

Math 5470/6440 Chaos theory

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1D: Linear stability analysis

Today we will discuss:

Linearization and linear stability analysis.

Growth / decay rate and characteristic time scale.

Existence / uniqueness of solution of one-dimensional system.

Impossibility of oscillations in 1D.

Potentials.

One-dimensional or first order system is:

$$\dot{x} = f(x), \quad x(0) = x_0$$

Here $x(t)$ is real valued function of time t ,

f is a smooth, real-valued function of x .

x is position of a particle, \dot{x} is velocity.

2.4

Linear Stability Analysis

? Rate of decay to the stable fixed pt?

To answer, linearize about the fixed pt

$\dot{x} = f(x)$, let x^* be a fixed pt,

$\eta(t) = x(t) - x^*$ is a small perturbation

Derive diff. eqn for η :

$$\dot{\eta} = \frac{d}{dt}(\eta(t)) = \frac{d}{dt}(x(t) - x^*) = \frac{d}{dt}x(t) = \dot{x} = f(x)$$

\uparrow
const
 $x = x^* + \eta$

$$\dot{\eta} = f(x^* + \eta) \quad - \text{use Taylor's expansion}$$

$$f(x^* + \eta) = f(x^*) + \eta f'(x^*) + O(\eta^2) \leftarrow \begin{cases} \text{quadratically} \\ \text{small terms in } \eta \end{cases}$$

" x^* is fixed pt

$$\dot{\eta} = \eta f'(x^*) + O(\eta^2)$$

If $f'(x^*) \neq 0$, $O(\eta^2)$ terms are negligible \Rightarrow

$$\dot{\eta} = \eta f'(x^*) \quad \left| \begin{array}{l} \text{linear eqn in } \eta, \\ \text{linearized about } x^* \end{array} \right.$$

$f'(x^*)$ - growth (or decay) rate

(as in $\eta = e^{kt}$, $k = f'(x^*)$,
- soln for $\dot{\eta} = k\eta$)

Linearized eqn for η : (if $f'(x^*) \neq 0$)

$$\boxed{\dot{\eta} = \eta f'(x^*)}$$

Approximation as we dropped quadratic terms in η

Solu: $\eta(t) = e^{kt}$ for $k = f'(x^*)$ - growth / decay rate

$f'(x^*) > 0 \Rightarrow \eta$ grows exponentially
 x^* is unstable

$f'(x^*) < 0 \Rightarrow \eta$ decays exponentially fast
 x^* is stable

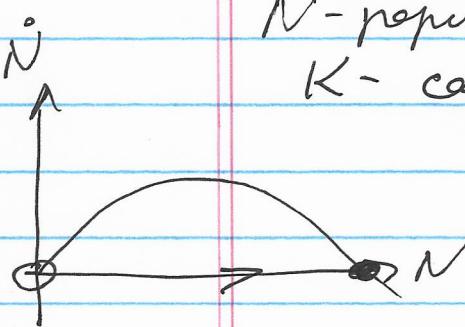
If $f'(x^*) = 0$ the nonlinear terms are not negligible, nonlinear analysis is required to determine stability

$$\frac{1}{|f''(x^*)|} \text{ - characteristic time scale}$$

Ex. Logistic eqn $\dot{N} = rN\left(1 - \frac{N}{K}\right)$

N -population size, $r > 0$ growth rate

K - carrying capacity



$$f(N) = rN\left(1 - \frac{N}{K}\right) = rN - \frac{r}{K}N^2$$

$$N^* = 0, N^* = K$$

$$f'(N) = r - 2r\frac{N}{K}$$

$f'(0) = r > 0$: $x^* = 0$ - unstable point, exponential growth

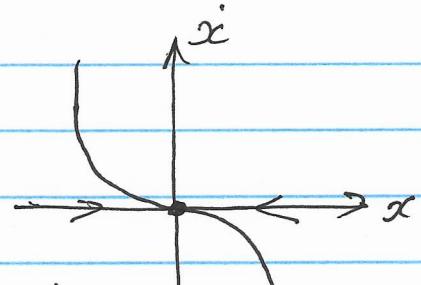
$f'(K) = -r < 0$: $x^* = K$ - stable fixed point with exponentially decaying perturbation

$$\frac{1}{|f''(N^*)|} = \frac{1}{r}$$

If $f'(x^*) = 0$, linear stability analysis does not work

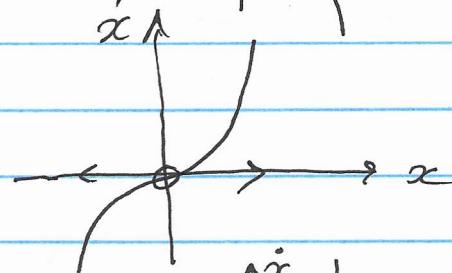
Ex. (a) $\dot{x} = -x^3$

$$x^* = 0, f'(x^*) = -3(x^*)^2 = 0$$



(b)

$$\dot{x} = x^3, x^* = 0, f'(x^*) = 0$$

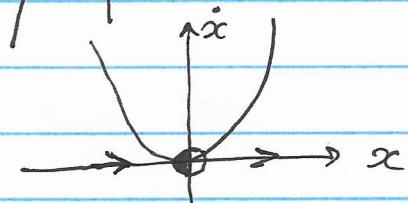


(c)

$$\dot{x} = x^2, x^* = 0, f'(x^*) = 0$$

$x^* = 0$ is half-stable fixed pt.

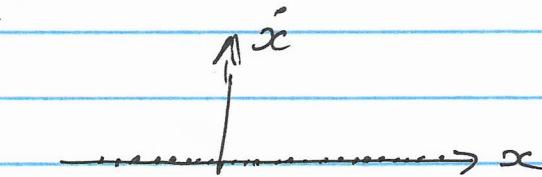
\Rightarrow attracting from left,
repelling on the right



(d)

$$\dot{x} = 0$$

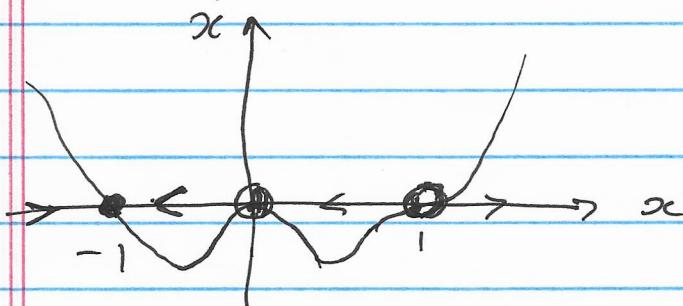
$$x^* = ?$$



Any pt in $(-\infty, \infty)$ is
stable fixed pt -
perturbations neither grow nor
decay

(e)

$$\dot{x} = (x+1)x^2(x-1)^3$$



$x_1^* = -1$ - stable

$x_2^* = 0$ - half-stable

$x_3^* = 1$ - unstable

Existence and Uniqueness: If $f(x)$ is smooth enough, $\exists!$ solution.

IVP: $\begin{cases} \dot{x} = f(x) \\ x(0) = x_0 \end{cases}$

Suppose $f(x), f'(x)$ are continuous on open interval R ,
(or f is cont. differentiable on R), and $x_0 \in R$. Then:

the IVP has a solution $x(t)$ on some time interval $(-\tau, \tau)$ about $t=0$, and the solution is unique.

① Does the solution exist for all times? - No

Ex. $\dot{x} = 1+x^2, x(0) = x_0 = 0$

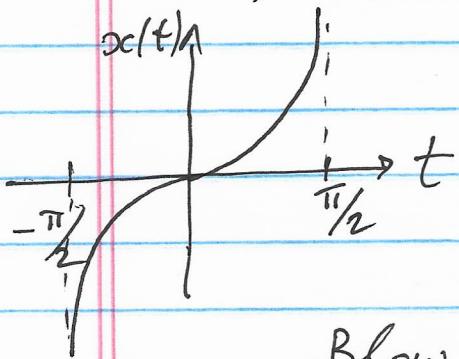
$f(x) = 1+x^2$ - cont., $f'(x)$ - cont. $\forall x \in R$

\Rightarrow Thm says: $\exists!$ soln $\forall x_0$.

let's integrate: $\frac{dx}{dt} = 1+x^2 \Rightarrow \int \frac{dx}{1+x^2} = \int dt \Rightarrow$

$\tan^{-1}(x) = t + C, \text{ as } x(0) = 0 \Rightarrow C = 0$

$\Rightarrow x(t) = \tan(t)$



- $x(t)$ exists only on interval $(-\frac{\pi}{2}, \frac{\pi}{2})$. There is no other

sln outside $(-\frac{\pi}{2}, \frac{\pi}{2})$ with $x(0) = 0$

when $t \rightarrow \pm \frac{\pi}{2}$, $x(t) \rightarrow \pm \infty$

Blow up: soln reaches ∞ at finite time!

* What if f is not cont. differentiable?

Ex. $\dot{x} = x^{2/3}$ $f(x) = x^{2/3}$ We will show that solution is not unique.

Use separation of variables:

$$\int \frac{dx}{x^{2/3}} = \int dt \quad (\text{use } \int u^n du = \frac{1}{n+1} u^{n+1} + C)$$

$$\int x^{-2/3} dx = 3x^{1/3} + C \Rightarrow 3x^{1/3} + C = t$$

$$x(t) = \frac{1}{3}(t - c)^3 = \frac{1}{27}(t - c)^3$$

$$\text{but } x(0) = 0 \Rightarrow x(t) = \frac{1}{27}t^3$$

$$\text{If } x(t_0) = 0 \Rightarrow c = t_0$$

$$x(t) = \frac{1}{27}(t - t_0)^3$$

Graphs are horizontal translation of $\frac{1}{27}t^3$

Missing solution:

$$x(t) \equiv 0$$

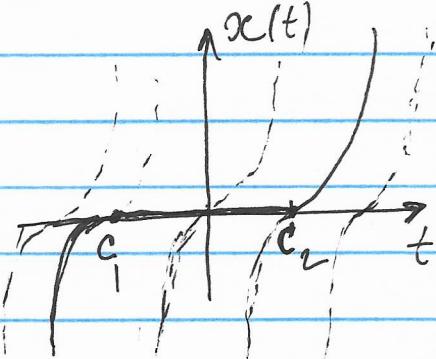
$x(t) = 0$ - singular solution which exists but cannot be constructed using separation of variables

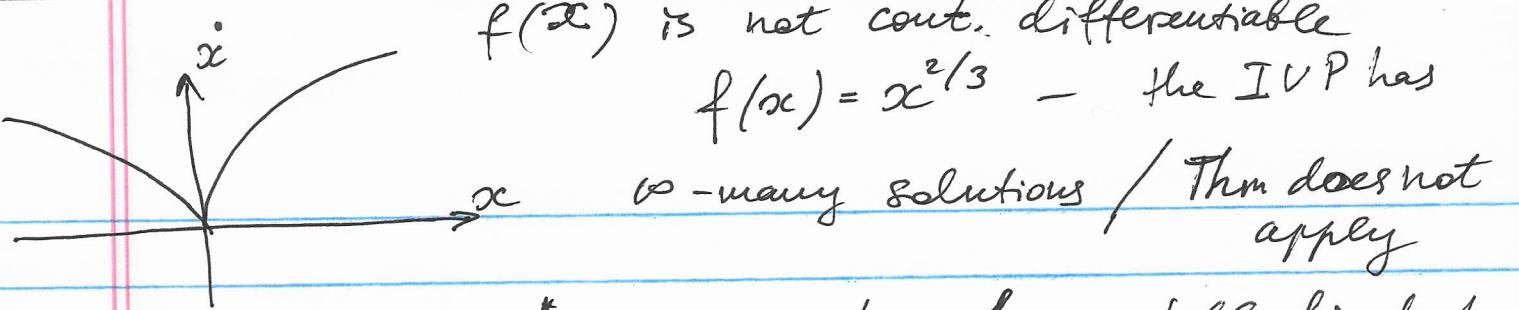
Why?

$f'(x)$ is not cont. at $x=0$; $f(x) = x^{2/3}$ is not cont. differentiable at $x=0$

\Rightarrow there are ∞ -many solns!

$$x(t) = \begin{cases} \frac{1}{27}(t - c_1), & -\infty < t < c_1 \\ 0, & c_1 \leq t \leq c_2 \\ \frac{1}{27}(t - c_2), & t > c_2 \end{cases}$$





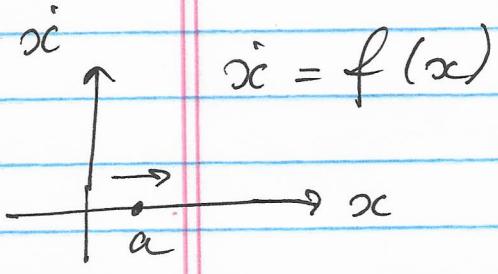
$x^* = 0$ is extremely unstable fixed pt.

$$f'(x) = (x^{2/3})' = \frac{2}{3} x^{-1/3} = \frac{2}{3} \frac{1}{x^{1/3}}$$

$$f'(0) = \frac{2}{3} \frac{1}{x^{1/3}} \Big|_{x=0} = \infty \Rightarrow \text{growth rate} = \infty$$

$$\text{Characteristic time} = \frac{1}{|f'(0)|} = 0 - \text{time scale}$$

Impossibility of oscillations in 1D



If for some point $x=a$,
particle passes it in positive direction,
velocity $\dot{x} > 0$.

To return back, particle would need to pass
this point in negative direction and to have
velocity $\dot{x} < 0$.

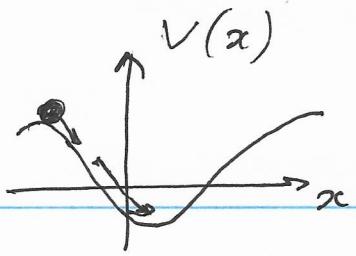
But particle cannot have both velocity $\dot{x} > 0$ and $\dot{x} < 0$
at the same point as $f(x)$ is single-valued
fn.

- ⇒ In 1D points move monotonically, never reverse direction
- ⇒ No periodic solutions!
- ? How about harmonic oscillator: $m\ddot{x} + kx = 0$?

2.7

Potentials

$$\dot{x} = f(x) \quad , \quad f(x) = -\frac{dV}{dx}$$



Particle sliding down the walls of potential well.

- Always downhill: to show this calculate

$$\frac{dV}{dt} = \frac{dV}{dx} \frac{dx}{dt} = \left(\frac{dV}{dx}\right)\dot{x} = \frac{dV}{dx} \cdot f(x) = -\left(\frac{dV}{dx}\right)^2 \leq 0$$

\Rightarrow Potential $V(t)$ decreases along the trajectory

\Rightarrow particle always moves toward lower potential.

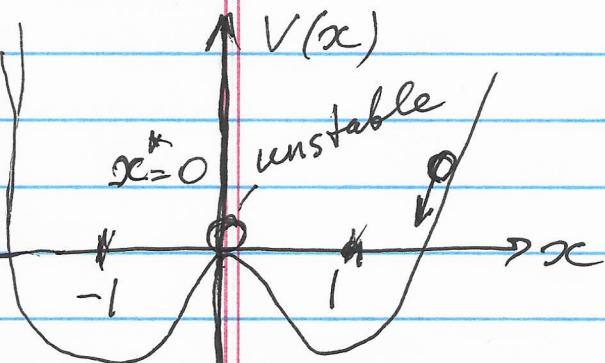
$$\frac{dV}{dx} = 0 \Leftrightarrow \dot{x} = 0 \quad - \text{equilibria points}$$

Min of $V(x)$ correspond to stable equilibria (fixed pts.)

Max of $V(x)$ \rightarrow unstable fixed pts.

$$\text{Ex. } \dot{x} = x - x^3$$

Graph the potential and find fixed pts.



$$x^* = -1 \quad \text{or} \quad x^* = 1$$

stable
fixed pts

$$f(x) = x - x^3 \quad - \text{Find } V(x)$$

$$-\frac{dV}{dx} = x - x^3 \quad \Rightarrow$$

$$V(x) = -\frac{1}{2}x^2 + \frac{1}{4}x^4 + C$$

put $C = 0$ (it doesn't matter)

Double-well potential

local minima: $x = \pm 1$ - stable
local max: $x = 0$ - unstable