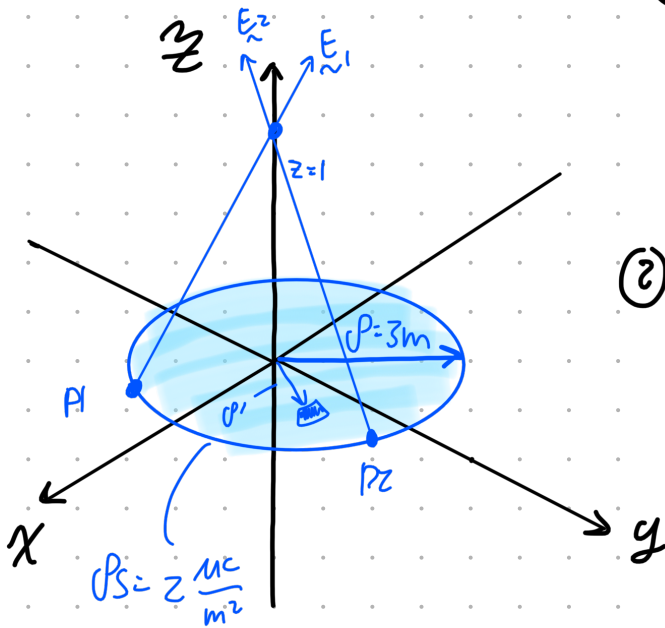


Stefan Tosti

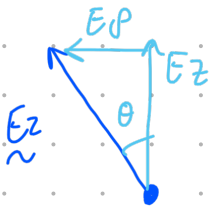
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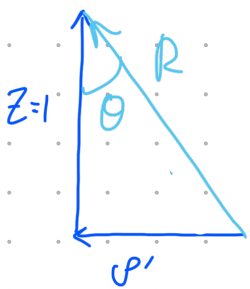
(1) Notice that due to the symmetry of the problem, the ρ component of all \vec{E} will cancel out, leaving only a z component of the field.

(2) The charge of an infinitesimal charge element, dq , can be written as $dq = \rho_s \cdot dS$, and we know that the given surface charges in the ρ and ϕ directions $\Rightarrow dq = \rho_s \cdot \rho' d\rho' d\phi'$



(3) The diagram to the left suggests that the z component of some force would be $E \cos \theta$.

$$\Rightarrow dE_z = dE \cos \theta$$



(4) We can represent θ using the diagram to the left, it will be a constant value, since $z=1$, $\rho=\rho'$ for some arbitrary observation point.

$$\Rightarrow R = \sqrt{(\rho')^2 + (1)^2}$$

$$\Rightarrow \cos \theta = \frac{1}{\sqrt{(\rho')^2 + (1)^2}}$$

(5) The \vec{E} from some arbitrary point is then...

$$dE_z = dE \cos \theta = \frac{dq}{4\pi\epsilon_0 R^2} \cdot \cos \theta = \frac{(\rho_s \cdot \rho' d\rho' d\phi')}{4\pi\epsilon_0 (\sqrt{(\rho')^2 + (1)^2})^2} \cdot \frac{1}{\sqrt{(\rho')^2 + (1)^2}}$$

$$\Rightarrow dE_z = \frac{\rho_s \rho'}{4\pi\epsilon_0 \sqrt{(\rho')^2 + (1)^2}} d\rho' d\phi'$$

⑥ Finally, we can integrate over all ρ and ϕ to determine the total field ---

$$E_z = \int_0^{2\pi} \int_0^1 \frac{\rho^S \rho'}{4\pi\epsilon_0 \sqrt{(\rho')^2 + (1)^2}}^3 d\rho' d\phi'$$

$$= \frac{\rho^S}{4\pi\epsilon_0} \int_0^{2\pi} d\phi \int_0^1 \frac{\rho'}{\sqrt{(\rho')^2 + (1)^2}}^3$$

$$= \left(\frac{2 \times 10^{-6}}{4\pi \left(\frac{10^{-9}}{36\pi} \right)} \right) (2\pi) \left(\frac{-1}{\sqrt{2}} + 1 \right)$$

$$= (0, 0, 3.31 \times 10^4) \text{ V/m}$$