

# How to Model Stars!

Note Title

1/9/2019

## Equations of Stellar Structure

$$(1) \quad \rho \frac{D\underline{u}}{Dt} = -\underline{\nabla} p - \rho \underline{\nabla} \phi \quad (\text{Euler}).$$

$$(2) \quad \frac{D\rho}{Dt} + \rho \underline{\nabla} \cdot \underline{u} = 0 \quad (\text{conservation of mass}).$$

$$(3) \quad \nabla^2 \phi = 4\pi G \rho \quad (\text{gravity})$$

$$(4) \quad \frac{DU}{Dt} - \frac{p}{\rho^2} \frac{D\rho}{Dt} = \epsilon - \frac{1}{\rho} \underline{\nabla} \cdot \underline{F} \quad (\text{energy}).$$

Some basic ("bad") assumptions:

- time independence. and static:  $\underline{u} = 0$ .
- spherical symmetry.

$$(1) \quad \frac{dp}{dr} = -\rho \frac{d\phi}{dr}$$

$$(2) \quad 0 = 0$$

$$(3) \quad \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{d\phi}{dr} \right) = 4\pi G \rho$$

$$(4) \quad \rightarrow \quad 0 = \epsilon - \frac{1}{\rho} \underline{\nabla} \cdot \underline{F}$$

Integrate (3):  $r^2 \frac{d\phi}{dr} = 4\pi G \int_0^r \rho(r') r'^2 dr'$

But the mass enclosed inside a radius  $r$  is

$$M = 4\pi \int_0^r \rho(r') r'^2 dr'$$

$$(3) \rightarrow \frac{d\phi}{dr} = \frac{GM(r)}{r^2}$$

$$(1) \rightarrow \boxed{\frac{dp}{dr} = -\frac{GM\rho}{r^2}} \quad \text{Hydrostatic equilibrium.}$$

But with another relationship between  $\rho$  and  $p$ , we can't solve this equation.

## Polytropic Model.

Assume  $p = k \rho^\gamma$   $k, \gamma$  are constants.

(e.g. for an adiabatic gas,  
 $p \rho^{-\gamma} = \text{constant}$ )

From ①,  $\frac{dp}{dr} = -\frac{GM\rho}{r^2}$

$$\Rightarrow \frac{r^2}{\rho} \frac{dp}{dr} = -GM$$

Take derivative of both sides:

$$\frac{d}{dr} \left( \frac{r^2}{\rho} \frac{dp}{dr} \right) = -G \frac{dM}{dr}$$

From ③,  $dM/dr = 4\pi r^2 \rho$

$$\frac{1}{r^2} \frac{d}{dr} \left( \frac{r^2}{\rho} \frac{dp}{dr} \right) = -4\pi G \rho$$

Use  $p = k \rho^\gamma \Rightarrow \frac{dp}{dr} = k\gamma \rho^{\gamma-1} \frac{d\rho}{dr}$

$$\frac{k\gamma}{r^2} \frac{d}{dr} \left( r^2 \rho^{\gamma-2} \frac{d\rho}{dr} \right) = -4\pi G \rho$$

This is an ODE involving just density.

Clear it up: • define a dimensionless parameter  

$$\rho(r) = \rho_c \rho_n^n(r)$$

• Define  $n$  such that

$$\gamma = \frac{n+1}{n}$$

$$\Rightarrow \left( \frac{(n+1)}{4\pi G} k \rho_c^{\frac{1-n}{n}} \right) \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{d\rho_n}{dr} \right) = -\rho_n^n$$

• Define  $\eta = \frac{r}{\lambda_n}$ ,  $\lambda_n = \left[ \frac{(n+1) k \rho_c^{\frac{1-n}{n}}}{4\pi G} \right]^{-\frac{1}{2}}$

$$\Rightarrow \boxed{\frac{1}{\eta^2} \frac{d}{d\eta} \left( \eta^2 \frac{d\rho_n}{d\eta} \right) = -\rho_n^n} \quad \text{Lane-Emden equation.}$$

Boundary Conditions:

• We want the density to go to zero at some radius - call it  $\eta_s$ .

$$\boxed{\rho_n(\eta_s) = 0}$$

• At the centre of the star, we want the density to be finite:

$$\boxed{\rho_n(0) = 1}$$

[So that  $\rho_c$  is the density at the centre of the star.]

- Consider a small volume at the centre with radius  $\delta$

$$V = \frac{4}{3} \pi \delta^3$$

The mass in this radius is

$$M = \frac{4\pi}{3} \bar{\rho} \delta^3 \quad \rightarrow \text{average density}$$

From ① we have

$$\frac{dp}{dr} = - \frac{GM_p}{r^2} = - \frac{4\pi}{3} G \bar{\rho}^2 \delta$$

$$\text{So as } \delta \rightarrow 0, \quad dp/dr \rightarrow 0$$

So

$$\frac{dp}{dr} \rightarrow 0 \quad \text{as } r \rightarrow 0$$

or

$$\boxed{\frac{dp_r}{d\eta} = 0 \quad \text{at} \quad \eta = 0.}$$

Analytic solutions:

- $n = 0, \quad L_0(\eta) = 1 - \frac{\eta^2}{6}$

The surface is at  $\eta_s = \sqrt{6}$

- $n = 1, \quad L_1(\eta) = \frac{\sin \eta}{\eta}, \quad \eta_s = \pi$

- $n = 5, \quad L_5(\eta) = \frac{1}{\sqrt{1 + \frac{1}{3}\eta^2}}, \quad \eta_s = \infty.$

## Solving ODEs (Press et al.)

Consider a general 2nd order ODE:

$$\frac{d^2 y}{dx^2} + g(x) \frac{dy}{dx} = r(x)$$

We can always turn this into 2 1st order ODEs:

$$\begin{cases} \frac{dy}{dx} = z(x) \\ \frac{dz}{dx} = r(x) - g(x)z(x) \end{cases}$$

So our general problem is to solve a set of coupled 1st order ODEs:

$$\frac{dy_i}{dx} = f_i(x, y_1, y_2, \dots, y_n), \quad i = 1 \dots n.$$

We'll also need BCs - but we'll assume we have an "initial value problem" - we have the  $y_i$ 's at some starting point, and are looking for  $y_i$ 's at some point.

## Euler's Method

Consider just 1 ODE,

$$\frac{dy}{dx} = f(x, y).$$

Write  $dy \rightarrow \Delta y$  ,  $dx \rightarrow \Delta x$

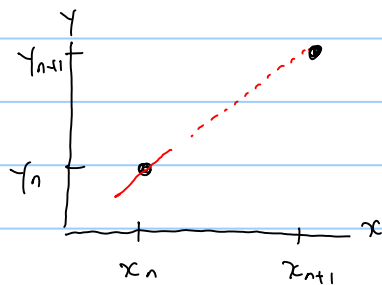
$$\frac{\Delta y}{\Delta x} = f(x, y)$$

$$\Delta y = f(x, y) \Delta x.$$

$$\Rightarrow y_{n+1} = y_n + \Delta y = y_n + f(x, y) \Delta x.$$

Write  $\Delta x = x_{n+1} - x_n = h$

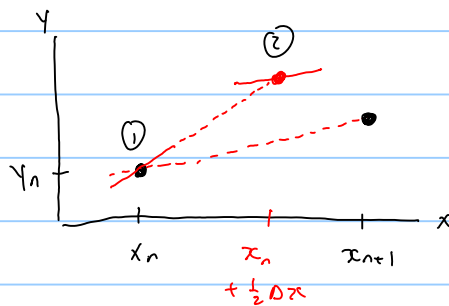
$$y_{n+1} = y_n + h f(x_n, y_n)$$



Error is  $\mathcal{O}(h^2)$

### MidPoint Method

Let's take a "trial" step to midpoint and use the derivative at that point :



- Compute the derivative at start ①  
 $k_1 = h f(x_n, y_n)$

- point (2) is at  $(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1)$   
 Compute the derivative at that point  
 $k_2 = h f(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1)$

So  $y_{n+1} = y_n + k_2$

Midpoint is  $\mathcal{O}(h^3)$

#### 4th Runge - Kutta

Why stop there? Use 4 derivatives:

- one at the start
- two at the midpoint
- one at the end

$$k_1 = h f(x_n, y_n)$$

$$k_2 = h f(x_n + \frac{h}{2}, y_n + \frac{k_1}{2})$$

$$k_3 = h f(x_n + \frac{h}{2}, y_n + \frac{k_2}{2})$$

$$k_4 = h f(x_n + h, y_n + k_3)$$

$$\Rightarrow y_{n+1} = y_n + \frac{k_1}{6} + \frac{k_2}{3} + \frac{k_3}{3} + \frac{k_4}{6}$$

∴ order  $\mathcal{O}(h^5)$