Madelung Representation

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```
[1]: __author__ = "@Tssp"
    __date__ = "17/03/2021"
    import sympy as sp
    import numpy as np
```

This notebook is aimed to obtain the superfluid hydrodynamic equations of motions through the Madelung representation of the Gross Pitaevskii equation.

```
[10]: x, t, m, hbar = sp.symbols('x t m \\hbar', real=True)
    psi = sp.Function('\\psi', real=False)(x, t)  # wavefunction
    rho = sp.Function('\\rho', real=True)(x, t)  # probability density
    V = sp.Function('V', real=True)(x, t)  # Trap Potential
    S = sp.Function('S', real=True)(x, t)  # Madelung Ansatz
    v = sp.Function('v', real=True)(x, t)  # fluid velocity
    Pq = sp.Function('P_Q', real=True)(x, t)  # Quantum Preassure
    g = sp.symbols('g', real=True)  # g constant
```

[11]:
$$Schr_{eq} = sp.I * hbar * sp.diff(psi, t) + (hbar**2/(2*m) * sp.diff(psi, x, 2) - V * psi - g*abs(psi)**2 * psi)$$

 $Schr_{eq}$

$$\frac{\hbar^2 \frac{\partial^2}{\partial x^2} \psi(x,t)}{2m} + i\hbar \frac{\partial}{\partial t} \psi(x,t) - g\psi(x,t) |\psi(x,t)|^2 - V(x,t) \psi(x,t)$$

Madelung transformation:

$$\psi(x,t) = \sqrt{\rho(x,t)}e^{imS(x,t)/\hbar}$$

$$\hbar^{2} \left(\frac{2 \frac{\partial^{2}}{\partial x^{2}} \rho(x,t) - \frac{\left(\frac{\partial}{\partial x} \rho(x,t)\right)^{2}}{\rho(x,t)}}{4 \sqrt{\rho(x,t)}} + \frac{m \left(i \frac{\partial^{2}}{\partial x^{2}} S(x,t) - \frac{m \left(\frac{\partial}{\partial x} S(x,t)\right)^{2}}{\hbar}\right) \sqrt{\rho(x,t)}}{\hbar} + \frac{i m \frac{\partial}{\partial x} S(x,t) \frac{\partial}{\partial x} \rho(x,t)}{\hbar \sqrt{\rho(x,t)}} \right) }{i \hbar \left(\frac{\partial}{\partial t} \rho(x,t)}{2 \sqrt{\rho(x,t)}} + \frac{i m \sqrt{\rho(x,t)} \frac{\partial}{\partial t} S(x,t)}{\hbar} \right) - g \sqrt{\rho(x,t)} \left| \sqrt{\rho(x,t)} \right|^{2} - V(x,t) \sqrt{\rho(x,t)}$$

$$\begin{split} &\left[\textbf{13} \right] : \frac{\hbar^2 \frac{\partial^2}{\partial x^2} \rho(x,t)}{4m\sqrt{\rho(x,t)}} - \frac{\hbar^2 \left(\frac{\partial}{\partial x} \rho(x,t) \right)^2}{8m\rho^{\frac{3}{2}}(x,t)} + \frac{i\hbar\sqrt{\rho(x,t)} \frac{\partial^2}{\partial x^2} S(x,t)}{2} + \frac{i\hbar\frac{\partial}{\partial x} S(x,t) \frac{\partial}{\partial x} \rho(x,t)}{2\sqrt{\rho(x,t)}} + \frac{i\hbar\frac{\partial}{\partial t} \rho(x,t)}{2\sqrt{\rho(x,t)}} - g\sqrt{\rho(x,t)} \left| \sqrt{\rho(x,t)} \right|^2 - m\sqrt{\rho(x,t)} \frac{\partial}{\partial t} S(x,t) - \frac{m\sqrt{\rho(x,t)} \left(\frac{\partial}{\partial x} S(x,t) \right)^2}{2} - V(x,t)\sqrt{\rho(x,t)} \end{split}$$

$$\frac{\lceil 14 \rceil :}{2} \frac{\hbar \sqrt{\rho(x,t)} \frac{\partial^2}{\partial x^2} S(x,t)}{2} + \frac{\hbar \frac{\partial}{\partial x} S(x,t) \frac{\partial}{\partial x} \rho(x,t)}{2 \sqrt{\rho(x,t)}} + \frac{\hbar \frac{\partial}{\partial t} \rho(x,t)}{2 \sqrt{\rho(x,t)}}$$

0.1 Imaginary part

Defining:

$$\vec{v} = \frac{\partial S(x,t)}{\partial x}$$

[15]:
$$m\rho(x,t)\frac{\partial^2}{\partial x^2}S(x,t) + m\frac{\partial}{\partial x}S(x,t)\frac{\partial}{\partial x}\rho(x,t) + m\frac{\partial}{\partial t}\rho(x,t)$$

[16]:
$$\rho(x,t)\frac{\partial^2}{\partial x^2}S(x,t) + \frac{\partial}{\partial x}S(x,t)\frac{\partial}{\partial x}\rho(x,t)$$

[17]:
$$Im_eq.subs(\{rho*sp.diff(S, x, 2): sp.diff(rho*sp.diff(S, x), x) - sp.diff(S, x)_{\sqcup} \rightarrow * sp.diff(rho, x)\})$$

[17]:
$$m\rho(x,t)\frac{\partial^2}{\partial x^2}S(x,t) + m\frac{\partial}{\partial x}S(x,t)\frac{\partial}{\partial x}\rho(x,t) + m\frac{\partial}{\partial t}\rho(x,t)$$

$$m\frac{\partial \rho}{\partial t} - \frac{\partial}{\partial x} \left((\rho \frac{\partial}{\partial x} S) - S - \rho + S + \rho \right) = 0$$

Or:

$$m\frac{\partial\rho}{\partial t}-\frac{\partial}{\partial x}\left((\rho\frac{\partial}{\partial x}S)\right)=0$$

$$\frac{\partial \rho}{\partial t} - \frac{\partial}{\partial x} \left(\rho v \right) = 0$$

A continuity equation, in 3D

$$\frac{\partial \rho}{\partial t} - \nabla \left(\rho v \right) = 0$$

0.2 Real part

In this part we are going to treat $\sqrt{\rho}$ as an unique variable in order to forbid Python to literally express the derivative and therefore simplify the equations

[19]:
$$\hbar^{2} \left(\frac{\partial^{2}}{\partial x^{2}} \sqrt{\rho}(x,t) + \frac{m \left(i \frac{\partial^{2}}{\partial x^{2}} S(x,t) - \frac{m \left(\frac{\partial}{\partial x} S(x,t) \right)^{2}}{\hbar} \right) \sqrt{\rho}(x,t)}{\hbar} + \frac{2im \frac{\partial}{\partial x} S(x,t) \frac{\partial}{\partial x} \sqrt{\rho}(x,t)}{\hbar} \right) + i\hbar \left(\frac{\partial}{\partial t} \sqrt{\rho}(x,t) + \frac{im \sqrt{\rho}(x,t) \frac{\partial}{\partial t} S(x,t)}{\hbar} \right) - g \sqrt{\rho^{3}}(x,t) - V(x,t) \sqrt{\rho}(x,t)$$

$$\frac{\hbar^2 \frac{\partial^2}{\partial x^2} \sqrt{\rho}(x,t)}{2m} + \frac{i\hbar \sqrt{\rho}(x,t) \frac{\partial^2}{\partial x^2} S(x,t)}{2} + i\hbar \frac{\partial}{\partial x} S(x,t) \frac{\partial}{\partial x} \sqrt{\rho}(x,t) + i\hbar \frac{\partial}{\partial t} \sqrt{\rho}(x,t) - g\sqrt{\rho}^3(x,t) - m\sqrt{\rho}(x,t) \frac{\partial}{\partial t} S(x,t) - \frac{m\sqrt{\rho}(x,t) \left(\frac{\partial}{\partial x} S(x,t)\right)^2}{2} - V(x,t)\sqrt{\rho}(x,t)$$

$$-\frac{\hbar^2\frac{\partial^2}{\partial x^2}\sqrt{\rho}(x,t)}{2m^2\sqrt{\rho}(x,t)} + \frac{g\sqrt{\rho}^2(x,t)}{m} + \frac{\partial}{\partial t}S(x,t) + \frac{\left(\frac{\partial}{\partial x}S(x,t)\right)^2}{2} + \frac{V(x,t)}{m}$$

The negative term corresponds with the quantum preassure contribution to the fluid. With some manipulation we get the superfluid hydrodynamic equations:

$$\frac{g\sqrt{\rho^2(x,t)}}{m} + P_{\mathcal{Q}}(x,t) + \frac{\partial}{\partial t}S(x,t) + \frac{\left(\frac{\partial}{\partial x}S(x,t)\right)^2}{2} + \frac{V(x,t)}{m}$$

[27]:
$$m \left(\frac{2g\sqrt{\rho}(x,t)\frac{\partial}{\partial x}\sqrt{\rho}(x,t)}{m} + v(x,t)\frac{\partial}{\partial x}v(x,t) + \frac{\partial}{\partial x}\operatorname{P}_{\mathbf{Q}}(x,t) + \frac{\partial}{\partial t}v(x,t) + \frac{\frac{\partial}{\partial x}V(x,t)}{m} \right) \rho(x,t)$$

or:

$$m\rho \left(\frac{\partial v}{\partial t} + (v \cdot \nabla)v\right) = -m\rho \nabla P_Q - \rho \nabla V - g\nabla \rho$$

Notice that the pressure depends only on the density. This property makes the condensate a barotropic fluid; as a consequence, surfaces of constant pres- sure are also surfaces of constant density. The quantum pressure is a pure quantum effect, and vanishes if we set Planck's constant equal to zero. It has the same origin as the quantum kinetic energy, i.e. zero point motion, which creates a pressure that opposes any `squashing' or `bending' of the condensate. In a uniform condensate the quantum pressure is zero because n is constant.

This equation is very similar to the classical Euler equation for an inviscid fluid. If in addition, the trapping potential is absent (V=0) then the equation reduces to the classical Euler equation, which describes the motion of a classical fluid without viscosity.