

# Nuclear Physics

A.I

Last updated: July 29, 2023

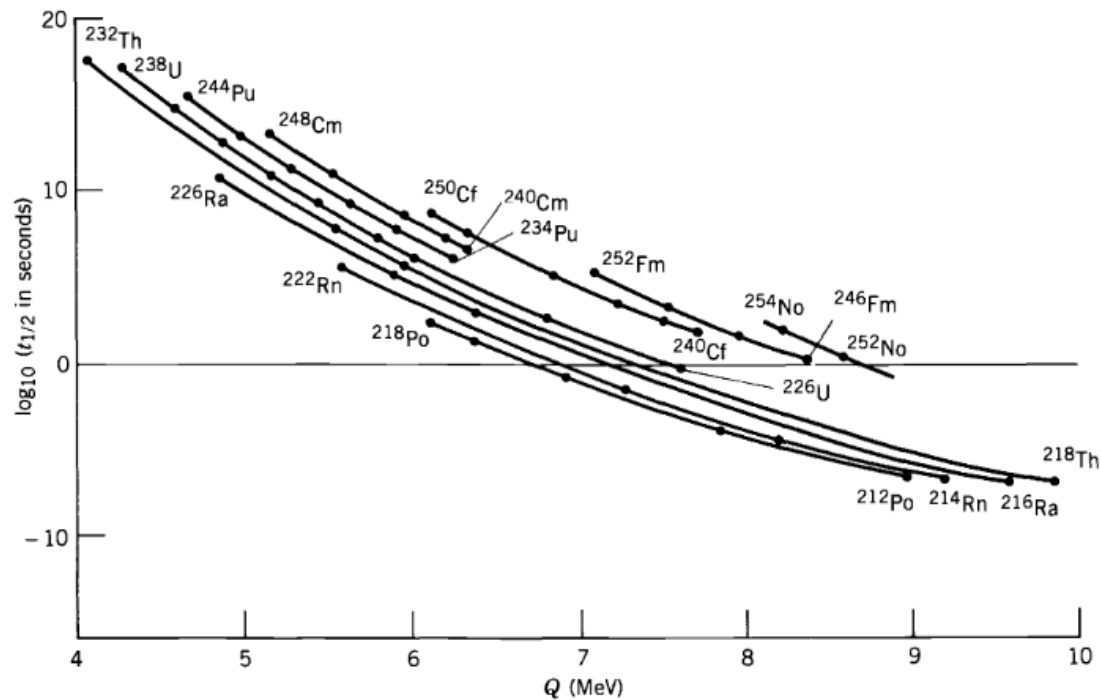
## Contents

<b>Lecture 1 – Alpha Particle</b>	<b>1</b>
1.1 Alpha Decay . . . . .	1
1.2 Second subsection . . . . .	4

## ❖ Lecture 1 (1/21)

### 1.1 Alpha Decay

Geiger and Nuttall noticed a trend where a slight difference in the disintegration energies or Q-values of radioactive samples brought about a huge difference in the half life.  $^{232}\text{Th}$  ( $1.4 \times 10^{10}$  y;  $Q = 4.08$  MeV)  $^{218}\text{Po}$  ( $3.1 \times 10^{-8}$  s;  $Q = 9.85$  MeV) so a factor of 2 in energy brings about a factor of  $10^{24}$  in half life.

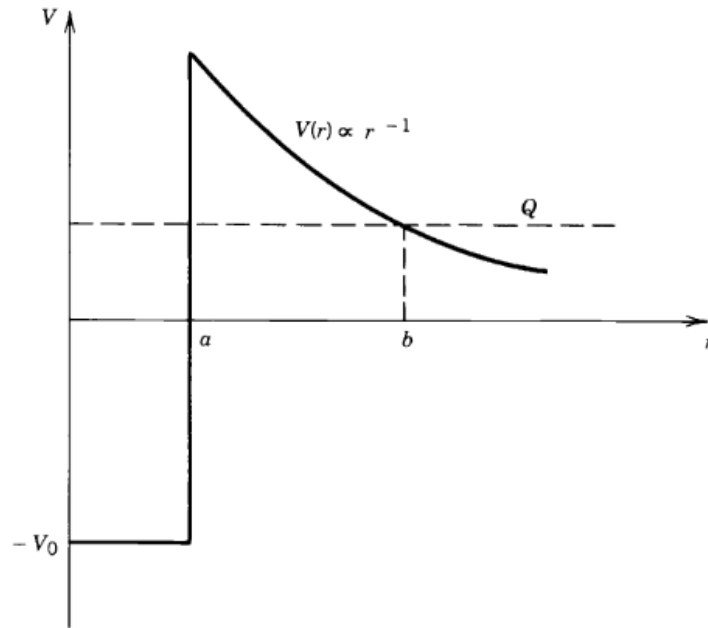


**Figure 8.1** The inverse relationship between  $\alpha$ -decay half-life and decay energy, called the Geiger-Nuttall rule. Only even- $Z$ , even- $N$  nuclei are shown. The solid lines connect the data points.

Even-odd, odd-even and odd-odd nuclei also show the same trend as the above even-even nuclei but don't result in lines as smooth. Their periods are 2-1000 times larger than the even-even types.

General features of the above plot can be explained by the Quantum Mechanical theory developed simultaneously in 1928 by Gamov and Churney and Condon. In this theory, the  $\alpha$  particle assumed to be formed inside the nucleus before emission.

Do keep in mind that just because this theory can explain nuclear behaviour does not mean that  $\alpha$  particle really is formed in the nucleus before emission, it just means that it behaves as if that was the case.



**Figure 8.3** Relative potential energy of  $\alpha$ -particle, daughter-nucleus system as a function of their separation. Inside the nuclear surface at  $r = a$ , the potential is represented as a square well; beyond the surface, only the Coulomb repulsion operates. The  $\alpha$  particle tunnels through the Coulomb barrier from  $a$  to  $b$ .

The above is a plot of potential between  $\alpha$  particle and residual nuclei ( $V_0$ ) against the distances between their centres ( $r$ ). The horizontal dotted line is the disintegration energy  $Q$  and ' $a$ ' can be taken as the sum of radii of  $\alpha$  particle and residual nucleus. There are three areas of interest, for  $r < a$ , we are inside the nucleus, with a potential well of  $-V_0$ , classically the particle can move here with potential of  $Q + V_0$ . The shell area  $a < r < b$  forms a potential barrier here as the potential energy here is higher than the available energy  $Q$ . So  $\alpha$  particle is not supposed to be able to enter this area through either side, classically that is.  $b < r$  is the classically allowed area outside the barrier. According to classical mechanics then,  $\alpha$  decay should not be possible at all, but quantum mechanically, there is a small possibility of an event called 'tunneling' to occur where the  $\alpha$  particle gets to the other side despite the barrier.

The probability of this event is extremely low, that's why  $\alpha$  unstable nuclei don't just decay all the time. For example, in  $^{238}\text{U}$ , the  $\alpha$  particle has to "hit" the potential barrier approximately  $10^{28}$  times before it can get out. The same principle follows in the case of an incident particle from the other side of the barrier (with lower energy than the potential well of course). Fusion reactions, such as those responsible for the energy released in stars are analyzed using the barrier penetration approach.

The disintegration constant of an  $\alpha$  emitter in one-body theory is given by

$$\lambda = fP$$

Where  $f$  is the frequency of the alpha particle "hitting" the barrier and  $P$  is the probability of transmission through the barrier.

## 1.2 Second subsection

## References

- [1] Kenneth Krane *Introductory Nuclear Physics*.