

# Gas dynamics

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# Notation

Eulerian description: one defines a fixed coordinate grid in space and follows how the gas quantities are changing at a given position. The

Eulerian time-derivative is  $\frac{\partial}{\partial t}$ .

Lagrangian description: one chooses a fluid element in the fluid and follows how its properties change. The Lagrangian time-derivative is

$\frac{D}{Dt}$ .

These points of view are related through:

$$\frac{DQ}{Dt} = \frac{\partial Q}{\partial t} + \mathbf{v} \cdot \nabla Q \quad (1)$$

# The basic equations of gas dynamisc

Continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0 \quad (2)$$

Euler equation

$$\rho \frac{\partial \mathbf{v}}{\partial t} + \rho \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla P + \mathbf{f} \quad (3)$$

Energy conservation

$$\frac{\partial}{\partial t} \left( \frac{1}{2} \rho v^2 + \rho \epsilon \right) + \nabla \cdot \left[ \left( \frac{1}{2} \rho v^2 + \rho \epsilon + P \right) \mathbf{v} \right] = \mathbf{f} \cdot \mathbf{v} - \nabla \mathbf{q} \quad (4)$$

# Continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0 \quad (5)$$

$\rho$  is the mass density (mass per unit volume),  $\rho \mathbf{v}$  is the mass flux density (mass flux per unit area.)

An alternative formula is:

$$\frac{D\rho}{Dt} = -\rho \nabla \cdot \mathbf{v} \quad (6)$$

# Equation of motion

Euler equation

$$\rho \left( \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla P + \mathbf{f} \quad (7)$$

$\rho \mathbf{v} \cdot \nabla \mathbf{v}$  - advection of momentum through the fluid by velocity gradients.

$\mathbf{f}$  - force density (force per unit volume)

gravity:  $\mathbf{f} = -\rho \mathbf{g} = \rho \nabla \Phi$ , where  $\Phi$  is gravitational potential

Poisson equation for self-gravity systems:

$$\nabla^2 \Phi = 4\pi G \rho \quad (8)$$

external magnetic field:  $(\nabla \times \mathbf{B}) \times \mathbf{B}$

viscous forces:  $\rho \nu \nabla^2 \mathbf{v}$ , here  $\nu$  is kinematic viscosity.

# Energy conservation

Energy = kinetic energy  $\frac{1}{2}\rho v^2$  + thermal energy  $\rho\epsilon$

Monoatomic gas:  $\epsilon = 3/2kT/(\mu m_H)$ ,  $\mu$  is the mean molecular weight,  
 $\mu = 1$  - neutral hydrogen,  $\mu = 0.6$  - Solar composition.

$$\frac{\partial E}{\partial t} + \nabla \cdot [(E + P) \mathbf{v}] = \mathbf{f} \cdot \mathbf{v} - \nabla \mathbf{q} \quad (9)$$

$$\frac{\partial}{\partial t} \left( \frac{1}{2} \rho v^2 + \rho \epsilon \right) + \nabla \cdot \left[ \left( \frac{1}{2} \rho v^2 + \rho \epsilon + P \right) \mathbf{v} \right] = \mathbf{f} \cdot \mathbf{v} - \nabla \mathbf{q} \quad (10)$$

here  $\mathbf{q} = -K\nabla T$  is the heat flux

# Steady flows

Continuity equation

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0 \quad (11)$$

Euler equation

$$\rho \frac{\partial \mathbf{v}}{\partial t} + \rho \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla P + \mathbf{f} \quad (12)$$

Energy conservation

$$\frac{\partial}{\partial t} \left( \frac{1}{2} \rho v^2 + \rho \epsilon \right) + \nabla \cdot \left[ \left( \frac{1}{2} \rho v^2 + \rho \epsilon + P \right) \mathbf{v} \right] = \mathbf{f} \cdot \mathbf{v} \quad (13)$$

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# Steady flows

$$\nabla(\rho \mathbf{v}) = 0 \quad (14)$$

$$\rho \mathbf{v} \cdot \nabla \mathbf{v} = -\nabla P + \mathbf{f} \quad (15)$$

$$\nabla \cdot \left[ \left( \frac{1}{2} \rho v^2 + \rho \epsilon + P \right) \mathbf{v} \right] = \mathbf{f} \cdot \mathbf{v} \quad (16)$$

Vector calculus:  $\nabla(\mathbf{a} \cdot \mathbf{k}) = \mathbf{a}(\nabla \cdot \mathbf{k}) + \mathbf{k}(\nabla \mathbf{a})$  Using this rule and equation (14) one can obtain:

$$\rho \mathbf{v} \cdot \nabla \left( \frac{1}{2} v^2 + \epsilon + \frac{P}{\rho} \right) = \mathbf{f} \cdot \mathbf{v} \quad (17)$$

Then multiply equation(15) on  $\mathbf{v}$ :

$$\mathbf{v} \rho \cdot (\mathbf{v} \cdot \nabla) \mathbf{v} = -\mathbf{v} \cdot \nabla P + \mathbf{f} \cdot \mathbf{v} \quad (18)$$

# Steady flows

$$\rho \mathbf{v} \cdot \nabla \left( \frac{1}{2} v^2 + \epsilon + \frac{P}{\rho} \right) = \mathbf{f} \cdot \mathbf{v} \quad (19)$$

Then multiply equation(15) on  $\mathbf{v}$ :

$$\mathbf{v} \rho \cdot (\mathbf{v} \cdot \nabla) \mathbf{v} = -\mathbf{v} \cdot \nabla P + \mathbf{f} \cdot \mathbf{v} \quad (20)$$

Here  $\mathbf{v} \rho \cdot (\mathbf{v} \cdot \nabla) \mathbf{v} = \rho \mathbf{v} \cdot \nabla \left( \frac{1}{2} v^2 \right)$  Then

$$\rho \mathbf{v} \cdot \nabla (\epsilon + P/\rho) = \mathbf{v} \cdot \nabla P \quad (21)$$

or

$$\mathbf{v} \cdot [\nabla \epsilon + P \nabla (1/\rho)] = 0 \quad (22)$$

$$\mathbf{v} \cdot [\nabla \epsilon + P \nabla(1/\rho)] = 0 \quad (23)$$

What is  $\mathbf{v} \cdot \nabla$ ? Remember equation(1):  $\frac{D}{Dt} = \frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla \equiv d$  So,

$$d\epsilon + Pd(1/\rho) = 0 \quad (24)$$

where for monoatomic gas enternal energy is  $\epsilon = \frac{3}{2} \frac{kT}{\mu m_H}$  and perfect gas law:  $P = \rho kT / \mu m_H$ :

$$\rho^{-1} T^{3/2} = \text{constant} \quad (25)$$

or

$$P \rho^{-5/3} = \text{constant} \quad (26)$$

This equation describes adiabatic flows.

# Adiabatic and isothermal flows

General case:

$$P\rho^\gamma = \text{constant} \quad (27)$$

here  $\gamma$  is adiabatic index, or the ratio of specific heats.

This condition is equivalent to setting the entropy of the gas constant.

Where? Very fast processes SN, thick disks

Isothermal flow with  $T$  is constant:  $P\rho^{-1} = \text{constant}$  Where?

Protoplanetary disks, HII clouds. Slow processes.

# Acoustic waves

Here neglect transport phenomena,  $\nabla q = 0$ . Consider a homogeneous perfect gas at rest:  $\rho = \rho_0$ ;  $p = p_0$ ;  $\mathbf{v} = \mathbf{v}_0 = 0$ .

Suppose that pressure is perturbed to  $p = p_0 + p_1$ ,  $\rho = \rho_0 + \rho_1$  and  $\mathbf{v} = \mathbf{v}_1$ , assume that perturbations are adiabatic, so:

$$0 = \frac{d}{dt} \left( \frac{p}{\rho^\gamma} \right) = \frac{1}{\rho^\gamma} \frac{dp_1}{dt} - \gamma \frac{p}{\rho^{\gamma+1}} \frac{d\rho_1}{dt} \quad (28)$$

from this  $dp_1 = c_s^2 d\rho_1$  where  $c_s^2 = \gamma p / \rho$ .

Assuming that the perturbations are small,  $\rho_1 \ll \rho_0$  and  $p_1 \ll p_0$ , we can replace the quantities in  $c_s^2$  by their equilibrium values,

$$c_s^2 = \frac{\gamma p_0}{\rho_0} \quad (29)$$

Linearizing continuity equation:

$$0 = \frac{\partial \rho_1}{\partial t} + \nabla \cdot [(\rho_0 + \rho_1)\mathbf{v}_1] \simeq \frac{\partial \rho_1}{\partial t} + \rho_0 \nabla \cdot \mathbf{v}_1 \quad (30)$$

then Euler equation:

$$0 = (\rho_0 + \rho_1) \left[ \frac{\partial \mathbf{v}_1}{\partial t} + (\mathbf{v}_1 \cdot \nabla) \mathbf{v}_1 \right] + \nabla p_1 \simeq \rho_0 \frac{\partial \mathbf{v}_1}{\partial t} + \nabla p_1 = \rho_0 \frac{\partial \mathbf{v}_1}{\partial t} + c_s^2 \nabla \rho_1 \quad (31)$$

Combining both these equation one can obtain:

$$\frac{\partial^2 \rho_1}{\partial t^2} - c_s^2 \nabla^2 \rho_1 = 0 \quad (32)$$

which is equation for acoustic waves showing that

$$c_s^2 = \sqrt{\frac{\gamma p_0}{\rho_0}} \quad (33)$$

is the sound speed. Assuming that  $p_1 = (dp/d\rho)_0 \rho_1$  give us:

$$c_s^2 = \left( \frac{dp}{d\rho} \right)_0 \quad (34)$$

$$c_s^2 = \left( \frac{dp}{d\rho} \right)_0 \quad (35)$$

adiabatic flow:

$$c_s^{\text{ad}} = \left( \frac{5p}{3\rho} \right)^{1/2} = \left( \frac{5kT}{3\mu m_H} \right)^{1/2} \quad (36)$$

isothermal flow:

$$c_s^{\text{iso}} = \left( \frac{p}{\rho} \right)^{1/2} = \left( \frac{kT}{\mu m_H} \right)^{1/2} \quad (37)$$

$c_s$  is the speed at which pressure disturbances travel through the gas, it limits the rapidly with which the gas can respond to pressure changes.



# Subsonic and supersonic flows

$c_s$  is the speed at which pressure disturbances travel through the gas, it limits the rapidly with which the gas can respond to pressure changes.

Supersonic flow:  $|\mathbf{v}| > c_s$  the pressure gradients has little effect on the flow.

Subsonic flow:  $|\mathbf{v}| < c_s$  to a first approximation the gas behaves as if in hydrostatic equilibrium.

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (38)$$

Keeping gravity + rotation

$$0 = -\nabla \Phi + \Omega^2 \mathbf{R} \quad (39)$$

Gives Thin Keplerian disks

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (40)$$

Keeping gravity + pressure

$$0 = -\nabla \Phi - \frac{1}{\rho} \nabla P \quad (41)$$

Gives Stars, stellar envelopes and atmospheres

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (42)$$

Keeping advection + gravity

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\nabla \Phi \quad (43)$$

Gives Gravitational collapse, dust in free-fall

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (44)$$

Keeping advection + gravity + rotation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\nabla \Phi + \Omega^2 \mathbf{R} \quad (45)$$

Gives Slim disks

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (46)$$

Keeping pressure + gravity + rotation

$$0 = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} \quad (47)$$

Gives Thick disks and tori

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (48)$$

Keeping advection + pressure + gravity

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi \quad (49)$$

Gives Bondi-Hoyle accretion

# Euler equation and the disks

Radial momentum equation

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} + \text{viscous terms} \quad (50)$$

Only viscous terms neglected

$$(\mathbf{v}_p \cdot \nabla) \mathbf{v}_p = -\frac{1}{\rho} \nabla P - \nabla \Phi + \Omega^2 \mathbf{R} \quad (51)$$

Gives Sun-Keplerian ADAFs, boundary layers



# What to read?

- ▶ Arnab Rai Choudhuri " The Physics of Fluids and Plasmas: An Introduction for Astrophysicists " + lecture notes by Rami Vainio
- ▶ Astrophysical Gasdynamics Lecture Notes by Garrelt Mellema
- ▶ Lecture notes "Astrophysical fluid dynamics" by Gordon Ogilvie, arxiv: 1604.03835
- ▶ Balbus, Potter "Surprises in astrophysical gasdynamics", arxiv: 1603.06489