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A Right and Left Truncated Gamma Distribution with Application to the Stars

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Abstract

The gamma density function is usually defined in interval between zero and infinity. This paper introduces an upper and a lower boundary to this distribution. The parameters which characterize the truncated gamma distribution are evaluated. A statistical test is performed on two samples of stars. A comparison with the lognormal and the four power law distribution is made.

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1 Introduction

A probability distribution function (PDF) which models a given physical variable is usually defined in the interval $0 \le x < \infty$. As an example the exponential, the gamma, the lognormal, the Pareto and the Weibull PDFs are defined in such interval, see [1]. We now briefly review the status of the research on the truncated gamma distribution (TG). A first attempt to deduce the parameters of a TG can be found in [2], [3] derived the minimum variance unbiased estimate of the reliability function associated with the TG distribution which is right truncated, [4, 5] estimated the parameters of a TG distribution over $0 \le x < t$, adopting the maximum likelihood estimator(MLE), [6] studied the properties of TG distributions and derived the simulation algorithms

which dominate the standard algorithms for these distributions, [7] considered a doubly-truncated gamma random variable restricted by both a lower (l) and upper (u) truncation.

On adopting an astronomical point of view the left truncation is connected with the minimum mass of a star, $\approx 0.02 M_{\odot}$ and the right truncation with the maximum mass of a star, $\approx 60 M_{\odot}$, see [8]. This paper first review the gamma PDF, introduces the right and left truncated gamma PDF and finally analyzes two samples of stars and brown dwarfs (BD).

2 The various gamma distributions

This Section reviews the gamma PDF, introduces the truncated gamma PDF and analyzes the data of two astronomical samples.

2.1 The gamma distribution

Let X be a random variable taking values x in the interval $[0, \infty]$; the gamma PDF is

$$f(x;b,c) = \frac{\left(\frac{x}{b}\right)^{c-1} e^{-\frac{x}{b}}}{b\Gamma(c)}$$
(1)

where

$$\Gamma(z) = \int_0^\infty e^{-t} t^{z-1} dt \quad , \tag{2}$$

is the gamma function, b > 0 is the scale and c > 0 is the shape, see formula (17.23) in [5]. Its expected value is

$$E(x;b,c) = bc \quad , \tag{3}$$

and its variance,

$$Var(x;b,c) = b^2c \quad . \tag{4}$$

The mode is at

$$m(x;b,c) = bc - b \quad when c > 1 \quad . \tag{5}$$

The distribution function (DF) is

$$DF(x;b,c) = \frac{\gamma(c,\frac{x}{b})}{\Gamma(c)} \quad , \tag{6}$$

where

$$\gamma(a,z) = \int_0^z t^{a-1} e^{-t} dt,$$
 (7)

is the lower incomplete gamma function, see [9, 10]. The two parameters can be estimated by matching the moments

$$b = \frac{s^2}{\bar{x}} \tag{8}$$

$$c = \left(\frac{\bar{x}}{s}\right)^2 \quad , \tag{9}$$

where s^2 and \bar{x} are the sample variance and the sample mean. More details can be found in [1].

2.2 The truncated gamma distribution

Let X be a random variable taking values x in the interval $[x_l, x_u]$; the truncated gamma (TG) PDF is

$$f(x;b,c,x_l,x_u) = k \left(\frac{x}{b}\right)^{c-1} e^{-\frac{x}{b}}$$

$$\tag{10}$$

where the constant k is

$$k = \frac{c}{b\Gamma\left(1 + c, \frac{x_l}{b}\right) - b\Gamma\left(1 + c, \frac{x_u}{b}\right) + e^{-\frac{x_u}{b}}b^{-c+1}x_u^c - e^{-\frac{x_l}{b}}b^{-c+1}x_l^c} \quad , \quad (11)$$

where

$$\Gamma(a,z) = \int_{z}^{\infty} t^{a-1} e^{-t} dt, \qquad (12)$$

is the upper incomplete gamma function, see [9, 10]. Its expected value is

$$E(b, c, x_l, x_u) = -b^2 k \left(-\Gamma \left(1 + c, \frac{x_l}{b} \right) + \Gamma \left(1 + c, \frac{x_u}{b} \right) \right) \quad . \tag{13}$$

The mode is at

$$m(x; b, c, x_l, x_u) = bc - b \quad when c > 1 \quad , \tag{14}$$

but in order to exist the inequality $x_l < m < x_u$ should be satisfied. The distribution function is

$$DF(x;b,c,x_l,x_u) = k\left(b\Gamma\left(1+c,\frac{x_l}{b}\right) - b\Gamma\left(1+c,\frac{x}{b}\right) + e^{-\frac{x}{b}}b^{-c+1}x^c - e^{-\frac{x_l}{b}}b^{-c+1}x_l^c\right) \quad . \tag{15}$$

A random number generation can be implemented by solving for x the following nonlinear equation

$$DF(x;b,c,x_l,x_u) - \mathbf{R} = 0 \quad , \tag{16}$$

where we have a pudendum number generator giving random numbers \mathbf{R} between zero and one, see [11]. A simple derivation of the lower and upper boundaries gives

$$\tilde{x_l} = minimum \ of \ sample \ \tilde{x_u} = maximum \ of \ sample \ .$$
 (17)

A first approximate derivation of \tilde{b} and \tilde{c} is through the standard estimation of parameters of the gamma distribution. We compute the χ^2 with these first values of \tilde{b} and \tilde{c} and we search a numerical couple which gives the minimum χ^2 . The χ^2 is computed according to the formula

$$\chi^2 = \sum_{i=1}^n \frac{(T_i - O_i)^2}{T_i},\tag{18}$$

where n is the number of bins, T_i is the theoretical value, and O_i is the experimental value represented by the frequencies. The merit function χ^2_{red} is evaluated by

$$\chi_{red}^2 = \chi^2 / NF \quad , \tag{19}$$

where NF = n - k is the number of degrees of freedom, n is the number of bins, and k is the number of parameters. The goodness of the fit can be expressed by the probability Q, see equation 15.2.12 in [12], which involves the degrees of freedom and the χ^2 . The Akaike information criterion (AIC), see [13], is defined by

$$AIC = 2k - 2ln(L) \quad , \tag{20}$$

where L is the likelihood function and k the number of free parameters in the model. We assume a Gaussian distribution for the errors and the likelihood function can be derived from the χ^2 statistic $L \propto \exp(-\frac{\chi^2}{2})$ where χ^2 has been computed by Equation (18), see [14], [15]. Now the AIC becomes

$$AIC = 2k + \chi^2 \quad . \tag{21}$$

2.3 Data analysis

A first test is performed on the low-mass initial mass function in the young cluster NGC 6611, see [16]. Table 1 shows the values of χ^2_{red} , the AIC, the probability Q, of the astrophysical fits and the results of the K-S test, the maximum distance, D, between the theoretical and the astronomical DF as well the significance level P_{KS} , see [17, 18, 19, 12]. Figure 1 shows the fit with the TG distribution of NGC 6611 and Figure 2 visually compares the three types of fits for NGC 6611.

A second test is performed on low-mass stars in NGC 2362, see [20]. Table 2 shows the statistical parameters which characterize the astrophysical fits. Figure 3 shows the fit with the TG distribution of NGC 2362 and Figure 4 visually compares the three types of fits for NGC 2362.

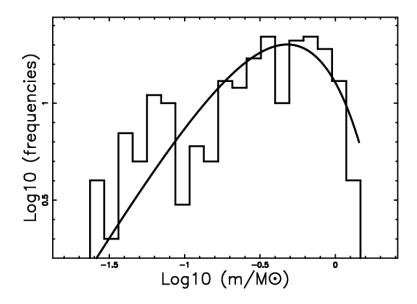


Figure 1: Logarithmic histogram of mass distribution as given by NGC 6611 cluster data (207 stars + BDs) with a superposition of the TG distribution when the number of bins, n, is 12, c = 1.287, b = 0.372, $x_l = 0.019$ and $x_u = 1.36$. Vertical and horizontal axes have logarithmic scales.

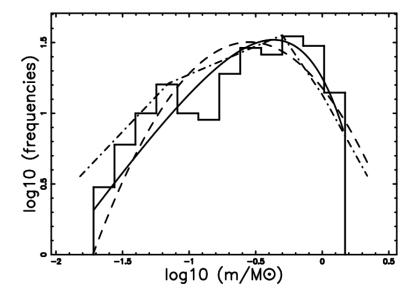


Figure 2: Histogram (step-diagram) of mass distribution as given by NGC 6611 cluster data (207 stars + BDs) with a superposition of the left TG distribution (full line), the lognormal (dashed), and the four power laws (dot-dash-dot-dash). Vertical and horizontal axes have logarithmic scales.

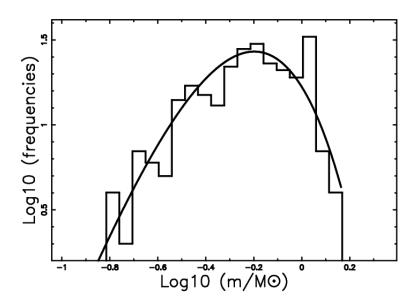


Figure 3: Logarithmic histogram of mass distribution as given by NGC 2362 cluster data (272 stars) with a superposition of the TG distribution when the number of bins, n, is 12, b=0.161, c=3.933, $x_l=0.12$ and $x_u=1.47$. Vertical and horizontal axes have logarithmic scales.

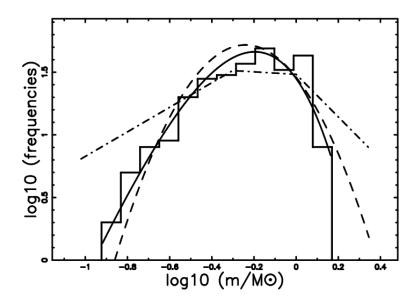


Figure 4: Histogram (step-diagram) of mass distribution as given by NGC 2362 cluster data (272 stars) with a superposition of the left TG distribution (full line), the lognornal (dashed), and the four power laws (dot-dash-dot-dash). Vertical and horizontal axes have logarithmic scales.

Table 1: Numerical values of NGC 6611 cluster data (207 stars + BDs). The number of linear bins, n, is 20.

PDF	parameters	AIC	χ^2_{red}	Q	D	P_{KS}
lognormal	$\sigma = 1.029, \mu_{LN} = -1.258$	71.24	3.73	1.310^{-7}	0.09366	0.04959
gamma	b=0.248, $c=1.717$	62.83	3.26	3.1510^6	0.109	0.0124
truncated gamma	b=0.372, $c=1.287$	52.34	2.77	0.00017	0.09	0.061
	$x_l = 0.019, x_u = 1.46$					
four	Eqn.(59)	81.39	5.18	2.4119^{-9}	0.12514	2.7210^{-3}
power laws	in Zaninetti 2013					

Table 2: Numerical values of the NGC 2362 cluster data (272 stars). The number of linear bins, n, is 20.

PDF	parameters	AIC	χ^2_{red}	Q	D	P_{KS}
lognormal	$\sigma = 0.5, \mu_{LN} = -0.55$	37.64	1.86	0.013	0.07305	0.10486
gamma	b=0.13 ,c $=4.955$	34.28	1.68	0.034	0.059	0.284
truncated gamma	b=0.161 ,c =3.933	33.88	1.61	0.055	0.071	0.122
	$x_l = 0.12, x_u = 1.47$					
four	Eqn.(58)	77.608	4.89	1.1710^{-8}	0.16941	2.610^{-7}
power laws	in Zaninetti 2013					

3 Conclusions

The right or left TG PDF has been extensively investigated in the field of mathematics , as an example [7] reports most of the mathematical details. The application of the TG PDF in astronomy represents conversely a new promising field. Here we have deduced the constant of normalization ,eqn.(11), the average value ,eqn.(13), the DF , eqn.(15), and presented an algorithm for the generation of the random numbers , (eqn.16). The application of the TG PDF to the IMF is positive and both the reduced χ^2 and the K-S test give better results in respect to the standard PDFs used by the astronomers which are the lognormal and the four power laws , see Tables 1 and 2. A comparison with the left truncated beta PDF , see Tables 1 and 2 in [21] allows to say that the left truncated beta PDF produces a better fit to the IMF in respect to the truncated gamma PDF here analyzed.

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