c} x \ y \ z \end{array}
ight], \quad oldsymbol{\eta} = \left[egin{array}{c} \phi \ heta \ \psi \end{array}
ight]$$

Figure 1 shows the frame of reference of the drone, referred to as body frame (B), with respect to the world frame.

Linear velocities (body frame): V_B Angular velocities (body frame): ν

$$oldsymbol{V}_B = \left[egin{array}{c} v_{x,B} \ v_{y,B} \ v_{z,B} \end{array}
ight], \quad oldsymbol{
u} = \left[egin{array}{c} p \ q \ r \end{array}
ight]$$

The rotation matrix from the body frame B to the world frame G is given by: $\mathbf{X}^B = \mathbf{R}_G^B \mathbf{X}^G = \mathbf{R}(\phi) \mathbf{R}(\theta) \mathbf{R}(\psi) \mathbf{X}^G$

$$\boldsymbol{R}_{B}^{G} = \boldsymbol{R} = \begin{bmatrix} C_{\psi}C_{\theta} & C_{\psi}S_{\theta}S_{\phi} - S_{\psi}C_{\phi} & C_{\psi}S_{\theta}C_{\phi} + S_{\psi}S_{\phi} \\ S_{\psi}C_{\theta} & S_{\psi}S_{\theta}S_{\phi} + C_{\psi}C_{\phi} & S_{\psi}S_{\theta}C_{\phi} - C_{\psi}S_{\phi} \\ -S_{\theta} & C_{\theta}S_{\phi} & C_{\theta}C_{\phi} \end{bmatrix}$$

where $S_x = \sin(x)$ and $C_x = \cos(x)$.

Time derivative of this rotation matrix provides us with the angular velocities. These are not simply the time derivatives of the independent Euler angles, instead:

$$\dot{m{\eta}} = m{W}_{\eta}^{-1} m{
u}, \quad egin{bmatrix} \dot{\phi} \ \dot{ heta} \ \dot{\psi} \end{bmatrix} = egin{bmatrix} 1 & S_{\phi} T_{ heta} & C_{\phi} T_{ heta} \ 0 & C_{\phi} & -S_{\phi} \ 0 & S_{\phi}/C_{ heta} & C_{\phi}/C_{ heta} \end{bmatrix} egin{bmatrix} p \ q \ r \end{bmatrix}$$

Conversely

$$oldsymbol{
u} = oldsymbol{W}_{\eta} \dot{oldsymbol{\eta}}, \qquad \left[egin{array}{ccc} p \\ q \\ r \end{array}
ight] = \left[egin{array}{ccc} 1 & 0 & -S_{ heta} \\ 0 & C_{\phi} & C_{ heta}S_{\phi} \\ 0 & -S_{\phi} & C_{ heta}C_{\phi} \end{array}
ight] \left[egin{array}{ccc} \dot{\phi} \\ \dot{ heta} \\ \dot{\psi} \end{array}
ight]$$

where $T_x = \tan(x)$. The matrix W_{η} is invertible if $\theta \neq (2k-1)\phi/2, (k \in \mathbb{Z})$

Symbol	Definition
ξ	Absolute position (inertial frame)
η	Attitude (inertial frame)
V_B	Linear velocities (body frame)
ν	Angular velocities (body frame)
\mathbf{R}	Rotation matrix from body to inertial frame
\mathbf{W}_{n}	Transformation matrix for angular velocities from inertial to
'	body frame
I	Inertia matrix
G	Gravity

TABLE I: Definitions and Notations

Assumptions:

• The quadrotor structure is rigid and symmetrical with a center of mass aligned with the center of the body frame of the vehicle. The four arms of the quadcopter are aligned with the body x- and y-axes. Therefore, the inertia matrix I is diagonal:

$$m{I} = \left[egin{array}{ccc} I_{xx} & 0 & 0 \ 0 & I_{yy} & 0 \ 0 & 0 & I_{zz} \end{array}
ight]$$

We further have that due to symmetry, that the inertial components in the x- and y-axes are equal: $I_{xx} = I_{yy}$.

- The thrust and drag of each motor is proportional to the square of the motor velocity.
- The propellers are considered to be rigid and therefore blade flapping is negligible (deformation of propeller blades due to high velocities and flexible material).
- The Earth is flat and non-rotating (difference of gravity by altitude or the spin of the earth is negligible).
- Ground effects that increase lift and decrease aerodynamic drag when flying close to the ground is considered negligible.

These assumptions and basic dynamics lead to the model used in this work.

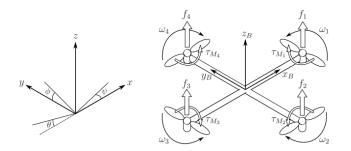


Fig. 1: The inertial and body frames of a quadcopter using the East North Up convention (ENU). Figure from [4].

B. Quadcopter Dynamics

The following derivations are a summary of the equations in [4].

• The angular velocity of rotor i, denoted with ω_i , creates force f_i in the direction of the rotor axis:

$$f_i = k\omega_i^2$$

where the lift constant is k. The combined forces of rotors create thrust T in the direction of the body z-axis (represented as T_B).

$$T = \sum_{i=1}^{4} f_i = k \sum_{i=1}^{4} \omega_i^2, \qquad T_B = \begin{bmatrix} 0 \\ 0 \\ T \end{bmatrix}$$

• The angular velocity and acceleration of the rotor also creates torque τ_{M_i} around the rotor axis:

$$I_M \dot{\omega}_i = \tau_{M_i} - b\omega_i^2$$

in which the drag constant is b and the inertia moment of the rotor is I_M . Usually the acceleration of the rotor $(\dot{\omega}_i)$ is considered small and thus it is omitted. This holds when we assume that the quadrotor is operating in stable flight and that the propellers are maintaining a constant thrust and not accelerating. This assumption results in the torque about the global z axis being equal to the torque due to drag.

Torque τ_B consists of the torques on the three axis $\tau_{\phi}, \tau_{\theta}, \tau_{\psi}$:

$$\begin{split} \boldsymbol{\tau}_{B} &= \left[\begin{array}{c} \boldsymbol{\tau}_{\phi} \\ \boldsymbol{\tau}_{\theta} \\ \boldsymbol{\tau}_{\psi} \end{array} \right] = \left[\begin{array}{c} lk \left(-\omega_{2}^{2} + \omega_{4}^{2} \right) \\ lk \left(-\omega_{1}^{2} + \omega_{3}^{2} \right) \\ \sum_{i=1}^{4} (-1)^{i+1} \boldsymbol{\tau}_{M_{i}} \end{array} \right] \\ &= \left[\begin{array}{c} lk \left(-\omega_{2}^{2} + \omega_{4}^{2} \right) \\ lk \left(-\omega_{1}^{2} + \omega_{3}^{2} \right) \\ b(\omega_{1}^{2} - \omega_{2}^{2} + \omega_{3}^{2} - \omega_{4}^{2}) \end{array} \right] \end{split}$$

where l is the distance between the rotor and the center of mass of the quadcopter.

1) Translational dynamics: Newton's second law:

$$m\ddot{\boldsymbol{\xi}} = \boldsymbol{G} + \boldsymbol{R}\boldsymbol{T}_{B}$$

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \end{bmatrix} = -g \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} + \frac{T}{m} \begin{bmatrix} C_{\psi} S_{\theta} C_{\phi} + S_{\psi} S_{\phi} \\ S_{\psi} S_{\theta} C_{\phi} - C_{\psi} S_{\phi} \\ C_{\theta} C_{\phi} \end{bmatrix}$$

2) Rotational dynamics: The rotational equations of motion are defined in the body frame so that the rotations can be computed about the quadrotor's center and not the center of the global coordinate frame. Applying Euler's second law:

$$I\dot{\nu} + \nu \times (I\nu) + \Gamma = \tau$$

where we have:

- Angular acceleration of the inertia: $I\dot{v}$
- Centripetal forces: $\nu \times (I\nu)$
- Gyroscopic forces: Γ
- External torque: τ_B

Replacing terms by their definitions, multiplying both sides by I^{-1} , and rearranging:

$$\dot{m{
u}} = m{I}^{-1} \left(- \left[egin{array}{c} p \\ q \\ r \end{array}
ight] imes \left[egin{array}{c} I_{xx}p \\ I_{yy}q \\ I_{zz}r \end{array}
ight] - I_r \left[egin{array}{c} p \\ q \\ r \end{array}
ight] imes \left[egin{array}{c} 0 \\ 0 \\ 1 \end{array}
ight] \omega_{\Gamma} + m{ au}_{B}
ight)$$

$$\begin{bmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{bmatrix} = \begin{bmatrix} (I_{yy} - I_{zz}) qr/I_{xx} \\ (I_{zz} - I_{xx}) pr/I_{yy} \\ (I_{xx} - I_{yy}) pq/I_{zz} \end{bmatrix} - I_r \begin{bmatrix} q/I_{xx} \\ -p/I_{yy} \\ 0 \end{bmatrix} \omega_{\Gamma}$$

$$+ \begin{bmatrix} \tau_{\phi}/I_{xx} \\ \tau_{\theta}/I_{yy} \\ \tau_{\psi}/I_{zz} \end{bmatrix}$$

where $\omega_{\Gamma} = \omega_1 - \omega_2 + \omega_3 - \omega_4$, and I_r is the rotor moment of inertia.

The angular accelerations $\ddot{\eta}$ in world frame are given by the time derivatives of the angular velocities $(\dot{\eta} = W_n^{-1} \nu)$,

$$\ddot{\boldsymbol{\eta}} = \frac{\mathrm{d}}{\mathrm{d}t} \left(\boldsymbol{W}_{\eta}^{-1} \boldsymbol{\nu} \right) = \frac{\mathrm{d}}{\mathrm{d}t} \left(\boldsymbol{W}_{\eta}^{-1} \right) \boldsymbol{\nu} + \boldsymbol{W}_{\eta}^{-1} \dot{\boldsymbol{\nu}} =$$

$$\begin{bmatrix} 0 & \dot{\phi}C_{\phi}T_{\theta} + \dot{\theta}S_{\phi}/C_{\theta}^{2} & -\dot{\phi}S_{\phi}C_{\theta} + \dot{\theta}C_{\phi}/C_{\theta}^{2} \\ 0 & -\dot{\phi}S_{\phi} & -\dot{\phi}C_{\phi} \\ 0 & \dot{\phi}C_{\phi}/C_{\theta} + \dot{\phi}S_{\phi}T_{\theta}/C_{\theta} & -\dot{\phi}S_{\phi}/C_{\theta} + \dot{\theta}C_{\phi}T_{\theta}/C_{\theta} \end{bmatrix} \boldsymbol{\nu} \\ + \boldsymbol{W}_{\boldsymbol{n}}^{-1}\dot{\boldsymbol{\nu}}.$$

At this point, it is common to do a simplification by setting $[\dot{\phi} \ \dot{\theta} \ \dot{\psi}]^T = \begin{bmatrix} p & q & r \end{bmatrix}^T$, which holds true for small angles of movement [5]. Nevertheless, in this work we omit such a simplification.

C. Aerodynamic Drag

We include the drag force generated by the air resistance. For this, we add a diagonal coefficient matrix that associates the linear velocities to the force slowing down the movement

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \end{bmatrix} = -g \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} + \frac{T}{m} \begin{bmatrix} C_{\psi} S_{\theta} C_{\phi} + S_{\psi} S_{\phi} \\ S_{\psi} S_{\theta} C_{\phi} - C_{\psi} S_{\phi} \\ C_{\theta} C_{\phi} \end{bmatrix} - \frac{1}{m} \begin{bmatrix} A_{x} & 0 & 0 \\ 0 & A_{y} & 0 \\ 0 & 0 & A_{z} \end{bmatrix} \begin{bmatrix} \dot{x} \\ \dot{y} \\ \dot{z} \end{bmatrix}$$

in which A_x, A_y and A_z are the drag force coefficients for velocities in the corresponding directions of the inertial frame. Several other aerodynamical effects could be included in the model. For example, dependence of thrust on angle of attack, blade flapping and airflow disruptions. We refrain from adding these for simplicity of the model. We also ignore rotational drag forces since we assume rotational velocities to be small. Alternatively, we could add the components $\tau_{\boldsymbol{w}} = \left[\begin{array}{ccc} \tau_{wx} & \tau_{wy} & \tau_{wz} \end{array}\right]^T$ to the overall torque.

D. State-space Model

We can write our nonlinear dynamics using the following state vector:

$$\mathbf{X}^T = \begin{bmatrix} x & y & z & \dot{x} & \dot{y} & \dot{z} & \phi & \theta & \psi & p & q & r \end{bmatrix}^T$$

Moreover, we define our inputs as follow:

- U_1 : the resulting thrust of the four rotors.
- U_2 : the difference of thrust between the motors on the x axis which results in roll angle changes and subsequent movement in the lateral x direction.

- U₃: the difference of thrust between the motors on the y axis which results in pitch angle changes and subsequent movement in the lateral y direction.
- U_4 : the difference of torque between the clockwise and counterclockwise rotors which results in a moment that rotates the quadrotor around the vertical z axis.

This results in the control vector U defined as:

$$U = \begin{bmatrix} U_1 \\ U_2 \\ U_3 \\ U_4 \end{bmatrix} = \begin{bmatrix} T \\ \tau_{\phi} \\ \tau_{\theta} \\ \tau_{\psi} \end{bmatrix} = \begin{bmatrix} k \sum_{i=1}^4 \omega_i^2 \\ lk \left(-\omega_2^2 + \omega_4^2\right) \\ lk \left(-\omega_1^2 + \omega_3^2\right) \\ b(\omega_1^2 - \omega_2^2 + \omega_3^2 - \omega_4^2) \end{bmatrix}.$$

Our nonlinear dynamics can be written as a nonlinear differential equation of the states X and control inputs U:

$$\dot{\mathbf{X}} = f(\mathbf{X}, \mathbf{U}). \tag{1}$$

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Or, more precisely:

E. Minimum Time Optimal Control Problem

Cost: our objective is to minimize the time taken for the drone to pass through all gates in a particular order. Therefore, the cost we try to minimize is simply the final time:

$$J = \int_{t_0}^{t_f} dt = t_f$$

where $t_0 = 0$ is fixed, but t_f is free. We consider that the drone has finished the race once it has reached the last gate.

Nonlinear Dynamics: are formulated using the generic form: $\dot{x}(t) = f(x(t), u(t), t)$.

As detailed in eq. (1), we make use of a nonlinear function of the state $\boldsymbol{X}(t)$ and the input $\boldsymbol{U}(t)$. Although the nonlinear function f may depend explicitly on time t, the drone dynamics do not depend directly on t, therefore the time dependency can be dropped.

Constraints: to simulate a realistic drone we must constraint the control inputs to a range of possible values.

• Generic input constraints can be formulated as:

$$M_i^- \le u_i(t) \le M_i^+,$$

where M_i^- is the lower bound of control input $u_i(t)$, while M_i^+ is its upper bound. In our case, we use the following bounds to the control inputs:

• Path constraints:

$$\forall t \quad t_0 \le t \le t_f, \qquad c(x(t), u(t), t) \le 0,$$

Boundary Conditions: Given

- Initial: $n(x(t_0), t_0) = 0$
- Final: $m(x(t_f), t_f) = 0$

III. EXPERIMENTAL RESULTS

IV. CONCLUSION

V. SUPPLEMENTARY MATERIAL

Euler angle rotation matrices:

$$\mathbf{R}(\psi) = \begin{bmatrix} \cos(\psi) & \sin(\psi) & 0 \\ -\sin(\psi) & \cos(\psi) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\mathbf{R}(\theta) = \begin{bmatrix} \cos(\theta) & 0 & -\sin(\theta) \\ 0 & 1 & 0 \\ \sin(\theta) & 0 & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R}(\phi) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\phi) & \sin(\phi) \\ 0 & -\sin(\phi) & \cos(\phi) \end{bmatrix}$$

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