

# Exercises for Computational Physics (physics760)

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Thomas Luu, Andreas Nogga, Marcus Petschlies, Andreas Wirzba

## Exercise 1

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### 1 In-class exercises

#### 1.1 Make buddies

If you haven't done so yet, use this opportunity to find a partner to do your homework with. We prefer teams of 2.

#### 1.2 Version control with git

Once you have your team, then please set up a `git` repo. Your tutor might do a hands-on, walk-through presentation about the basic usage of `git` if there is enough demand by the class. Either way, make sure you:

- Can use of `git` on your own local machine
- Understand and can use the following `git` commands
  - `init`, `status`, `branch`, `log`, `add`, `commit`, `show`, `help`, `diff`
- Commit a change **with an accompanying message of your change** (ie. the heart of version control)
- *Provide access to your tutor!*

### 2 Homework (due Nov.11 at 18:00)

Please note the due date of this homework. Submission of homework *requires* submitting your solutions (e.g. answers to questions, graphs, results in tabular form) in the form of a brief report (please no 100 page submissions!) **AND** a copy of your code that you used to do the simulations. Please address all questions and requirements written below in *italics* in your homework report.

#### Simulation of the 1-D Ising model

The Ising model is a wonderful sandbox to study critical phenomena with numerical methods. Our main strategy in the coming weeks will be to investigate the Ising model in 2 dimensions, which already shows plenty of interesting behavior. To uncover these features

via non-perturbative simulation techniques is our principal goal in this tutorial series.

As a warm up we start with the simpler Ising model in 1 dimension. The analytic solution is known ( you might have seen this as an exercise in some solid-state physics lecture ), which conveniently allows us to implement, verify and test our code by the exact solution.

## 2.1 The model

Consider a chain of spins ( “1d lattice” ) with sites labeled by  $x \in \{0, \dots, N-1\}$ . To each site  $x$  we attach a spin variable  $s_x \in \{\pm 1\}$ . One specific choice of a spin for each site we call a spin configuration  $\mathbf{s} = \{s_x : \forall x\}$ .

This system of spins is immersed in a heat bath of constant temperature  $T$  and in an external magnetic field  $h$ . Its dynamics are governed by the Hamiltonian

$$\mathcal{H}(\mathbf{s}) = -J \sum_{\langle x,y \rangle} s_x s_y - h \sum_x s_x . \quad (1)$$

$\langle x, y \rangle$  denotes the nearest-neighbor pair  $x$  and  $y$  on the chain,  $J$  is a real number and we assume periodic boundary conditions. In the heat bath environment the probability for finding the spin system at spin configuration  $\mathbf{s} = (s_0, s_1, \dots, s_{N-1})$  is given by the Boltzmann distribution

$$P(\mathbf{s}) = \exp\left(-\frac{\mathcal{H}(\mathbf{s})}{k_B T}\right) / \sum_{\mathbf{s}'} \exp\left(-\frac{\mathcal{H}(\mathbf{s}')}{k_B T}\right) \equiv \frac{1}{Z} \exp\left(-\frac{\mathcal{H}(\mathbf{s})}{k_B T}\right) , \quad (2)$$

where we define

$$Z = \sum_{\mathbf{s}'} \exp\left(-\frac{\mathcal{H}(\mathbf{s}')}{k_B T}\right) .$$

as the partition function. The sum in the partition function is over all possible spin field  $\mathbf{s} = (s_0, s_1, \dots, s_{N-1})$  configurations. To simplify notation, let's switch to a system of units such that  $k_B = 1$ .

In 1d the partition function can be determined analytically as a function of  $N$ ,  $J/T$ , and  $h/T$ . We won't derive this result here, but only state it,

$$Z = \lambda_+^N + \lambda_-^N \quad ; \quad \lambda_{\pm} = e^{\frac{J}{T}} \left( \cosh\left(\frac{h}{T}\right) \pm \sqrt{\sinh^2\left(\frac{h}{T}\right) + e^{-4\frac{J}{T}}} \right) . \quad (3)$$

For those who are interested, the derivation can be found here:

[https://en.wikipedia.org/wiki/Ising\\_model](https://en.wikipedia.org/wiki/Ising_model)

You'll find this expression useful later when comparing your numerical results with exact results.

**1:** Discuss the physical meaning of  $J$ , in particular the sign of  $J$ , and the role it plays in magnets, for example.

**2:** Clarify what it means to have periodic boundary conditions (nearest neighbors).

**3:** Implement the Ising 1d simulation: determine an estimate for magnetization per spin

$$\langle m \rangle = -\frac{T}{N} \frac{\partial \log Z}{\partial h} , \quad (4)$$

and attempt to estimate the error of your estimate.

*Since we work with units where  $k_B = 1$ , what are the relevant dimensionless ratios in this problem?*

*Use your simulation code to study the dependency of  $\langle m \rangle$  on*

- *the external field strength  $h$  for fixed  $N$*
- *the number of spins  $N$  for fixed  $h$*

*Compare your numerical results to the exact solution at finite  $N$ , as well as to the “infinite volume” (also known as the “thermodynamic limit”) solution, i.e. for  $N \rightarrow \infty$ . Provide graphs of your comparisons. Use values of  $h \in [-1, 1]$  and  $N$  up to  $N \sim 20$  (but no larger). Describe your findings on the dependency of  $\langle m \rangle$  with  $N$ . You can fix  $J = 1$  for all simulations.*