

## Visual Computing - Lecture notes week 13

- Author: Ruben Schenk
- Date: 20.12.2021
- Contact: ruben.schenk@inf.ethz.ch

# 8.4 Physics-Based Animation

## 8.4.1 Introduction

We first differentiate between two important terms in the field of physics and physics-based animation:

- **Kinematics:** The branch of mechanics concerned with the motion of objects without reference to the forces which cause the motion.
- **Dynamics:** The branch of mechanics concerned with the motion of bodies under the action of forces.

## 8.4.2 The Animation Equation

We have already seen the *rendering equation* which is concerned with rasterization and path tracing which give approximate solutions to the rendering equation.

The **animation equation** is concerned with the large spectrum of physical systems and phenomena, such as solids, fluids, elasticity, etc. For animations, the connection between force and motion is essential:

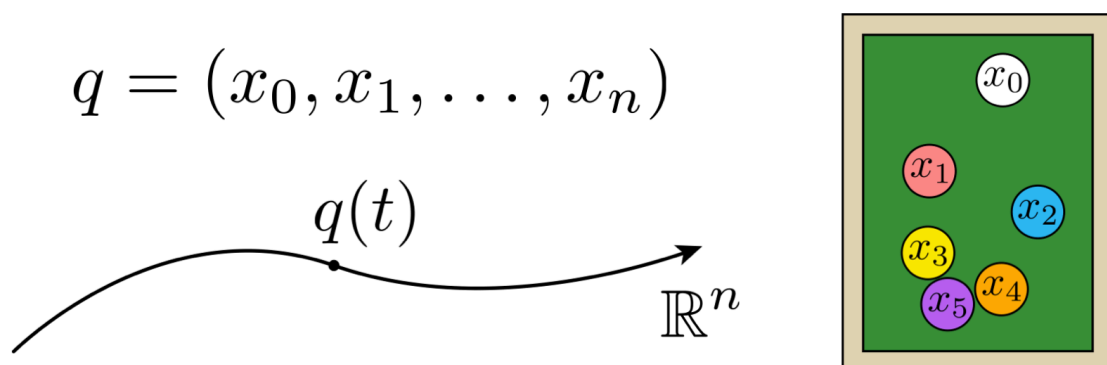
A change in motion is proportional to the motive force impressed and takes place along the straight line in which that force is impressed. – Sir Isaac Newton, 1687

However, there is more to be said than  $F = ma$ :

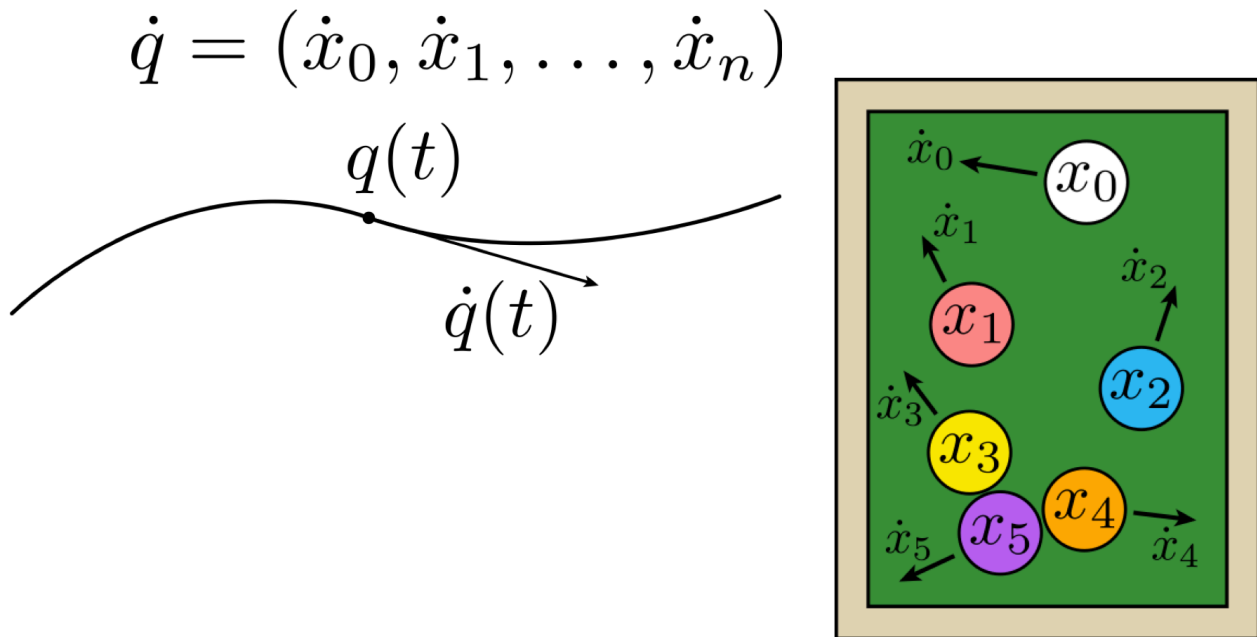
- Every system has a *configuration*  $q(t)$
- It also has a *velocity*  $\dot{q} := \frac{d}{dt} q$
- It has some kind of *mass*  $M$
- There are *forces*  $F$  acting on the system

## 8.4.3 Generalized Coordinates

In physics, we often need to describe a system with many moving parts, e.g. a collection of billiard balls, each with position  $x_i$ . We usually collect them all into a single vector of **generalized coordinates**.



We can think of  $q$  as a single point moving along some trajectory in  $\mathbb{R}^n$ . If we take the time derivative of the generalized coordinates, we get **generalized velocity**:



### 8.4.4 Ordinary Differential Equations

Many dynamical systems can be described via an **ordinary differential equation**:

$$\frac{d}{dt}q = f(q, \dot{q}, t),$$

where  $\frac{d}{dt}$  is the change in configuration over time and  $f$  is the velocity function.

*Example:* Assume we have a function where the rate of growth is proportional to the value, i.e.

$$\frac{d}{dt}u(t) = au,$$

then our solution is given by  $u(t) = be^{at}$ .

Note that Newton's 2nd law is an ODE as well, i.e.  $\ddot{q} = F/m$ . We can also write this as a system of two first order ODEs, by introducing new variables for velocity:

$$\begin{aligned} \dot{q} &= v, & \dot{v} &= \frac{F}{m} \\ \frac{d}{dt} \begin{bmatrix} q \\ v \end{bmatrix} &= \begin{bmatrix} v \\ F/m \end{bmatrix}. \end{aligned}$$

### 8.4.5 Solving ODEs Numerically

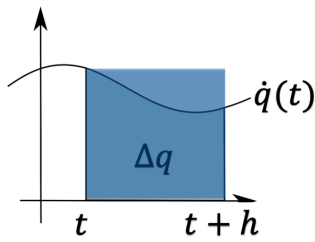
When we are talking about solving ODEs we mean that given some initial conditions  $q(0)$  and  $\dot{q}(0)$ , we want to find the function  $q(t)$ . Solving ODEs *numerically* means solving numerical time integration:

$$q(t+h) = q(t) + \int_t^{t+h} \dot{q}(t) dt$$

We use some discrete approximation of the form

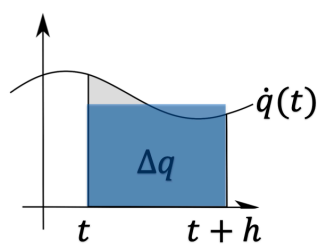
$$\Delta q_i \simeq \int_t^{t+h} \dot{q}(t) dt,$$

and then apply the following **numerical integration rules**:



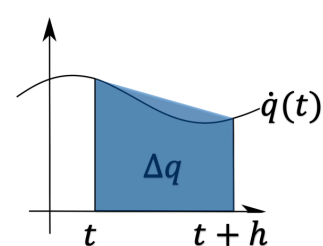
**Rectangle rule**

$$\Delta q_i \approx \dot{q}(t) \cdot h$$



**Midpoint rule**

$$\Delta q_i \approx \dot{q}(t + h/2) \cdot h$$



**Trapezoid rule**

$$\Delta q_i \approx \frac{\dot{q}(t) + \dot{q}(t+h)}{2} \cdot h$$

**Configuration update (rectangle rule):**

$$q_{i+1} = q_i + h \cdot \dot{q}_i$$

## 8.4.6 Forward Euler

**Forward Euler** describes a simple scheme: We evaluate the derivative at the current configuration and write the new state explicitly in terms of known data:

$$\begin{aligned} q_{i+1} &= q_i + h \cdot \dot{q}_i \\ \dot{q}_{i+1} &= \dot{q}_i + h \cdot \ddot{q}_i = \dot{q}_i + hM^{-1}F(q_i, \dot{q}_i) \end{aligned}$$

*Example:* Assume some simple linear ODE, i.e.  $\dot{u} = -au$ ,  $a > 0$ . The exact solution to this ODE would be  $u(t) = u(0)e^{-at}$ , so  $u_k \rightarrow 0$  as  $k \rightarrow \infty$ . The forward Euler approximation is given by:

$$u_n = (1 - ha)^n u_0,$$

from where we can derive that this decays only if  $|1 - ha| < 1$ , or equivalently  $h < 2/a$ , so in practice we may need very small time-steps!

## 8.4.7 Backward Euler

We might try something else and evaluate the velocity at some new configuration. This scheme is also known as **backward Euler**. The new configuration is then implicit, and we must solve for it:

$$\begin{aligned} q_{i+1} &= q_i + h \cdot \dot{q}_{i+1} \\ \dot{q}_{i+1} &= \dot{q}_i + h \cdot \ddot{q}_{i+1} = \dot{q}_i + hM^{-1}F(q_{i+1}, \dot{q}_{i+1}) \end{aligned}$$

We can again observe the stability of the backward Euler with our previous example, i.e.  $\dot{u} = -au$ ,  $a > 0$ . The backward Euler approximation is given by:

$$u_n = \left( \frac{1}{1 + ha} \right)^n u_0,$$

which decays if  $|1 + ha| > 1$ , which is always true! Backward Euler is *unconditionally stable* for linear ODEs.