

**DIFFERENTIAL JET PRODUCTION CROSS SECTION MEASUREMENT
IN Z + JET EVENTS FROM PROTON - PROTON COLLISIONS AT $\sqrt{S} = 13$
TEV USING THE CMS DETECTOR AT LHC**

by

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Abstract

The standard model of particle physics, while describing our universe well on many scales, has yet to be precisely measured in all energy regimes. Recent theoretical advances in higher order QCD calculations have provided a way to compare the standard model's predictions to precision measurements of data and simulation. Within this dissertation, I present a measurement of the double differential jet production cross section as a function of the jet mass and transverse momentum. The measurement is presented in events with a Z + Jet topology, with and without soft radiation included in the jet. The removal of the soft radiation is done using a jet "grooming" algorithm.

a jet grooming algorithm applied.

Studying Z + jet events will yield a light quark enriched jet sample, which has not yet been studied at $\sqrt{s} = 13$ TeV. Comparing groomed and ungroomed jets will allow us to better understand the jet mass in all energy regimes since the groomed jets will have varying amounts of soft and collinear radiation with respect to the ungroomed counterpart. For ungroomed jets, leading-order and next-to-leading order QCD Monte Carlo programs are found to predict the jet mass spectrum in the data reasonably well, with some disagreement at very low masses. For groomed jets, the agreement between the Monte Carlo programs and the data improves overall, and extends lower in jet mass due to the

removal of soft and colinear portions of the jet. First-principles theoretical calculations of the groomed jet mass are being prepared currently by colleagues in order to make comparisons with the measurement presented herein. Ultimately these measurements will be used to tune Monte Carlo generators, producing more accurate parton showering simulations, leading to tighter constraint of backgrounds in future searches for new physics.

Chapter 1

Introduction

This thesis presents a measurement of the differential production cross section of Z+jets events as a function of the jet mass and transverse momentum (p_T). The cross section is presented for events before and after the jets are groomed with the “soft drop” procedure [4], using multiple values of the tunable parameters β and z_{cut} . Each unique combination of the tunable parameters β and z_{cut} leads to a jet with less soft and collinear radiation with respect to the ungroomed counterpart in varying degrees depending on the choice. The soft drop grooming algorithm is described in more detail in 3.2.

Comparing the production cross section, described in ?? , for groomed and ungroomed jets separately allows us to, in the language of Soft-Collinear Effective Theory (SCET) described in 3.5, gain sensitivity to both the “hard” and “soft” jet physics. The groomed cross section can be directly compared to theoretical calculations of the jet mass now and in the future, which is a very active area of theoretical research at this time [7, 8, 9, 10, 11, 12, 13, 14, 15]. Furthermore, separating the hard and soft jet physics allows a deeper understanding of the various effects involved in QCD radiation. In particular, Ref. [14] calculates

the groomed jet mass at next-to-next-to-leading order using soft colinear effective theory, matched to a parton shower at leading order using MCFM [16, 17], and the authors of Ref. [5] are preparing a next-to-leading logarithm calculation with traditional perturbative QCD, matched to a parton shower at leading order, also using MCFM. In the near future we will compare these theoretical predictions to our data measurement. Both CMS and ATLAS have similar measurements in a dijet sample at Ref. [18, 19].

The analysis strategy is similar to that of Ref. [18]. However, there are several differences. As in that paper, the cross section is now unfolded in both jet mass and p_T . However, while the previous measurement considered only one value for the soft drop parameter β , this analysis considers several. We apply the soft drop algorithm to compare directly to theoretical computations. Additionally, we not only measure the cross section as a function of mass, but also as a function of dimensionless mass, $\rho = 2\log(m/(ptR))$, as is also done in the previously mentioned ATLAS measurement. The dimensionless mass ρ only weakly depends on p_T , unlike mass, which is highly correlated. Additionally, the use of this variable aids in the separation of fixed order, perturbative and non-perturbative effects. We present the normalized double differential jet production cross section with respect to jet mass and transverse momentum (p_T) as well as with respect to jet dimensionless mass and p_T . We compute the cross sections normalized per p_T bin (the “normalized” cross section) with respect to the jet p_T and jet mass by unfolding a binned two-dimensional distribution in p_T and mass with widths Δp_T and Δm , respectively.

The normalized differential cross section

$$\frac{1}{d\sigma/dpt} \frac{d^2\sigma}{dpt dm} = \frac{1}{N/\Delta pt} R\left(\frac{N_{ij}}{\Delta pt \Delta m}\right) \quad (1.1)$$

where N is the total number of Z +jets events in our selection, N_{ij} is the number of such events in pt bin i and mass bin j , and $R(\alpha)$ is the unfolding procedure applied to the two-dimensional distribution α .

Chapter 2

Theoretical Framework

2.1 Introduction To The Standard Model

The Standard Model (SM) of particle physics is a quantum field theory (QFT) description of the strong, weak and electromagnetic forces of nature. The known particles of the SM are: 1 scalar Higgs boson, 4 gauge bosons, 6 types of quarks and 6 types of leptons.

The quarks and leptons are fermions that constitute matter and so obey Fermi-Dirac statistics due to their half-integer spin. In contrast, the bosons have integer spin and obey Bose-Einstein statistics. Gauge bosons mediate the 3 fundamental forces and the Higg's boson is responsible for the electro-weak symmetry breaking which gives mass to the other particles [20].

The fermions are arranged into 3 generations, arranged in columns from left to right on Figure 2.1

The SM is humanity's most rigorous theory of our universe, providing predictions of observables that have since been measured, in the case of Quantum Electrodynamics, QED, to the highest precision of any scientific theory. Despite

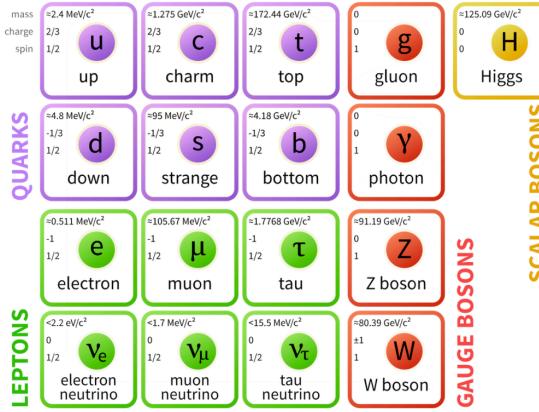


Figure 2.1: Fundamental particles of the Standard Model [1].

the impressive predictions, the gravitational force and more subtle phenomena, such as flavor oscillation of neutrinos [21], indicate the existence of physics beyond the standard model, BSM.

Various attempts have been made to unify the fundamental forces under one theory, thus far the electromagnetic and weak interactions have been united by electro-weak theory.

The Standard Electroweak Model can be described $SU(2) \times U(1)$ mathematically. The $SU(2) \times U(1)$ gauge group is a unification of the special unitary symmetry group $SU(2)$ describing 3 mixed massive vector bosons, (W_- W_+ Z_0), as carriers of the weak nuclear force, and the unitary gauge group $U(1)$, describ-

ing the massless chargeless photon, of the electromagnetic interaction.

The standard model of the strong interaction is known as Quantum Chromodynamics, QCD, a non-Abelian gauge theory described by the special unitary group ($SU(3)_c$), where the flavours of quark are the physical manifestation of the symmetry group. This force is mediated by the 8 massless gluons that carry color charge, making QCD more complicated mathematically than QED.

The SM also contains a Higgs boson, an excitation of a scalar Higg's field, which gives rise to spontaneous symmetry breaking of the electroweak theory, providing the particles with mass, but that won't be discussed herein.

The quarks and leptons are arranged in generations according to their relative masses, as shown in Figure 2.1. The table also shows the spins of the particles, the leptons and quarks have half-integer spin, fermions, that obey the fermi exclusion principle. Conversely, the bosons have half integer spin and therefore obey bose-einstein statistics. Through the SM we interpret the observed hadronic particles, mesons (baryons), as 2 quark (3 quark) bound states. The existence of spin $\frac{3}{2}$ baryons, which are symmetric bound states in space, spin and flavour, and the need to obey Fermi-Dirac statistics, by maintaining total assymmetry of the wave function, implies there is another degree of freedom, called color, so that each quark is either red, green or blue. Granted only color singlet states exist. Furthermore there exists a property of asymptotic freedom where the QCD coupling between quarks and gluons increases as they asymptotically approach one another. There exist a wealth of experimental data to support the concept of asymptotic freedom. Asymptotic freedom is a useful property as it allows for perturbative calculations of QCD observables, such as the jet mass, this is discussed in Section 3.4.3. The running of the QCD (strong) coupling constant as measured by CMS experiment can be seen in Figure 2.2

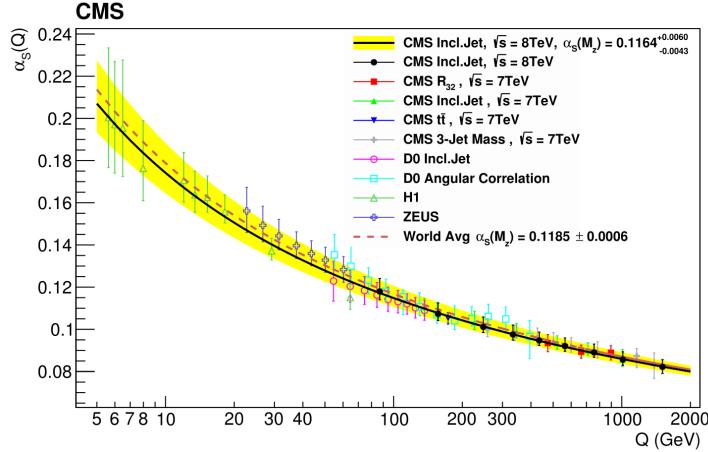


Figure 2.2: The running of the strong coupling constant as compiled by CMS including measurements from CMS and HERA among others [2].

Nuclei in ordinary matter are composed solely of 1st generation particles, up and down quarks, bound by gluons. Neutral atoms contain an equal number of protons (composed of 2 up quarks and a down quark) and electrons, 1st generation leptons. The main distinction between leptons and quarks, both fermions (particles of $\frac{1}{2}$ integer spin), being that leptons do not experience the color interaction ($SU(3)_c$) like their quark friends. In each generation there is a quark with charge $Q = +\frac{2}{3}$ (up, charm, top) and another of charge $Q = -\frac{1}{3}$ (down, strange, bottom).

2.1.1 Quantum Chromodynamics

QCD is a quantum field theory that describes the color force, experienced by quarks and mediated by gluons. The quarks each possess a color charge ; red, green or blue (anti-red, anti-green or anti-blue). In contrast the gluons are "bicolored", each carrying one unit of negative and one of positive charge. [20]. As previously mentioned, QCD a non-Abelian gauge theory described by the special

unitary group ($SU(3)_c$ (color charge) and has a Lagrangian that can be written as follows:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^A F_A^{\mu\nu} + \sum_{\text{flavours}} \bar{\psi}_a (i\gamma_\mu D^\mu - m)_{ab} \psi_b \quad (2.1)$$

The sum is over each of the generators of the ($SU(3)_c$ gauge group, and the quark field fermion multiplets, ψ , belong to its irreducible representation. [22]

The gluon fields A_ν^a of spin 1 have a field strength tensor, $F_{\mu\nu}^A$, given below:

$$F_{\mu\nu}^A = \partial_\mu A_\nu^A - \partial_\nu A_\mu^A + g_s f^{ABC} A_\mu^B A_\nu^C \quad (2.2)$$

The structure functions are denoted as f^{ABC} and their indices run over all of the gluon color degrees of freedom. It is noteable to mention that the third term in the above equation is what gives rise to asymptotic freedom through the gluon quartic and triple self-interactions it induces [22].

The covariant derivative is defined as :

$$(D_\mu)_{ab} = \partial_\mu \delta_{ab} - i g_s A_\mu^A t_{ab}^A \quad (2.3)$$

Where the t_a are the matrices of the fundamental representation of ($SU(3)$).

Lastly, for completeness, I mention the final gauge invariant term of the QCD Lagrangian below, where θ_{QCD} is a free parameter of QCD known as the vacuum angle parameter.

$$\mathcal{L}_\theta = \theta_{\text{QCD}} \frac{\alpha_s^2}{64\pi^2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^a F_{\rho\sigma}^a \quad (2.4)$$

QCD is discussed from a phenomenological perspective in Section 3.6.

Chapter 3

Jets in Proton-Proton Collisions

A jet is a collimated grouping of hadrons usually associated with the LO production of a parton, quark or gluon. In this case the quark (gluon) could be initiated by the hard scatter of 2 constituent partons from protons or be the decay product of another particle that was produced in the hard scatter, such as a W boson, which can then decay to 2 quarks each forming a jet. The initial parton radiates other quarks and gluons, called the "Parton Shower" and all color charged particles fragment into hadrons, mainly pions and kaons, before reaching the detector as mentioned in 5.

Studies of jets at LHC are complicated by experimental complexities such as "underlying event", other partons from the same proton interacting and depositing energy in the same region of the detector. "Pileup" is also increasingly relevant, like "underlying event" but initiated from other proton interactions from this or a previous bunch, since the LHC crosses bunches with 10^{11} protons every 25 nanoseconds.

Any measurement is limited by the resolution of the measurement device and any detector effect. In this thesis the results are disentangled from the re-

constructed data by "unfolding" the reconstructed distributions back to generator level.

While jets are often used as simple proxies for the quark or gluon from which they originated, the structure of the radiation pattern of the hard scatter is encoded within the jet's constituent particles [23]. Jet studies are essential for a complete understanding of proton-proton interactions since the majority of interesting physics signatures contain a color charged parton in the final state. This chapter covers the basics of jet physics at LHC, from the algorithms used for clustering the constituent particles in experimental data to the language and calculations of the theory.

3.1 Jet Clustering Algorithms

At LO, a jet represents a quark [gluon], however realistically that is not a well defined concept and a more precise definition is useful. This thesis adopts the definition presented at the Les Houches conference in 2015 :

"A phase space region (as defined by an unambiguous hadronic fiducial cross section measurement) that yields an enriched sample of quarks [gluons] (as interpreted by some suitable, though fundamentally ambiguous, criterion)" [24].

I will discuss one class of "suitable, though fundamentally ambiguous" criteria for defining jets, known as sequential recombination algorithms. These algorithms take pairs of particles and successively combine them into 1 particle, in a way which is intended to reconstruct the successive branchings of partons within the jet as described by perturbative QCD [25].

In sequential recombination algorithms a distance metric, d_{ij} , is defined between all particle pairs. These pairs are then sequentially combined in order of increasing distance. Three popular algorithms in this class can be described by the equation 3.1:

$$\begin{aligned} d_{ij} &= \min \left(p_{ti}^{2p}, p_{tj}^{2p} \right) \frac{\Delta R_{ij}^2}{R^2} \quad \Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2 \\ d_{iB} &= p_{ti}^{2p} \end{aligned} \quad (3.1)$$

Depending on the value chosen for p , this equation can produce a variety of clusterings, herein I discuss the 3 popular choices $p = [1, 0, -1]$ referring to them by their names: KT [26], Cambridge/Aachen [27] and Antik_T algorithms [28] respectively.

Each algorithm has its own specific use cases and drawbacks. For example, the Antik_T algorithm produces circular jets of constant area as seen in 3.1.

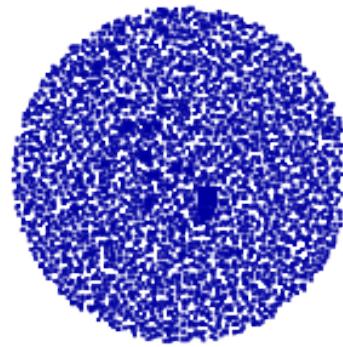


Figure 3.1: For jets clustered using the Antik_T algorithm a circular pattern emerges, making pileup and underlying event subtraction simpler for experimentalists [3].

This constant area is useful for experimental studies as it allows for simple removal of underlying event energy, which is evenly distributed throughout the jet area [3]. In contrast to the other 2 algorithms described, the C/A algorithm does not contain any momentum weighting, this leads to jets which are not circular, depicted in 3.3, and instead follow the radiation pattern of the original jet constituent particles as seen in 3.3. The variation in the way these algorithms cluster the underlying radiation is depicted in 3.2. In this thesis we have chosen to recluster the Antik_T jets with C/A in order to preserve the original radiation pattern.

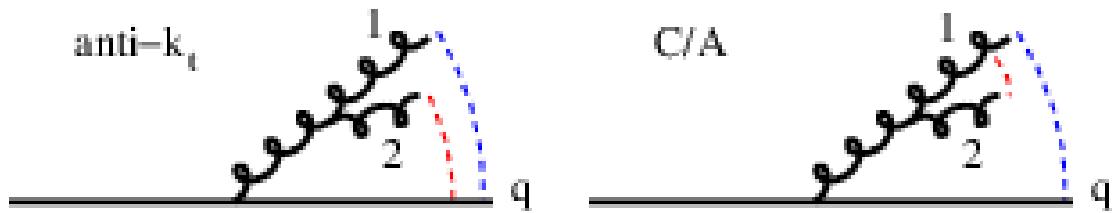


Figure 3.2: Notice that the clustering follows the radiation pattern for C/A while the softer emissions are clustered with the jet axis in the case of Antik_T [3].

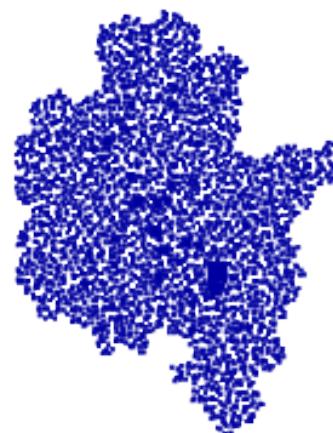


Figure 3.3: Jets clustered using C/A algorithm are not circular, instead they are shaped like the radiation pattern [3].

3.2 Jet Grooming

Jet grooming is a broad term to describe a number of different algorithms intended to remove some portion of the soft and collinear radiation within a jet. This analysis uses the Soft-Drop (SD) algorithm [4] which removes soft and collinear radiation in a theoretically controlled manner. The SD procedure is as follows: Begin with an Antik_T [28] clustered jet composed of particle flow candidates, reclustered with Cambridge-Aachen algorithm [27] to remove the p_T weighting dependance of the clustering. At this stage we regress through the clustering history, keeping a subjet if it meets of SD criterion and otherwise "dropping" it.

SD iteratively declusters a jet j with distance parameter R into two subjets, j_1 and j_2 . If the softdrop condition

$$\frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}} > z_{cut} \cdot \left(\frac{\Delta R_{12}}{R}\right)^\beta \quad (3.2)$$

is met, then the procedure terminates and j is the final state jet. Otherwise, the declustering continues - the higher p_T subjet is relabeled as j and the lower p_T ("softer") subjet is dropped, hence the name. By design, this condition fails for wide-angle soft radiation, which is therefore removed by the soft drop procedure. The tunable parameters, β and z_{cut} , control the degree of jet grooming: β tunes the algorithm's sensitivity to wide-angle radiation, while z_{cut} sets the energy scale of the grooming. In the case of $\beta \rightarrow \infty$, an ungroomed jet is returned. In the $\beta = 0$ case, the soft drop procedure is identical to the "modified mass drop tagger" (MMDT) from Ref. [5]. For $\beta > 0$ soft emissions are removed from the jet while the majority of the collinear emissions remain, as illustrated in 3.4.

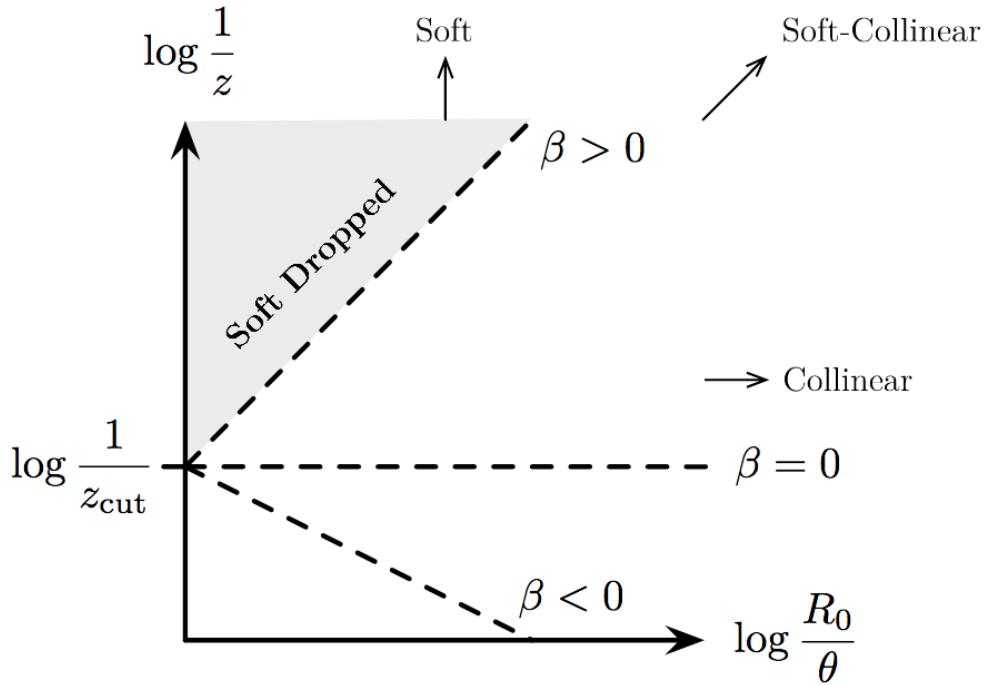


Figure 3.4: An illustration of the phase space for jet emissions outlining which would be groomed away by the SD procedure [4].

The soft drop algorithm removes soft and wide-angle radiation from jets in a very theoretically controlled manner, making it suitable to separate the “hard” and “soft” parts of the jet. Specifically, the soft drop algorithm can remove non-global logarithms from correlations of radiation within and between jets, which are extremely difficult to compute theoretically [29, 5, 4, 30, 31, 32].

3.3 Lund Jet Plane

A Lund kinematic Diagram, shown in 3.5, is an illustration of the kinematics of emissions within a jet in terms of two variables: horizontally, the logarithm of the inverse of the angle of a given emission with respect to the jet axis, and

vertically the logarithm of the emission's transverse momentum k_t . The line of constant jet mass and the shaded region constitute the part of the phase space in which emissions are vetoed, this leads to a Sudakov form factor [5].

The Lund kinematic diagram, or Lund jet plane, is a theoretical representation of the phase space within jets, and has been a useful tool for studying parton showers in the past, and has now been extended to individual jets through repeated Cambridge/Aachen declustering [3].

The primary Lund plane is shown on the bottom of Figure 3.6 [3] where the black particle (a) "initiated" the jet and is described by the grey shaded region and particles (b) and (c) are marked with dots and constitute emissions within the jet. In contrast, the Lund diagram shows the phase space encompassed by each subsequent emission, not just the emission (a). [3]

3.4 Jet Mass

The calculations for the jet mass in perturbative QCD are described in this section. Jet invariant mass is a simple observable, useful for probing QCD, it is defined as:

$$m^2 = \left(\sum_{i \in \text{jet}} k_i \right)^2 \quad (3.3)$$

the summation includes all particles i clustered into the jet. In the case of this thesis the particles are particle-flow candidates, Section 5.2 contains information on how the events for this thesis were reconstructed.

The plain jet mass is discussed in the first subsection while a later subsec-

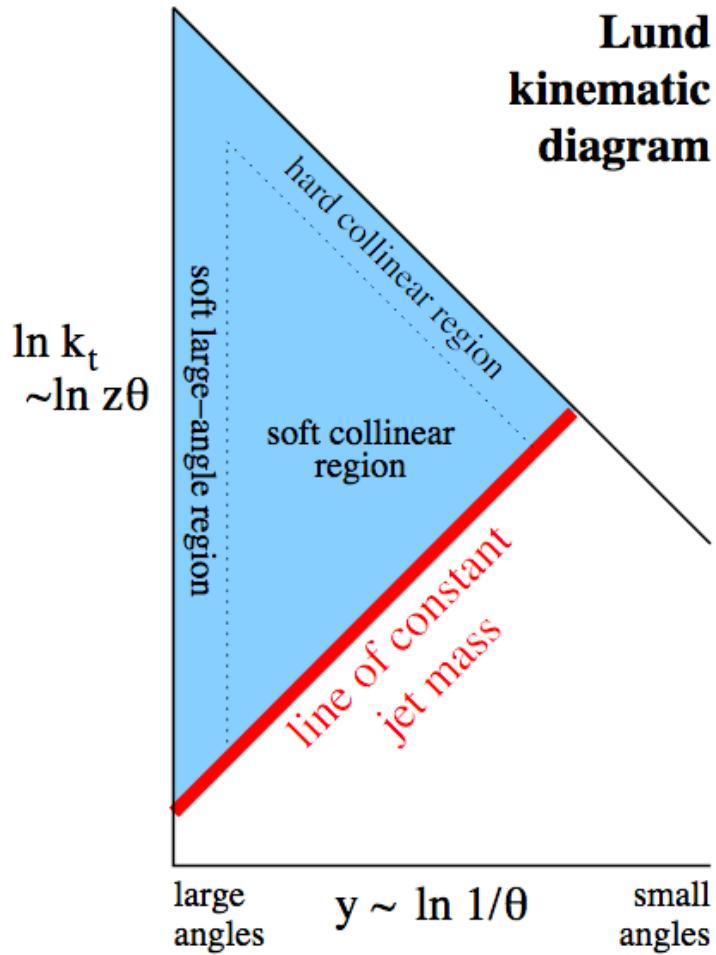


Figure 3.5: A Lund kinematic Diagram. The line of constant jet mass is shown in red [5].

tion shows how the jet mass distribution is calculated for jets modified by a “grooming” algorithm, in this case soft-drop [4], as are the jets in the groomed measurement presented herein.

This discussion centers on “QCD jets”, initiated by a hard parton and evolv-

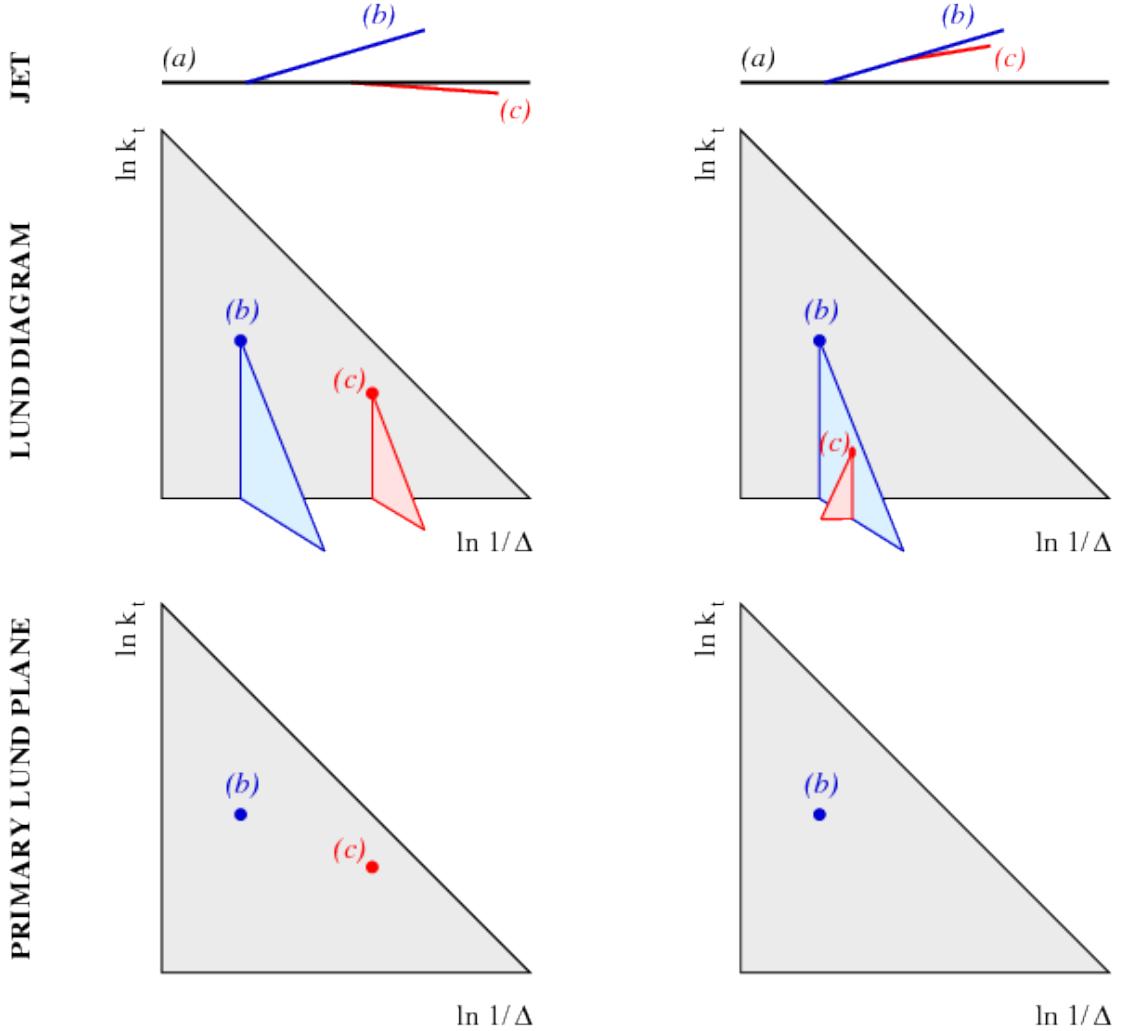


Figure 3.6: The primary and secondary lund planes for 2 example jets [3].

ing via a parton shower but stopping at parton level as perturbation theory is not able to describe the transition to particle level [25]. Event generators employ different hadronization techniques to describe the parton-to-hadron transition, those used in this thesis are discussed in Section 5.1.

Within the valid phase space of perturbative QCD one finds that the fixed-order approach to jet mass is inadequate, most noticeably in the boosted regime as each emission contributes a large logarithm [25]. All-order resummation techniques are then exploited to theoretically describe the jet mass observable

more accurately than is possible with fixed-order.

3.4.1 Plain Jet Mass

In this section quark jet mass calculations are discussed at LO. It is first useful to define the dimensionless jet mass, ρ , as the plain mass is p_T dependant.

$$\rho = m^2 / (E^2 R^2) \quad (3.4)$$

Using the small-angle approximation, ρ is invariant under Lorentz boosts along the jet axis direction, since boosts vary the jet p_T up by some factor (say κ) and scale its opening angle by the inverse factor ($\frac{1}{\kappa}$) while leaving the mass unchanged. Due to this invariance, the analytical results are often simplest when expressed in terms of ρ [5].

Now the integrated jet mass distribution can be defined as :

$$D(\rho) = \int_{\rho}^1 \frac{d\rho'}{\rho'} \int_{\rho'}^1 dz p_{gq}(z) \frac{\alpha_s(z\rho'R^2 p_t^2) C_F}{\pi} \quad (3.5)$$

Above the quark-gluon splitting function is $p_{gq} = \frac{1+(1-z)^2}{2z}$. This can be elucidating if we take the fixed coupling approximation to reveal the below relationship. [5]

$$D(\rho) \simeq \frac{\alpha_s C_F}{\pi} \left[\frac{1}{2} \ln^2 \frac{1}{\rho} - \frac{3}{4} \ln \frac{1}{\rho} + \mathcal{O}(1) \right], \quad (\text{fixed coupling approx.}) \quad (3.6)$$

In terms of the Lund diagram, discussed in 3.3, the above equation can be described simply, as the $\frac{1}{2} \ln^2 \frac{1}{\rho}$ term represents the bulk of the area of 3.5 and

the $-\frac{3}{4} \ln \frac{1}{\rho}$ term is due to the hard collinear region where z is finite.

The Next-to-Leading Logarithm (NLL) approximation gives the integrated jet mass distribution :

$$\Sigma(\rho) = e^{-D(\rho)} \cdot \frac{e^{-\gamma_E D'(\rho)}}{\Gamma(1 + D'(\rho))} \cdot \mathcal{N}(\rho) \quad (3.7)$$

In the above equation, the first term accounts for the Sudakov Suppression of emissions that would induce a jet mass greater than ρ . To restate this in the language of the Lund diagram, this is the probability that there are no emissions in the shaded region of 3.5. The central term encodes the information of single-logarithmic corrections from multiple emissions, generally near the line of constant jet mass in 3.5, sum together to obtain the total jet mass. The final factor accounts for variations in the jets radiation pattern (non-global logarithms) and boundaries of the jet (clustering logarithms) induced by soft radiation near the jet's edge, i.e. near the left-hand, vertical edge of the shaded region in 3.5 [5]. Note that the non-global logarithms are problematic and that is the main reason why there does not exist a full resummation of the standard jet mass beyond NLL accuracy [5].

In order to better visualize the jet mass distribution, one can take the fixed coupling approximation then look at only the leading order differential jet mass distribution. [5]

$$\frac{\rho}{\sigma} \frac{d\sigma}{d\rho} \simeq \frac{\alpha_s C_F}{\pi} \left(\ln \frac{1}{\rho} - \frac{3}{4} \right) e^{-\frac{\alpha_s C_F}{2\pi} \left(\ln^2 \frac{1}{\rho} - \frac{3}{2} \ln \frac{1}{\rho} + \mathcal{O}(1) \right)} \quad (3.8)$$

The above distribution allows one to see why the red line for plain jet mass looks as it does in the figure 3.7.

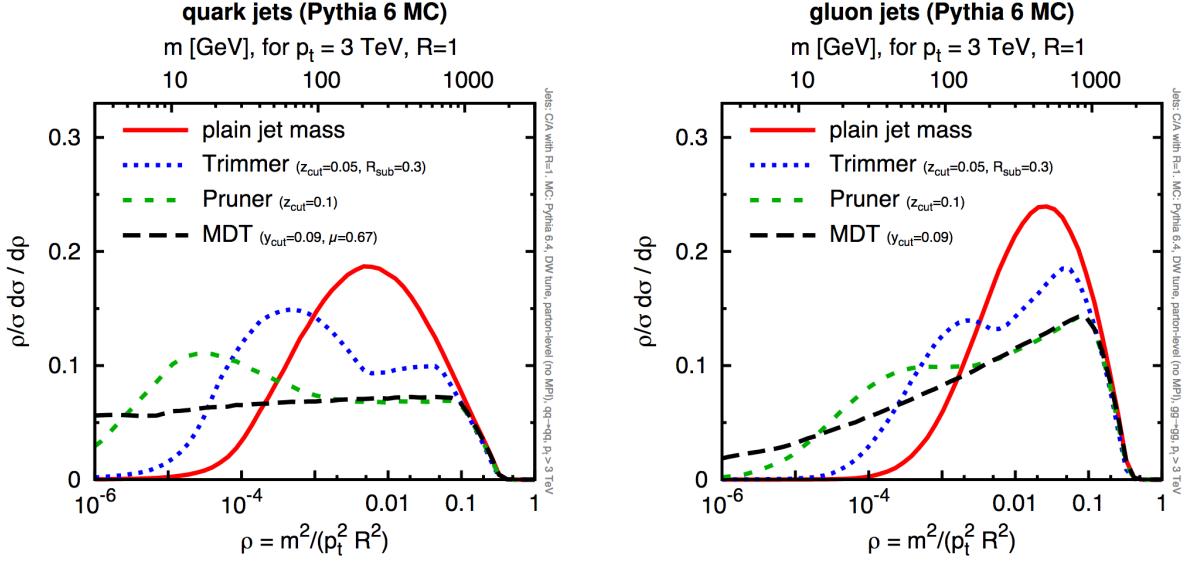


Figure 3.7: (Left) The distribution of ρ for tagged jets, with three taggers/groomers: pruning , trimming and the mass-drop tagger (MDT). (Right) Analogous distribution for gluon jets [5].

Starting from the far right of the distributions in 3.7 one can see that as ρ decreases the plain jet mass (red line) grows linearly in $\ln \frac{1}{\rho}$ only to be cut off by Sudakov suppression (The exponential term in the jet mass equation at LO). Notice that the peak of the distributions can be calculated using [5]:

$$L_{\text{peak}} = 1/\sqrt{\alpha_s} + \mathcal{O}(1) \quad (3.9)$$

Higher order corrections to the hard process are necessary in order to properly predict measurements in the non-perturbative regime. This can be seen clearly in the dijet dimensionless mass spectrum measured by the ATLAS collaboration (Maybe replace with our result but I am not sure which plot to use as ours seem to agree even at low mass[presumably bc we only showed the more

accurate predictions by Marzani and Frye rather than all of the others]) 3.8

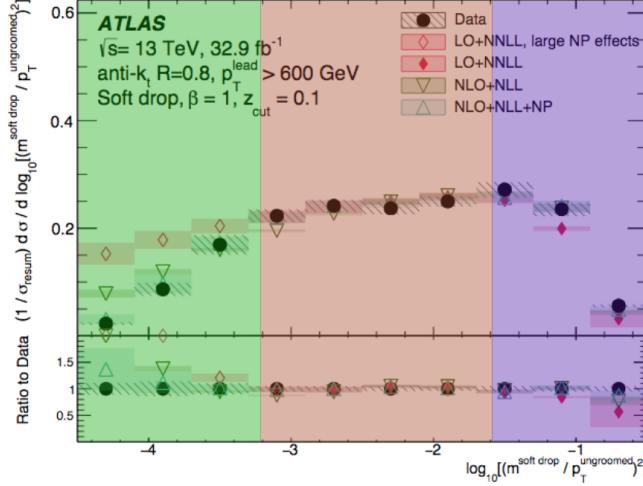


Figure 3.8: Above is the ATLAS Dijets dimensionless mass result [3]. Next to Leading Order + Next to Leading Log + NonPerturbative corrections were required in order to match data in non-perturbative regime (Far left of this plot, shaded green column).

3.4.2 Initial-State Radiation Contributions to Jet Mass

To expand upon the simple calculation above to include additional emissions, one can consider the case of emission of a soft gluon from a dipole formed by 2 hard quarks. This process is illustrated in Figure 3.9.

If the soft gluon from the kinematic diagram is clustered with the jet then it will make the following contribution to the jet mass [25]:

$$m^2 = (p_3 + k)^2 = 2p_3 \cdot k = 2p_t k_t (\cosh(\eta - y) - \cos \phi) \quad (3.10)$$

Given the above, one can write the contribution to the cumulative distribution as :

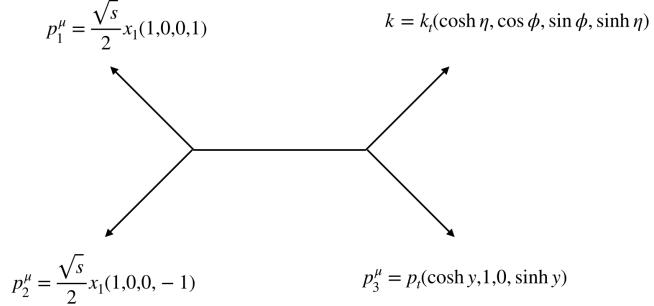


Figure 3.9: Above the incoming hard partons, of momenta p_1 and p_2 , form a jet of momentum p_3 and a soft gluon of momentum k_t is radiated as initial state radiation. For simplicity the assumption is made that the jet was produced at $\phi = 0$

$$\alpha_s \Sigma_{12}^{(1)} = C_{12} \int k_t dk_t d\eta \frac{d\phi}{2\pi} \frac{\alpha_s (\kappa_{12})}{2\pi} \frac{(p_1 \cdot p_2)}{(p_1 \cdot k)(p_2 \cdot k)} \Theta((\eta - y)^2 + \phi^2 < R^2) \cdot \left(\frac{2k_t}{p_t R^2} (\cosh(\eta - y) - \cos \phi) > \rho \right) \quad (3.11)$$

Θ is the jet clustering condition [25]. Integrating and simplifying one finds:

$$\alpha_s \Sigma_{12}^{(1)} = C_{12} R^2 \int_{\rho p_t}^{p_t} \frac{\alpha_s (k_t)}{2\pi} \frac{dk_t}{k_t} \quad (3.12)$$

The equation tells us that the soft gluon will not give rise to collinear enhancements since it is away from the hard partons in the process. The lower limit on the integration comes from the jet mass constraint.

3.4.3 Jet Mass in Proton-Proton $\rightarrow Z + \text{jet}$

The specific jet production channel studied in this thesis is $Z + \text{jet}$. This process is useful to study as it produces a sample of light quark enriched jets for the measurement as opposed to the dijet measurements previously performed by CMS and ATLAS, those studied a quark-gluon admixture enriched jet sample. In principle, the interaction is simpler than the dijet case having only 3 coloured hard legs as opposed to 4 [25]. Additionally, the $Z + \text{jets}$ channel is interesting to understand in the boosted regime, $p_T >> m$, as it is the main background in searches for a boosted Higgs boson recoiling off a Z boson [25].

The Figure ??, shows calculations at Born level for 2 partonic processes: $qg \rightarrow Zq$ and $q\bar{q} \rightarrow Zg$. The Feynman diagrams for these processes are shown on the bottom and top respectively, of Figure ??.

In Figure ?? it is noteworthy to realize that the bulk of the order R^2 corrections (black line) are composed of contributions from the simple ISR gluon as illustrated in Section 3.4.2

3.4.4 Groomed Jet Mass

The jet mass distribution for soft-dropped jets differs from the plain jet because it lacks soft and collinear radiation contributions. It is shown in the Lund diagram in Figure 3.11 at LL accuracy, where the shaded red area is that vetoed and associated with Sudakov suppression and the red line is a line of constant jet ρ .

For the soft-dropped mass the LO cumulative distribution can be written as :

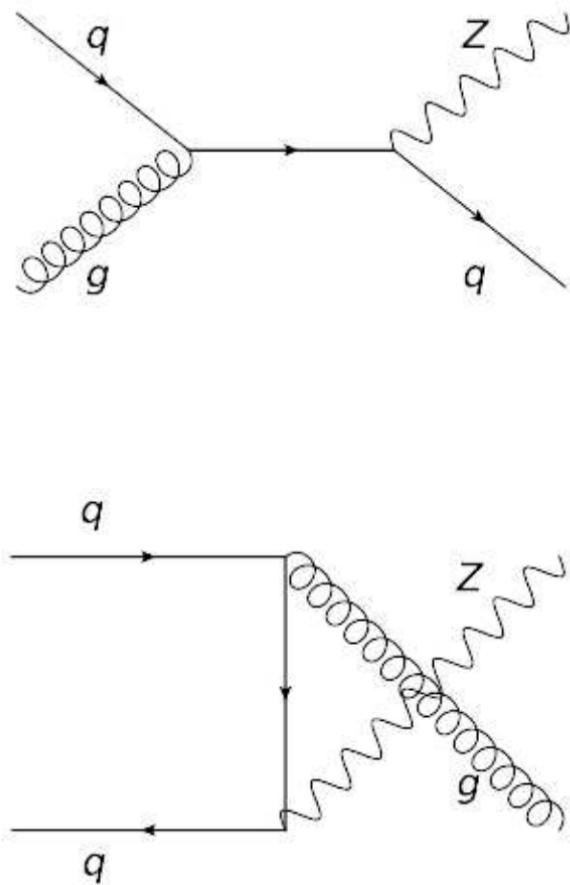


Figure 3.10: The LO Feynman diagrams for (Top) The Compton Process (Bottom) Quark - Anti-Quark Annihilation, the top being far more likely to occur at LHC.

$$\begin{aligned} \Sigma_{\text{SD}}^{(\text{LO})}(\rho) &= \frac{1}{\sigma_0} \int_0^\rho d\rho' \frac{d\sigma}{d\rho'} = 1 - \frac{1}{\sigma_0} \int_\rho^1 d\rho' \frac{d\sigma}{d\rho'} \\ &\quad \left(1 - \frac{\alpha_s C_F}{\pi} \left[\frac{1}{2} \log^2 \left(\frac{1}{\rho} \right) - \frac{3}{4} \log \left(\frac{1}{\rho} \right) \right] \right), \quad \text{if } \rho > z_{\text{cut}} \\ &1 - \frac{\alpha_s C_F}{\pi} \left[\frac{1}{2} \log^2 \left(\frac{1}{\rho} \right) - \frac{1}{2+\beta} \log^2 \left(\frac{z_{\text{cut}}}{\rho} \right) - \frac{3}{4} \log \left(\frac{1}{\rho} \right) \right], \quad \text{if } \rho < z_{\text{cut}} \end{aligned} \quad (3.13)$$

at order α_s [25].

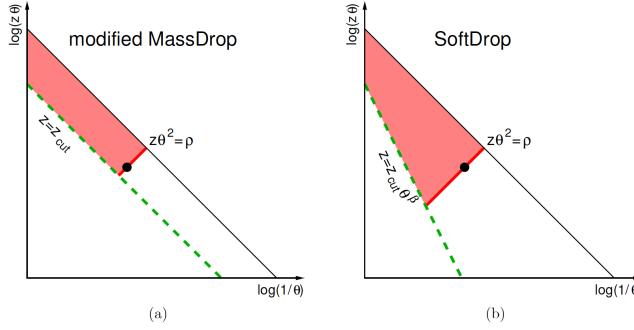


Figure 3.11: Lund diagrams for : (a) The modified mass drop tagger, (b) The Soft-Drop procedure . The dotted green line is the edge of the dropped region in both cases.

Since many of the complications of the calculations, such as non-global logarithms, are eliminated by dropping the soft and collinear contributions theorists are able to analytically calculate the soft-drop groomed jet mass and perform an all-order resummation of the logarithms.

$$\Sigma_{SD}(\rho) = \exp [-R_{SD}(\rho)] \quad (3.14)$$

$$R_{SD}(\rho) = \int_0^1 \frac{d\theta^2}{\theta^2} dz P_i(z) \frac{\alpha_s(z\theta p_t R)}{2\pi} \Theta(z\theta^2 > \rho) \Theta(z > z_{cut}\theta^\beta) \quad (3.15)$$

Including the folling term accounts for a running coupling correction [25]

$$\alpha_s(z\theta p_t R) = \frac{\alpha_s(p_t R)}{1 + 2\alpha_s\beta_0 \log(z\theta)} \quad (3.16)$$

Performing the integration retaining only hard-collinear branchings, Equation 3.4.4, and only the leading double-logarithm contributions from soft and collinear emissions [25], Equation 3.4.4.

$$R_{\text{SD}}^{(\text{LL})}(\rho) = \frac{C_i}{2\pi\alpha_s\beta_0^2} \left[\frac{2+\beta}{1+\beta} W \left(1 - \frac{\lambda_c + (1+\beta)\lambda_\rho}{2+\beta} \right) - \frac{W(1-\lambda_c)}{1+\beta} - 2W \left(1 - \frac{\lambda_\rho}{2} \right) \right. \quad (3.17)$$

$$\left. - 2\alpha_s\beta_0 B_i \log \left(1 - \frac{\lambda_\rho}{2} \right) \right] \quad (3.18)$$

where

$$\lambda_\rho = 2\alpha_s\beta_0 \log(1/\rho), \quad \lambda_c = 2\alpha_s\beta_0 \log(1/z_{\text{cut}}), \quad \text{and} \quad W(x) = x \log(x) \quad (3.19)$$

Lastly, the expression can be expanded up to NNLL corrections [25]:

$$R_{\text{SD}}^{(\text{LL})}(\rho) = \frac{C_i}{2\pi\alpha_s\beta_0^2} \left[\frac{2+\beta}{1+\beta} W \left(1 - \frac{\lambda_c + (1+\beta)\lambda_\rho}{2+\beta} \right) - \frac{W(1-\lambda_c)}{1+\beta} - 2W \left(1 - \frac{\lambda_\rho + \lambda_B}{2} \right) \right. \quad (3.20)$$

$$\left. + W(1 - \lambda_B) \right] \quad (3.21)$$

Define

$$\lambda_B = -2\alpha_s\beta_0 B_i \quad (3.22)$$

3.5 Factorization in Soft-Collinear Effective Theory

Factorization allows one to separate QCD calculations into parts, which can be calculated individually, using different methods and then summed for a complete understanding of the event at in an energy regimes.

Below the factorization structure of the double differential jet production cross section is displayed in the context of Soft-Collinear Effective Theory, SCET, following the framework for inclusive jet production $pp \rightarrow jet + X$ developed for jets of radius $R \ll 1$ the triple differential cross section can be defined [33] as :

$$\frac{d\sigma}{d\eta dp_T d\tau_a} = \sum_{abc} f_a(x_a, \mu) \otimes f_b(x_b, \mu) \otimes H_{ab}^c(x_a, x_b, \eta, p_T/z, \mu) \otimes G_c(z, p_T R, \tau_a, \mu, z_{cut}, \beta) \quad (3.23)$$

In the above cross section equation:

$f_{a,b}$ represent the parton distribution functions that give the probability of a given parton, of momentum fraction $x_{a,b}$, from a proton [33].

H_{ab}^c is the hard function describing the $ab \rightarrow c$ process, where c is the final state parton, possessing momentum $\frac{p_T}{z}$, that fragments to become the observed jet.

G_c being the semi-inclusive jet function, which can be refactored as follows

:

$$G_c(z, p_T R, \tau_a, \mu, z_{cut}, \beta) = \sum_i \mathcal{H}_{c \rightarrow i}(z, p_T R, \mu) S_i^{tg(z_{cut} p_T R, \beta, \mu)} \quad (3.24)$$

$$\times \int d\tau_a^{C_i} d\tau_a^{S_i} \delta(\tau_a - \tau_a^{C_i} - \tau_a^{S_i}) C_i(\tau_a^{C_i}, p_T, \mu) S_i^{\text{gr}}(\tau_a^{S_i}, p_T, R, \mu, z_{\text{cut}}, \beta) \quad (3.25)$$

In the above factorization the term S_i^{tgr} represents soft emissions that will always fail the soft drop criterion while the latter soft function accounts for soft emissions which may or may not pass depending on the values of the tunable parameters [33].

3.6 Jets Initiated by Quarks and Gluons

In order to emphasize the relevance to the measurement presented herein, the phenomenology of Quantum Chromodynamics is discussed rather than the Lagrangian perspective. This is useful as jet studies help probe QCD in the soft and collinear limits. Jets are formed by the hadronization of quarks and gluons. In this thesis I present a measurement of a light quark enriched jet sample.

Consider the simplest process that could produce a quark initiated jet, a quark of energy E_q emitting a gluon of energy E_g . The probability that this will occur is a function of the gluon's energy fraction, z , and the emission angle $, \theta$ [34].

$$z = \frac{E_g}{E_q + E_g}$$

$$1 - \cos\theta = \frac{m^2}{2E_q E_g}$$

Then the probability of gluon emission from the quark is :

$$P_q(z, \cos\theta) dz d\cos\theta = \frac{\alpha_s C_F}{\pi} \frac{dz}{z} \frac{d\cos\theta}{1-\cos\theta}$$

It is useful to assume the small angle approximation, $\theta \ll 1$, giving:

$$P_q(z, \theta^2) dz d\theta^2 = \frac{\alpha_s C_F}{\pi} \frac{dz}{z} \frac{d\theta^2}{\theta^2}$$

Notice that the probability of emission diverges for very soft (small z) or very collinear (small θ) gluons. In the soft and collinear limits the probability can be interpreted as an expectation value for the number of soft/collinear gluons [34].

It is elucidating to rewrite the probability in terms of inverse logarithms and introduce the "Lund Diagram" in order to visualize the uniform distribution of soft and collinear gluons in the $\log \frac{1}{\theta^2}, \log \frac{1}{z}$ space. Recall that the Lund plane is described in 3.3.

$$P_q(z, \theta^2) dz d\theta^2 = \frac{\alpha_s C_F}{\pi} d(\log \frac{1}{z}) d(\log \frac{1}{\theta^2})$$

At leading order (LO) jets can also be "initiated" by gluons and this probability is incredibly similar :

$$P_g(z, \theta^2) dz d\theta^2 = \frac{\alpha_s C_A}{\pi} d(\log \frac{1}{z}) d(\log \frac{1}{\theta^2})$$

This similarity allows us to interpret the variations in quark enriched and gluon enriched jet samples in terms of the fundamental C_F and adjoint C_A casimirs, in $SU(3)$, $C_F = \frac{4}{3}$ and $C_A = 3$. This is also the only difference between quark and gluon jet masses at LO as shown in 3.4.3.

Comparing the probability of a quark to emit a gluon and that of a gluon to emit a gluon, we can see the ratio is simply $\frac{C_A}{C_F} = \frac{9}{4}$. This has strong experimental implications since it implies gluon jets will on average be composed of

about twice as many constituent particles as quark jets, and will be broader. This will cause gluon jets to have higher mass, since the mass of a $1 \rightarrow 2$ splitting is given by $\theta \simeq \frac{2}{\gamma}$, this is derived in A..

CMS Experiment at LHC

4.1 The Large Hadron Collider

The LHC is the largest machine created by mankind to date and currently the world's highest energy particle accelerator, it accelerates and collides bunches of 10^{11} protons at a time which collide at combined center-of-mass energy of 13 TeV.

4.2 The CMS Detector

The Compact Muon Solenoid (CMS) detector was used to collect the data presented in this thesis, it is one of two large general purpose detectors at the LHC. CMS experiment has recorded 162 fb^{-1} integrated luminosity in the dataset presented in this thesis, collected during Run 2 of LHC.

CMS is one of 4 detectors that measure collisions of protons and lead ions produced by the Large Hadron Collider, LHC, at CERN. CMS is the smaller, overall length of 22m, a diameter of 15m, and weighs 14 000 tonnes, of the 2 large general-purpose detectors, the other being ATLAS. The most notable fea-

ture of the detector is its powerful 3.8 Tesla solenoid magnet, the largest superconducting magnet ever built, as of the year 2011.

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Forward calorimeters, made of steel and quartz-fibres, extend the pseudorapidity coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid.

4.2.1 Calorimeter Energy Resolution

In the barrel section of the ECAL, an energy resolution of about 1% is achieved for unconverted or late-converting photons that have energies in the range of tens of GeV. The remaining barrel photons have a resolution of about 1.3% up to a pseudorapidity of $|\eta| = 1$, rising to about 2.5% at $|\eta| = 1.4$. In the endcaps, the resolution of unconverted or late-converting photons is about 2.5%, while the remaining endcap photons have a resolution between 3 and 4% [35]. The lead tungstate crystals are $25.8X_0$ thick in the barrel and $24.7X_0$ thick in the endcaps. When combining information from the entire detector, the jet energy resolution amounts typically to 15% at 10 GeV, 8% at 100 GeV, and 4% at 1 TeV, to be compared to about 40%, 12%, and 5% obtained when the ECAL and HCAL calorimeters alone are used.

4.3 From Calorimeter Energy Deposits to Jets

In the region $|\eta| < 1.74$, the HCAL cells have widths of 0.087 in pseudorapidity and 0.087 in azimuth (ϕ). In the η - ϕ plane, and for $|\eta| < 1.48$, the HCAL cells map on to 5×5 arrays of ECAL crystals to form calorimeter towers projecting radially outwards from close to the nominal interaction point. For $|\eta| \gtrsim 1.74$, the coverage of the towers increases progressively to a maximum of 0.174 in $\Delta\eta$ and $\Delta\phi$. Within each tower, the energy deposits in ECAL and HCAL cells are summed to define the calorimeter tower energies, subsequently used to provide the energies and directions of hadronic jets.

Jets are reconstructed offline from the energy deposits in the calorimeter towers, clustered using the anti- k_T algorithm [28, 36] with a distance parameter of 0.4. In this process, the contribution from each calorimeter tower is assigned a momentum, the absolute value and the direction of which are given by the energy measured in the tower, and the coordinates of the tower. The raw jet energy is obtained from the sum of the tower energies, and the raw jet momentum by the vectorial sum of the tower momenta, which results in a nonzero jet mass. The raw jet energies are then corrected to establish a relative uniform response of the calorimeter in η and a calibrated absolute response in transverse momentum p_T .

4.4 Muon Reconstruction

Muons are measured in the pseudorapidity range $|\eta| < 2.4$, with detection planes made using three technologies: drift tubes, cathode strip chambers, and resistive plate chambers. The single muon trigger efficiency exceeds 90% over

the full η range, and the efficiency to reconstruct and identify muons is greater than 96%. Matching muons to tracks measured in the silicon tracker results in a relative transverse momentum resolution, for muons with p_T up to 100 GeV, of 1% in the barrel and 3% in the endcaps. The p_T resolution in the barrel is better than 7% for muons with p_T up to 1 TeV [37].

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [38]. The global event reconstruction (also called particle-flow event reconstruction [39]) is described in 5.2.

Event Generation and Reconstruction

The events used for this measurement were reconstructed from data acquired by the CMS detector in the case of the data and generated using monte carlo generators PYTHIA and HERWIG in the case of the generated data.

5.1 Brief Introduction to Monte Carlo Event Generators

Monte Carlo (MC) event generators are tools used by both experimental and theoretical physicists to simulate different physical processes in order to make predictions and prepare future experiments. The main tasks of such generators are to calculate matrix elements of the relevant hard processes but they must also describe parton showering, hadronization and underlying event. MC can be utilized to extrapolate data measurements beyond the acceptance of the detector or in the case of this thesis it is used in the unfolding process to correct

the data for detector efficiency and resolution.

MC generators provide an ensemble of generated events that simulate the physics process in question. Each individual generator implements a slightly different scheme in order to approximate the necessary calculations for the factorization and renormalization scales relevant to a process. These variations mean that the choice of MC generator will have a slight effect on the generated distributions. For this reason, in the analysis described herein we compared results from 2 different generators: PYTHIA and HERWIG.

All event generators break the calculations into separate parts as depicted by the different colored objects in Figure 5.1

PYTHIA [40] is a very commonly used general purpose event generator which uses the parton shower approach for higher order corrections to the hard scattering matrix element.

HERWIG [41] is another commonly used event generator, incredibly similar to PYTHIA, differing mainly in hadronization and parton showering behaviors.

5.2 Event Reconstruction with Particle Flow

The data from proton-proton collisions at LHC are detected by the CMS experiment and then the Particle Flow (PF) event reconstruction [6] is applied to these raw detector outputs in order to construct "particle flow objects" that contain information from multiple CMS detector subsystems and constitute a global event description. More details about the CMS apparatus can be found in 4. The particle flow objects are given defined object categories based on which subsystems they are measured in. This requires the use of a "link algorithm" which combines subdetector information together into a single object [6]. For example,

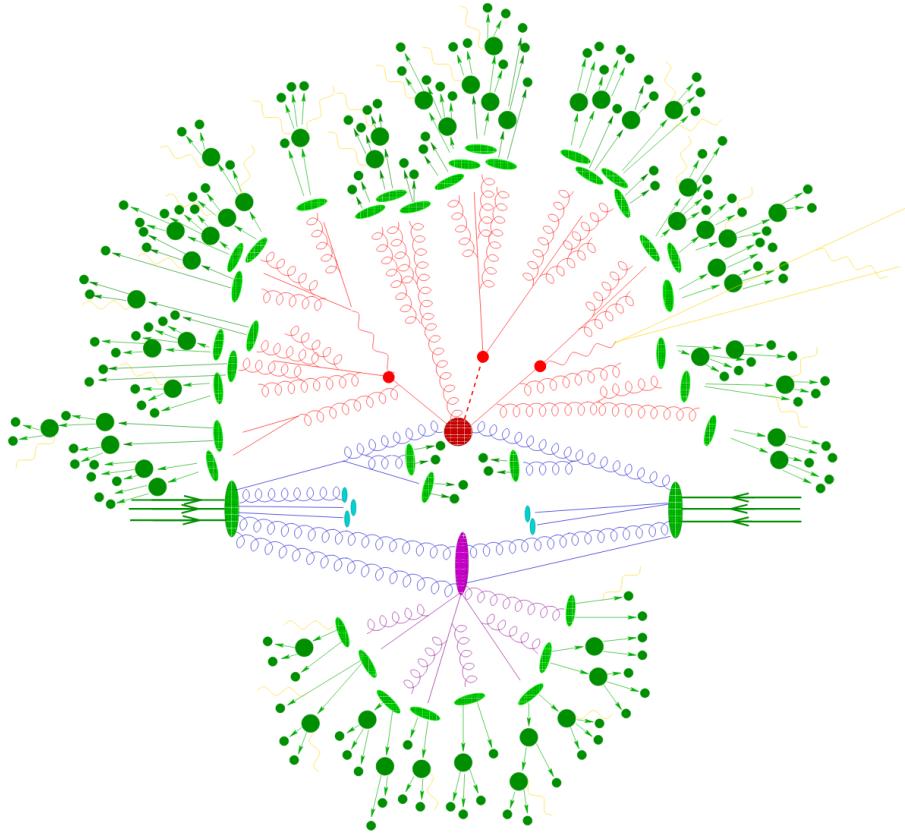


Figure 5.1: A pictoral representation of the way the MC generators simulate hadron-hadron collisions [?].

one can see that the solid, light blue, line in figure 5.2 corresponds to a muon PF object as it left hits in the tracker then traversed all of the calorimeters only to deposit its energy in the muon chambers.

PF aims to reconstruct and identify each individual particle in an event, with an optimized combination of all subdetector information. In doing so PF characterizes the particles into 5 types: photon, electron, muon, charged hadron, neutral hadron. The type determines which sub-detector information will be combined to determine the energy and direction of that particle. Photons (e.g. coming from electron bremsstrahlung) are identified as ECAL energy clusters

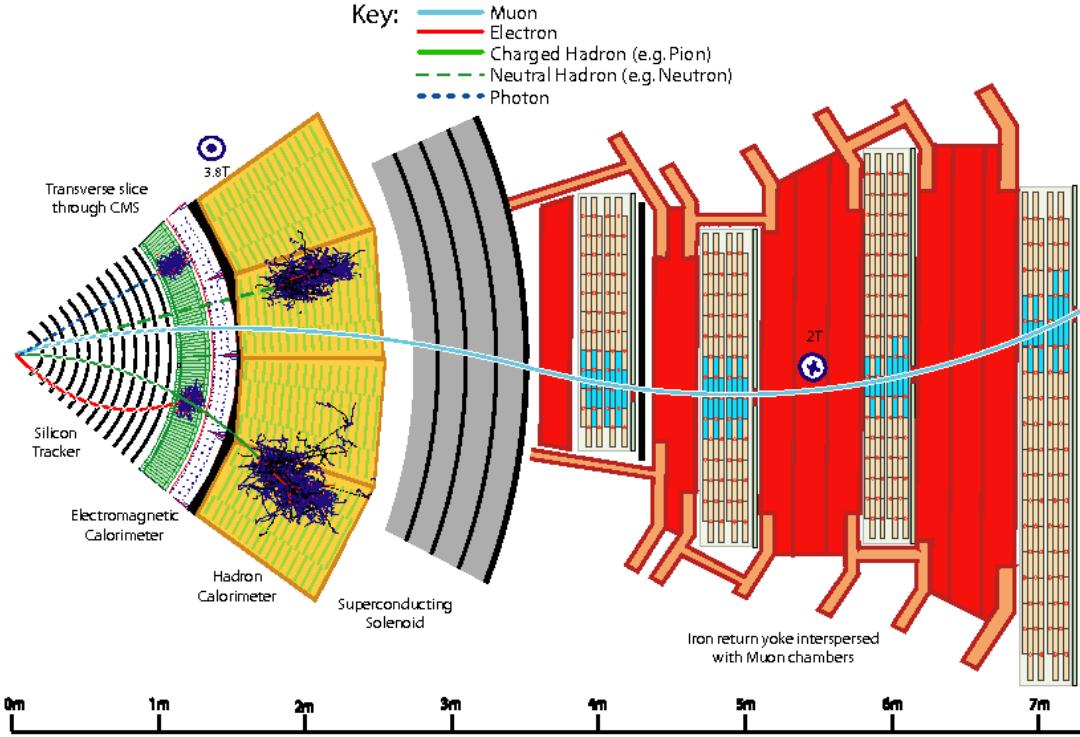


Figure 5.2: A pictoral representation of the way the Particle Flow algorithm determines which objects correspond to which particles based on an optimal combination of sub-detector information [6].

that are not linked to the extrapolation of any charged particle trajectory to the ECAL. Electrons (e.g. coming from photon conversions in the tracker material or from semileptonic decays of hadrons) are identified as a primary charged particle track and potentially many ECAL energy clusters corresponding to this track extrapolation to the ECAL and to any bremsstrahlung photons emitted along the way within the tracker. Muons (e.g. from hadron semileptonic decays) are identified as tracks in the central tracker consistent with either a track or several hits in the muon system, and associated with calorimeter deposits compatible with the muon hypothesis. Charged hadrons are identified as charged particle tracks neither identified as electrons, nor as muons. Finally, neutral hadrons are

identified as HCAL energy clusters not linked to any charged hadron trajectory, or as a combined ECAL and HCAL energy excess with respect to the expected charged hadron energy deposit.

The ECAL is used to obtain the energy of photons. The energy of electrons is more complex, determined from a combination of the track momentum at the main interaction vertex, the corresponding ECAL cluster energy, and the energy sum of all bremsstrahlung photons originating from the track. The energy of muons is obtained from the corresponding track momentum. The energy of charged hadrons is determined from a combination of the track momentum and the corresponding ECAL and HCAL energies, corrected for zero-suppression effects and for the response function of the calorimeters to hadronic showers. Finally, the energy of neutral hadrons is obtained from the corresponding corrected ECAL and HCAL energies.

After particle flow is used to determine the particle’s type, the particle flow objects which are not categorized as muons, electrons and isolated photons are then described as hadrons and are clustered into “jets”. An example PF jet with 5 constituent PF objects is depicted in 5.3.

For each event, hadronic jets are clustered from these reconstructed particles using the infrared and collinear safe anti- k_T algorithm [28, 36] with a distance parameter of 0.4. Jet momentum is determined as the vectorial sum of all particle momenta in the jet, and is found from simulation to be, on average, within 5 to 10% of the true momentum over the whole p_T spectrum and detector acceptance. Additional proton-proton interactions within the same or nearby bunch crossings (pileup) can contribute additional tracks and calorimetric energy depositions to the jet momentum.

The pileup per particle identification (PUPPI) algorithm [42] is used to mit-

igate the effect of pileup at the reconstructed particle level, making use of local shape information, event pileup properties and tracking information. Charged particles identified to be originating from pileup vertices are discarded. For each neutral particle, a local shape variable is computed using the surrounding charged particles compatible with the primary vertex within the tracker acceptance ($|\eta| < 2.5$), and using both charged and neutral particles in the region outside of the tracker coverage. The momenta of the neutral particles are then rescaled according to their probability to originate from the primary interaction vertex deduced from the local shape variable, superseding the need for jet-based pileup corrections [43].

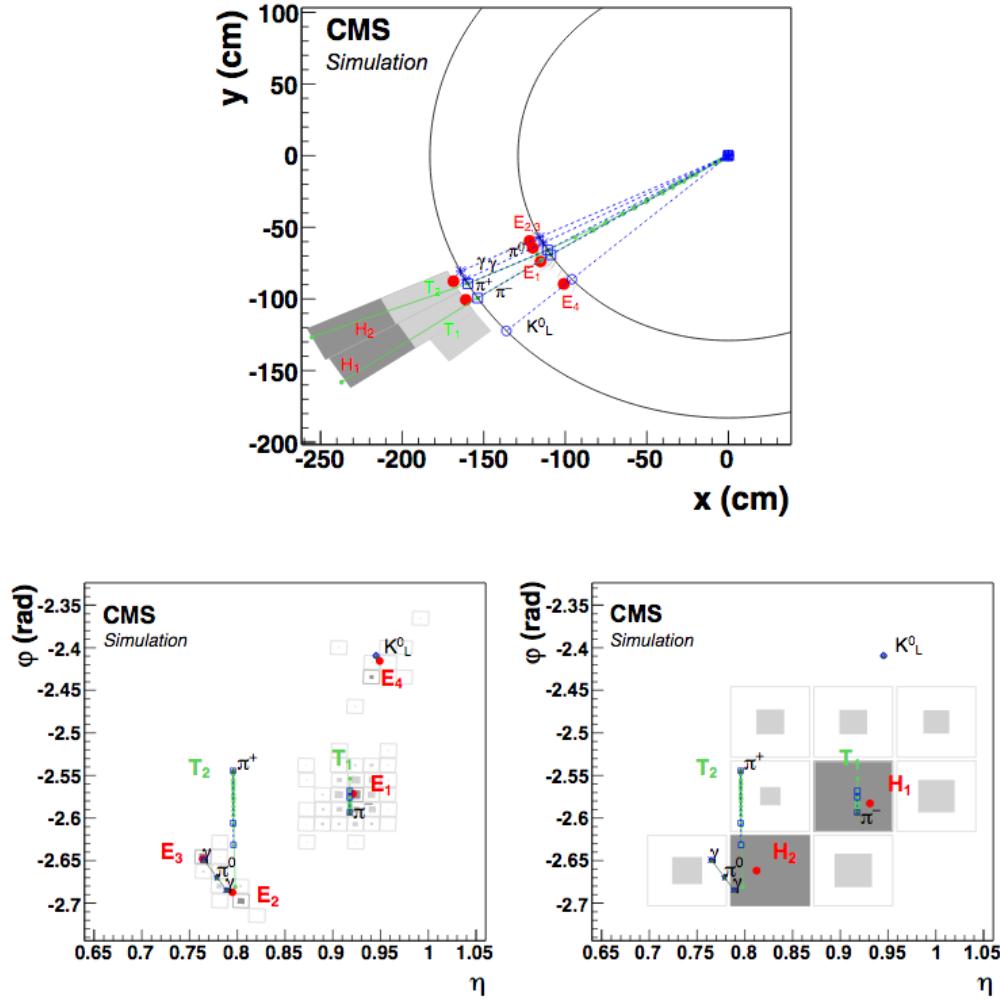


Figure 5.3: A pictorial representation of the way the Particle Flow objects are clustered into jets [6]. The top image shows a jet composed of 5 PF candidates in the x,y plane and the concentric circles describe the surfaces of the ECAL and HCAL detectors respectively. The lower left image is the PF candidates as measured by the ECAL and the image on the lower right shows the same for the HCAL surface. The dotted lines represent generated particles and the solid boxes represent energy deposits in the detector.

Measurement of the differential jet production cross section with respect to jet mass and transverse momentum in Z + Jet events from pp collisions at $\sqrt{s} = 13$ TeV

6.1 Introduction

The analysis strategy is similar to that of Ref. [18]. However, there are several differences. As in that paper, the cross section is now unfolded in both jet mass and p_T . However, while the previous measurement considered only one value for the soft drop parameter β , this analysis considers several. We apply the soft drop algorithm to compare directly to theoretical computations. Additionally, we not only measure the cross section as a function of mass, but also as

a function of dimensionless mass, $\rho = 2\log(m/(ptR))$, as is also done in the previously mentioned ATLAS measurement. The dimensionless mass ρ only weakly depends on p_T , unlike mass, which is highly correlated. Additionally, the use of this variable aids in the separation of fixed order, perturbative and non-perturbative effects. We present the normalized double differential jet production cross section with respect to jet mass and transverse momentum (p_T) as well as with respect to jet dimensionless mass and p_T . We compute the cross sections normalized per p_T bin (the “normalized” cross section) with respect to the jet p_T and jet mass by unfolding a binned two-dimensional distribution in p_T and mass with widths Δpt and Δm , respectively.

The normalized differential cross section

$$\frac{1}{d\sigma/dpt} \frac{d^2\sigma}{dpt dm} = \frac{1}{N/\Delta pt} R\left(\frac{N_{ij}}{\Delta pt \Delta m}\right) \quad (6.1)$$

where N is the total number of $Z+jets$ events in our selection, N_{ij} is the number of such events in pt bin i and mass bin j , and $R(\alpha)$ is the unfolding procedure applied to the two-dimensional distribution α .

6.2 Data and MC

The data used are from proton-proton collisions in 2016, 2017, and 2018 with 25 ns bunch spacing. The datasets are summarized in Table ??.

The approved luminosity sections used for 2016, 2017, 2018 Runs A – C, and 2018 Run D, respectively, are listed below:

- Cert_271036-284044_13TeV_23Sep2016ReReco_Collisions16_JSON.txt
- Cert_294927-306462_13TeV_EOY2017ReReco_Collisions17_JSON.txt

Data sample	Cross section (pb^{-1})
2016	
/SingleMuon/Run2016*-17Jul2018-v1/MINIAOD	35.9
/SingleElectron/Run2016*-17Jul2018-v1/MINIAOD	35.9
2017	
/SingleMuon/Run2017*-31Mar2018-v1/MINIAOD	41.5
/SingleElectron/Run2017*-31Mar2018-v1/MINIAOD	41.5
2018 Runs A – C	
/SingleMuon/Run2018*-17Sep2018-v1/MINIAOD	27.9
/EGamma/Run2018*-17Sep2018-v1/MINIAOD	27.9
2018 Run D	
/SingleMuon/Run2018D-PromptReco-v2/MINIAOD	31.3
/EGamma/Run2018D-PromptReco-v2/MINIAOD	31.1

Table 6.1: Data samples used in the analysis.

- Cert_314472-325175_13TeV_PromptReco_Collisions18_JSON.txt
- Cert_314472-325175_13TeV_PromptReco_Collisions18_JSON.txt

In this analysis, Drell-Yan Monte Carlo (MC) events were simulated with MADGRAPH [44] and showered with two different generators.

The first generator is PYTHIA 8 [45], for which inclusive as well as (\hat{p}_T) binned samples were used.

An ensemble of samples, each varying the hard-scatter transverse momentum (\hat{p}_T) from 70 GeV to Infinity, was used to fill the response matrices for the measurement. The samples are :

- 2016: /DYJetsToLL_M-50_HT-*_TuneCUETP8M1_13TeV-madgraphMLM-pythia8/
- 2017 and 2018: /DYJetsToLL_M-50_HT-*_TuneCP5_13TeV-madgraphMLM-pythia8/

The underlying event tune is CUETP8M1 [46] for 2016 and CP5 [47] for 2017 and 2018. These samples were generated in the

- RunIISummer16MiniAODv3-PUMoriond17_94X_mcRun2_asymptotic_v3
- RunIIFall17MiniAODv2-PU2017_12Apr2018_94X_mc2017_realistic_v14
- RunIIAutumn18MiniAOD-102X_upgrade2018_realistic_v15

MC campaigns for 2016, 2017, and 2018, respectively. A second inclusive PYTHIA8 sample was used to evaluate the physics model uncertainty, for which the underlying event tune is CUETP8M1. It generated with CMSSW 9.4.X in the RunIISummer16MiniAODv3-PUMoriond17_94X_mcRun2_asymptotic_v3 MC campaign.

The second generator is HERWIG++ [48] with tune CUETP8M1, which is also used to generate a single sample which was also used to evaluate the physics model uncertainty. This sample was generated with CMSSW 9.4.X in the RunIISummer16MiniAODv3-PUMoriond17_94X_mcRun2_asymptotic_v3 MC campaign [This was used because there was only a 2016 version available].

Table 6.2 summarizes the Drell-Yan samples, the number of events, and the cross sections.

The backgrounds considered in this analysis comprise $t\bar{t} + jets$ and diboson (WW , WZ , and ZZ) samples. The $t\bar{t} + jets$ samples were simulated with MADGRAPH and showered with PYTHIA 8, while the diboson samples with simulated and showered with PYTHIA 8. The tunes and MC campaigns used were the same as for the Drell-Yan samples. Table 6.3 summarizes the background samples, the number of events, and the cross sections.

The reconstruction used for the data and MC was CMSSW 9.4.X (CMSSW 10.0.X) for 2016 and 2017 (2018), with private extended NANOAOD based on central NANOAODv4. This corresponds to the following MINIAOD versions:

- 2016: v3

MC sample	Number of Events	Cross section (pb), LO
2016		
HERWIG++	29883521	1091.1
PYTHIA 8 Tune CUETP8M1	96531428	14420
PYTHIA 8 Tune CUETP8M1 70-100	9691660	175.3
PYTHIA 8 Tune CUETP8M1 100-200	2751187	147.4
PYTHIA 8 Tune CUETP8M1 100-200 extension	8265899	147.4
PYTHIA 8 Tune CUETP8M1 200-400	962195	40.99
PYTHIA 8 Tune CUETP8M1 200-400 extension	8646942	40.99
PYTHIA 8 Tune CUETP8M1 400-600	1070454	5.678
PYTHIA 8 Tune CUETP8M1 400-600 extension	8655207	5.678
PYTHIA 8 Tune CUETP8M1 600-800	8292957	1.363
PYTHIA 8 Tune CUETP8M1 800-1200	2673066	0.6759
PYTHIA 8 Tune CUETP8M1 1200-2500	596079	0.116
PYTHIA 8 Tune CUETP8M1 2500-Inf	399492	0.002592
2017		
PYTHIA 8 Tune CP5 70-100	9344037	143
PYTHIA 8 Tune CP5 100-200	10235418	164.4
PYTHIA 8 Tune CP5 100-200 extension	3950339	164.4
PYTHIA 8 Tune CP5 200-400	10728447	49.8
PYTHIA 8 Tune CP5 200-400 extension	1200863	49.8
PYTHIA 8 Tune CP5 400-600	9533635	6.8
PYTHIA 8 Tune CP5 400-600 extension	1124294	6.8
PYTHIA 8 Tune CP5 600-800	8153358	1.8
PYTHIA 8 Tune CP5 800-1200	3089861	0.85
PYTHIA 8 Tune CP5 1200-2500	625517	0.018
PYTHIA 8 Tune CP5 2500-Inf	404986	0.0036
2018		
PYTHIA 8 Tune CP5 70-100	10019684	143
PYTHIA 8 Tune CP5 100-200	11530510	164.4
PYTHIA 8 Tune CP5 200-400	11210867	49.8
PYTHIA 8 Tune CP5 400-600	9697098	6.8
PYTHIA 8 Tune CP5 400-600 extension	9358053	6.8
PYTHIA 8 Tune CP5 600-800	8862104	1.8
PYTHIA 8 Tune CP5 800-1200	3138129	0.85
PYTHIA 8 Tune CP5 1200-2500	536416	0.018
PYTHIA 8 Tune CP5 2500-Inf	427051	0.0036

Table 6.2: List of Monte Carlo samples used. The number of generated events and the total cross section are also provided for each subsample.

MC sample	Number of Events	Cross section (pb)
2016		
TTJets PYTHIA 8 Tune CUETP8M1	10199051	831.76, NNLO
WW PYTHIA 8 Tune CUETP8M1	6988168	118.7, NNLO
WZ PYTHIA 8 Tune CUETP8M1	2997571	47.13, NLO
ZZ PYTHIA 8 Tune CUETP8M	998034	16.523, NLO
2017		
TTJets PYTHIA 8 Tune CUETP8M1	8026103	831.76, NNLO
WW PYTHIA 8 Tune CUETP8M1	7765828	118.7, NNLO
WZ PYTHIA 8 Tune CUETP8M1	3928630	47.13, NLO
ZZ PYTHIA 8 Tune CUETP8M	1925931	16.523, NLO
2018		
TTJets PYTHIA 8 Tune CUETP8M1	10244307	831.76, NNLO
WW PYTHIA 8 Tune CUETP8M1	7850000	118.7, NNLO
WZ PYTHIA 8 Tune CUETP8M1	3885000	47.13, NLO
ZZ PYTHIA 8 Tune CUETP8M	1979000	16.523, NLO

Table 6.3: List of Monte Carlo samples used. The number of generated events and the total cross section are also provided for each subsample.

- 2017: v2
- 2018: v1

6.3 Trigger

The data are collected with single-lepton triggers, with muons (electrons) of $p_t > 29(37)$ GeV using the Isolated muon above 27 GeV (Tight working point electron above 35 GeV GSF || Photon above 200 GeV) triggers, as recommended by the muon (egamma) POG. All triggers used for this analysis are prescaled to 1.

6.4 Reconstruction and Selection

The event reconstruction is based on the CMS Particle Flow (PF) algorithm [49, 50, 6], which takes into account information from all subdetectors, including charged particle tracks from the tracking system, energy deposits in the electromagnetic and hadronic calorimeters and tracks reconstructed in the muon chambers. Given this information, all particles in the event are reconstructed as electrons, muons, photons, charged hadrons or neutral hadrons. Charged hadrons associated with pileup vertices are removed from consideration, referred to as “charged hadron subtraction” (CHS).

Hadronic jets are clustered from particle flow inputs using the anti- k_T algorithm [28] using the FASTJET 3.0 software package [51] with $R = 0.8$ (AK8 jets). Corrections based on the jet area [52] are applied to the jets to remove the energy contribution of neutral hadrons arising from pileup collisions. Further corrections are used to account for the nonlinear calorimetric response, as a function of η and p_T [53], derived from simulation and data-to-simulation correction factors. L1FastJet, L2, L3, and L2L3Residual AK8PFchs corrections from version Fall115_25nsV2 were used. For the jet energy, all corrections are applied. For the jet mass, only L2L3+Residual corrections are applied, and the L1 correction is not applied. This is to ensure that theoretical calculations can be directly compared. It is also the recommendation from JMAR.

The constituents of the AK8 jets are reclustered using the Cambridge-Aachen algorithm [54, 55]. The jet grooming algorithm used is the modified mass drop tagger (MMDT) algorithm [5], also known as the “soft drop” algorithm with angular exponents $\beta = 0, 1, -1$, soft thresholds $z_{cut} < 0.05, 0.1, 0.15$, all with characteristic radius $R_0 = 0.8$ [4], giving a total of 9 different groomings. This

algorithm is also used to identify two subjets within the AK8 jet.

The same reconstruction criteria are also applied to stable particles at the generator level (`GenJets`), removing neutrinos.

All reconstructed AK8 jets are required to have $p_T > 200$ GeV and rapidity $|y| < 2.4$, as well as to satisfy “tight” jet identification requirements to remove detector noise: CHECK 2017 TIGHT JET requirements

- Neutral hadron fraction < 0.90 ,
- Neutral EM fraction < 0.90 ,
- Charged hadron fraction > 0.00 ,
- Charged EM fraction < 0.99 ,
- At least two constituents,
- At least one charged hadron.

To account for the worse jet energy resolution in the data than the MC, the MC are smeared in the nominal case by the amounts listed in Ref. [53]. The typical values are around 5-8%. These are applied to the entire jet four-vector, including p_T and mass. Currently the Fall17 V3 JER files are being used.

The jet mass can also have a different scale and resolution in data and MC. These are determined by fitting the W mass in semileptonic $t\bar{t}$ events. The ratio of the fitted means (data/MC) is the jet mass scale, JMS, and the ratio of the fitted widths (data/MC) is the jet mass resolution, JMR. Traditionally, this procedure is performed on events where the W decay products merge forming a single AK8 Jet. At high momentum, above 400 GeV, this method breaks down since the b quark decay products merge into the same AK8 as the W. Using

the highest mass soft-drop subjet as the W candidate, we are able to more accurately determine JMR and JMS for higher transverse momentum jets. These studies are ongoing, currently, we are using dummy values from the nanoAOD-tools modules JMR [0.1, 0.2, 0.0] and JMS [1. , 0.99, 1.01].

The jet mass resolution for ungroomed jets varies from XXX% depending on p_T and mass. The ungroomed jet mass resolution below XX GeV is very non-Gaussian, thus poorly estimated (see Figures ??-??). FIX THIS SENTENCE For this reason, the jet mass measurement is deemed to be invalid for ungroomed masses below 20 GeV for $p_T < 760$ GeV and below 40 GeV for $p_T > 760$ GeV. This is not very problematic because there are very few events for ungroomed masses in this regime.

The groomed jets, on the other hand, can access jet masses very close to zero. The jet mass scale in MC is between XXX 0.8-1.0 for the groomed jets for all jet p_T and masses. Small deviations from unity are accounted for in the response matrix. However, the jet mass resolution for very closely-separated jet constituents becomes extremely poor. This is because of the finite spatial resolution of the detector. Since the opening angle between the constituents scales like m/p_T , there will be minimum masses at a given p_T , below which the measurement is invalid. The criterion chosen here is to require that the uncertainty of the fits to the jet mass scale be lower than 50%. This is shown as the uncertainty bands in Figs ??-?? and then again as the central values in Figs. ??-??. With this criterion in mind, the measurement of the groomed masses are restricted to $m > 10$ GeV.

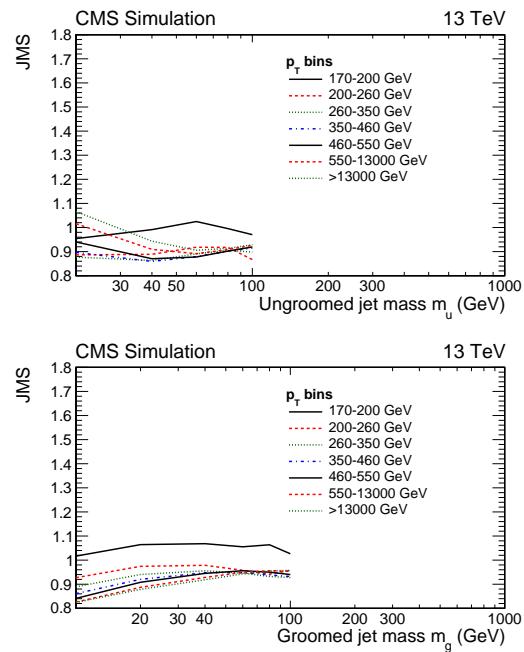


Figure 6.1: Jet mass scale in MC (mean of fits to m_{reco}/m_{gen} as a function of $p_{T_{gen}}$) for ungroomed (top) and groomed (bottom) jets in different p_T bins as a function of mass.

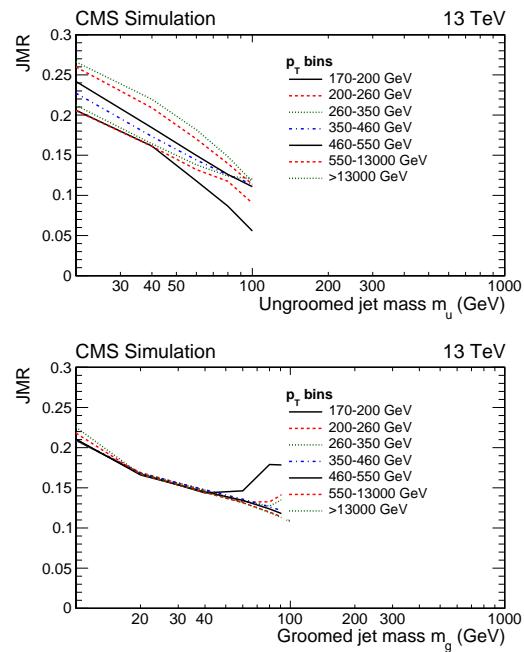


Figure 6.2: Jet mass resolution in MC (widths of fits to m_{reco}/m_{gen} as a function of $p_{T_{gen}}$) for ungroomed (top) and groomed (bottom) jets in different p_T bins as a function of mass.

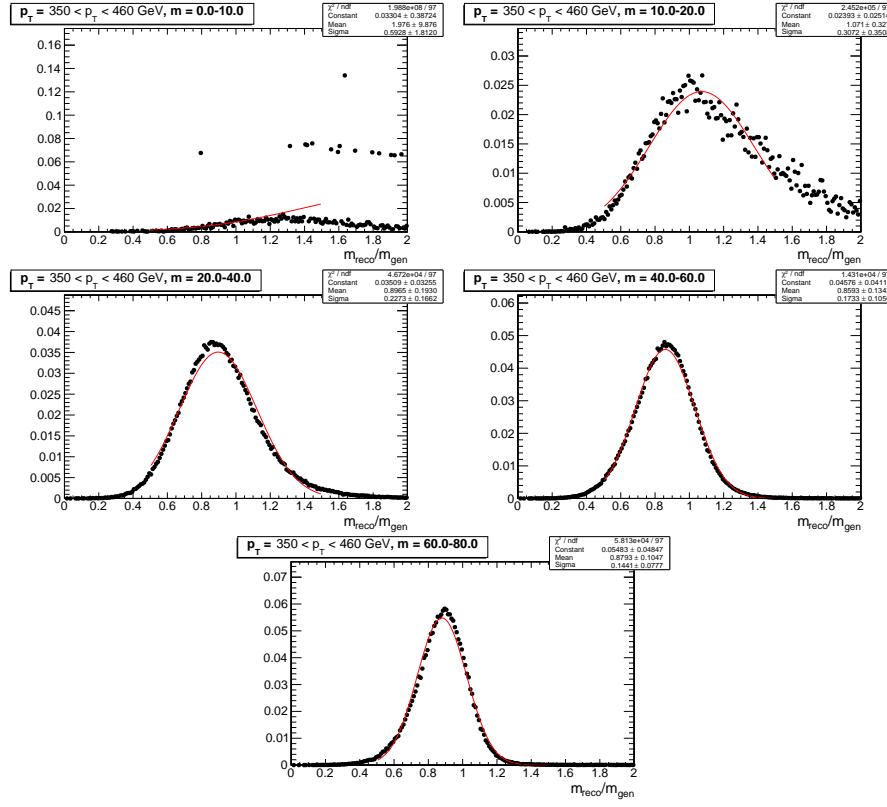


Figure 6.3: Fits to $m_{\text{reco}}/m_{\text{gen}}$ for different m_{gen} bins are shown for ungroomed jets of p_T 350-460 GeV . The full fits are in Appendix ?? for separate p_T, m_{gen} bins.

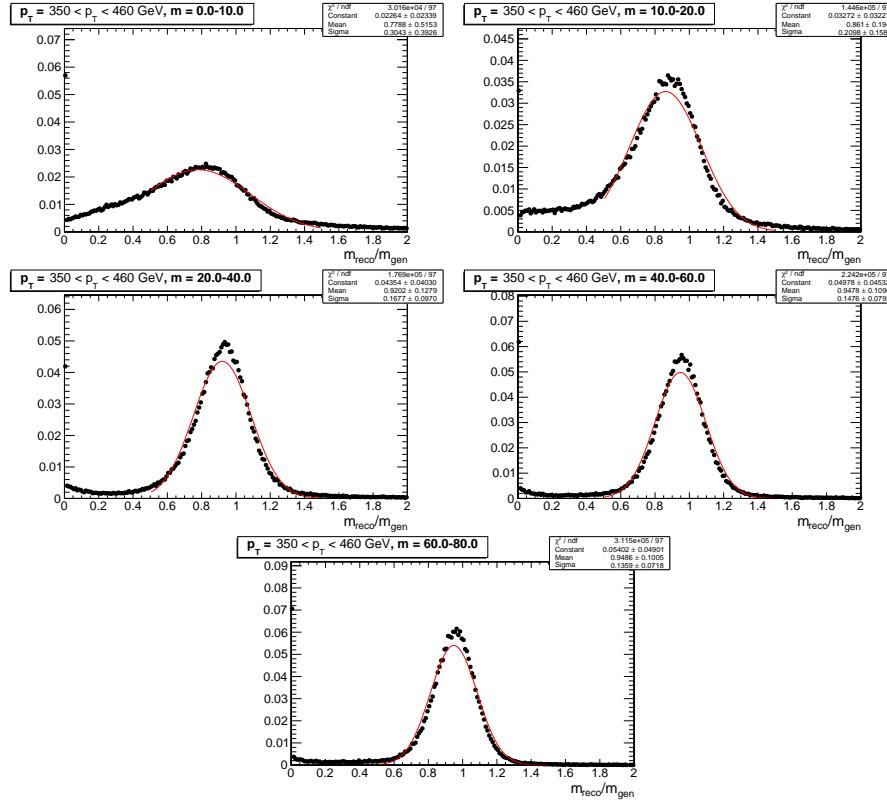


Figure 6.4: Fits to $m_{\text{reco}}/m_{\text{gen}}$ for different m_{gen} bins are shown for groomed jets of p_T 350-460 GeV . The full fits are in Appendix ?? for separate p_T, m_{gen} bins.

6.5 Data to MC Comparisons

The jet p_T of the leading jet, this is the only jet considered in the event, after luminosity scaling is shown for muon events which passed any of the 3 triggers in Fig. 6.5.

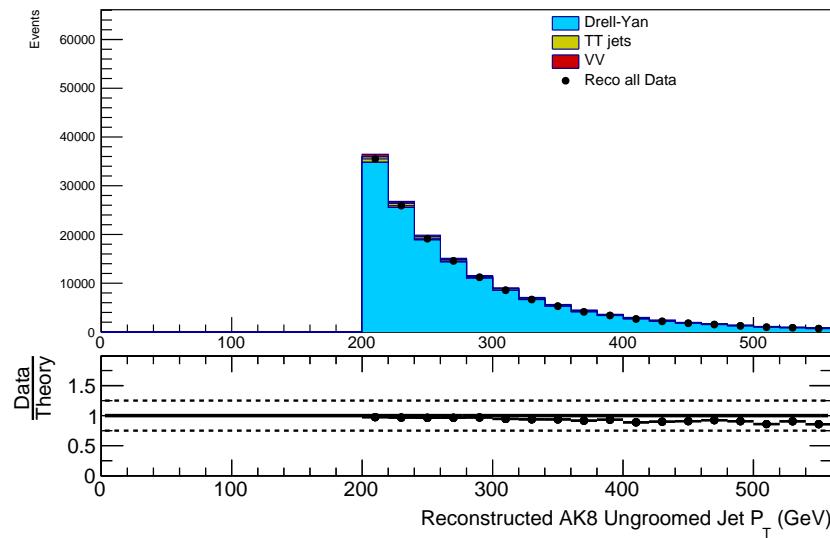


Figure 6.5: Jet p_T of leading jet of all data compared with PYTHIA 8 monte carlo simulated events, with luminosity scaling for 2017 data applied to the simulation.

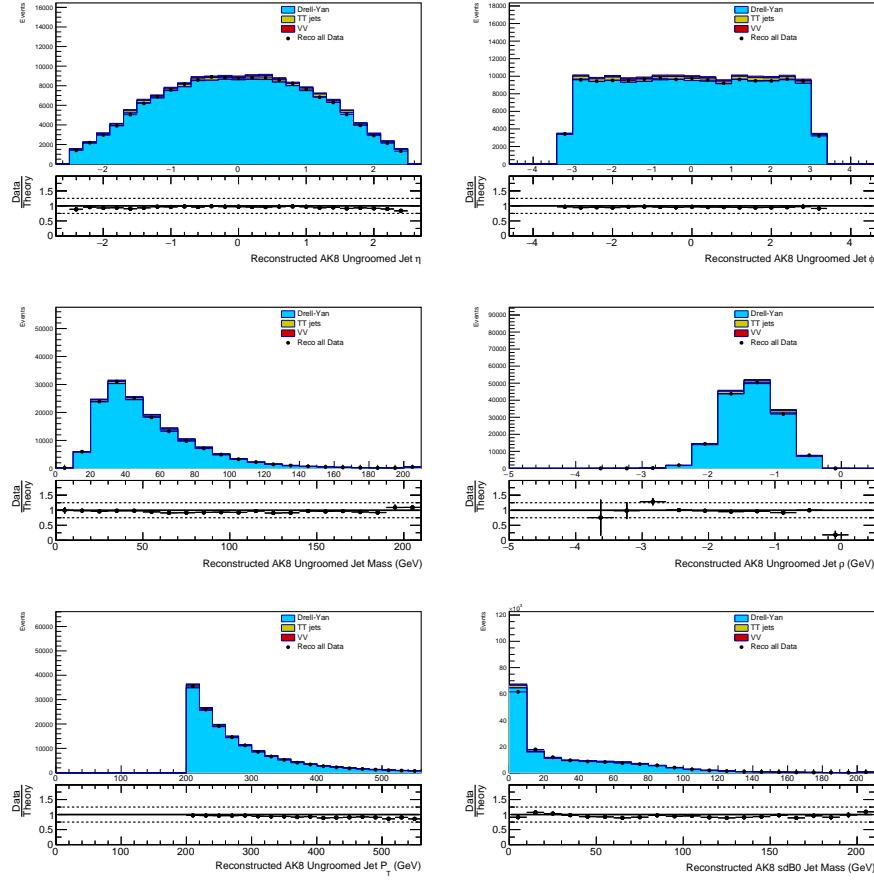


Figure 6.6: Jet pseudorapidity, azimuthal angle, mass, ρ [perhaps groomed ρ is better or we should show both side by side.] , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC, inclusive in jet p_T .

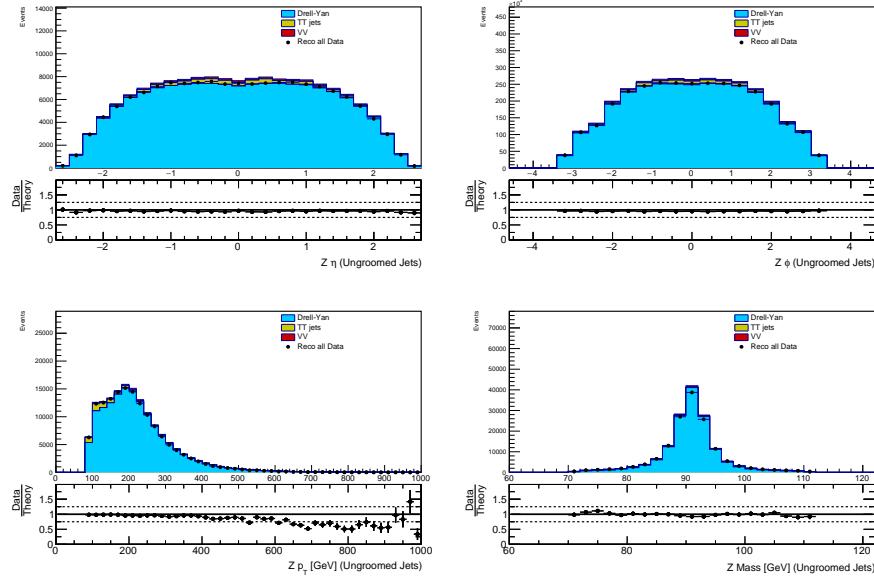


Figure 6.7: Z candidate pseudorapidity, azimuthal angle, mass, p_T , inclusive in jet p_T .

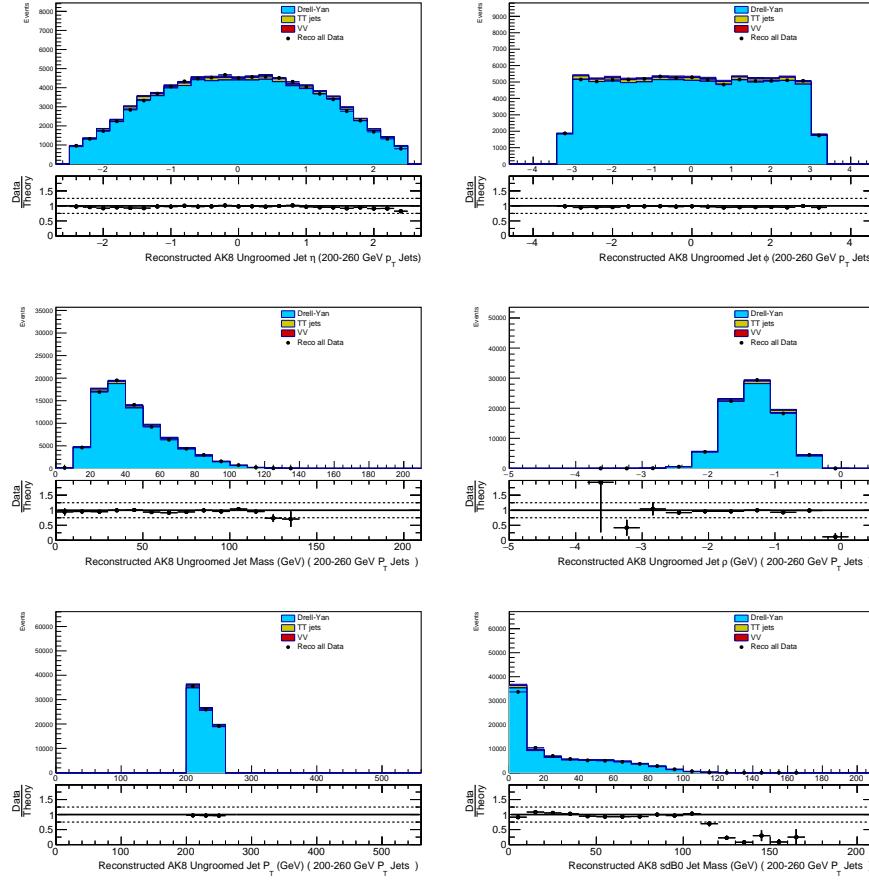


Figure 6.8: Jet pseudorapidity, azimuthal angle, mass, ρ , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC for $p_T = 200 - 260$ GeV.

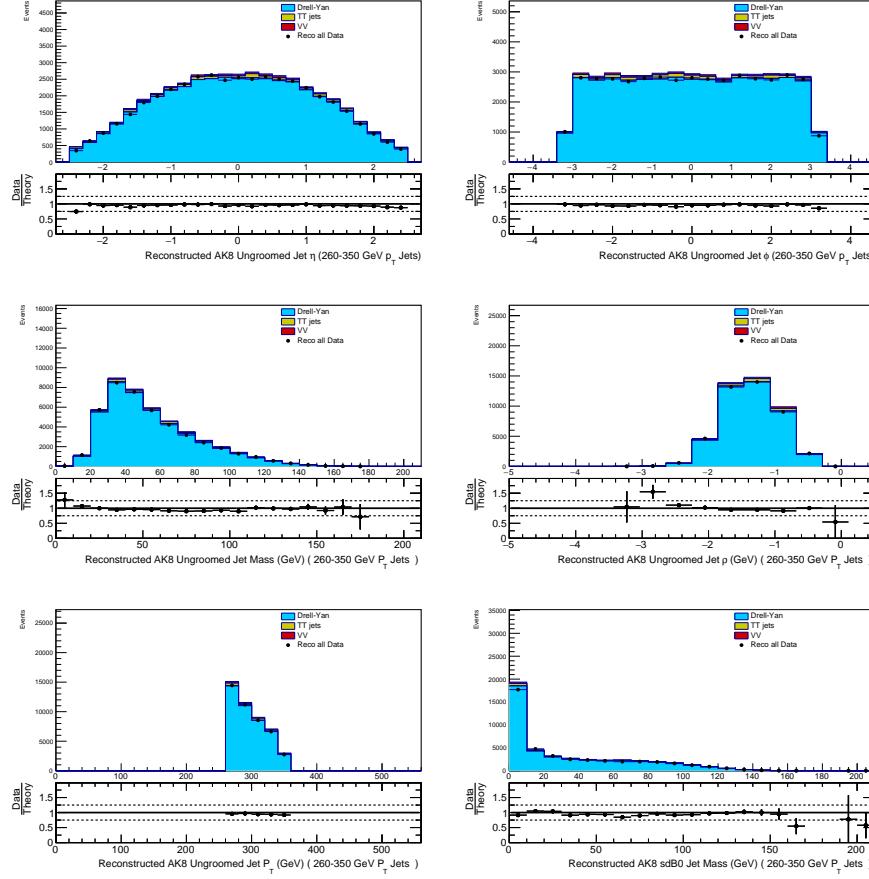


Figure 6.9: Jet pseudorapidity, azimuthal angle, mass, ρ , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC for $p_T = 260 - 350$ GeV.

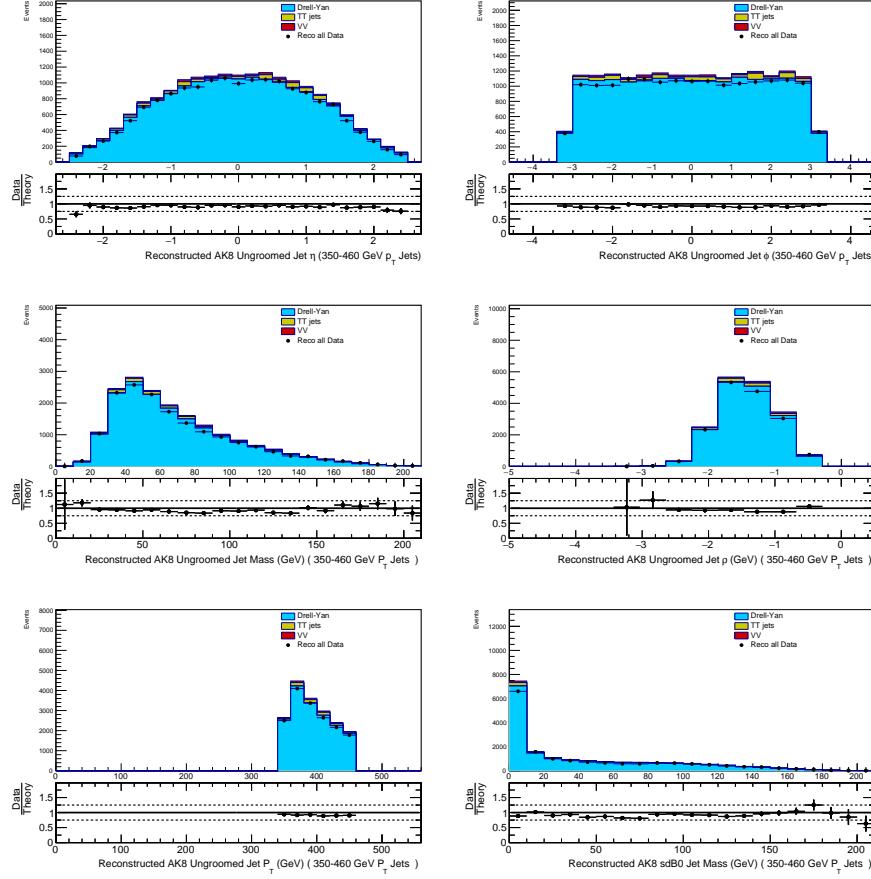


Figure 6.10: Jet pseudorapidity, azimuthal angle, mass, ρ , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC for $p_T = 350 - 460$ GeV.

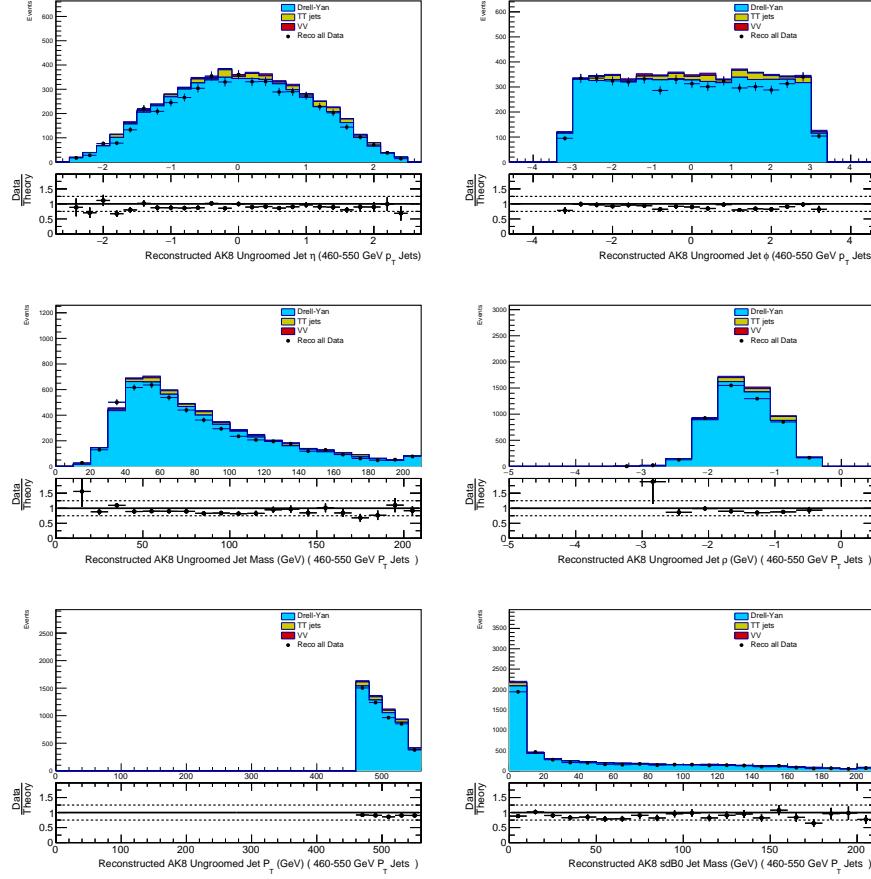


Figure 6.11: Jet pseudorapidity, azimuthal angle, mass, ρ , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC for $p_T = 460 - 550$ GeV.

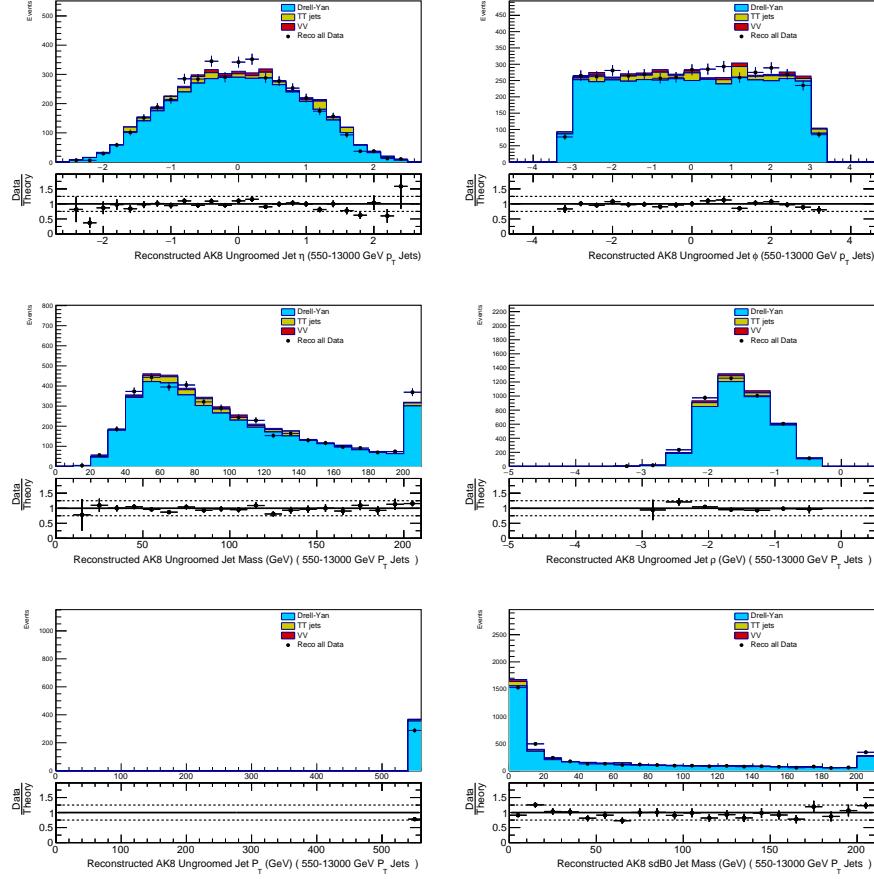


Figure 6.12: Jet pseudorapidity, azimuthal angle, mass, ρ , p_T , and soft-drop mass (shown here for the selected case $\beta = 0$ and $Z_{cut} = 0.1$ with plots of the other groomed masses in the appendix "Data and MC comparisons for soft-dropped jet masses") for data and PYTHIA8 MC for $p_T = 550 - 1300$ GeV.

6.6 Detector Response

The two-dimensional detector response is generated from MC simulation in bins of jet p_T and mass (Later we will include these distributions for jet p_T and ρ). For the groomed jet measurement, the ungroomed jet p_T and the groomed jet mass are used (here results for $\beta = 0$ and $z_{cut} = 0.1$ are shown while alternatively groomed distributions are given in the appendix). This is because the

theoretical predictions we will compare to use the ungroomed jet p_T , which is a Sudakov safe observable unlike the soft-drop groomed pt . The central value is constructed using PYTHIA8 with all systematic variations at their nominal values. The systematic variations are treated by unfolding the data or MC with new responses with the systematic variations applied, then taking the full difference from the nominal value. The response matrix uncertainties are symmetrized. The matrix is visualized by normalizing each row to unity.

The p_T bins are 200, 260, 350, 460, 550, 13000 GeV, they were chosen to match those of the previous measurement in the dijet channel in order to enable later comparisons of jets with different flavor compositions, here we present a like-quark enriched sample. We did not have sufficient statistics at high Pt to maintain the same binning as the dijet case above 550 GeV jet Pt. The response matrix is shown for ungroomed jets and groomed jets in Figures ?? and ??, respectively. The jet mass bins are chosen as a trade-off between the approximate detector resolution of ~ 10 GeV and the purity and stability. To avoid many bins below 10%, the binning chosen is initially finely binned and then becomes broader.

The purity is defined as the fraction of events in a reconstructed bin that are generated from the same bin. (That is, the diagonal value divided by the row it is in). The stability is defined as the fraction of events in a generated bin that are reconstructed in the same bin. (That is, the diagonal value divided by the column it is in). To avoid unphysical regions (where the jet mass is above what can occur given a jet p_T), the purity and stability are restricted to the regions where $m < p_T * R / \sqrt{2}$. These are shown in Figures 6.17 and Figures ?? for ungroomed and groomed jets, respectively.

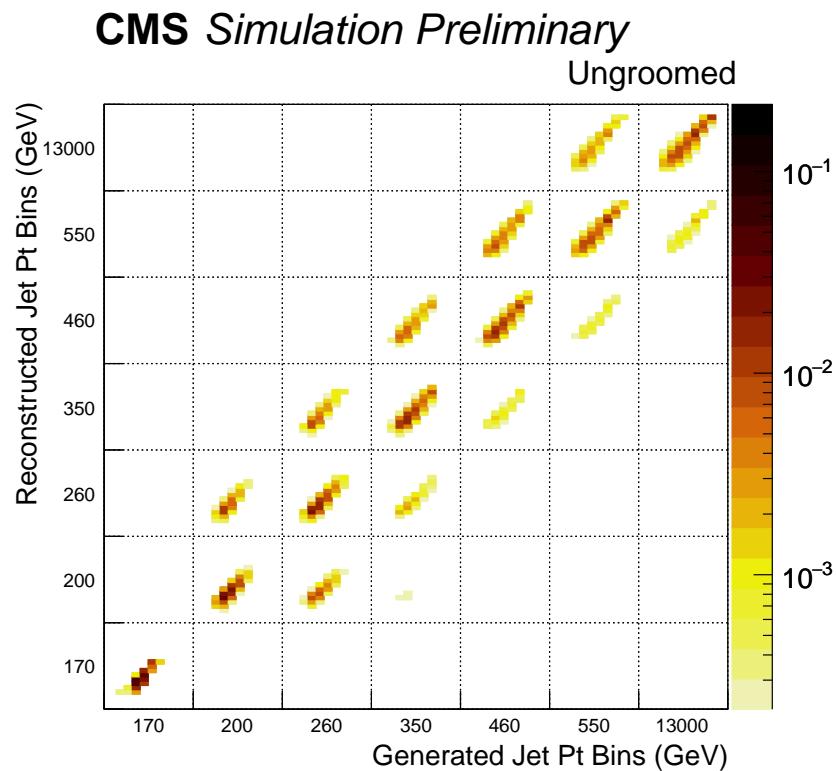


Figure 6.13: Two-dimensional response matrix for ungroomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). While the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

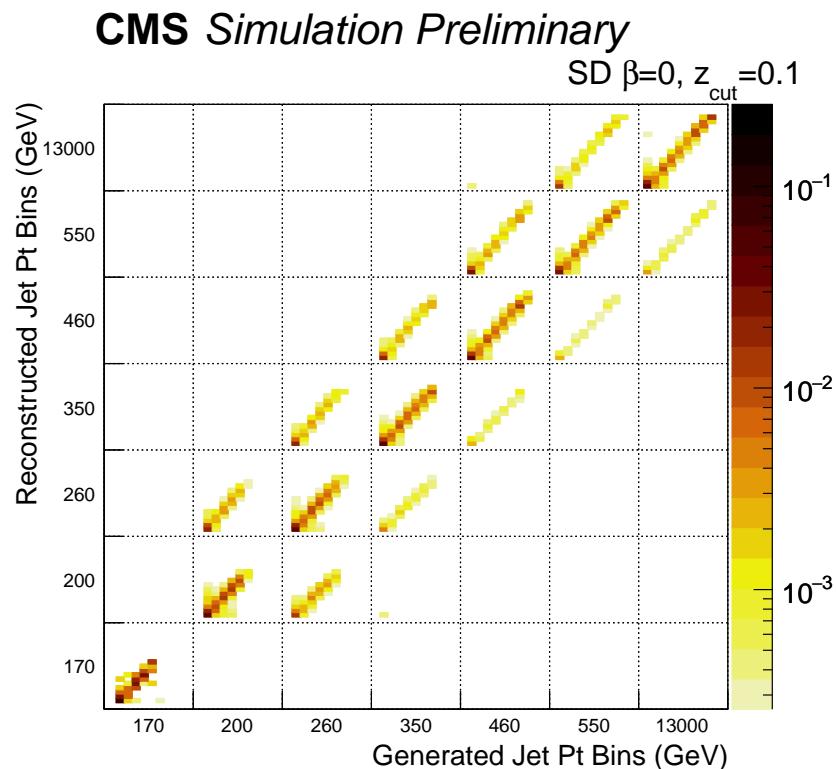


Figure 6.14: Two-dimensional response matrix for groomed jets $\beta = 0, z_{cut} = 0.1$. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV) while the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

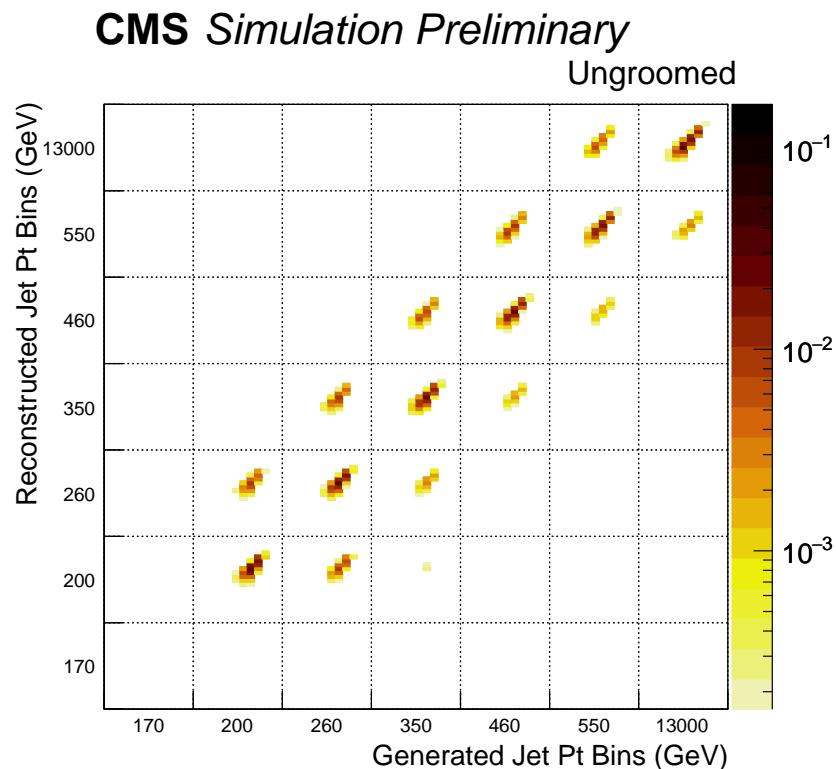


Figure 6.15: Two-dimensional response matrix for ungroomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). while the jet dimensionless mass is indexed by the minor blocks on the x (y) axis -5.0, -3.5, -3.0, -2.5, -2.0, -1.5, -1.0, -0.5, 0.0 (-5.0, -4.0, -3.5, -3.25, -3.0, -2.75, -2.5, -2.25, -2.0, -1.75, -1.5, -1.25, -1.0, -0.75, -0.5, -0.25, 0.0).

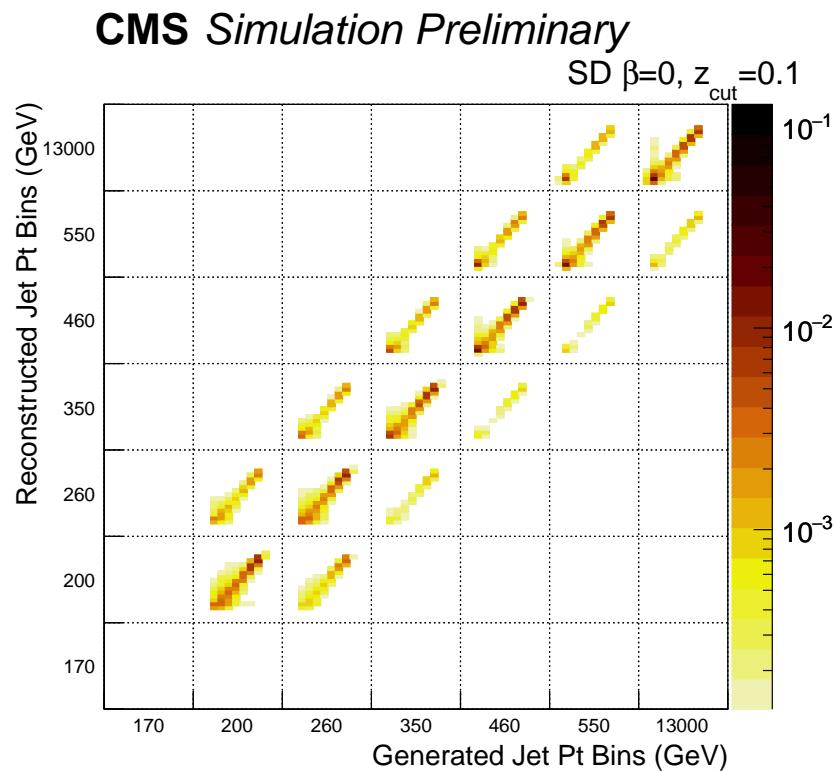


Figure 6.16: Two-dimensional response matrix for groomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). while the jet dimensionless mass is indexed by the minor blocks on the x (y) axis -5.0, -3.5, -3.0, -2.5, -2.0, -1.5, -1.0, -0.5, 0.0 (-5.0, -4.0, -3.5, -3.25, -3.0, -2.75, -2.5, -2.25, -2.0, -1.75, -1.5, -1.25, -1.0, -0.75, -0.5, -0.25, 0.0).

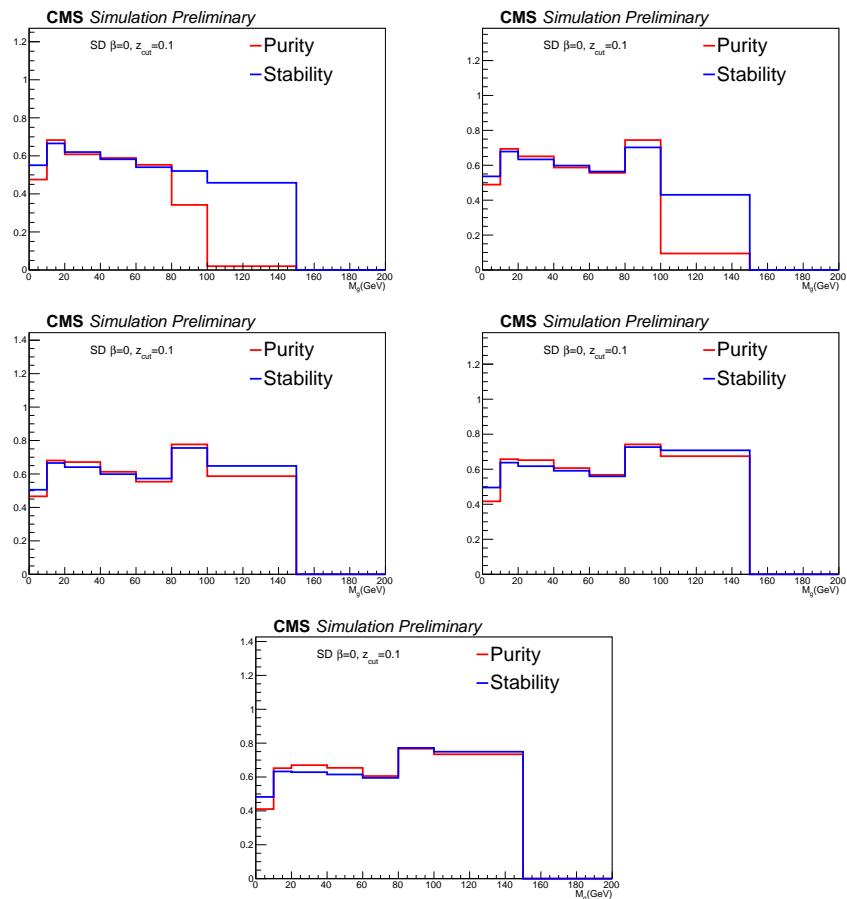


Figure 6.17: Purity and stability for ungroomed jets in various p_T bins.

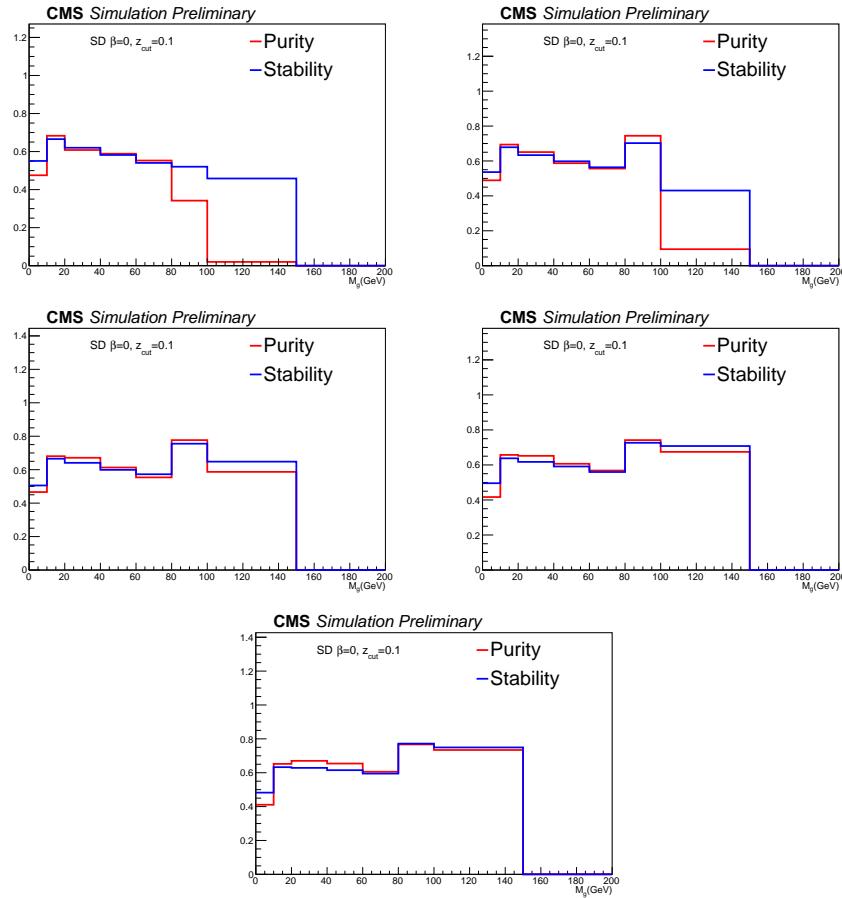


Figure 6.18: Purity and stability for groomed jets, where the soft-drop criterion was applied with $\beta = 0$ and $z_{cut} = 0.1$, in various p_T bins.

6.7 Uncertainties

The statistical uncertainties from the MC are non-trivial to estimate due to the unfolding procedure. Instead of standard methods, we employ a “jackknife resampling” technique (a linear bootstrapping method), where we create 10 response matrices that each remove an exclusive 10% of the sample. Each event is removed exactly once. The RMS of this ensemble of distributions, multiplied by 10/9ths, is taken as the statistical uncertainty [56].

There are both theoretical and experimental systematic uncertainties that affect this measurement, along with uncertainties from the unfolding procedure. Unless otherwise specified, the uncertainties are estimated by creating response matrices with the systematic uncertainty applied, and unfolding the reconstructed data or MC. The difference between the systematically varied output and the nominal output is taken to be the systematic uncertainty of the measurement. The response matrix uncertainties are symmetrized.

For the normalized cross section, the differences are taken after normalizing for each individual p_T bin. This removes much of the p_T dependence of the measurement since it cancels in the ratio. For the full cross section, the differences are taken without normalizing each p_T bin, so include these effects. As such, the normalized cross section has very small jet energy uncertainties, whereas the absolute cross section has large jet energy uncertainties. The luminosity uncertainty also cancels in the ratio.

The predominant experimental uncertainties that affect the normalized cross section measurement are the jet mass scale, the jet mass resolution, the parton shower uncertainty, the pileup reweighting scheme, the jet energy scale, and the jet energy resolution. The predominant theoretical uncertainty is the physics

model.

In addition to the previous effects, the absolute cross section also has sizeable uncertainties from the jet energy scale, jet energy resolution, and overall luminosity measurement.

- **Jet energy scale:** The JES uncertainty is evaluated by varying the L1L2L3 corrections up and down by the appropriate uncertainties as described in Ref [53].
- **Jet energy resolution:** The JER uncertainty is evaluated by smearing the jet energy by an additional factor as described in Ref [53]. The central value is determined by smearing to the JER nominal value, and the uncertainties are evaluated by smearing to the JER systematic variations.
- **Jet mass scale:** The JMS in Run 1 was equal to unity within experimental uncertainties, derived from investigating the mass peak of the W boson within $t\bar{t}$ events, and comparing the results in data and MC. In Run 2, this is still effectively unity within statistical uncertainties [57]. The uncertainty on this is less than 1%, and is thus neglected.
- **Jet mass resolution:** The JMR in Run 1 was equal to 1.11 ± 0.09 . In Run 2 we now observe 1.07 ± 0.12 . The JMR is evaluated by stretching the jet mass by 7% relative to the generator-level jets [57]. The uncertainty is set to be the uncertainty in this measurement, 12%.
- **Pileup:** The pileup uncertainty is estimated by reweighting the MC, assuming the minimum bias cross section of 69 mb is scaled up and down by 1.4%.

- **Luminosity:** The CMS measurement of the luminosity for the 2015 data-taking period is 2.7% [58].
- **Parton distribution functions:** The PDF uncertainty is evaluated by using the PDF4LHC15 meta-PDF set.
- **Physics model uncertainty:** The PS uncertainty, which includes the parton shower (PS), generator tune, and physics showering model, is evaluated by comparing a response matrix derived from PYTHIA8 to a response matrix derived from HERWIG++. The difference between the response matrices for the two cases is taken as a shifted systematic uncertainty.

Figure-6.19 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

Figure-6.20 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

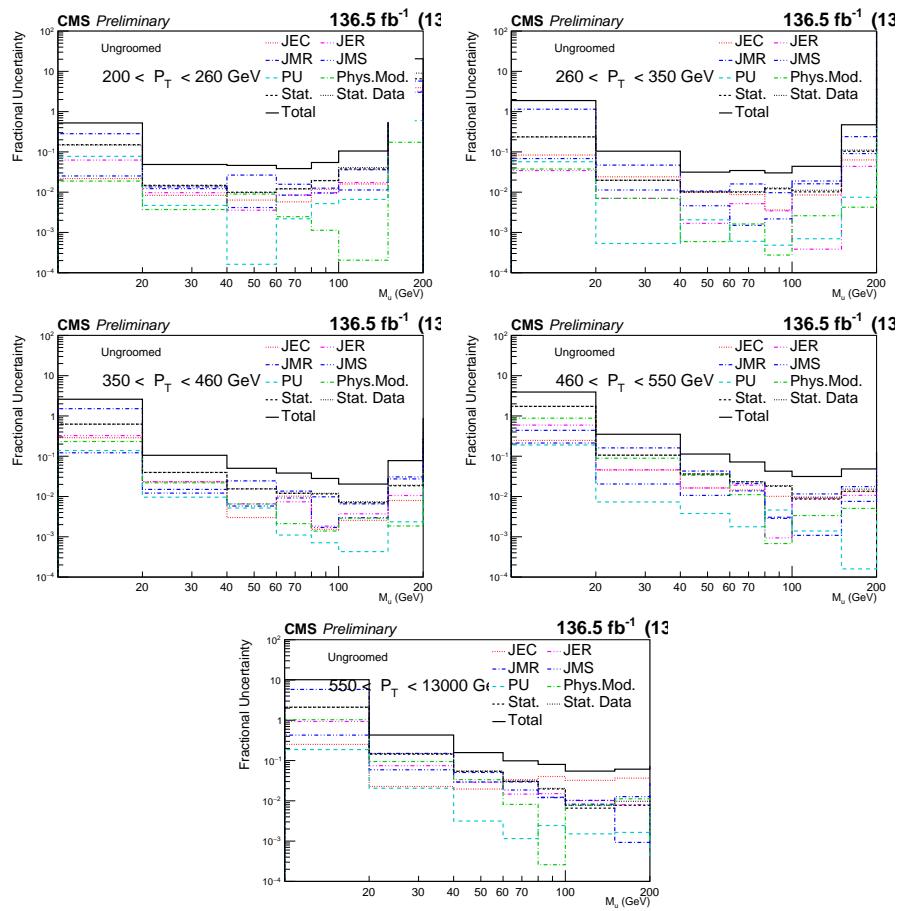


Figure 6.19: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw jet mass.

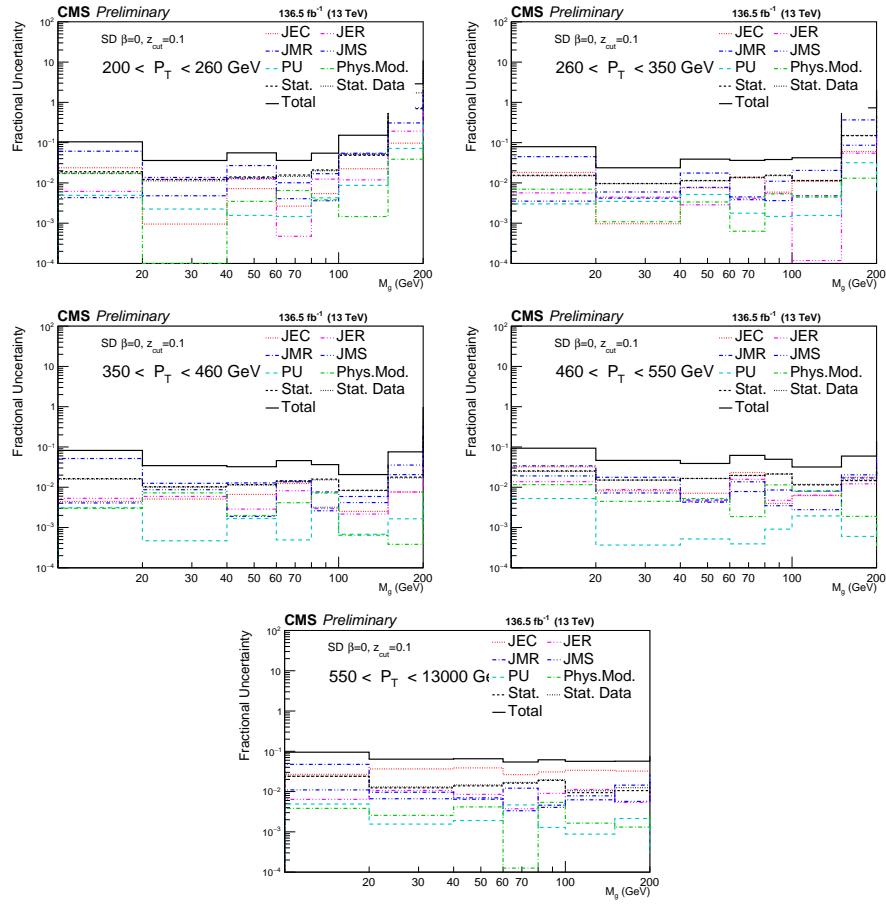


Figure 6.20: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

6.8 Unfolding

The unfolding procedure in this analysis is a least squares fitting with optional Tikhinov regularization, implemented using the TUnfold algorithm [59]. This procedure was used rather than the D'Agostini [60] method employed in previous measurements as it has the capacity to accept variant truth MC vectors as input, not simply the same one used to fill the response matrix as required in the current RooUnfold version. This issue was brought to our attention due to the recommendation of the CMS statistics committee and their claim that TUnfold is better documented. We have found this claim to be unfounded in the two dimensional case however we have developed some informational materials outlining TUnfold in 2D using the algorithm's built-in systematic uncertainty handling. Regardless, moving to TUnfold from RooUnfold prevents a bias towards the input distribution since TUnfold allows the option of using different input MC truth vectors, not necessarily filled with the same events as the response matrix, which is lacking in RooUnfold and causing regularization to bias the results towards the input distributions. This is currently the first analysis that we are aware of on CMS, which makes use of TUnfold in 2D using the built-in systematic uncertainty handling.

The response matrix was used to unfold the MADGRAPH + PYTHIA8 MC with and without Tikhinov regularization .It was determined that regularization is not necessary for the current binning scheme due to the high purity and stability of the binning scheme. Furthermore, the condition number was determined to be 1.43 which was considered small enough, less than 10 is suggested by the statistics committee currently, such that regularization was not required. A simple matrix multiplication was then used to re-fold the unfolded MC with-

out regularization. The resulting re-folded distribution was compared to the original reconstructed distribution.

Several validations were performed to ensure that the unfolding procedure was accurately reproducing input distributions.

1. **Closure test in MC:** The PYTHIA8 MC was used to construct both the response matrix as well as the reconstructed distribution. By construction, this should reproduce exactly the input given, since that is where it was derived from. Figure ?? showing the closure tests for ungroomed jets for all p_T bins, and Figures 6.31-6.34 show the same for groomed jets.
2. **Physical bias test in MC:** To investigate the capability of the response matrix to correctly unfold a spectrum different from the inputs, we have performed another closure test unfolding the 2016 inclusive HERWIG++ sample as the input distribution with the 2016 inclusive PYTHIA8 response matrix . This is not the same as the systematic variation due to the parton shower uncertainty, where the difference between the PYTHIA8 response matrix and the HERWIG++ response matrix is treated as a shifted uncertainty, using TUnfold's internal systematic uncertainty handling, applied to the unfolding of the nominal PYTHIA8 reconstructed response matrix.

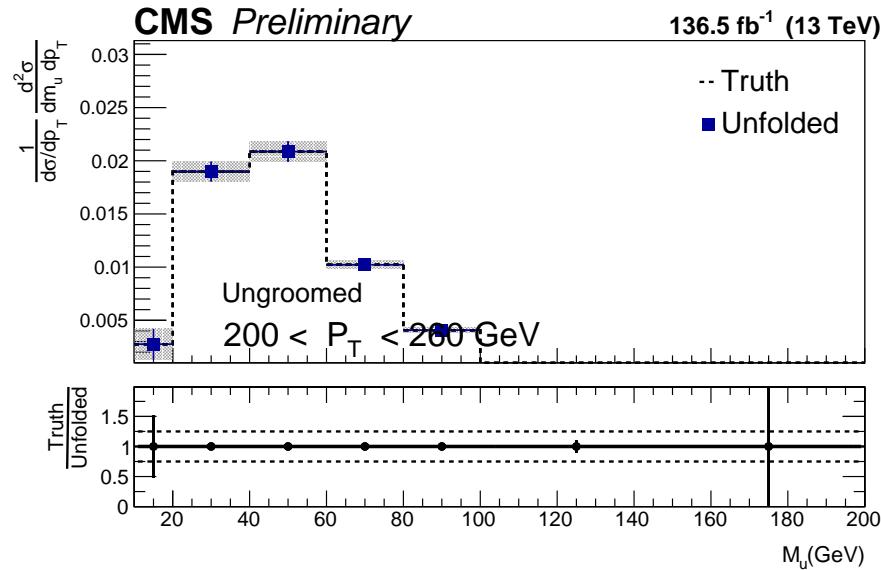


Figure 6.21: Closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

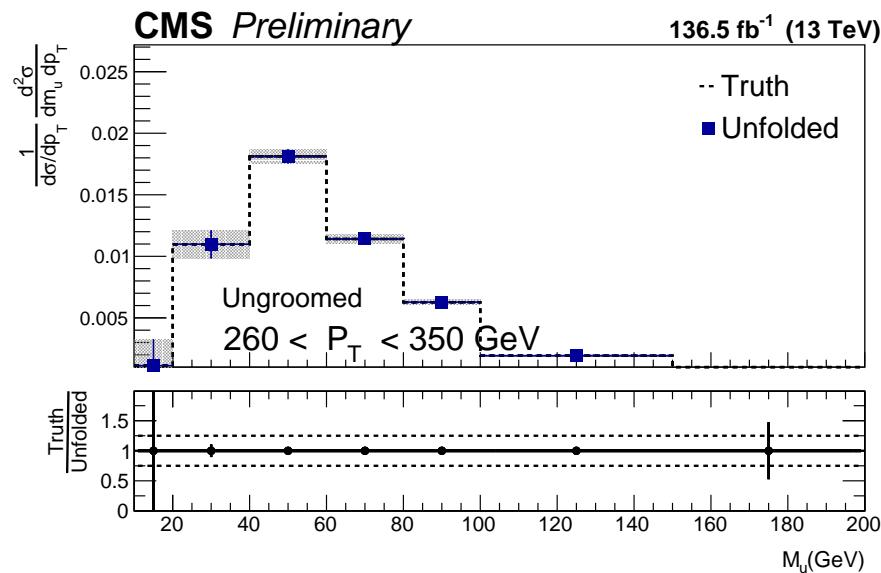


Figure 6.22: Closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

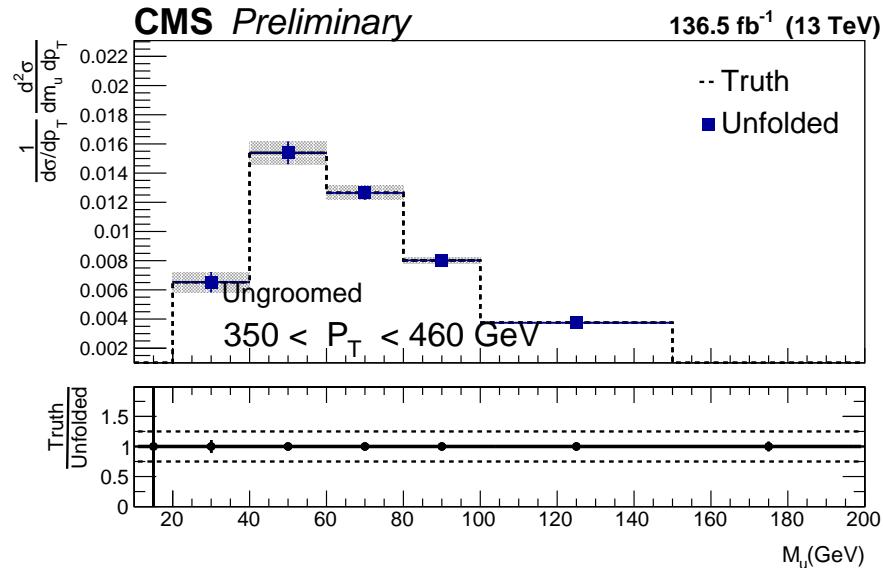


Figure 6.23: Closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

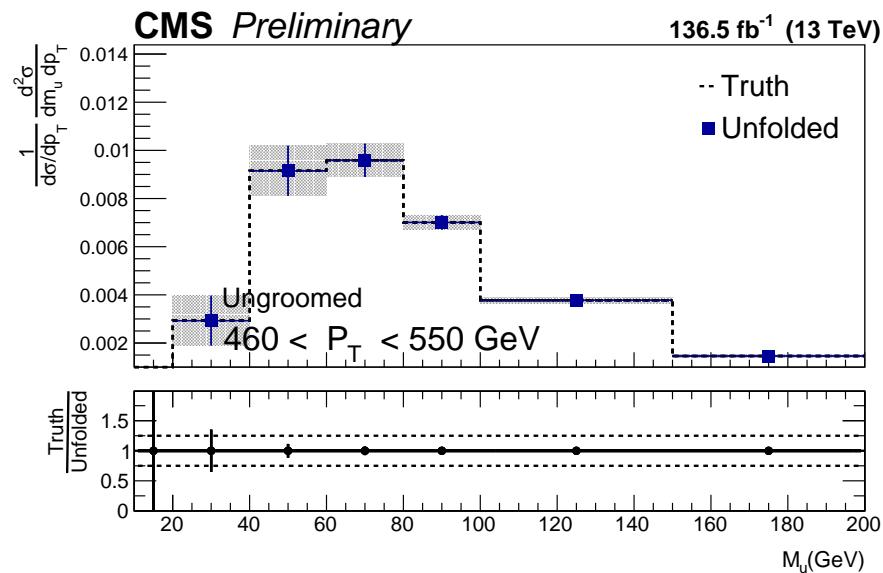


Figure 6.24: Closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

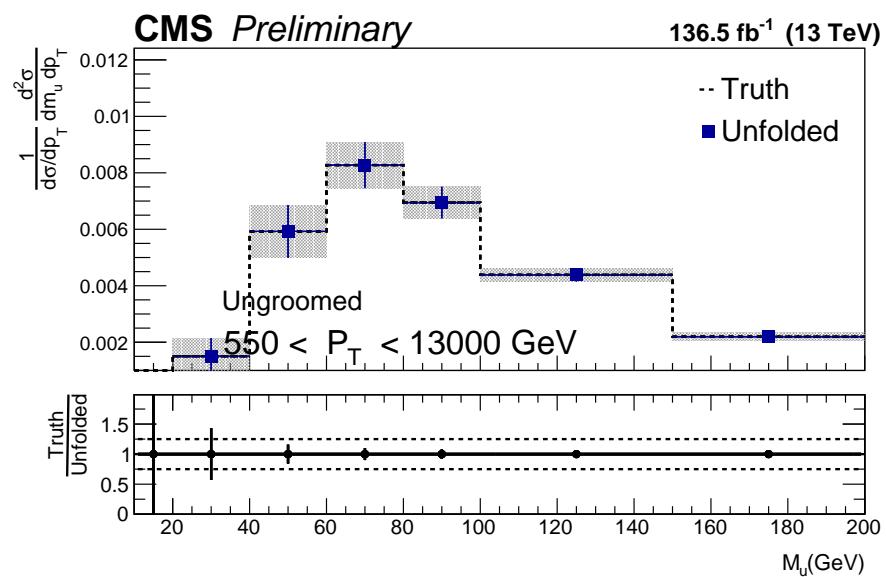


Figure 6.25: Closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

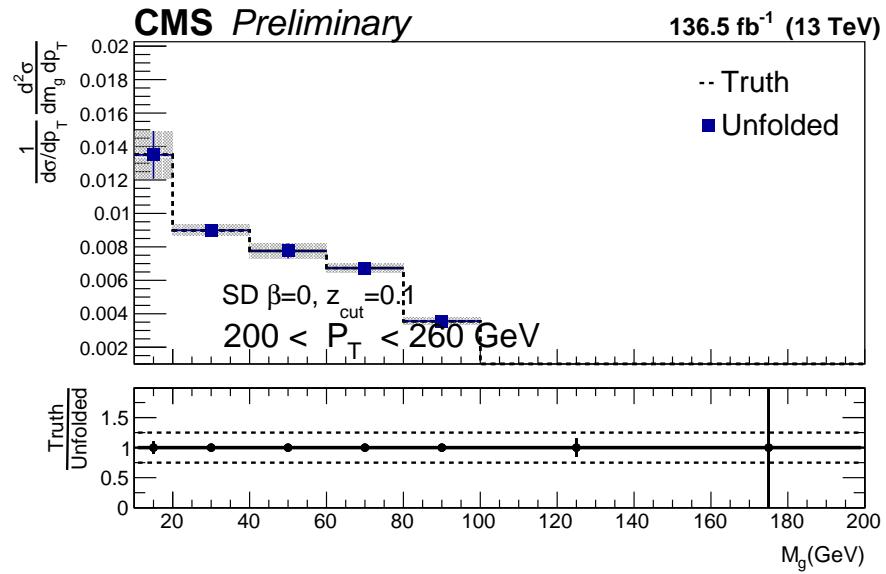


Figure 6.26: Closure test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

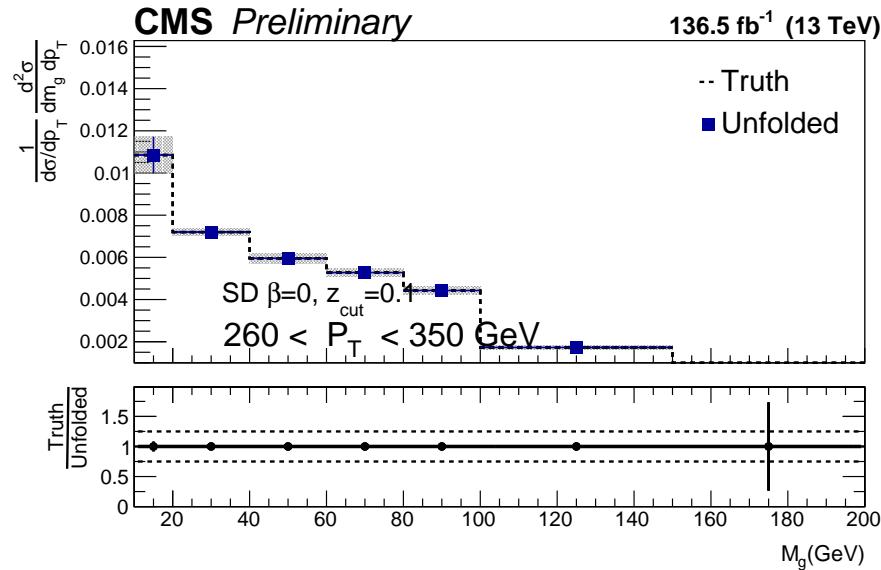


Figure 6.27: Closure test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

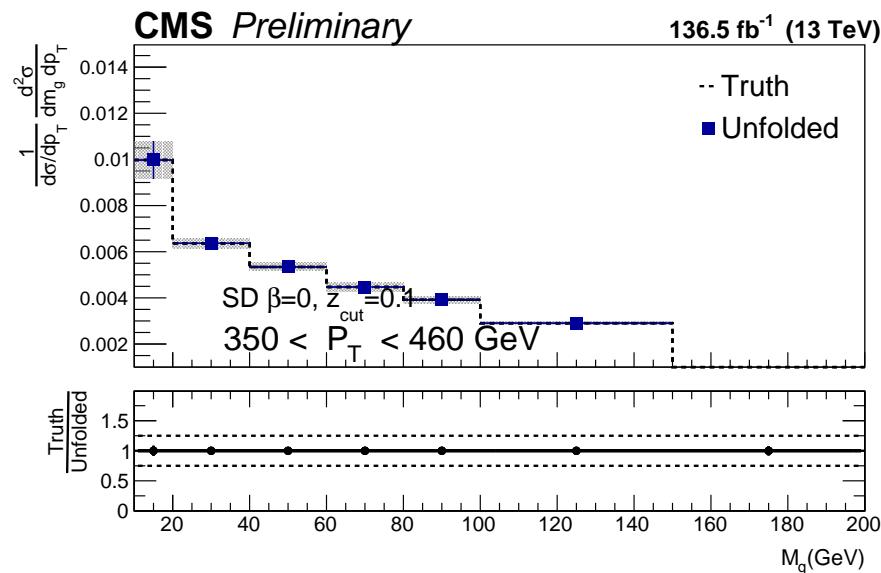


Figure 6.28: Closure test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

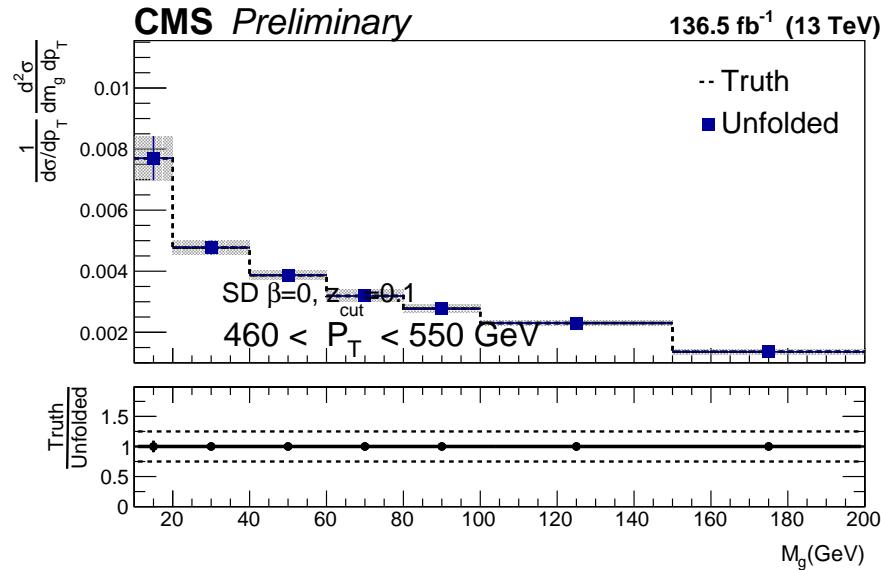


Figure 6.29: Closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

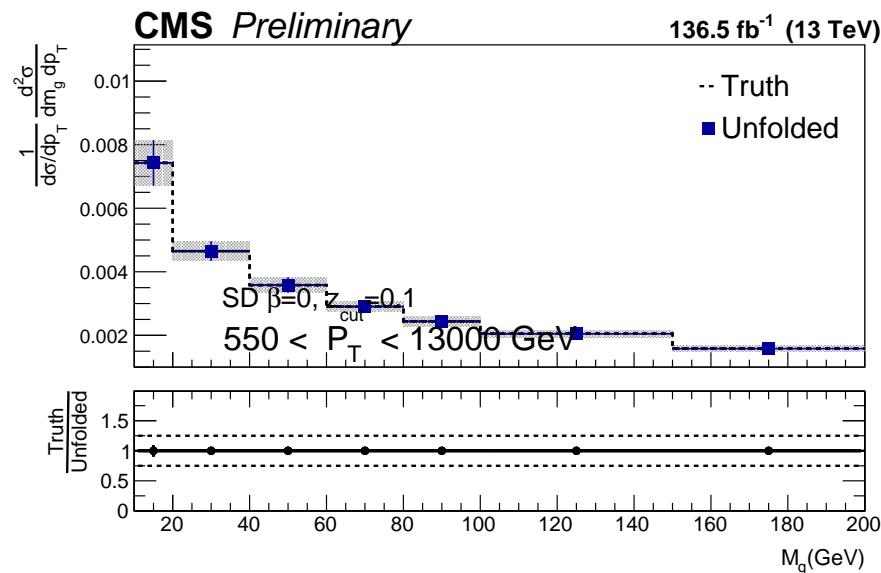


Figure 6.30: Closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

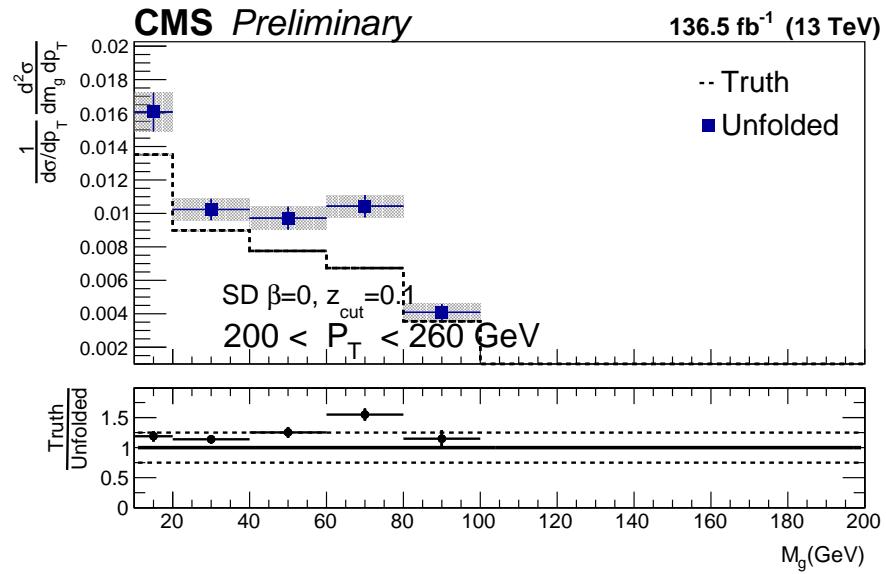


Figure 6.31: Physical bias test of groomed reconstructed Monte Carlo, p_T 200–260 GeV.

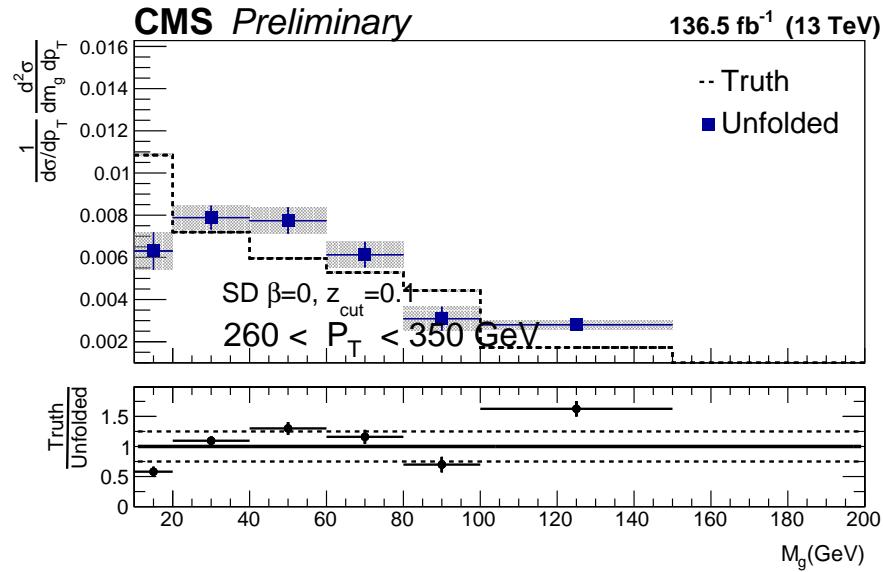


Figure 6.32: Physical bias test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

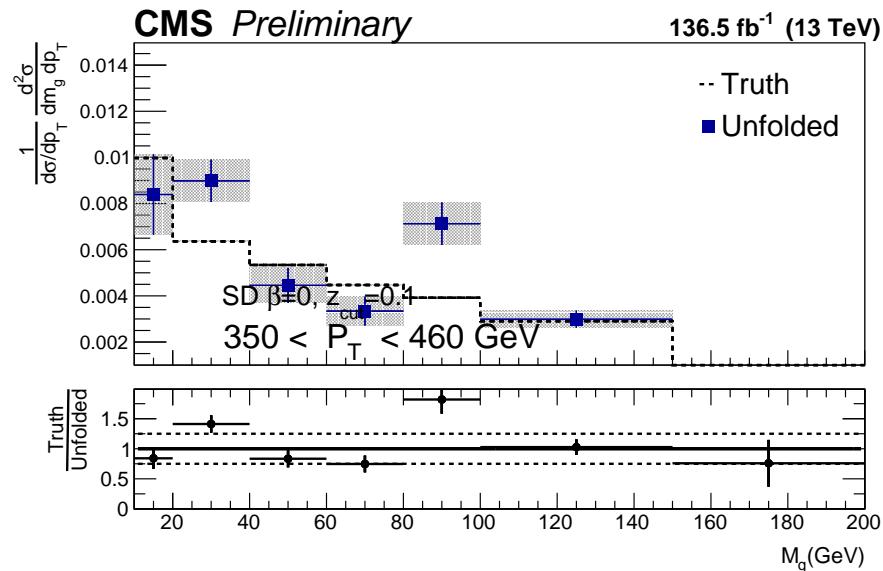


Figure 6.33: Physical bias test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

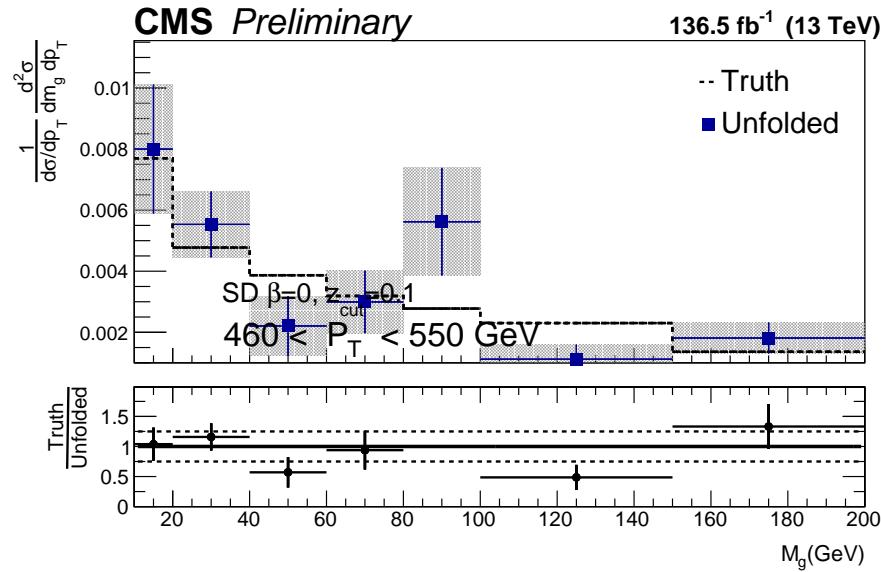


Figure 6.34: Physical bias test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

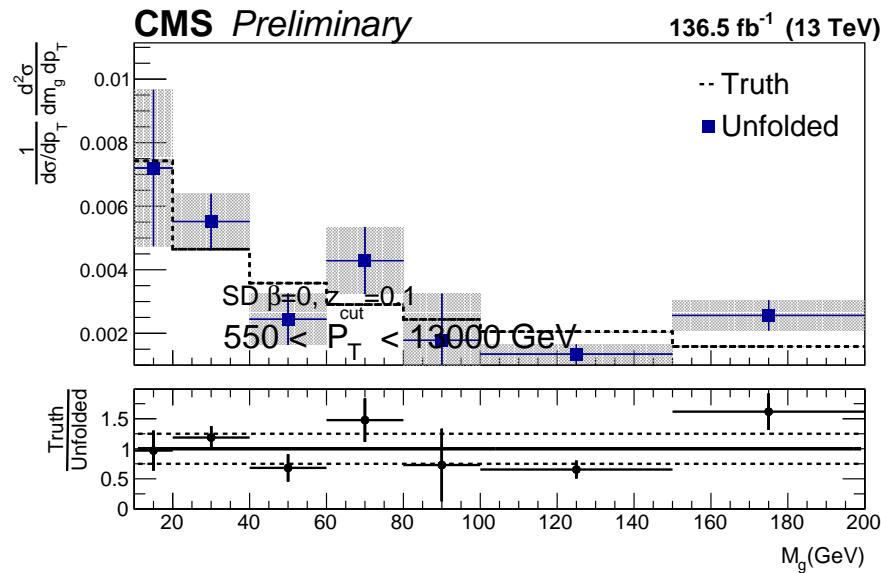


Figure 6.35: Physical bias test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

6.9 Results

The unfolded jet p_T spectrum is shown in Fig. ???. The dominant uncertainty is the jet energy scale.

The normalized cross section results for the ungroomed jets are shown in Figures ??-??, and for the groomed jets are shown in Figures ??-???. The distributions are shown normalized per p_T bin to compare only the shapes. All p_T bins are shown in double-logarithmic plots in Figures 6.46 and 6.47.

In order to visualize the results in a slightly different way, we also present the normalized cross section with respect to $\log(m)$. This is equivalent to multiplying the unfolded cross section by the mass. We are approximating this by multiplying each bin in the final result by the average mass in that bin. Thus, we present

$$\frac{m}{d\sigma/dp_T} \frac{d^2\sigma}{dm dp_T}. \quad (6.2)$$

These results are shown for ungroomed jets in Figures ??-???, and for the groomed jets are shown in Figures 6.41-6.43. The distributions are normalized per p_T bin and multiplied by the center of the mass bin.

The results are currently presented compared to MADGRAPH + PYTHIA8 and will soon be compared to HERWIG++ when the 2017 samples are ready, and, for the normalized groomed case, we will eventually compare the calculation of the jet mass from Marzani et al. in progress now, expected in early 2020. The normalized results are scaled to unity for each individual jet p_T bin for the unfolded data, the MADGRAPH + PYTHIA8, and results from the calculation from the authors of Ref. [5]. The calculation from Ref. [14], on the other hand, is normalized to match the data at 50 GeV, since the calculation does not consider

nonperturbative corrections that are important below that value. Furthermore, at very high masses (relative to the p_T), the LO matrix element matching is also insufficient to capture the true dynamics, since at LO there is a kinematic turn-off at $p_T R/2$ (where R is the distance parameter for the jet clustering), from the relativistic kinematics of a $1 \rightarrow 2$ decay. However, for real jets the turn-off is closer to $p_T R/\sqrt{2}$ due to stochastic effects. To see this, consider a particle of energy E and mass m decaying to two massless particles, each with an energy $E/2$ and separated by an angle θ . The mass must satisfy $m^2 < \frac{E^2}{2} (1 - \cos \theta)$. In the small angle limit, this would be $m^2 < E^2 \theta^2 / 4$, or $m < E\theta/2$. With more particles, the stochastic nature of the shower increases this to $m < E\theta/\sqrt{2}$. Thus, a leading-order ($1 \rightarrow 2$) decay will have a faster kinematic turn-off than an all-orders ($1 \rightarrow \text{many}$) decay. This is described in A.

The data are reproduced reasonably well by the simulations, with poorer agreement below the Sudakov peak. The matrix element does not have a large impact on the results due to the excellent agreement between Madgraph+PYTHIA8 and HERWIG++ alone, however the parton shower program chosen can have larger effects. Values of the χ^2 probabilities for the data-to-MC and data-to-theory comparisons are shown in Tables ?? and ?? . Table ?? shows the probabilities if we restrict the range of the comparison to accommodate known deficiencies in the predictions. At very low masses, nonperturbative effects are large, and at very large masses, the fixed order matching is insufficient to capture the kinematics. As such, the comparisons are restricted to a range of intermediate jet masses. In Table ?? the ranges are unrestricted, for comparison. In addition to these effects, at very low groomed masses, the resolution of the detector itself spoils the agreement between simulation and theory. For these reasons, the probability after the range restriction is more appropriate for comparisons.

The soft drop algorithm considerably lowers the jet mass distribution, as observed in Figs. ??-?? and Fig. 6.47. The soft and collinear parts of the jet are removed, leaving the hard jet function, which tends to have low mass. The precision improves after the grooming algorithm is applied, and the physics modeling (primarily parton shower) uncertainty is reduced.

It is also worth noting that the groomed jet distributions exhibit falling spectra for nearly all jet p_T bins, whereas the ungroomed jets have a peaked structure, making unfolding more difficult. The groomed jets are able to be nicely unfolded with much lower uncertainties. In addition to this feature, the parton shower uncertainty, which is one of the largest uncertainties for the ungroomed jets below the Sudakov peak, is dramatically reduced after grooming.

Checks on the unfolding are shown in Figures 6.52-?. Figures 6.52-? and Figures 6.57-? show the ratio of the unfolded to raw data and MC for ungroomed and groomed jets, respectively. Figures 6.62-? and Figures 6.67-? show the generator truth over the unfolded data, divided by the raw MC over the raw data. The behaviors observed are well predicted by the MC. There are some features at very low masses (below the Sudakov peak) that the unfolding introduces, but this is reproduced by the MC.

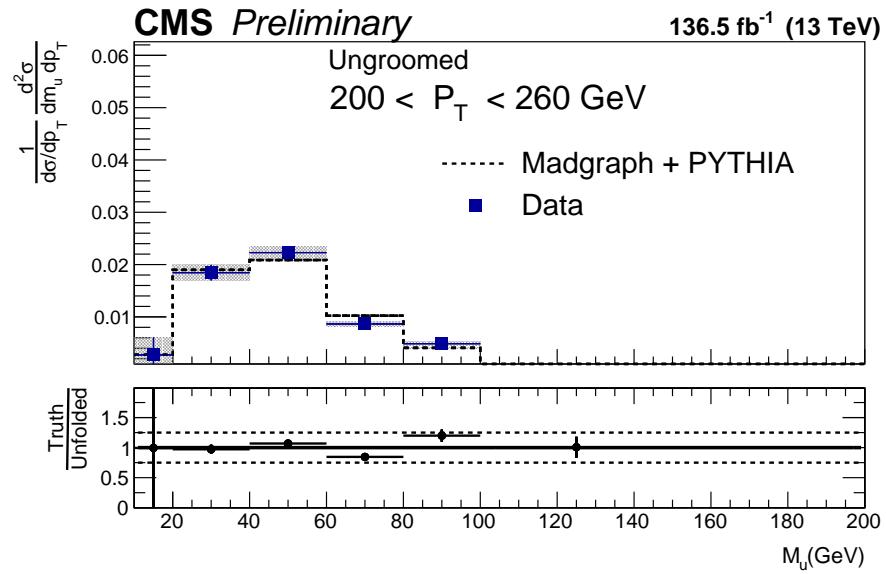


Figure 6.36: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 200-260 GeV.

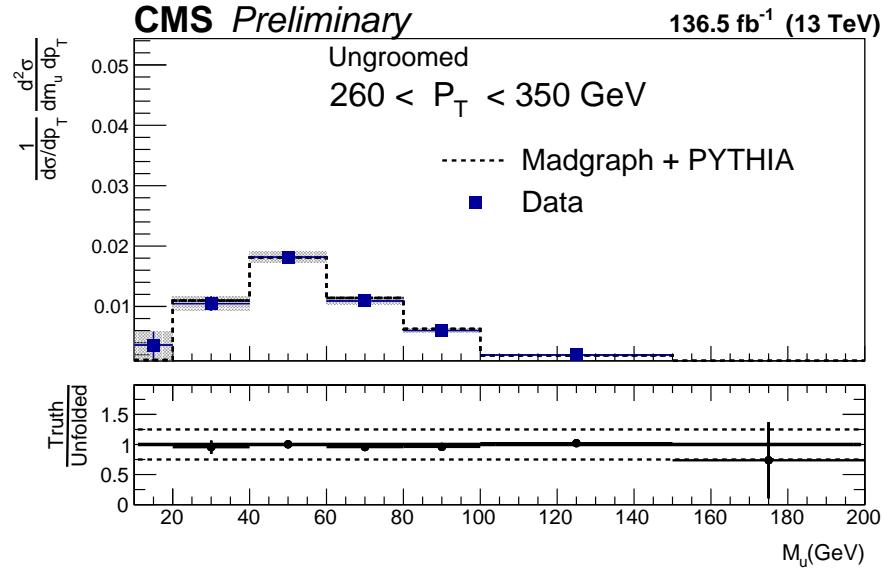


Figure 6.37: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 260–350 GeV.

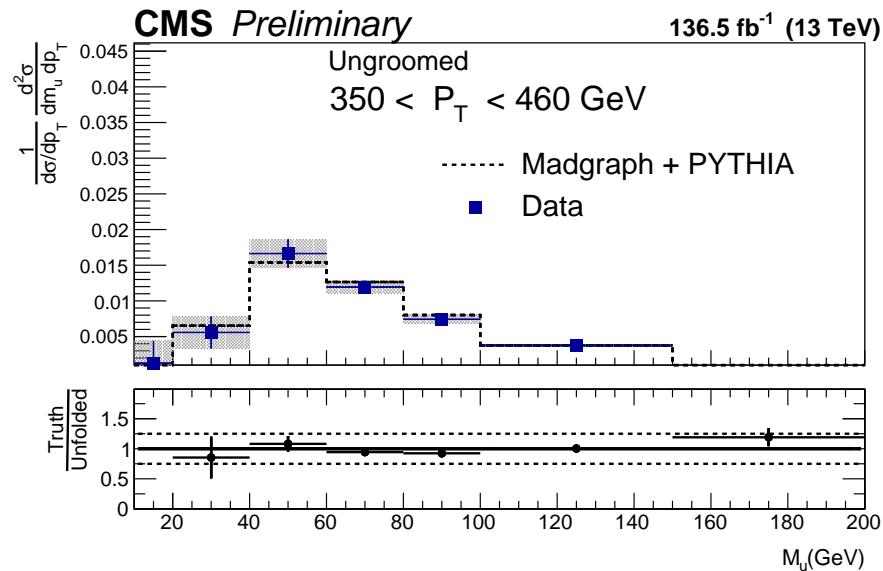


Figure 6.38: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 350–460 GeV.

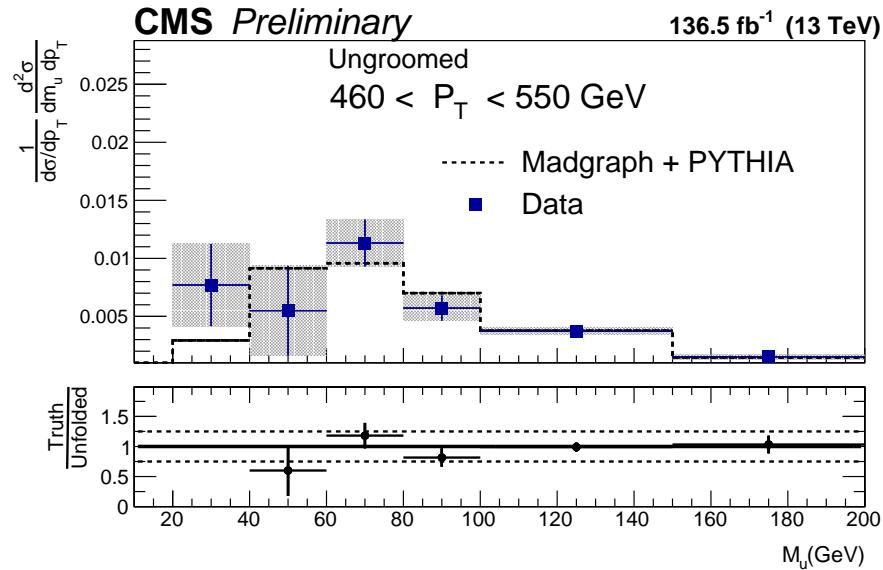


Figure 6.39: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 460-550 GeV.

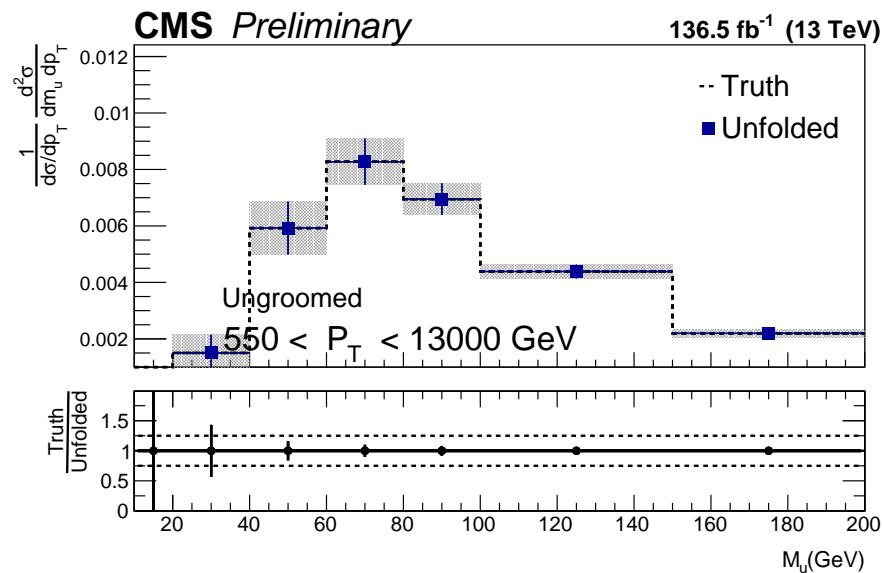


Figure 6.40: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 550-13000 GeV.

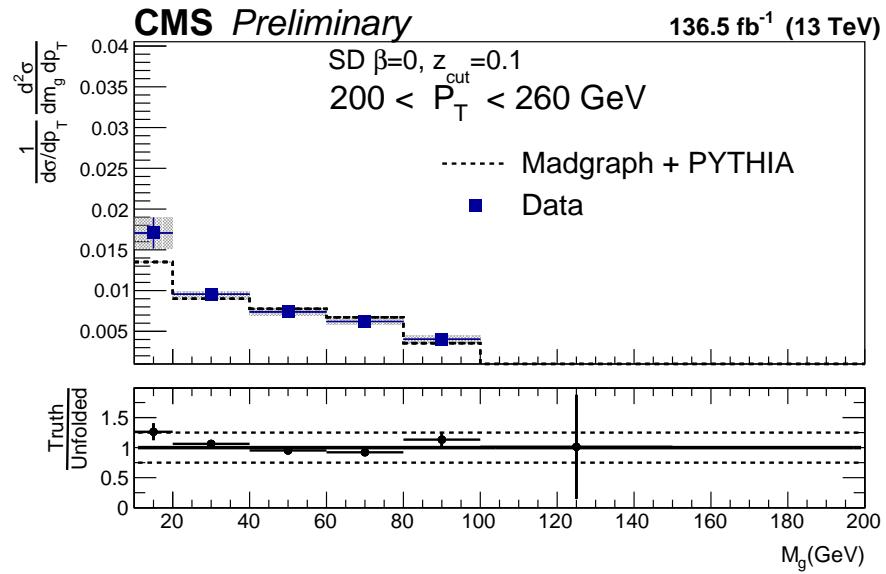


Figure 6.41: Normalized cross section results with respect to jet mass for groomed jets, p_T 200-260 GeV.

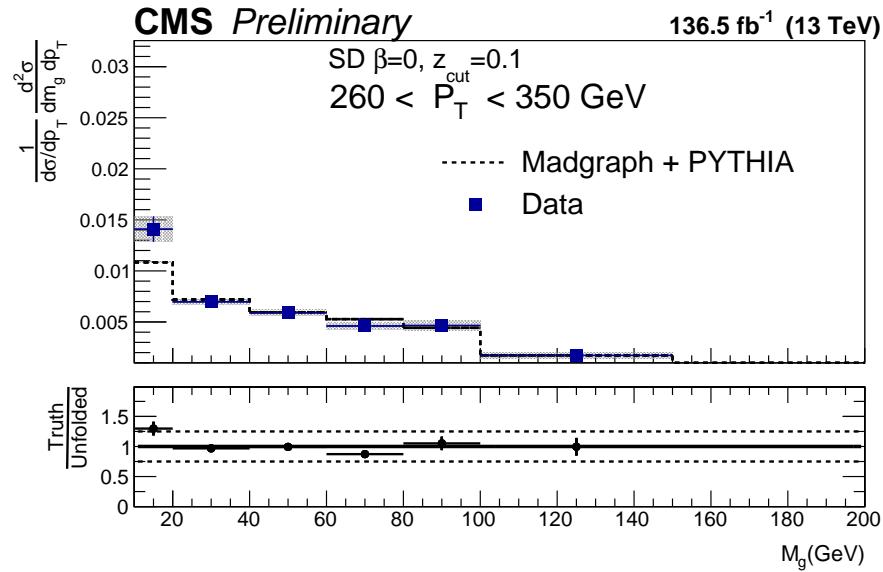


Figure 6.42: Normalized cross section results with respect to jet mass for groomed jets, p_T 260-350 GeV.

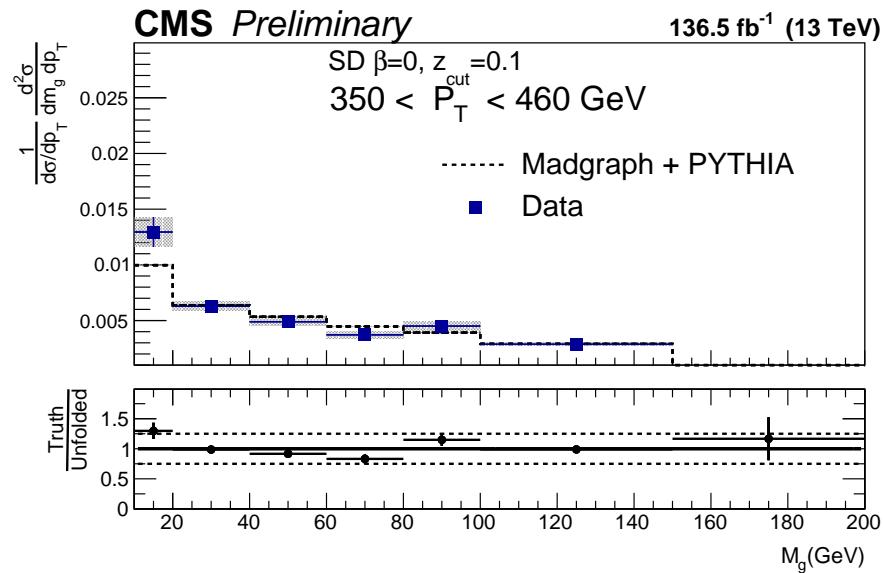


Figure 6.43: Normalized cross section results with respect to jet mass for groomed jets, p_T 350-460 GeV.

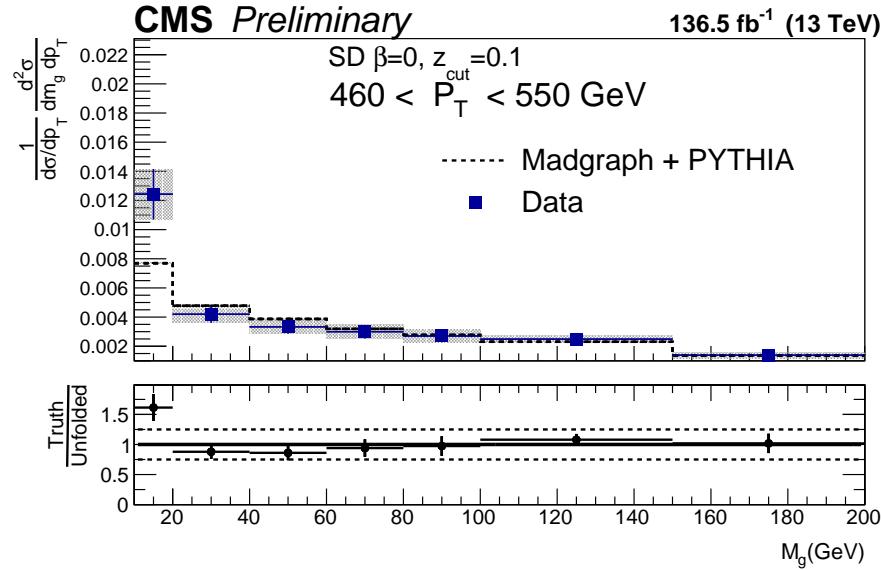


Figure 6.44: Normalized cross section results with respect to jet mass for groomed jets, p_T 460-550 GeV.

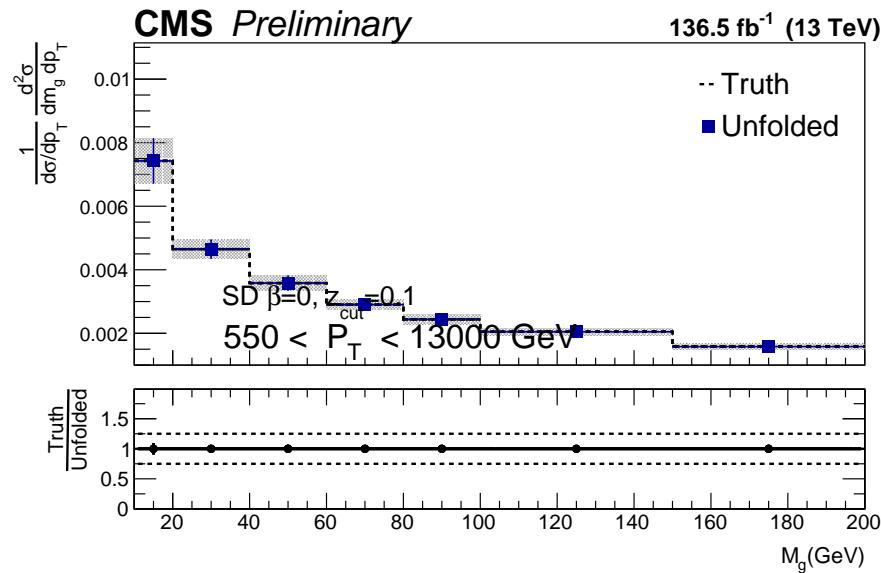


Figure 6.45: Normalized cross section results with respect to jet mass for groomed jets, p_T 550-13000 GeV.

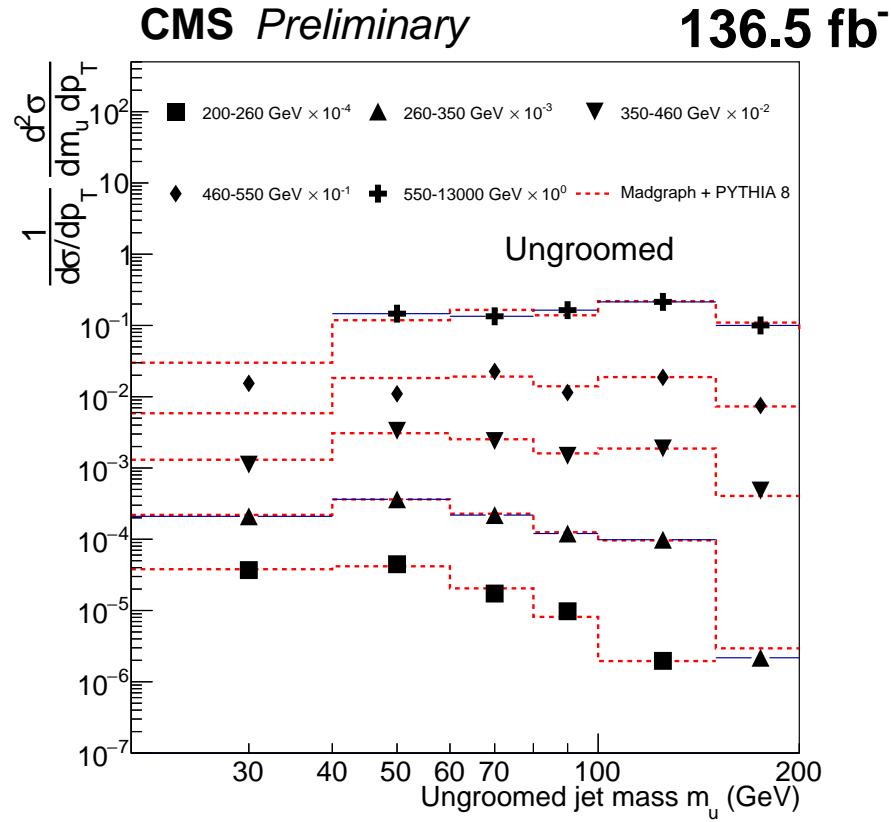


Figure 6.46: Results for ungroomed reconstructed unfolding with jet mass.

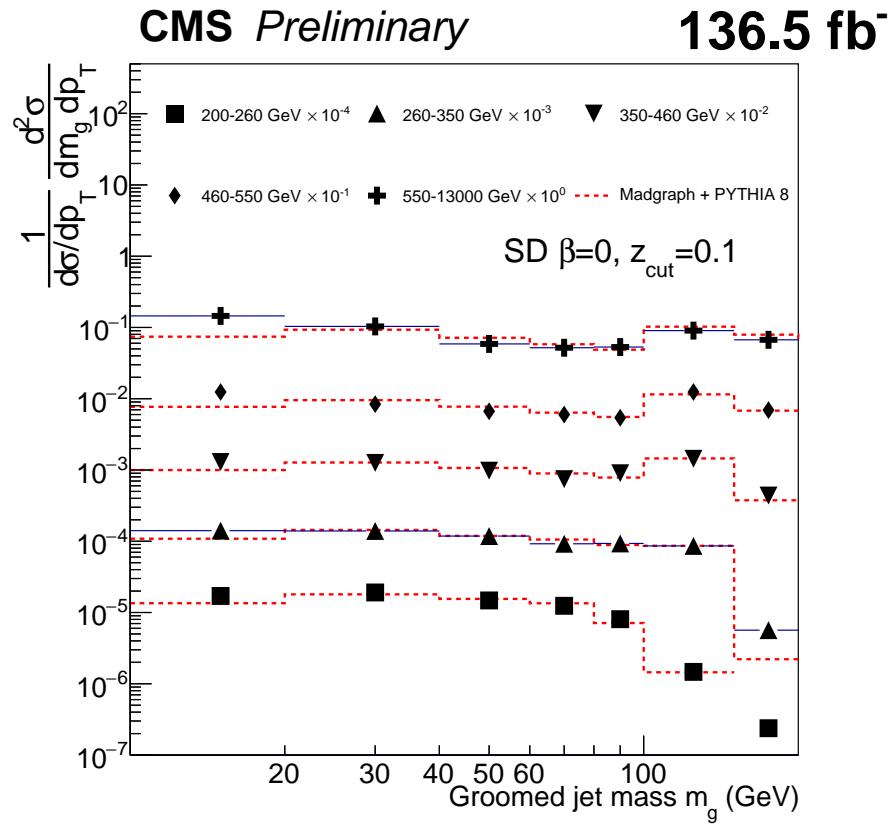


Figure 6.47: Results for groomed reconstructed unfolding with jet mass.

The resulting correlation matrix without systematic uncertainties for the ungroomed jets is shown in Fig. 6.48, and for the groomed jets is shown in Fig. 6.49. The same figures with systematic uncertainties are shown for the ungroomed jets in Fig. 6.50, and for the groomed jets is shown in Fig. 6.51.

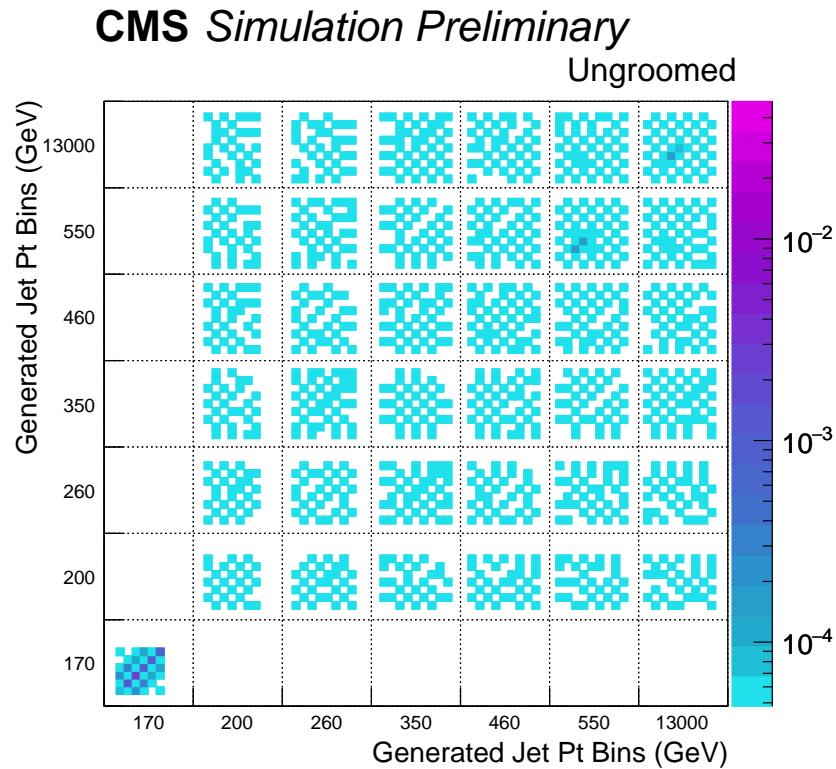


Figure 6.48: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 10, 20, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 350, \dots, 13000$ GeV).

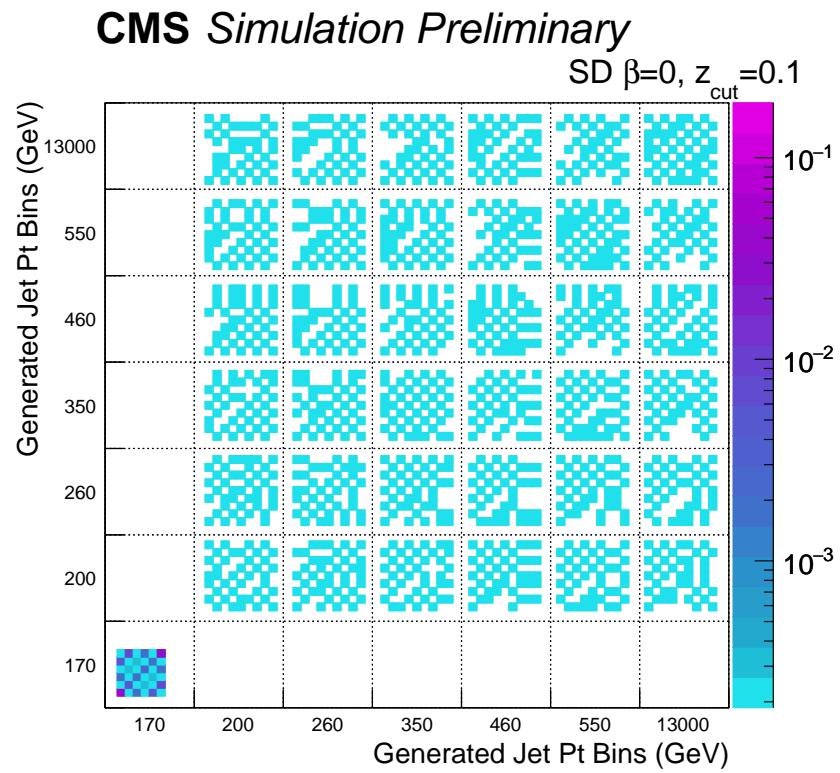


Figure 6.49: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...760 GeV).

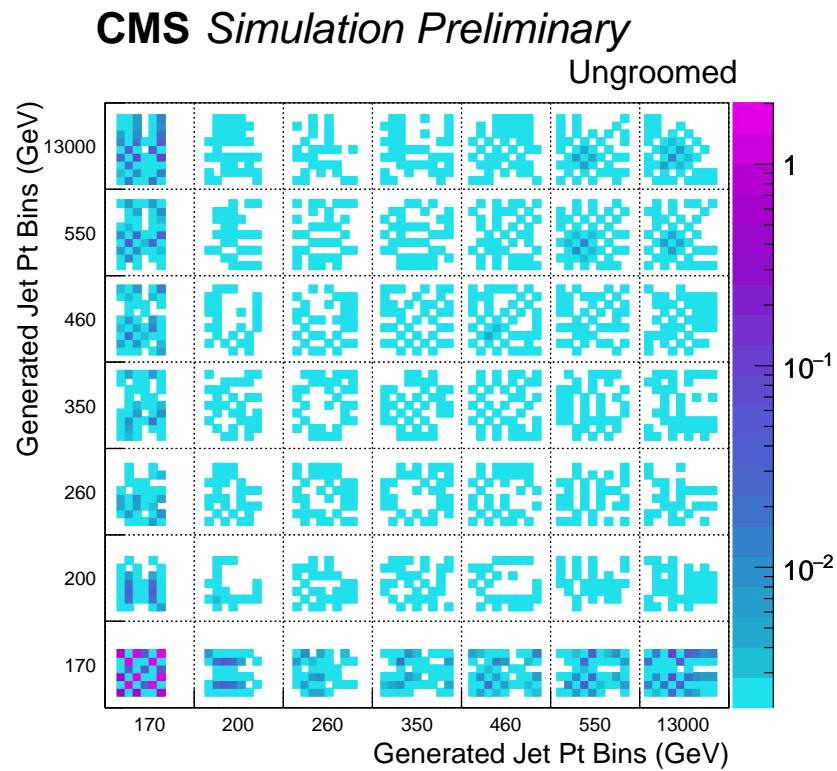


Figure 6.50: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,260,...13000 GeV).

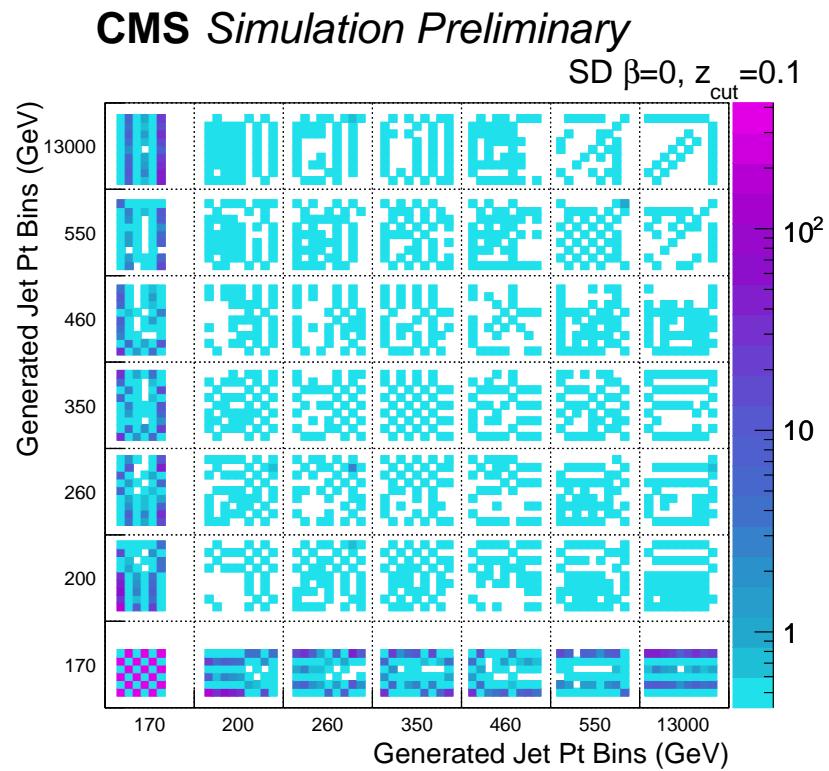


Figure 6.51: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...0 GeV).

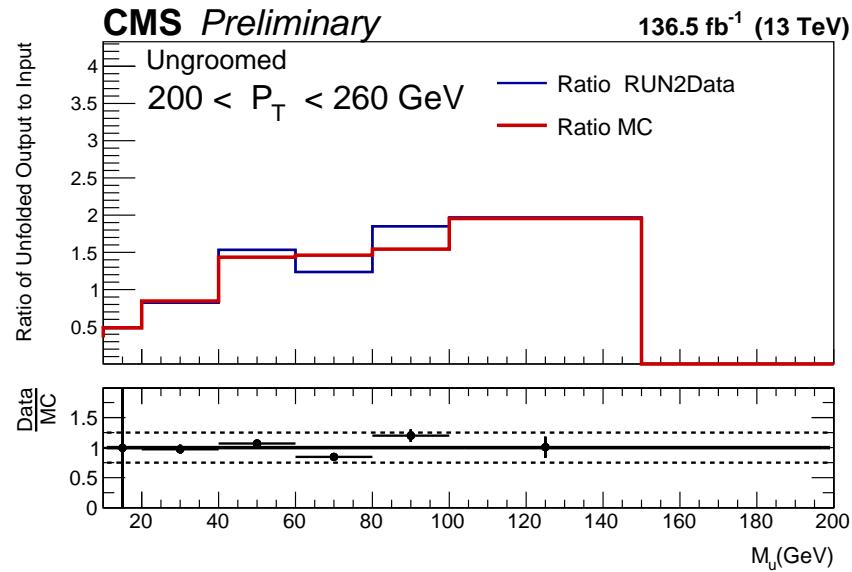


Figure 6.52: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 200-260 GeV.

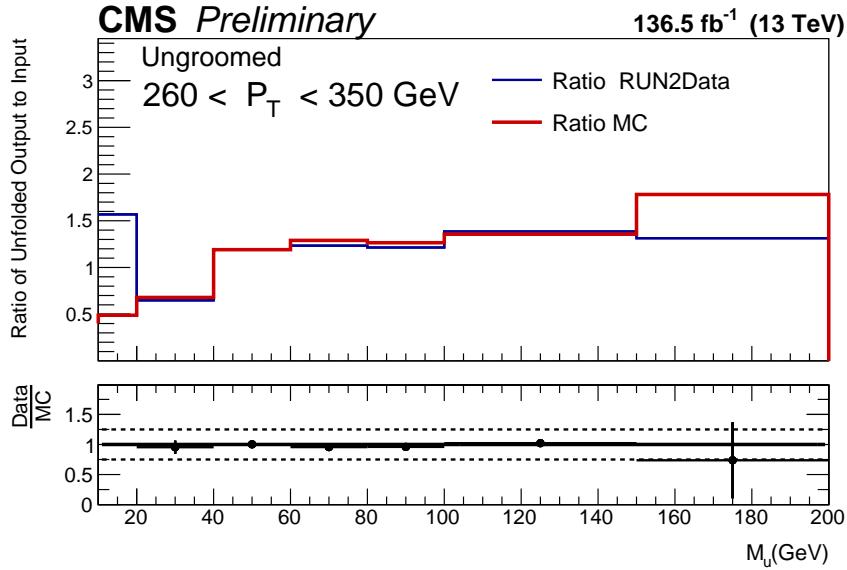


Figure 6.53: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 260-350 GeV.

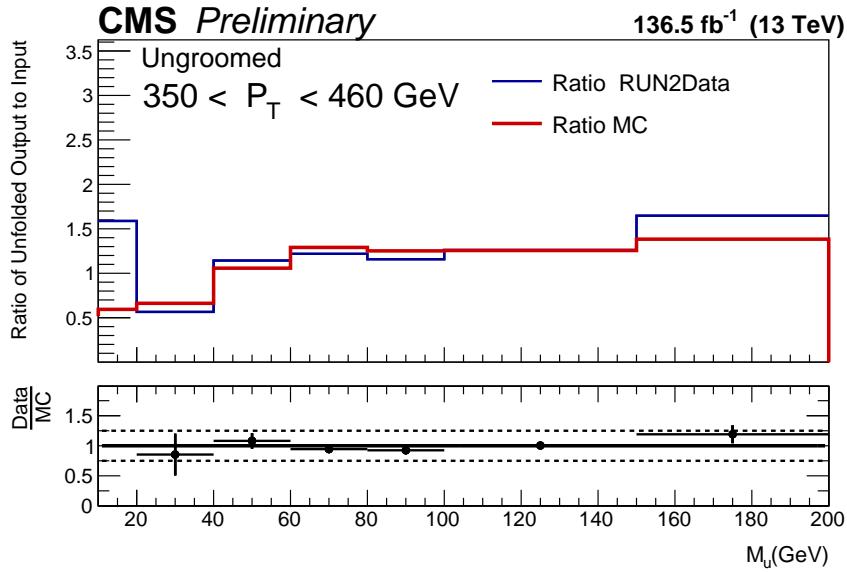


Figure 6.54: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 350-460 GeV.

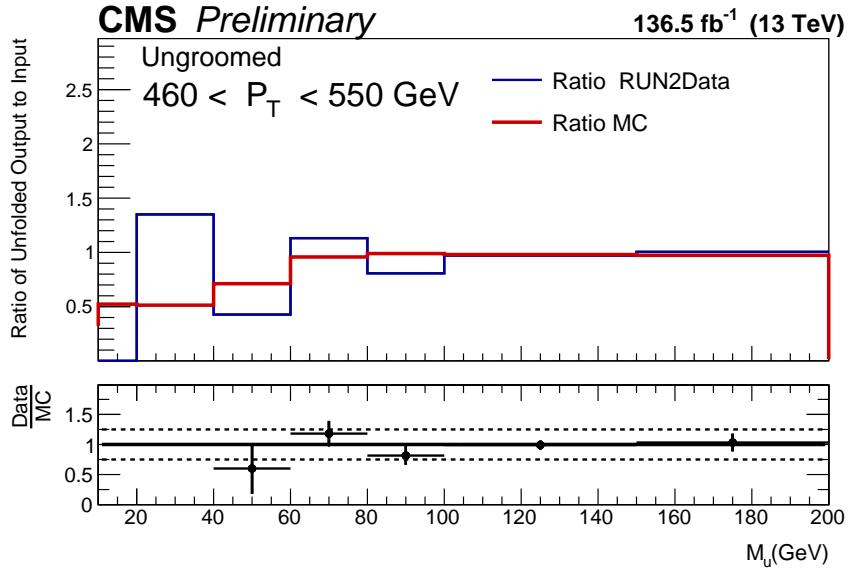


Figure 6.55: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 460-550 GeV.

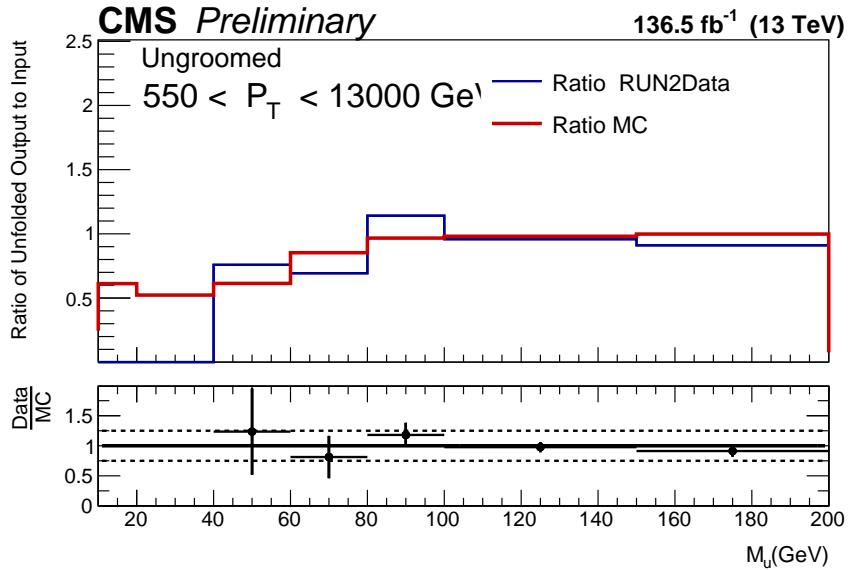


Figure 6.56: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 550-13000 GeV.

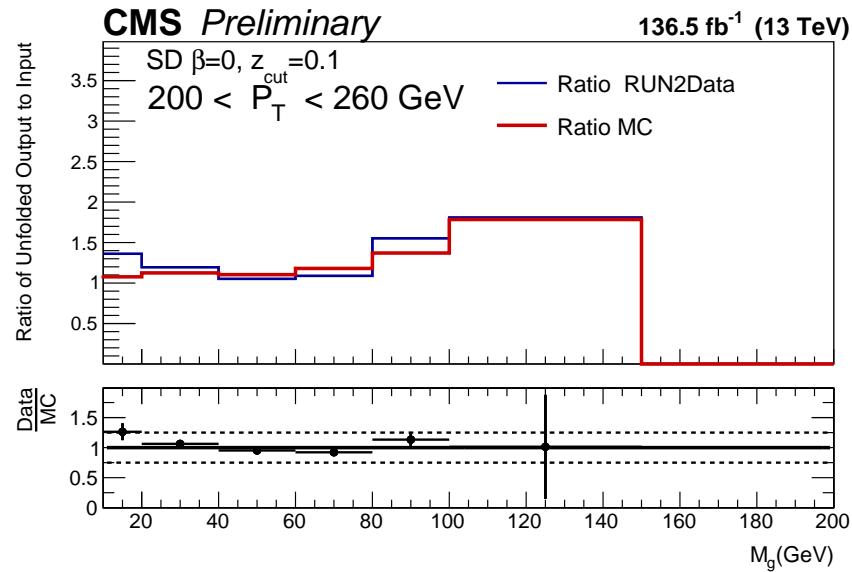


Figure 6.57: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 200-260 GeV.

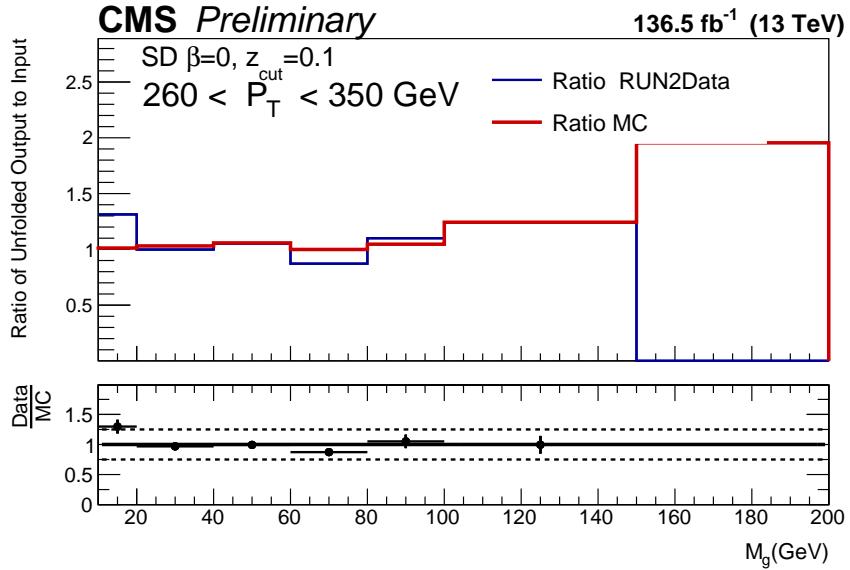


Figure 6.58: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 260-350 GeV.

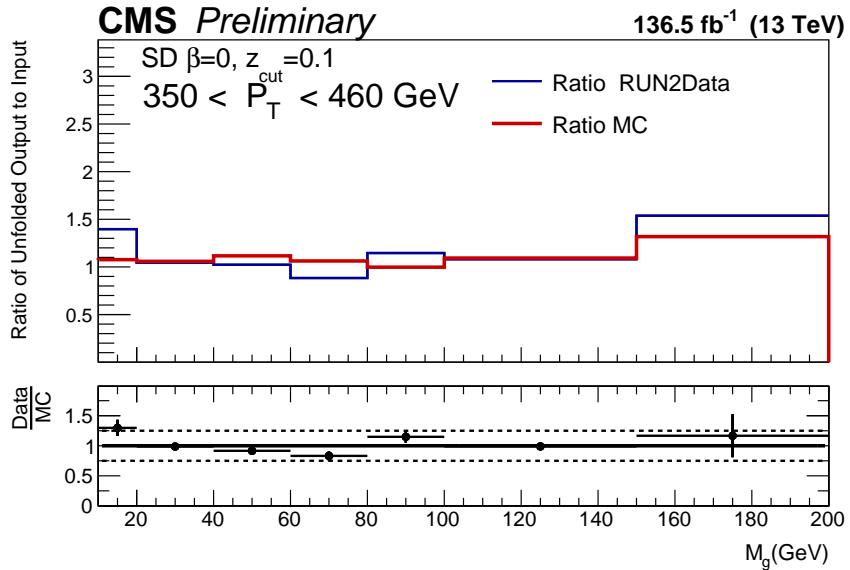


Figure 6.59: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 350-460 GeV.

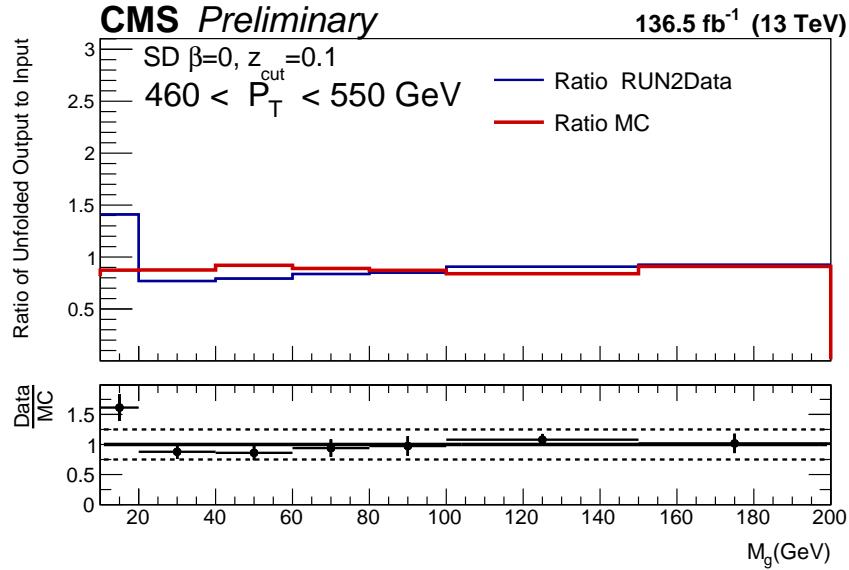


Figure 6.60: Ratio of unfolded over raw data and MC for groomed jets, p_T 460-550 GeV.

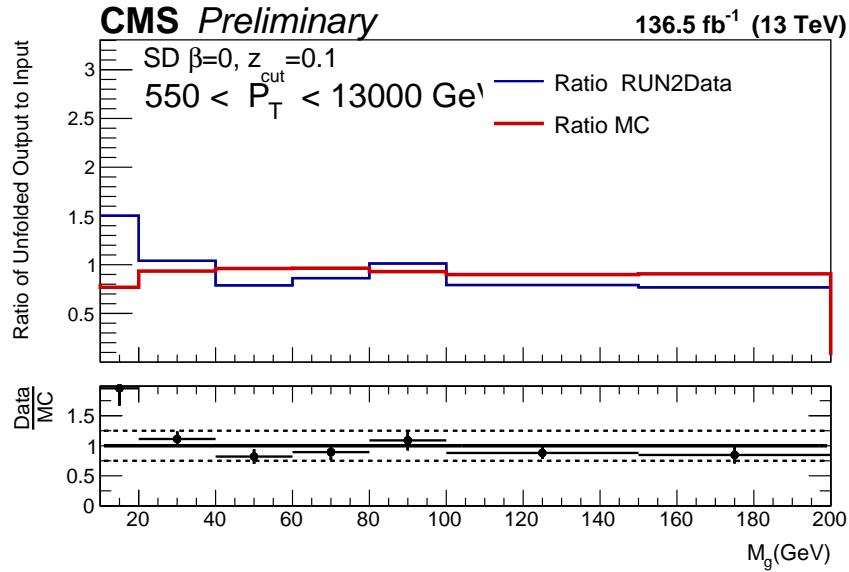


Figure 6.61: Ratio of unfolded over raw data and MC for groomed jets, p_T 550-13000 GeV.

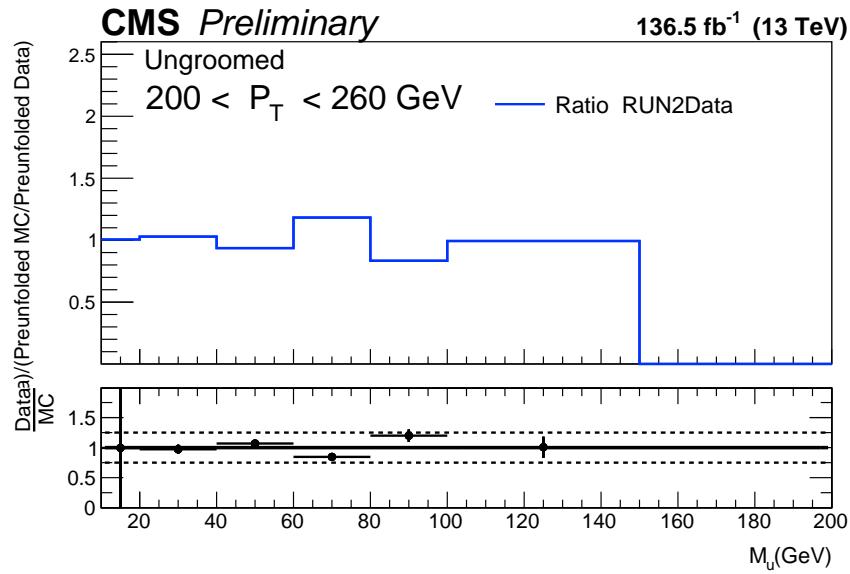


Figure 6.62: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 200-260 GeV.

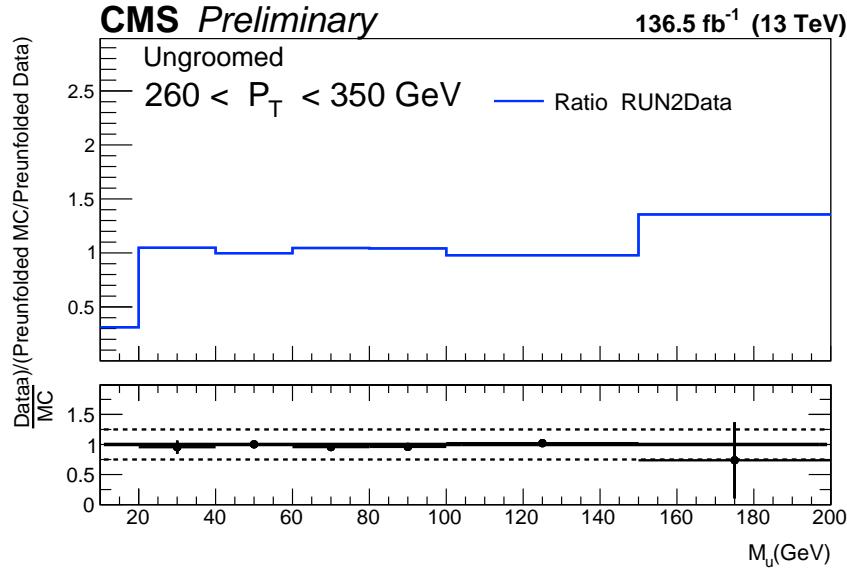


Figure 6.63: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 260-350 GeV.

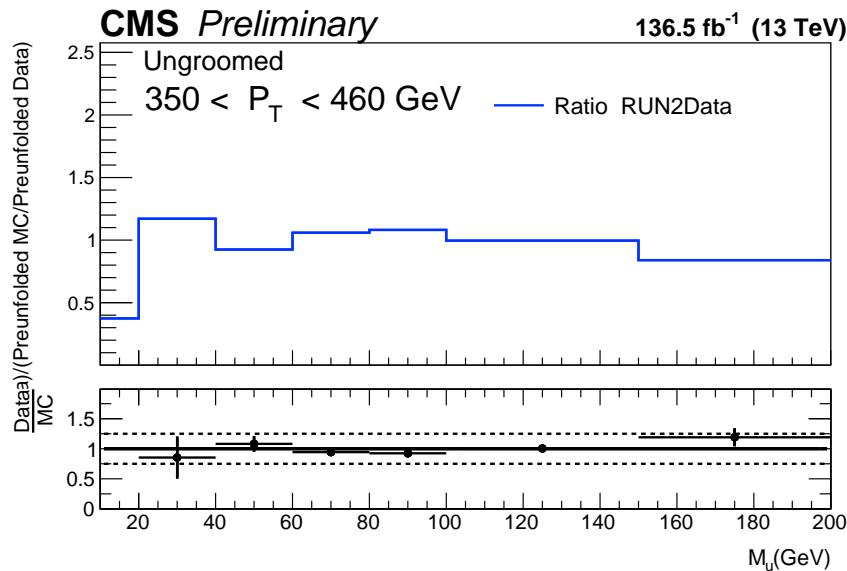


Figure 6.64: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 350-460 GeV.

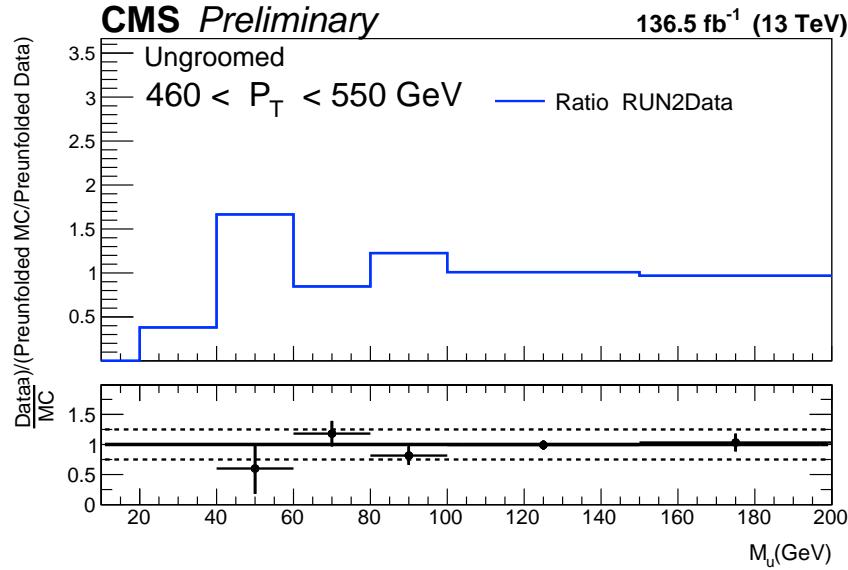


Figure 6.65: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 460-550 GeV.

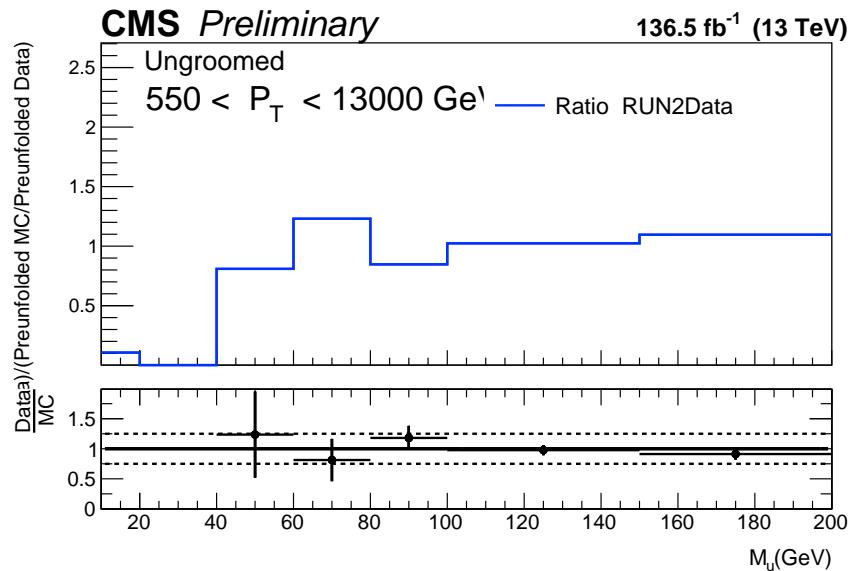


Figure 6.66: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 550-13000 GeV.

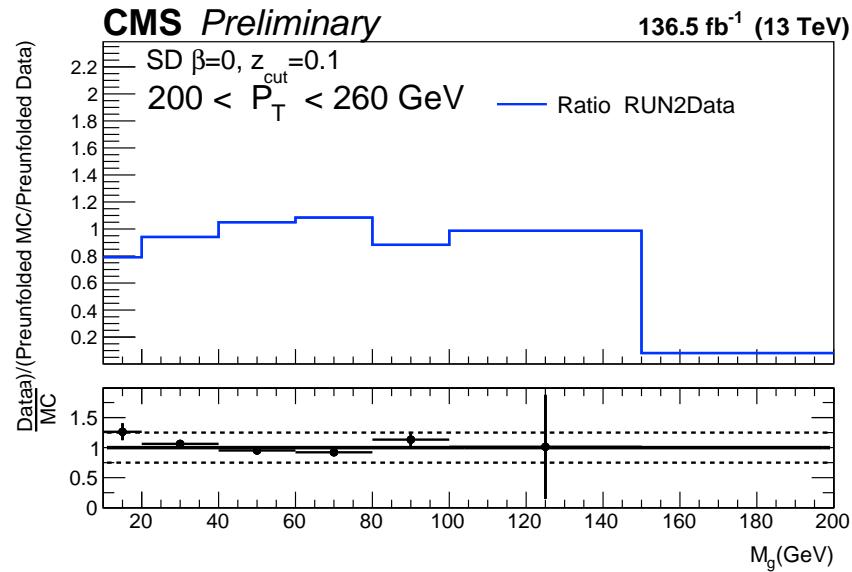


Figure 6.67: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_{T} 200-260 GeV.

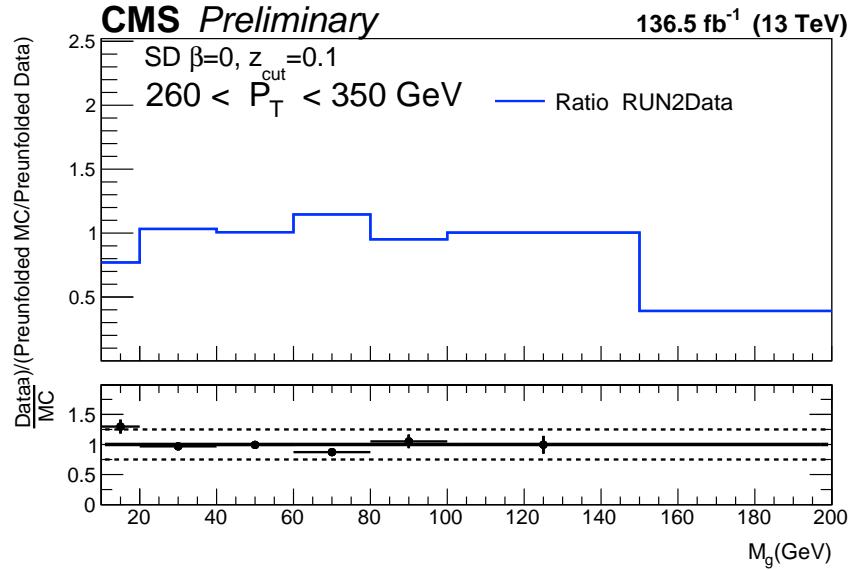


Figure 6.68: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 260-350 GeV.

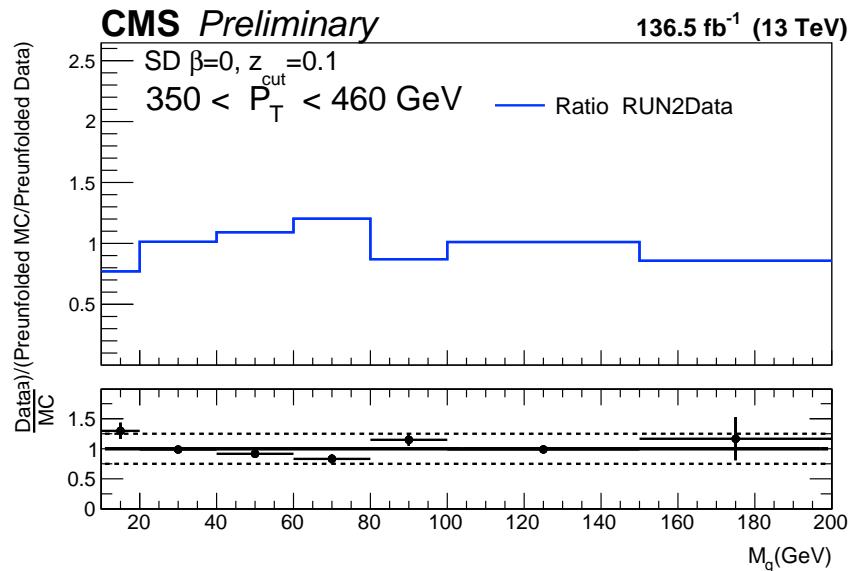


Figure 6.69: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 350-460 GeV.

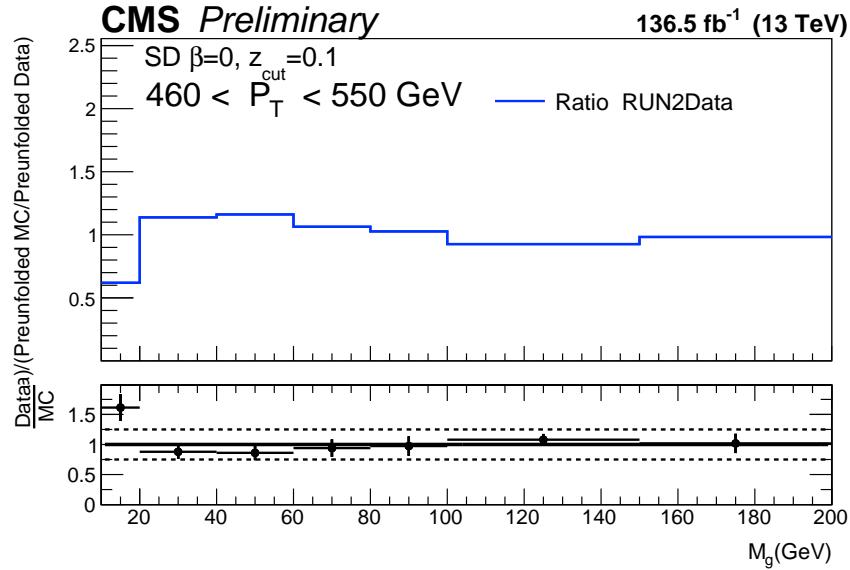


Figure 6.70: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 460-550 GeV.

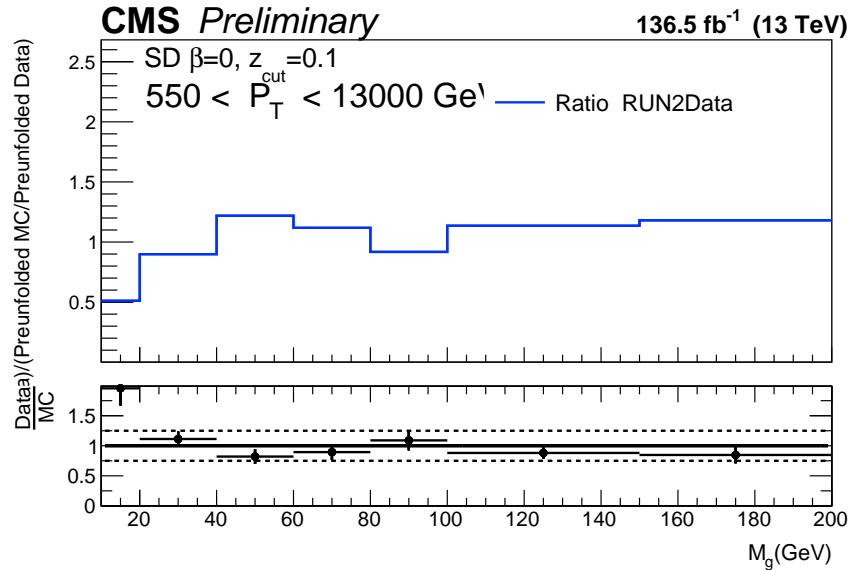


Figure 6.71: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 550-13000 GeV.

6.10 Summary

In conclusion, we have presented a differential jet cross section measured in $Z + \text{Jet}$ events in bins of the ungroomed jet p_T in conjunction with the ungroomed and groomed jet mass (as well as dimensionless mass) using the “soft drop” (a.k.a. “modified mass drop tagger”) algorithm with 9 different combinations of parameters. The results are presented as the normalized cross section, normalized per reconstructed jet p_T bin. Overall leading-order MC simulation agrees reasonably well with the data within our uncertainties. Agreement below the Sudakov peak is slightly worse than above. The application of a grooming algorithm improves the overall precision, with larger improvement at low jet masses. This analysis improves over previous iterations by using various parameter values for the “soft drop” jet grooming algorithm, as well as by including an additional unfolding in both transverse momentum and dimensionless mass, as was done by the ATLAS collaboration [61].

.1 2016 data results

This Appendix shows the distributions from the “Detector Response” through the “Results” sections of the main analysis note with only the 2016 data rather than the full Run 2 statistics seen in the main body of the note.

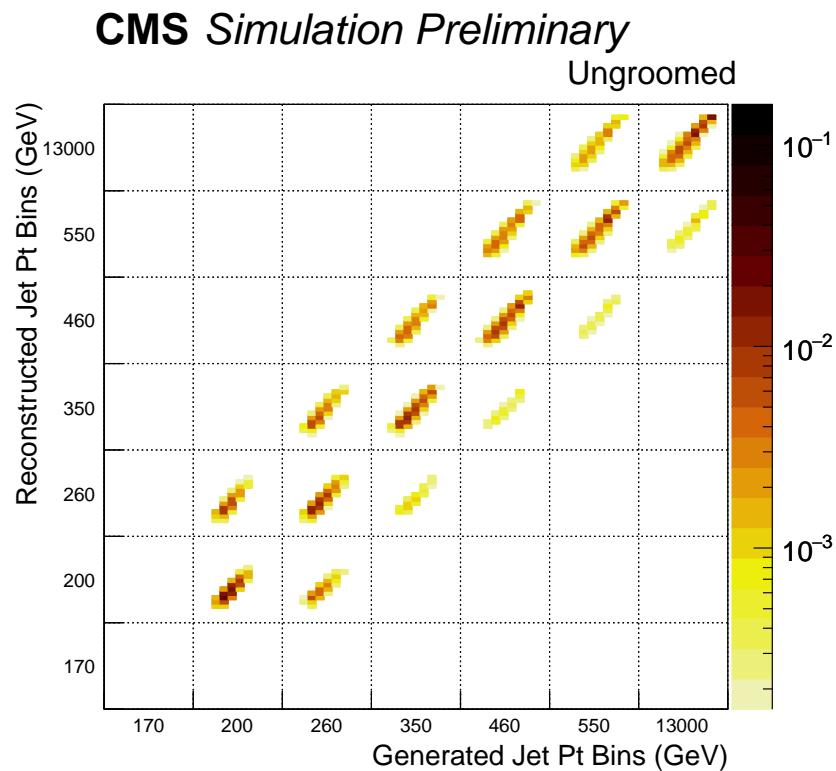


Figure 72: Two-dimensional response matrix for ungroomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). While the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

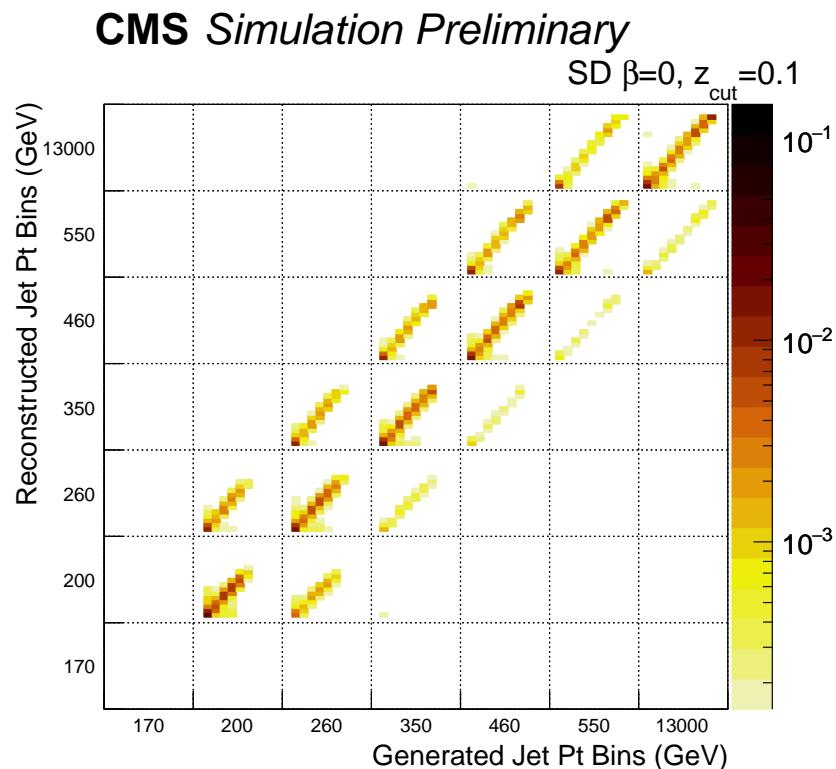


Figure 73: Two-dimensional response matrix for groomed jets $\beta = 0, z_{cut} = 0.1$. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV) while the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

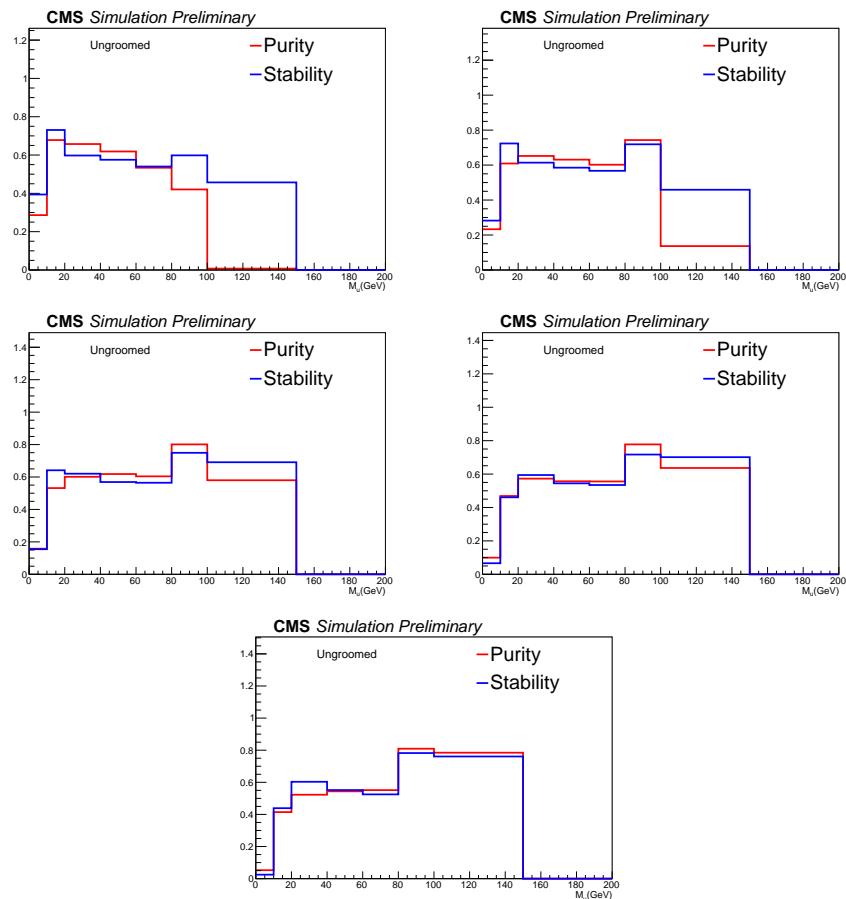


Figure 74: Purity and stability for ungroomed jets in various p_T bins.

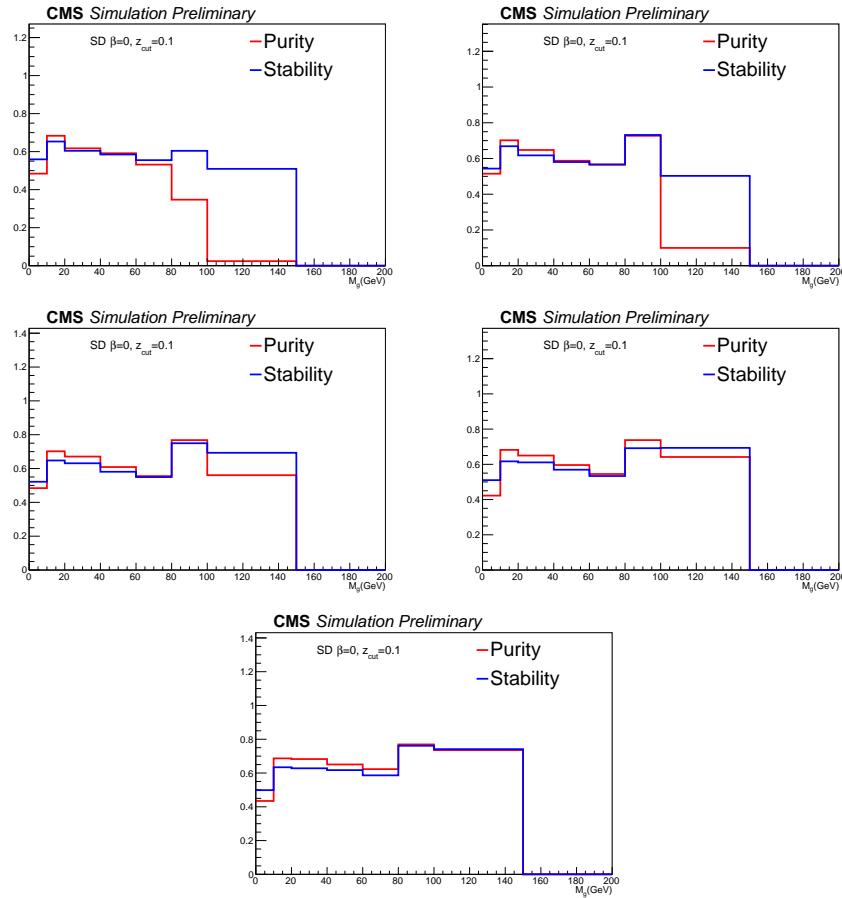


Figure 75: Purity and stability for groomed jets, where the soft-drop criterion was applied with $\beta = 0$ and $z_{cut} = 0.1$, in various p_T bins.

Figure-76 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

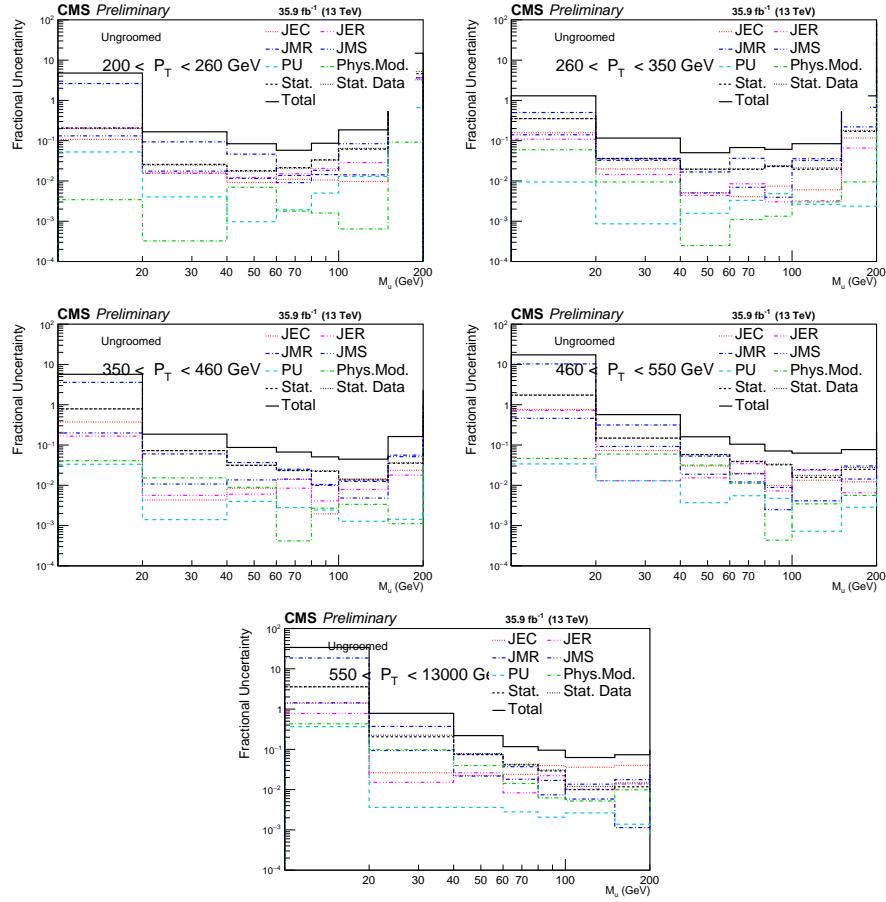


Figure 76: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw jet mass.

Figure-78 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

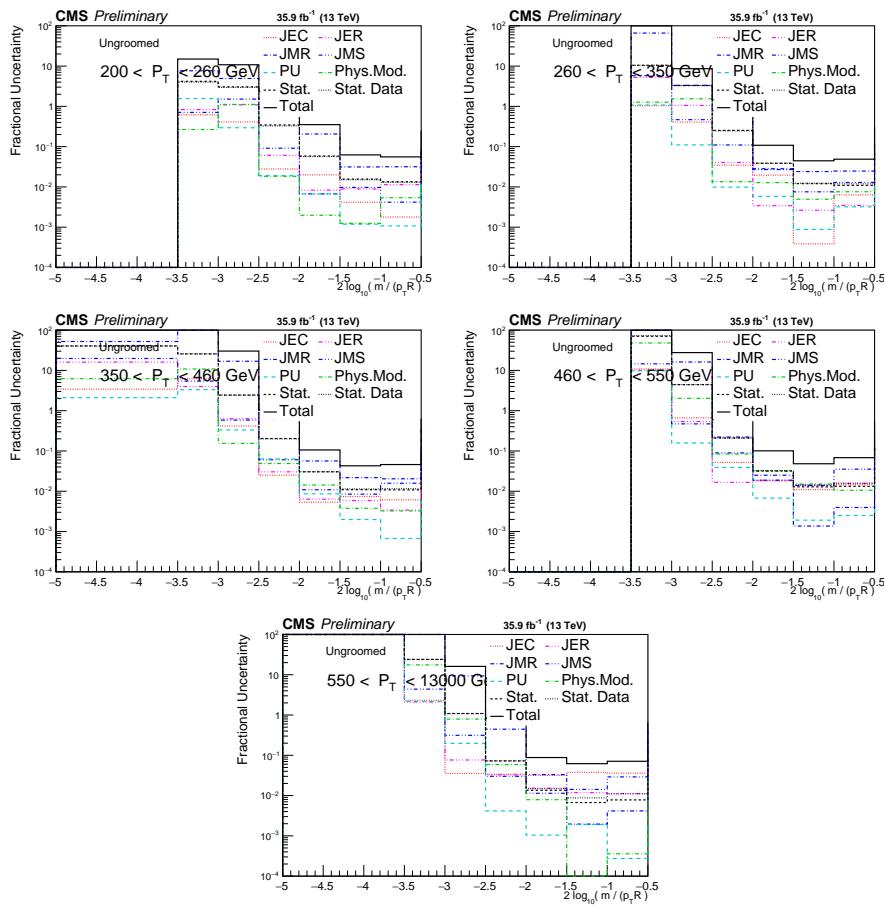


Figure 77: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw dimensionless jet mass.

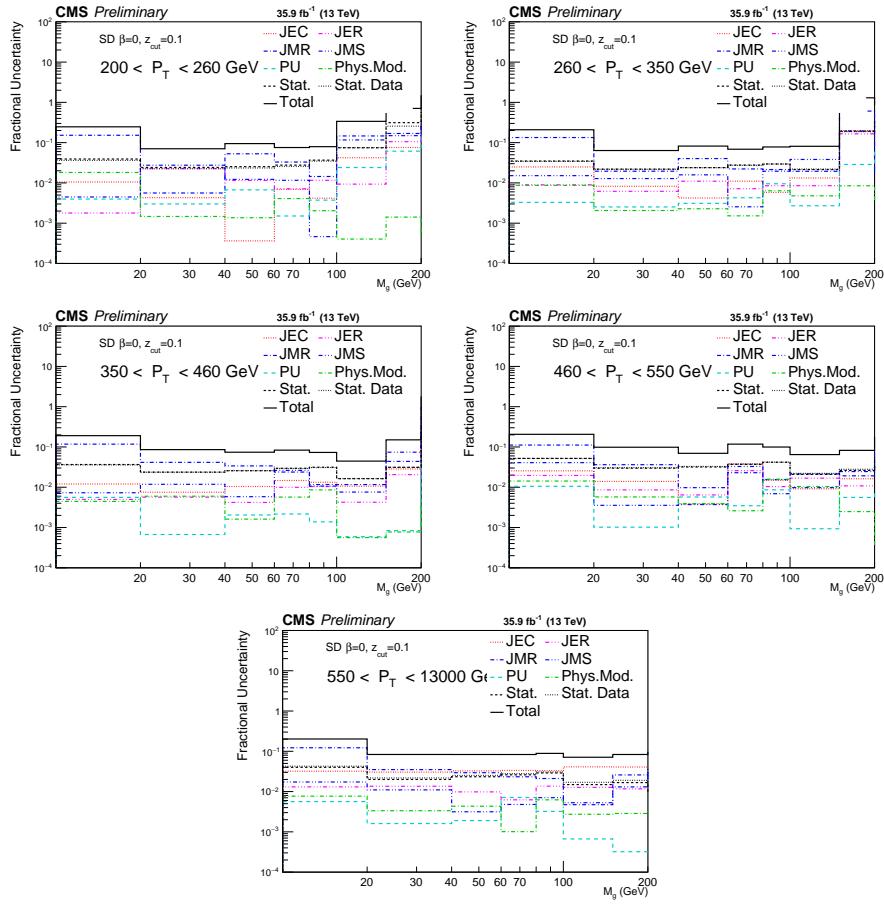


Figure 78: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

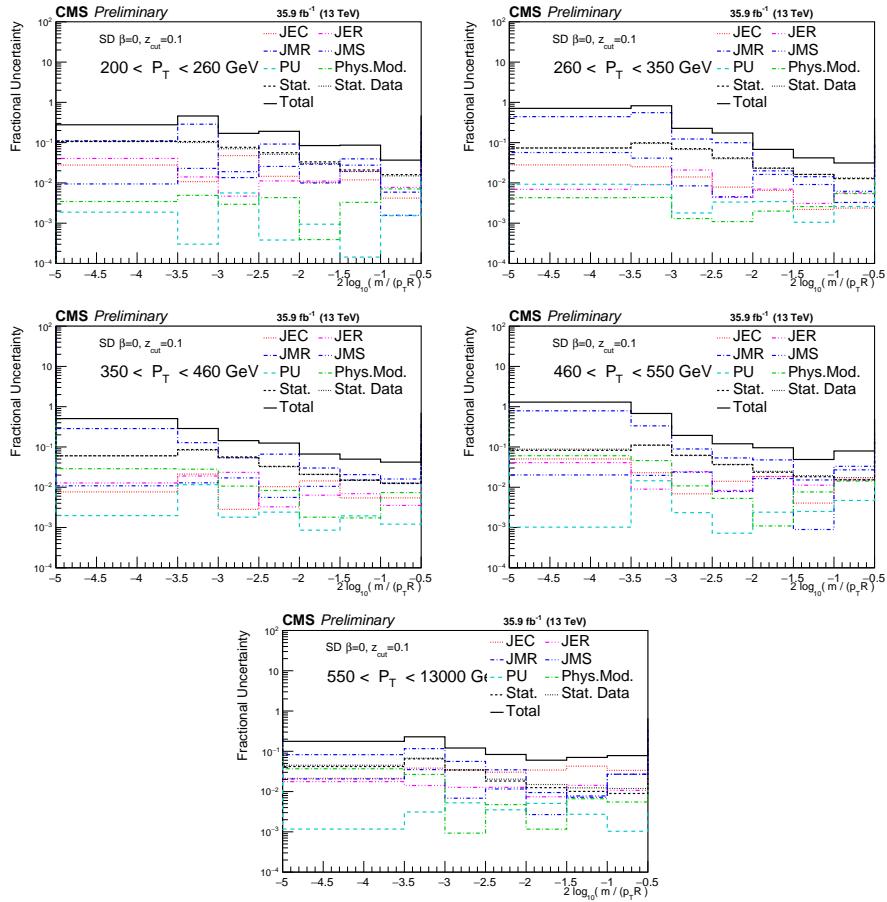


Figure 79: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

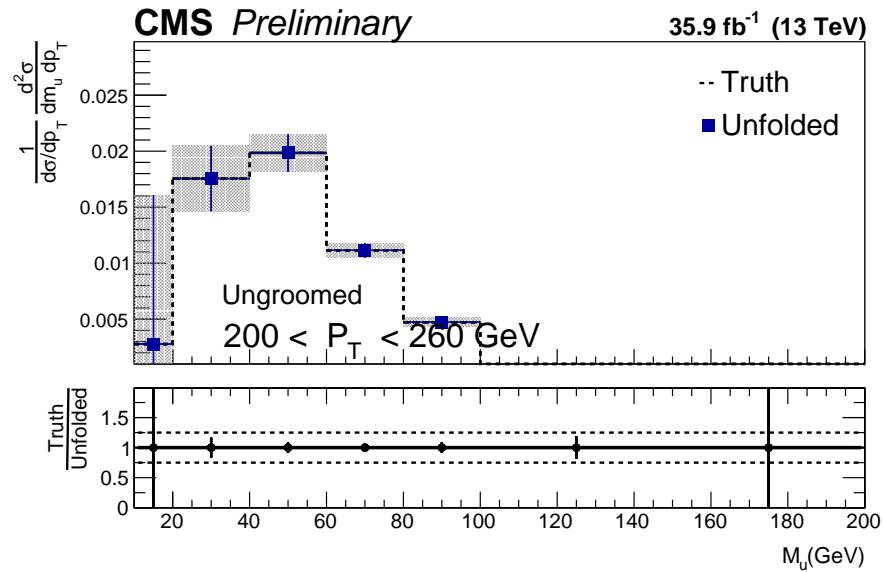


Figure 80: Closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

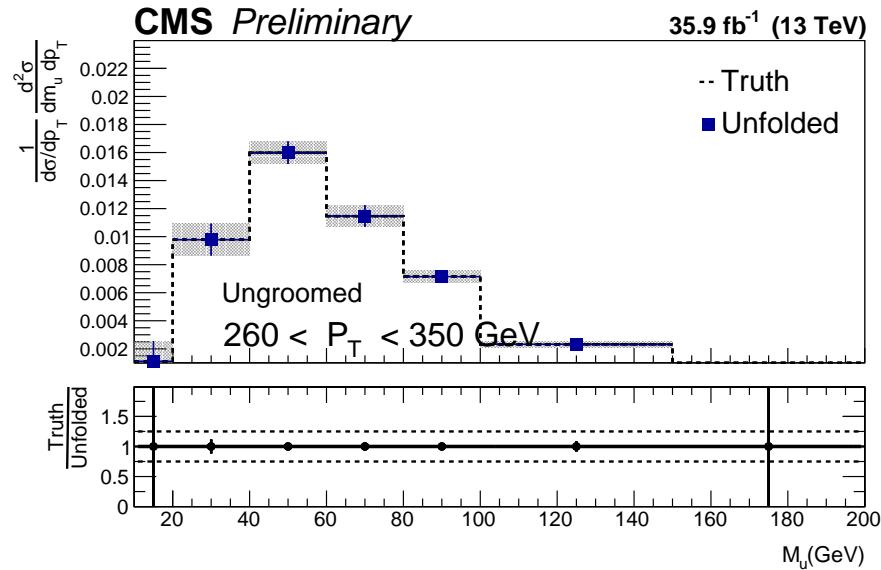


Figure 81: Closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

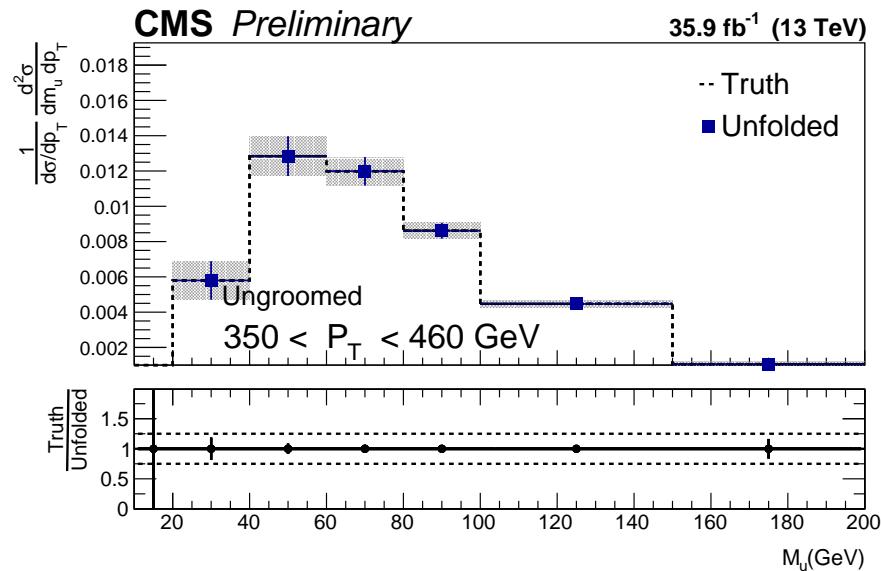


Figure 82: Closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

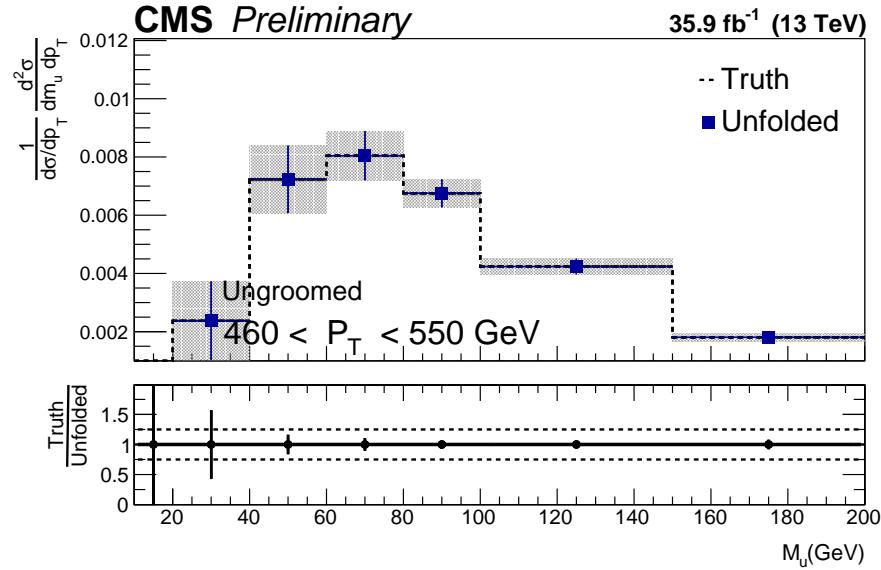


Figure 83: Closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

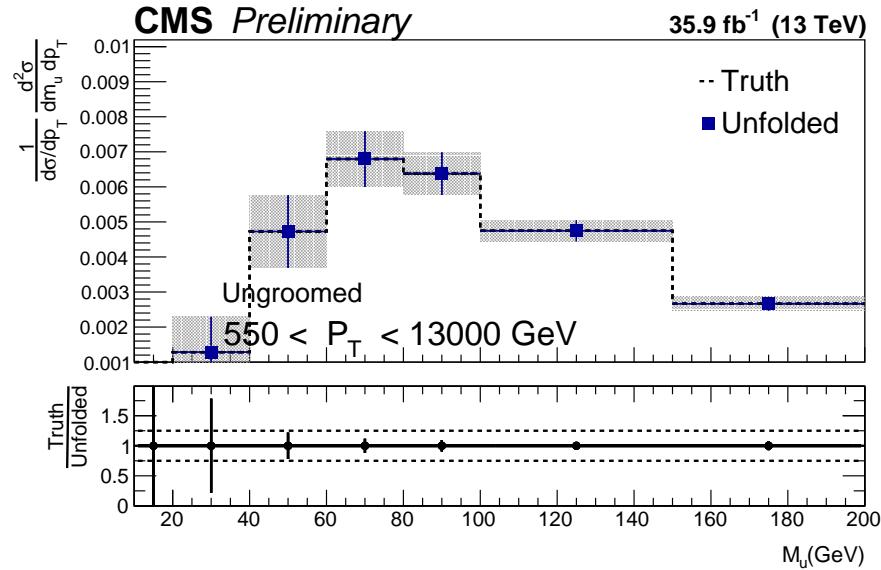


Figure 84: Closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

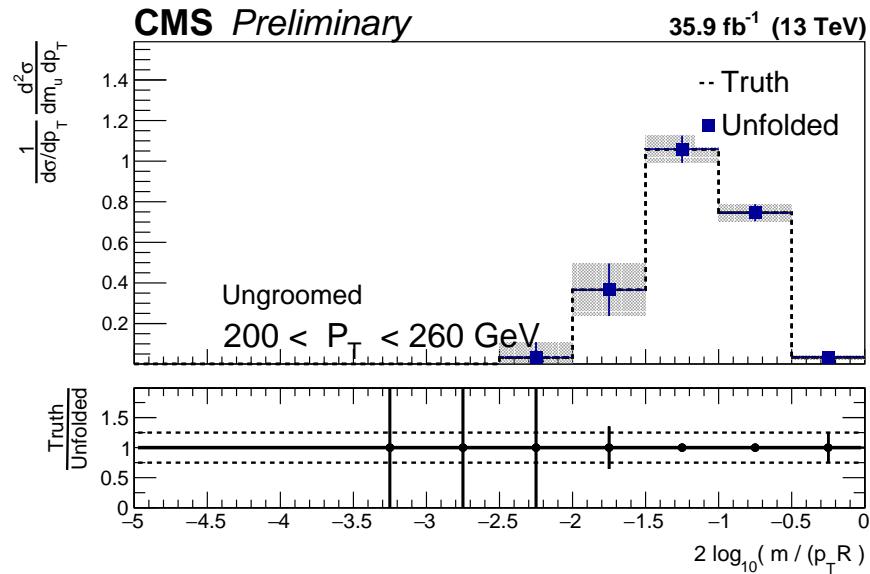


Figure 85: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

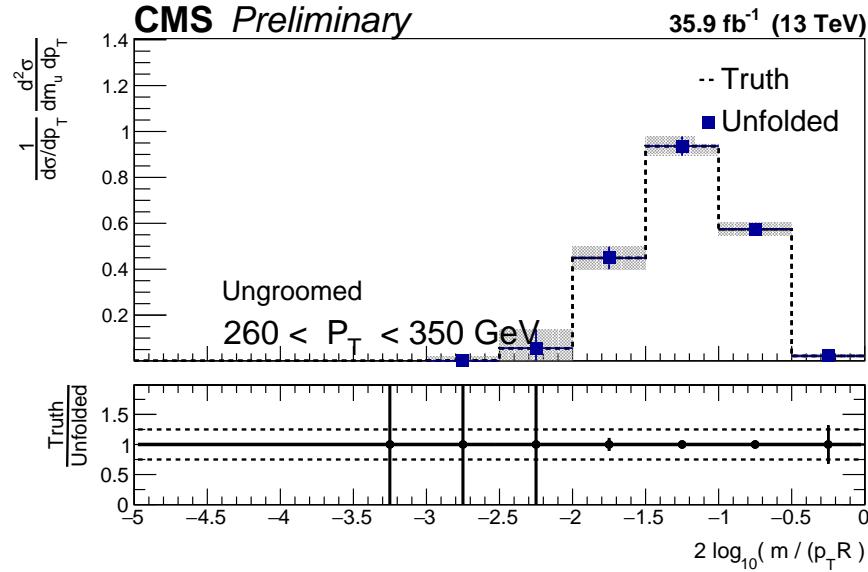


Figure 86: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

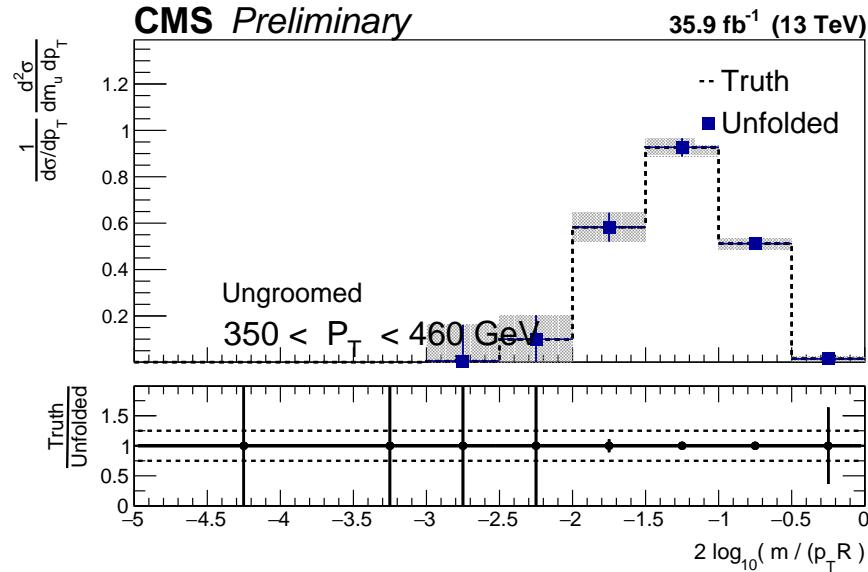


Figure 87: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

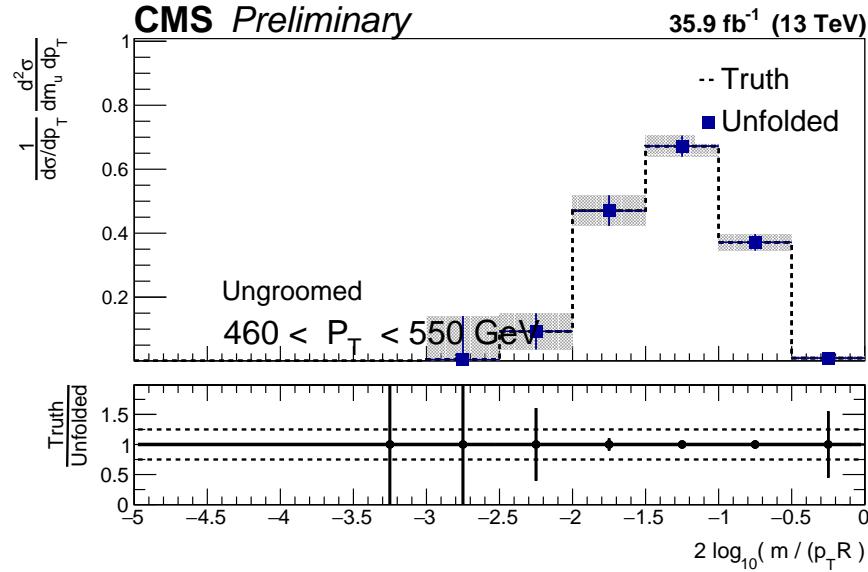


Figure 88: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

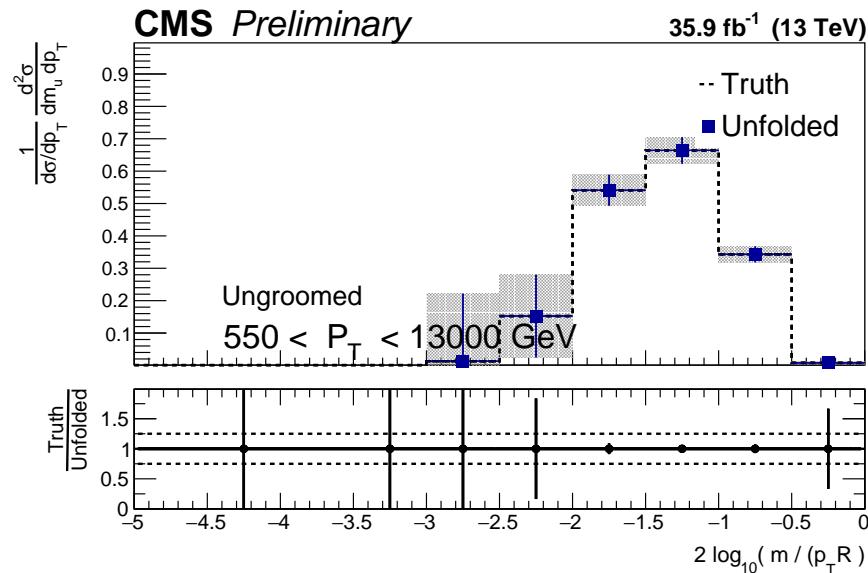


Figure 89: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

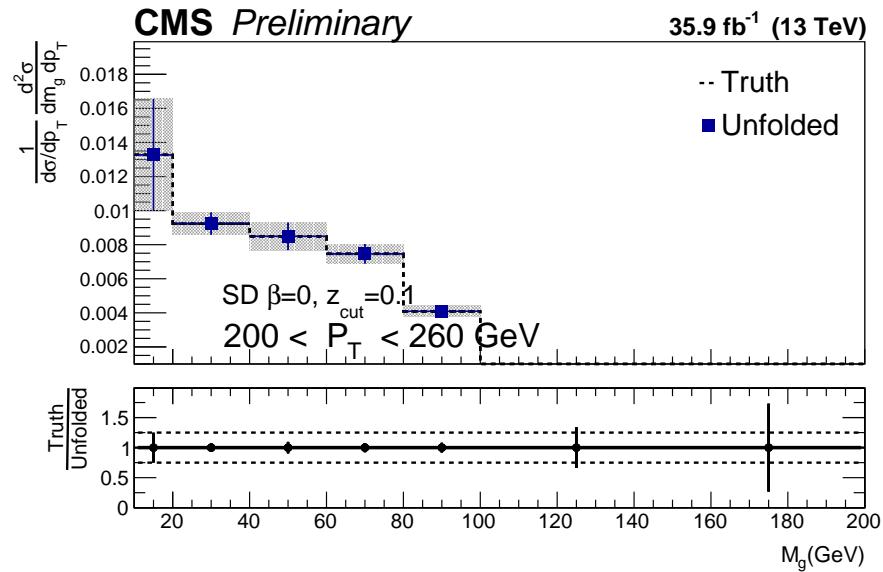


Figure 90: Closure test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

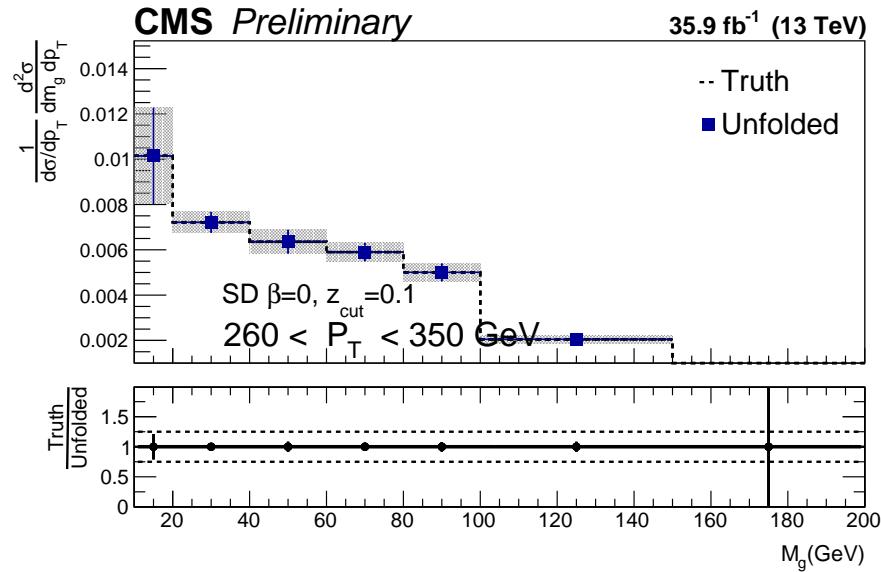


Figure 91: Closure test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

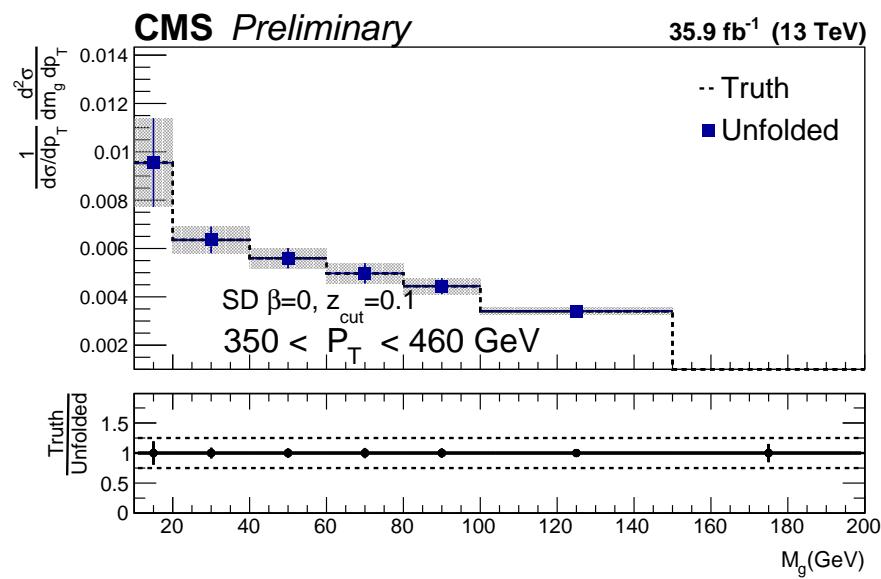


Figure 92: Closure test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

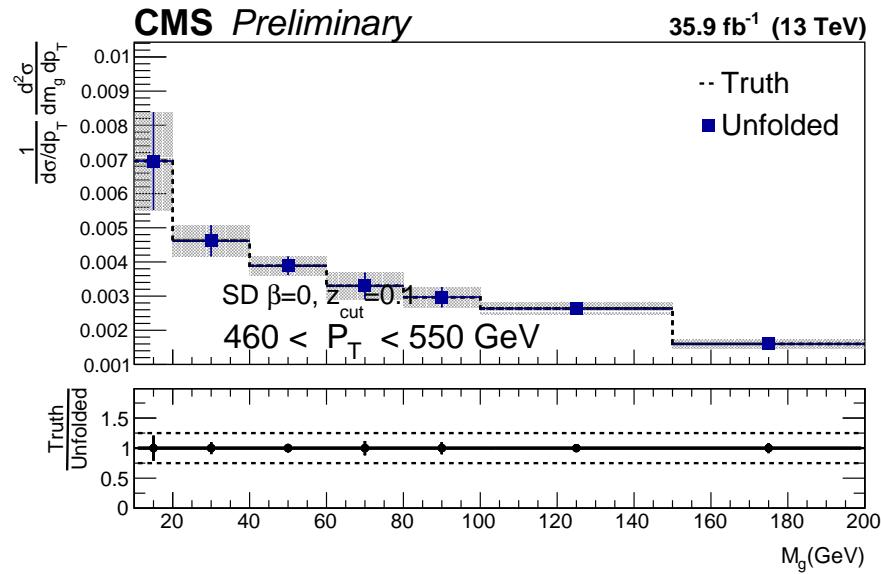


Figure 93: Closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

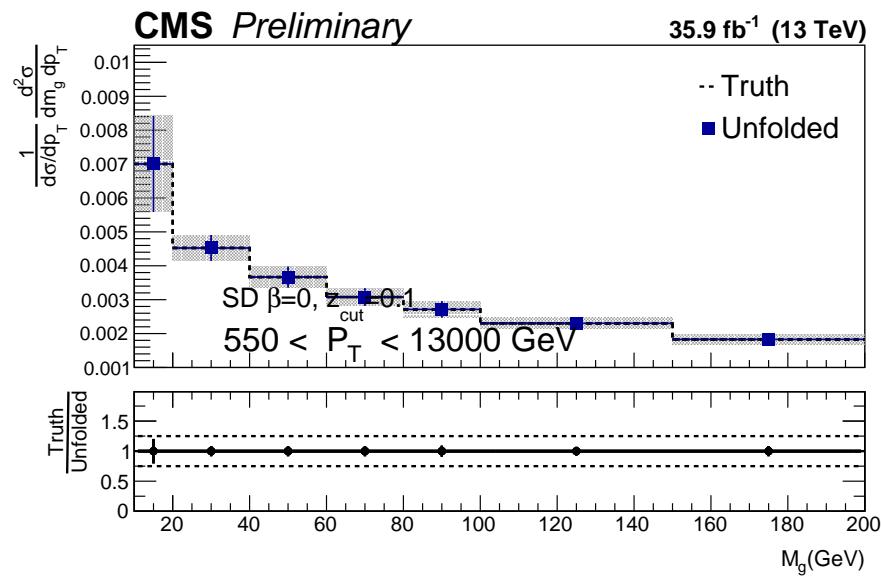


Figure 94: Closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

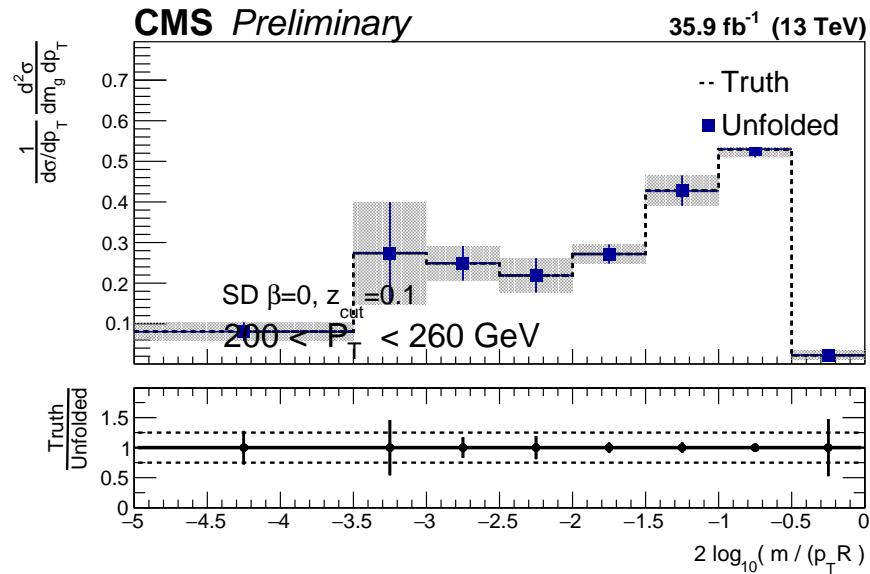


Figure 95: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

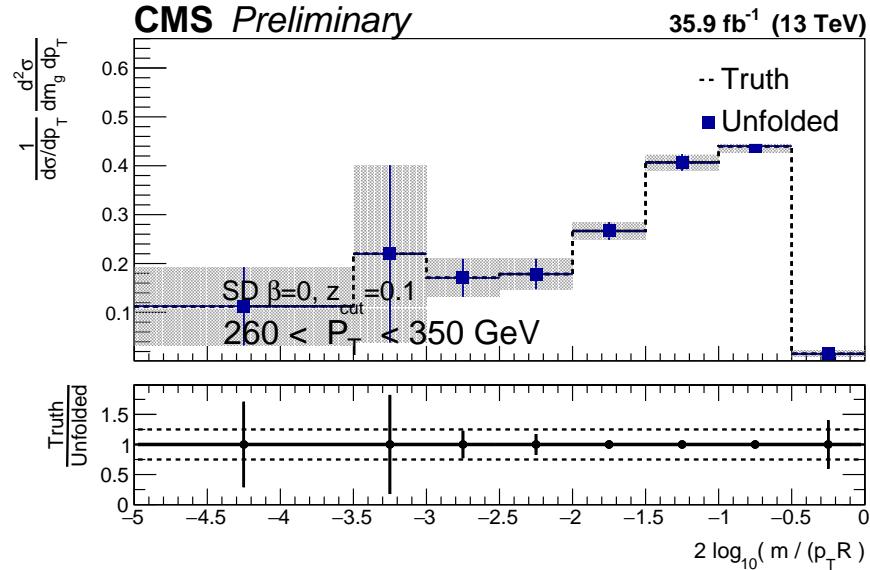


Figure 96: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

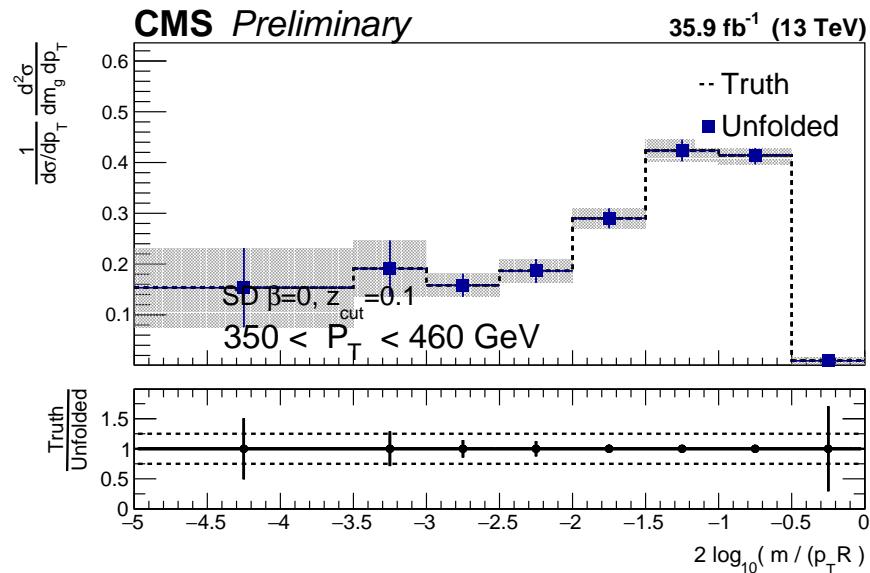


Figure 97: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

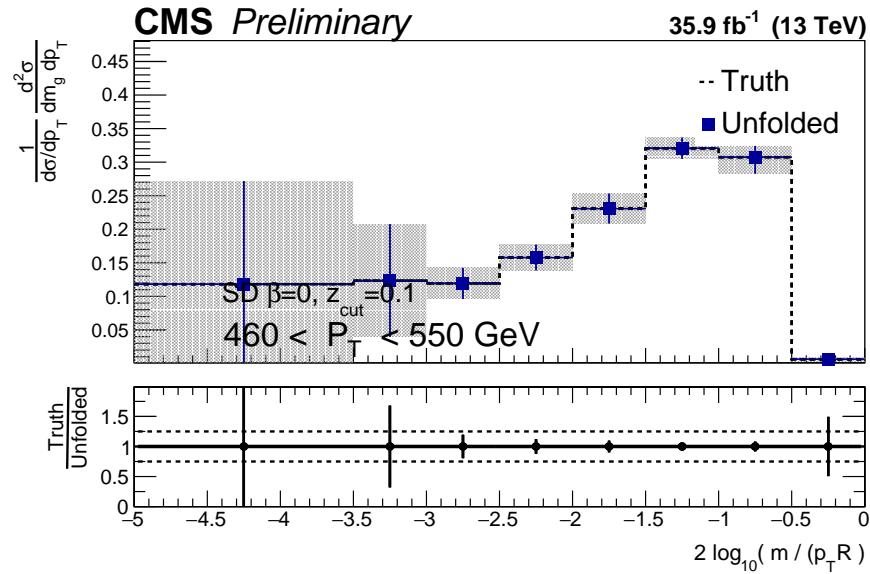


Figure 98: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

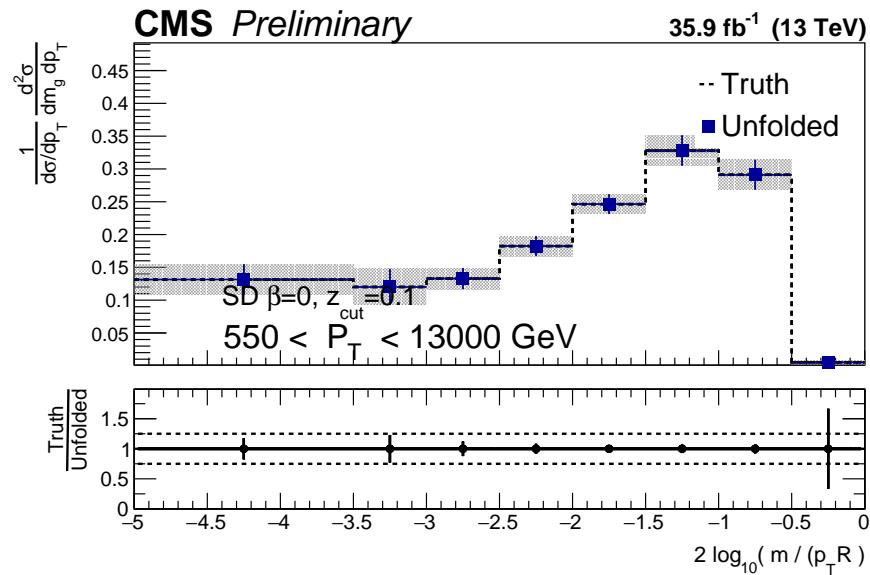


Figure 99: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

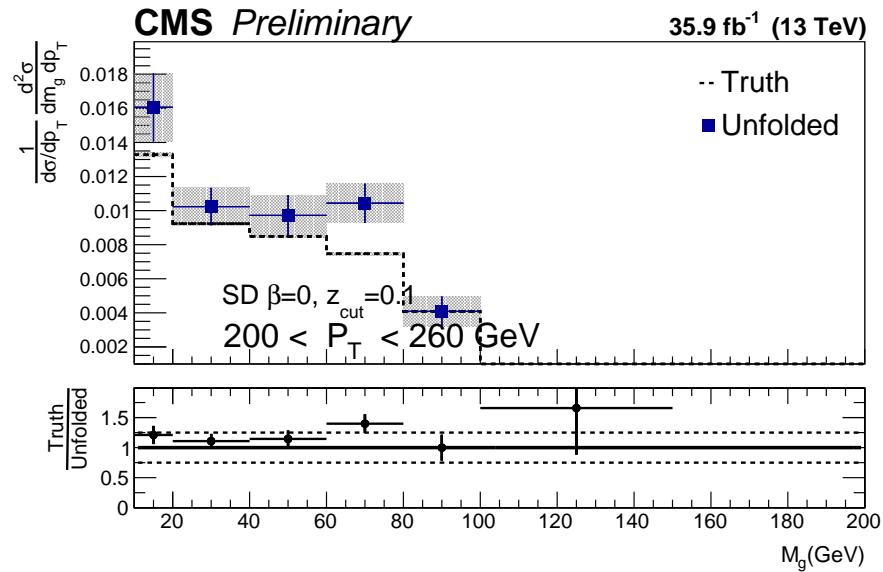


Figure 100: Physical bias test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

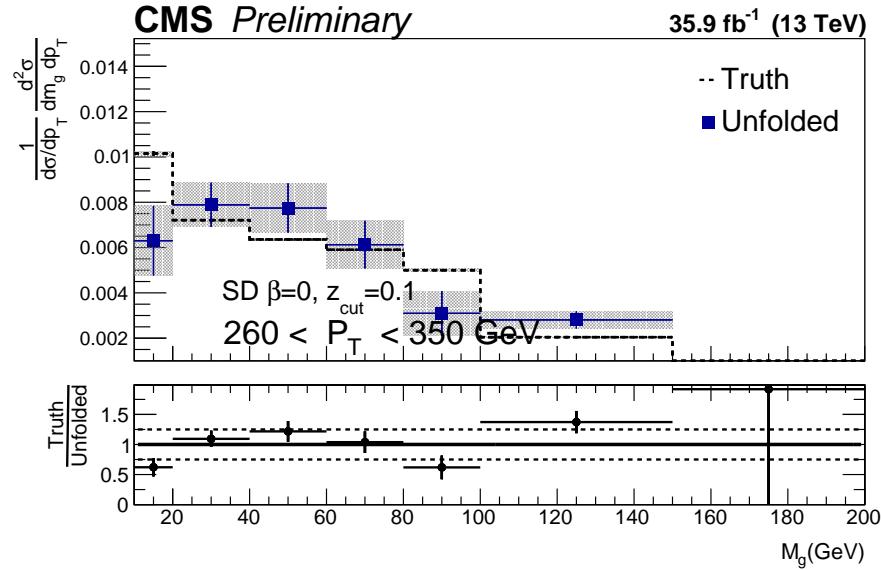


Figure 101: Physical bias test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

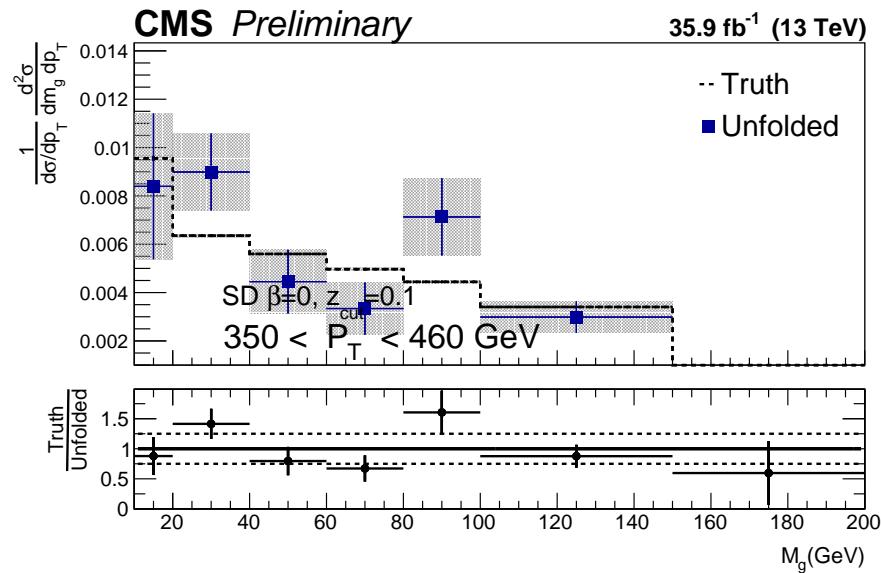


Figure 102: Physical bias test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

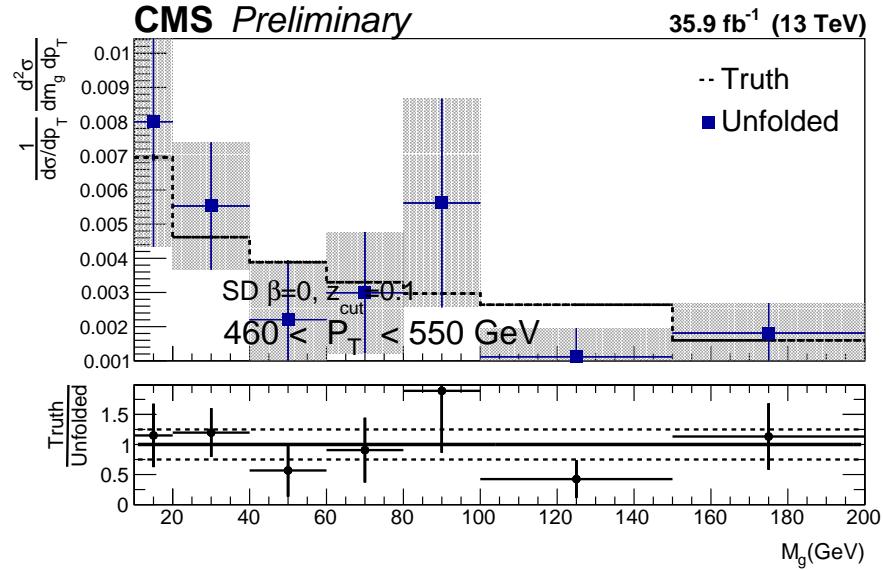


Figure 103: Physical bias test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

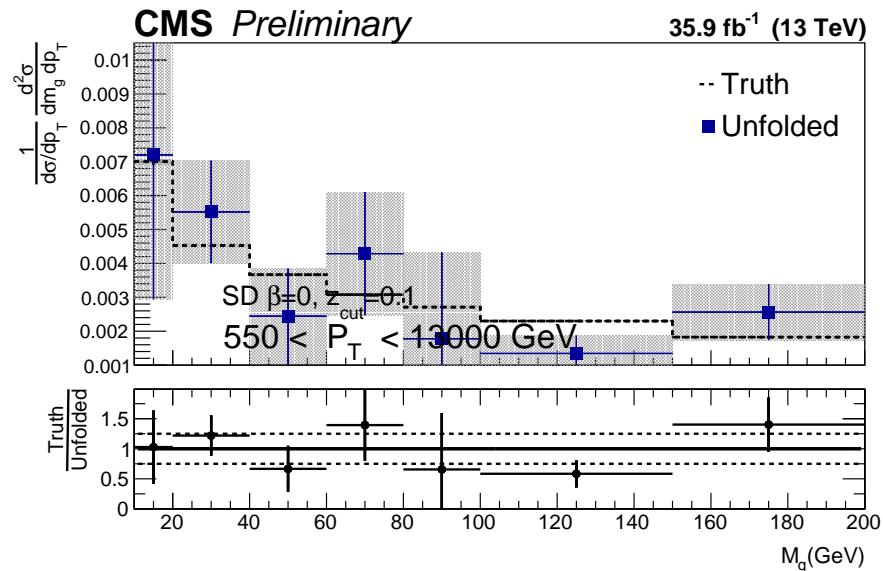


Figure 104: Physical bias test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

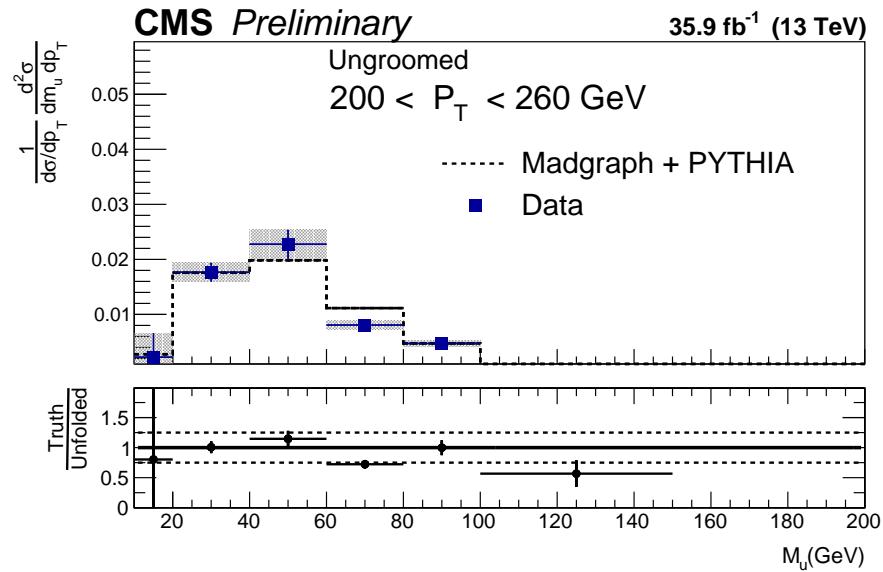


Figure 105: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 200-260 GeV.

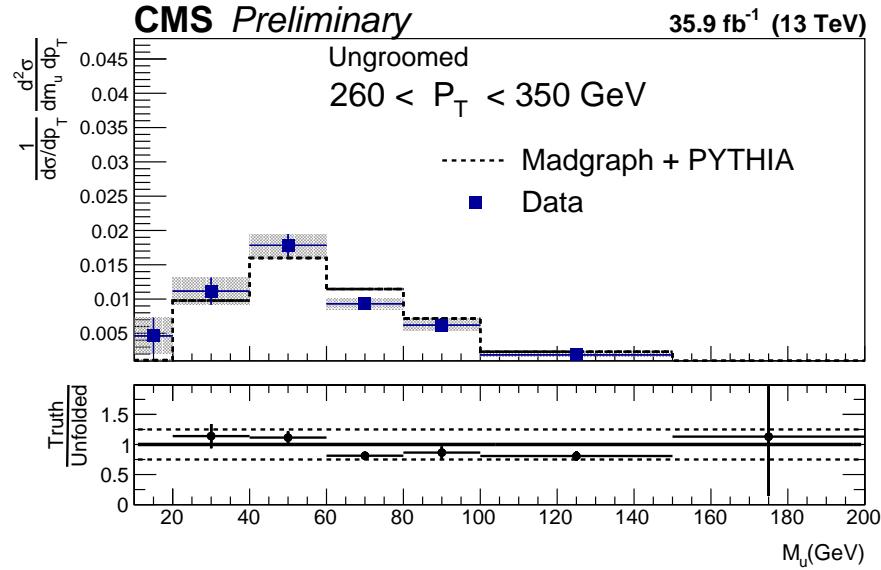


Figure 106: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 260–350 GeV.

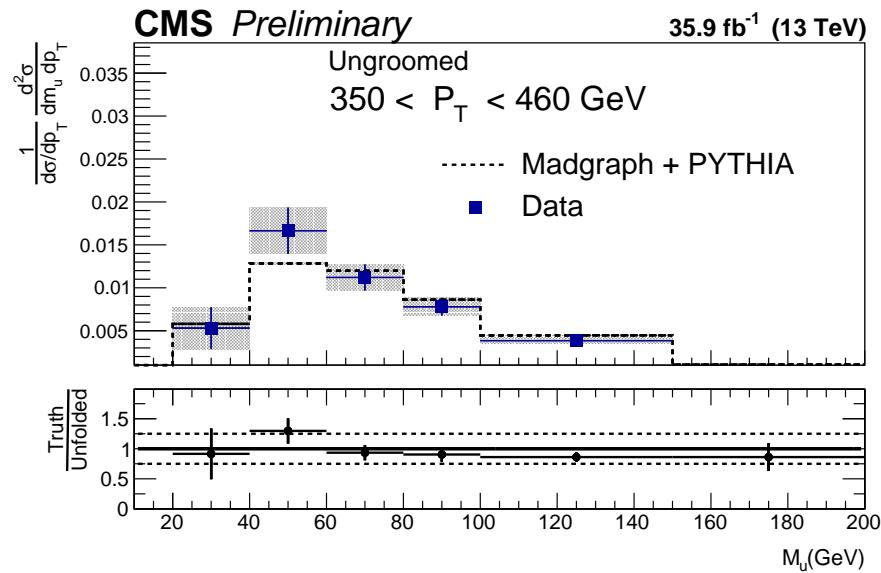


Figure 107: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 350–460 GeV.

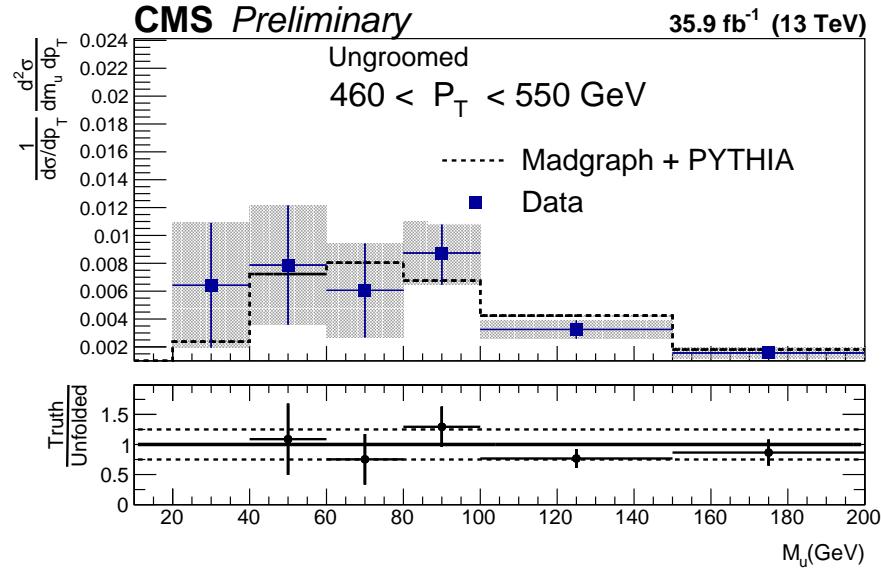


Figure 108: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 460-550 GeV.

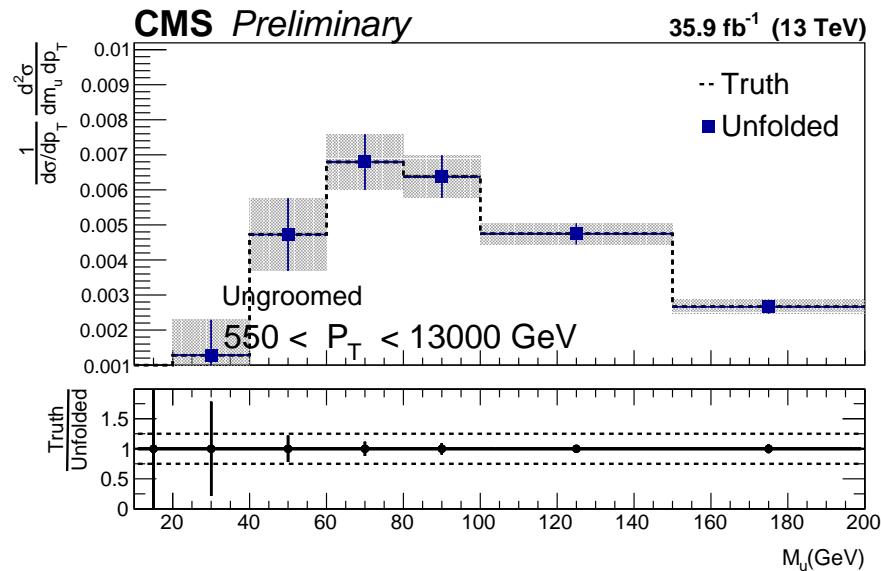


Figure 109: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 550-13000 GeV.

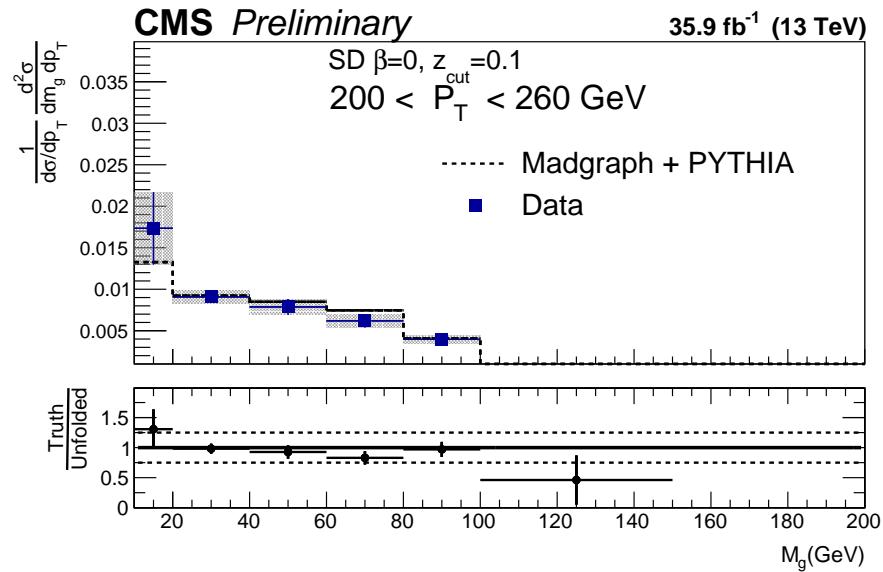


Figure 110: Normalized cross section results with respect to jet mass for groomed jets, p_T 200-260 GeV.

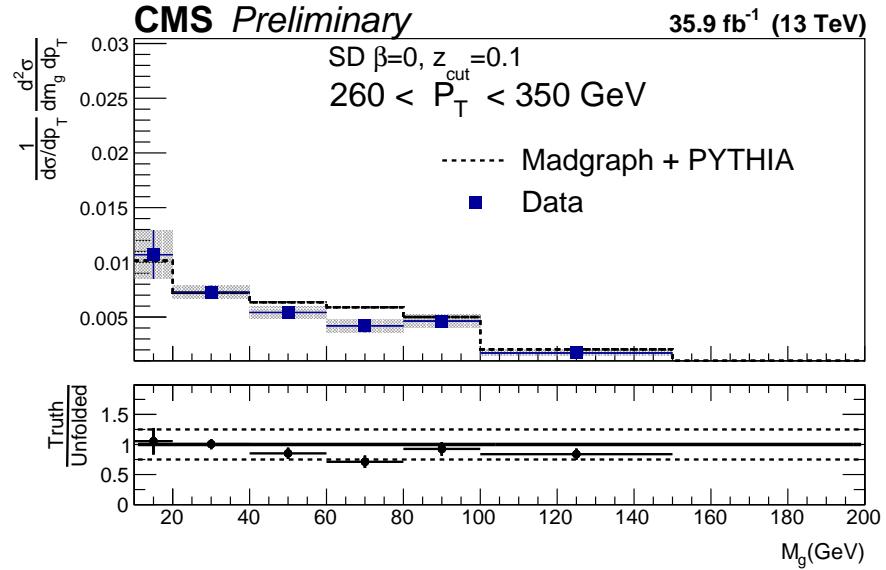


Figure 111: Normalized cross section results with respect to jet mass for groomed jets, p_T 260-350 GeV.

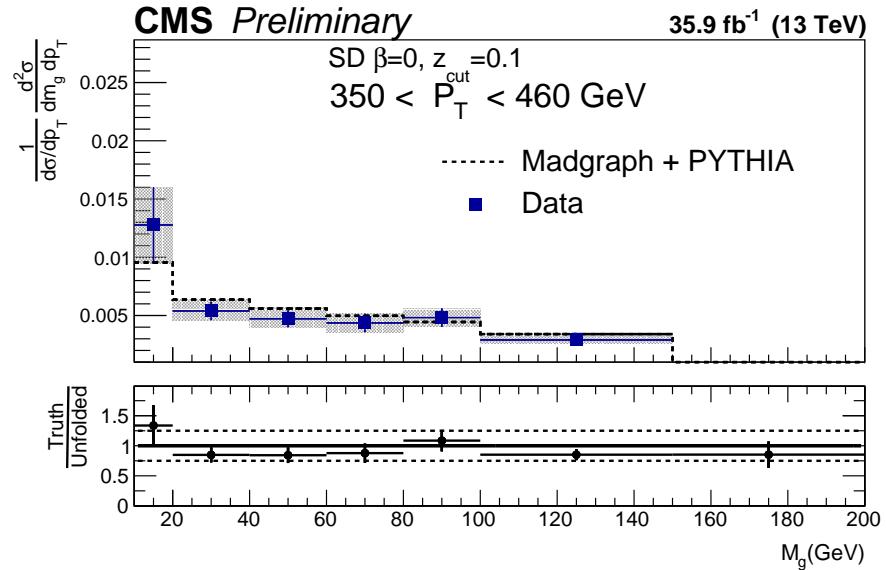


Figure 112: Normalized cross section results with respect to jet mass for groomed jets, p_T 350-460 GeV.

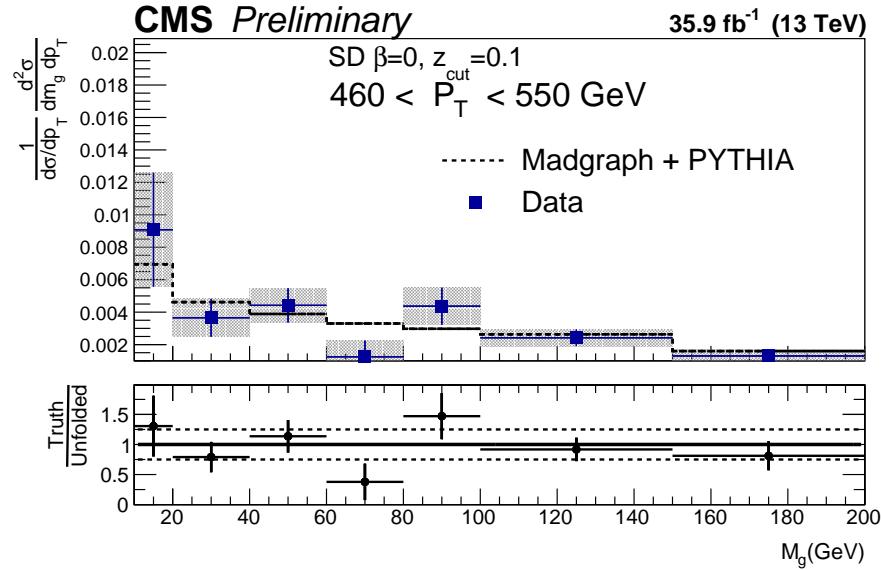


Figure 113: Normalized cross section results with respect to jet mass for groomed jets, p_T 460-550 GeV.

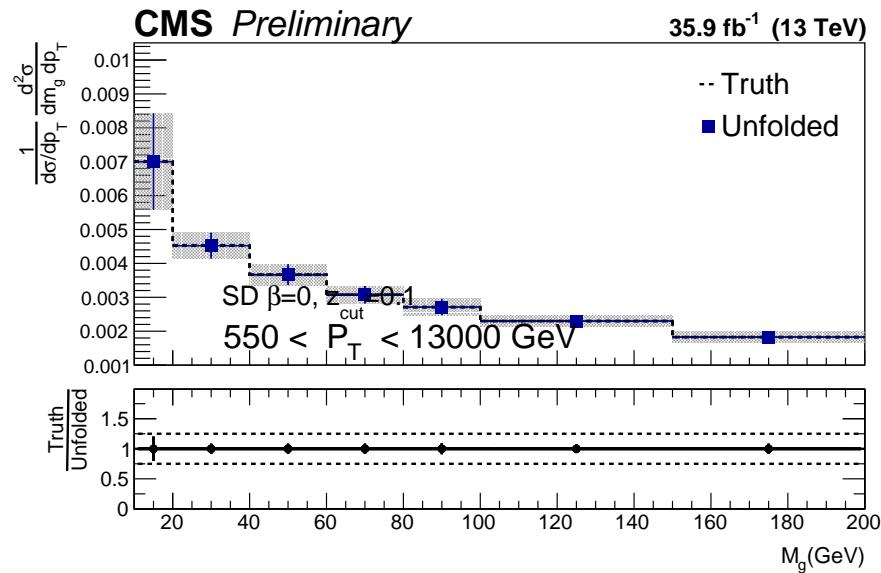


Figure 114: Normalized cross section results with respect to jet mass for groomed jets, p_T 550-13000 GeV.

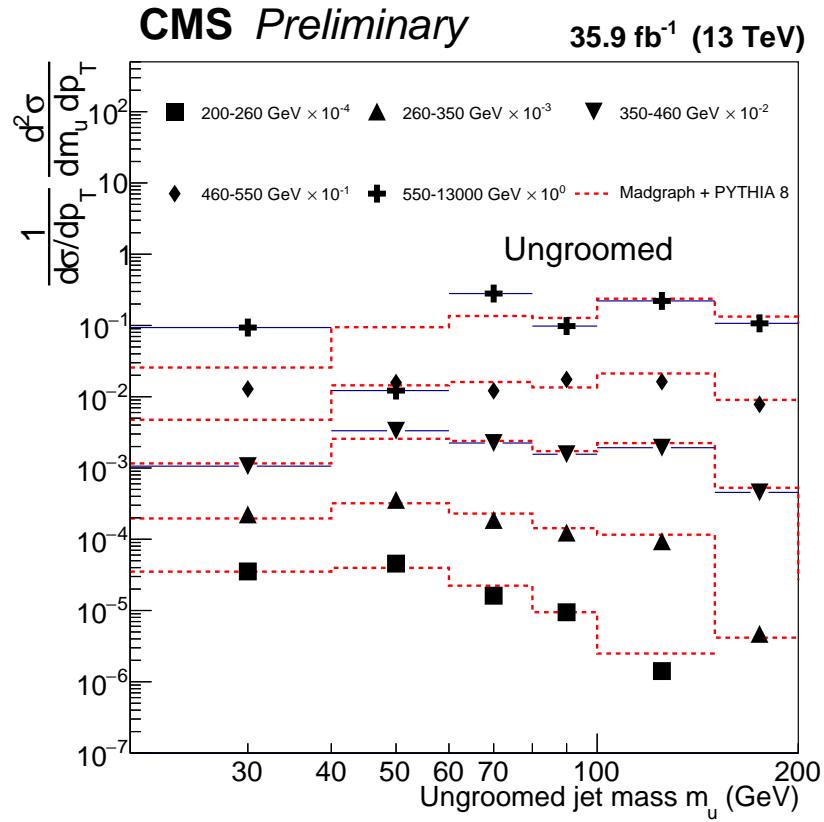


Figure 115: Results for ungroomed reconstructed unfolding with jet mass.

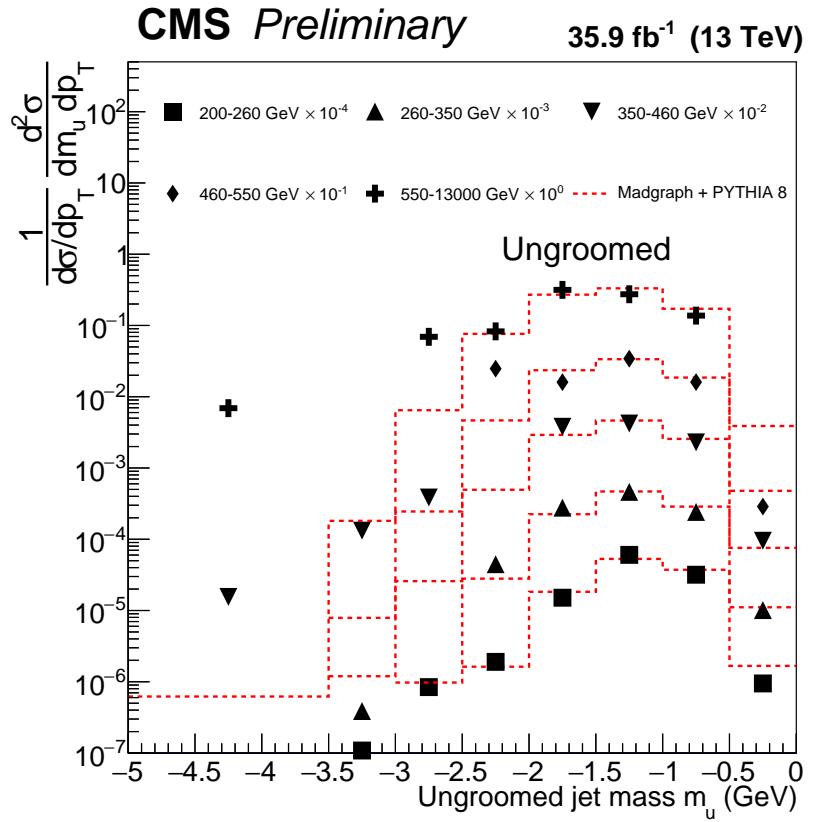


Figure 116: Results for ungroomed reconstructed unfolding with jet dimensionless mass.

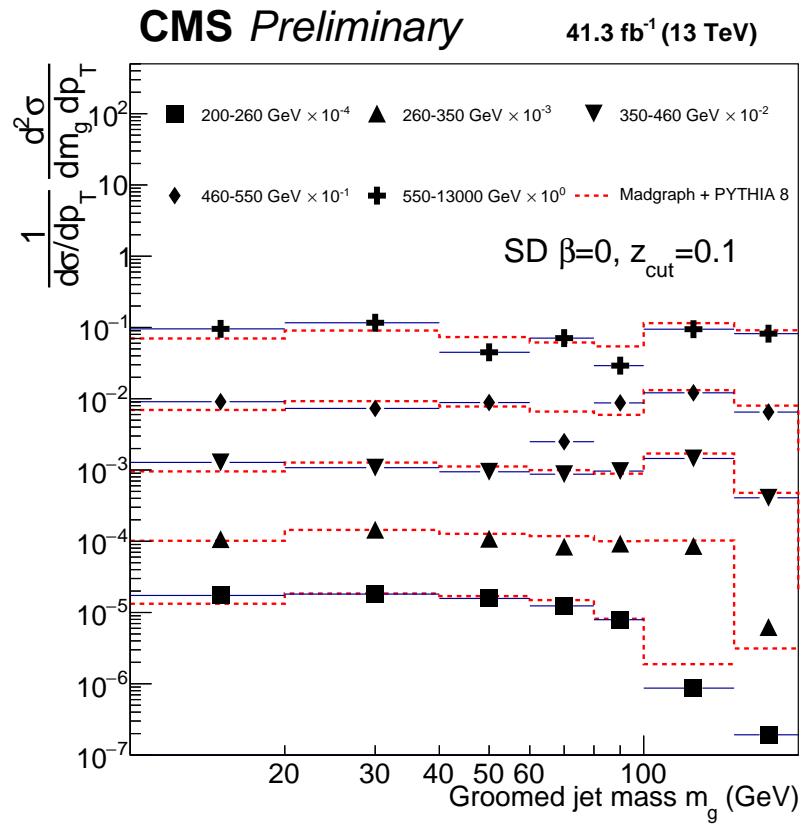


Figure 117: Results for groomed reconstructed unfolding with jet mass.

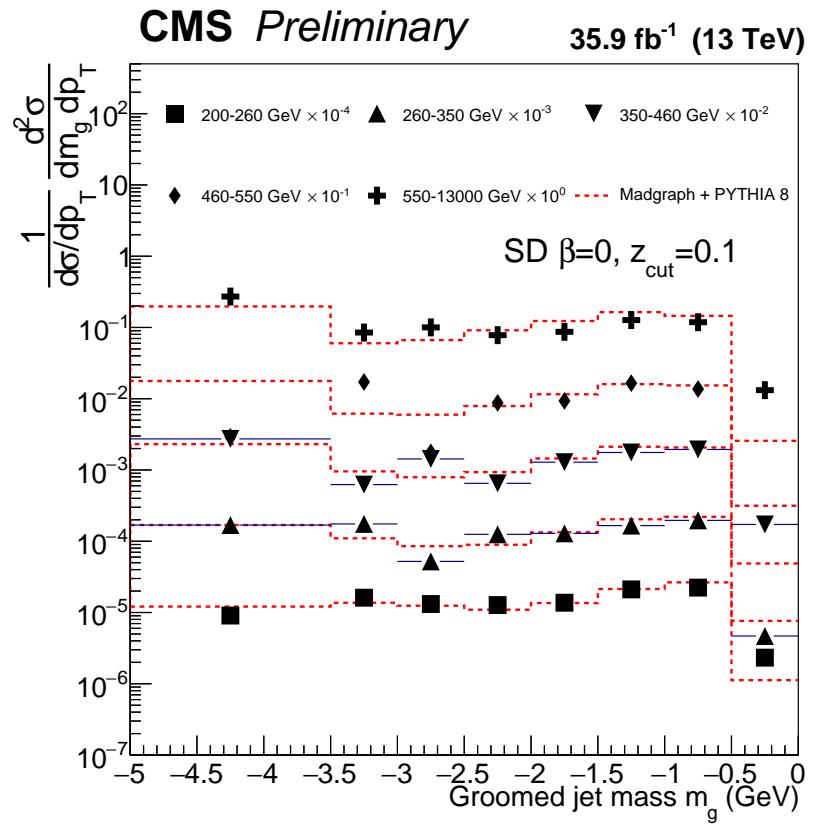


Figure 118: Results for groomed reconstructed unfolding with jet dimensionless mass.

The resulting correlation matrix without systematic uncertainties for the ungroomed jets is shown in Fig. 119, and for the groomed jets is shown in Fig. 121. The same figures with systematic uncertainties are shown for the ungroomed jets in Fig. 123, and for the groomed jets is shown in Fig. 125.

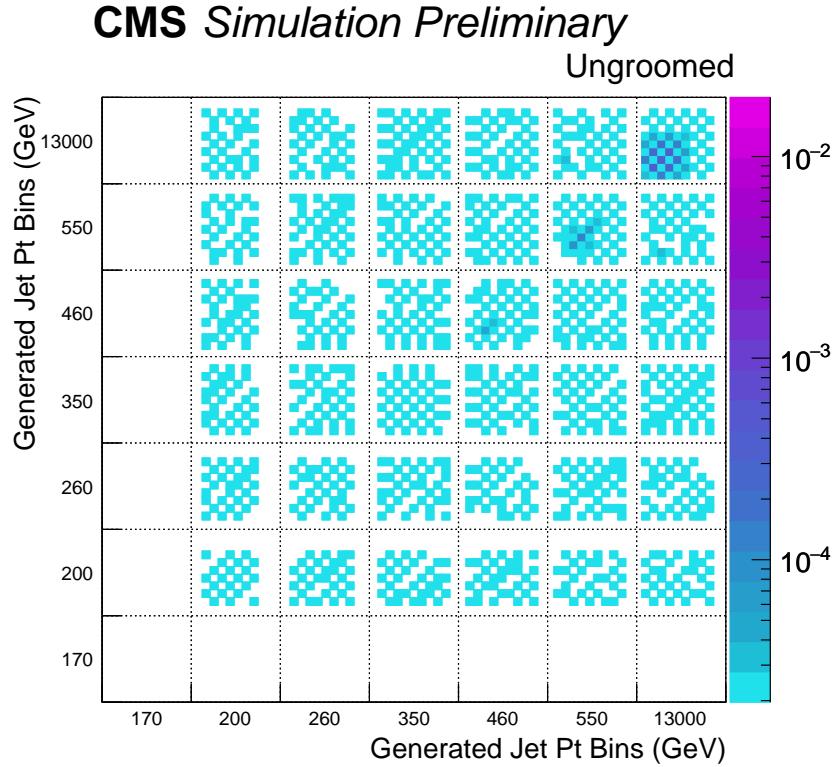


Figure 119: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 10, 20, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 350, \dots, 13000$ GeV).

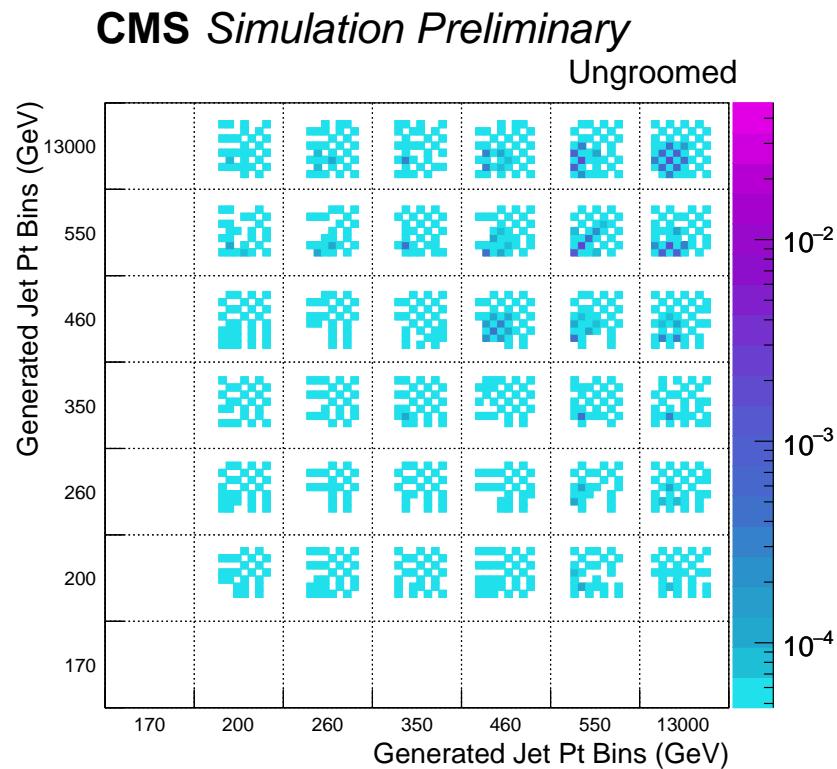


Figure 120: Correlation matrix for ungroomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 350, ... 13000 GeV).

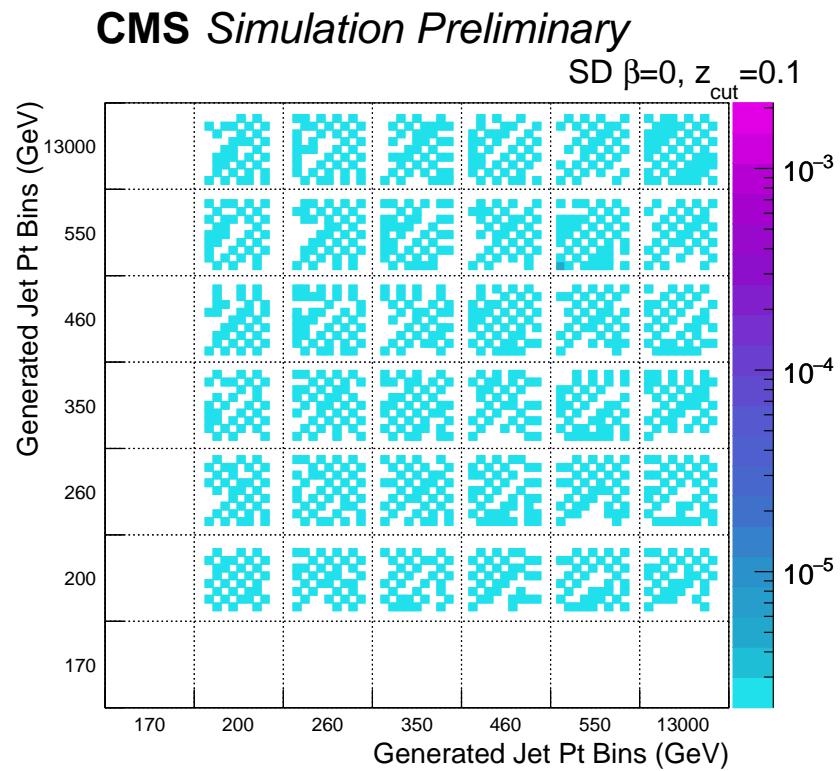


Figure 121: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...760 GeV).

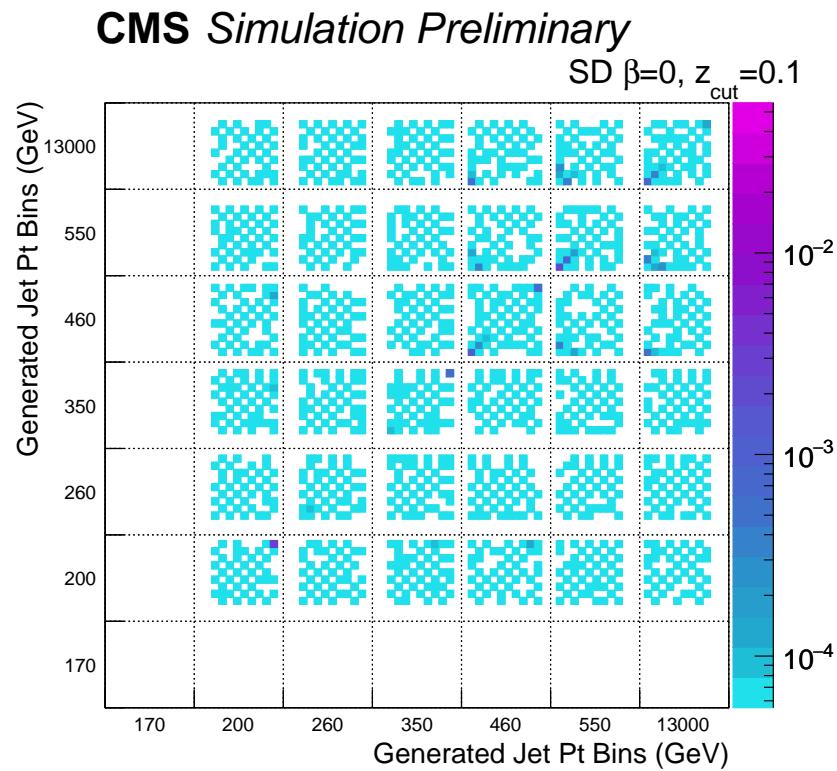


Figure 122: Correlation matrix for groomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 340, ... 760 GeV).

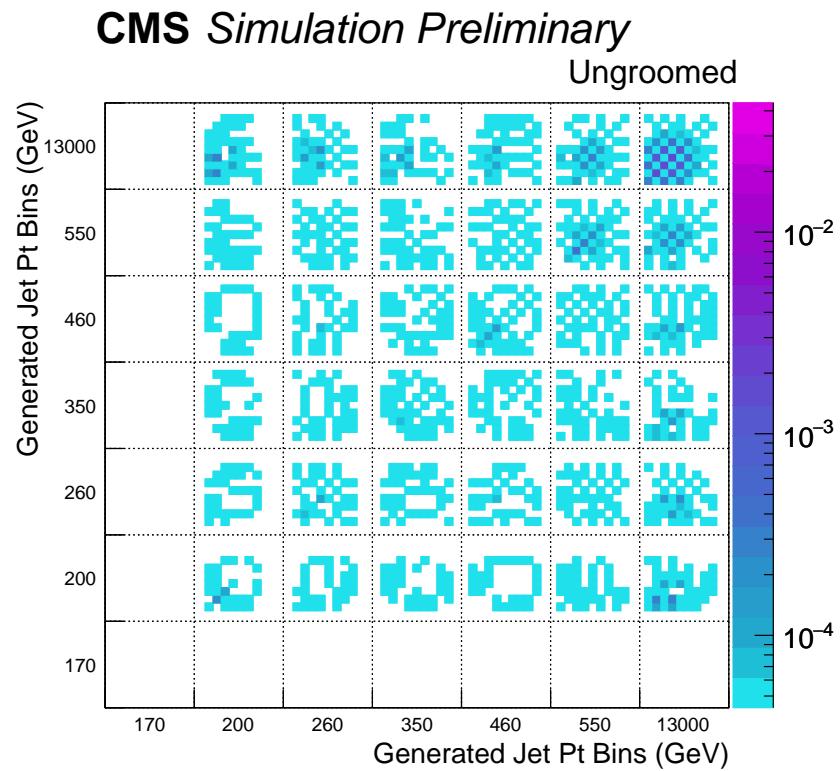


Figure 123: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 20, 40, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 260, \dots, 13000$ GeV).

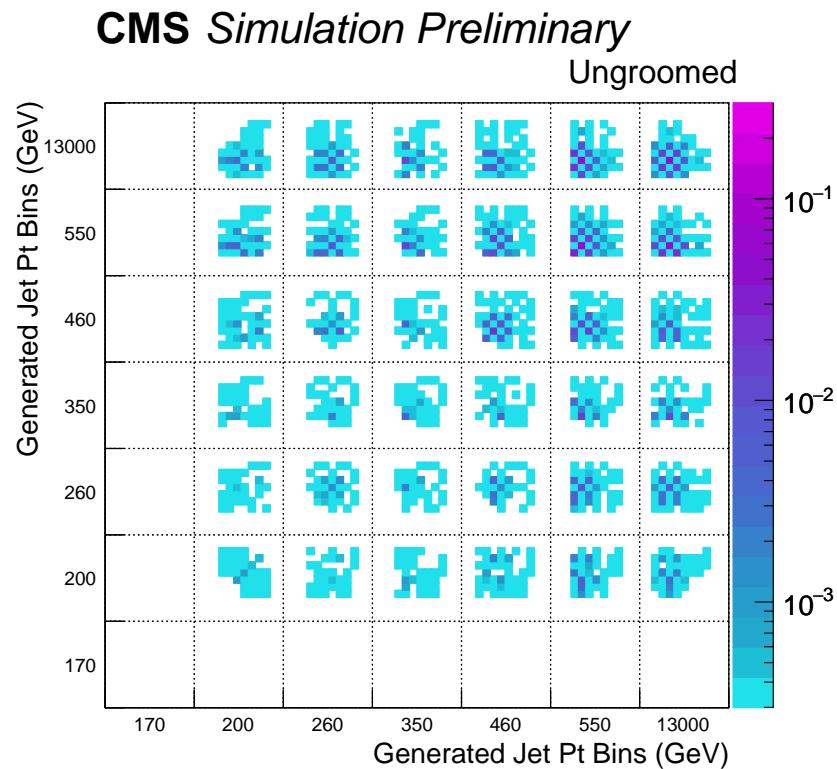


Figure 124: Correlation matrix for ungroomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 260, ... 13000 GeV).

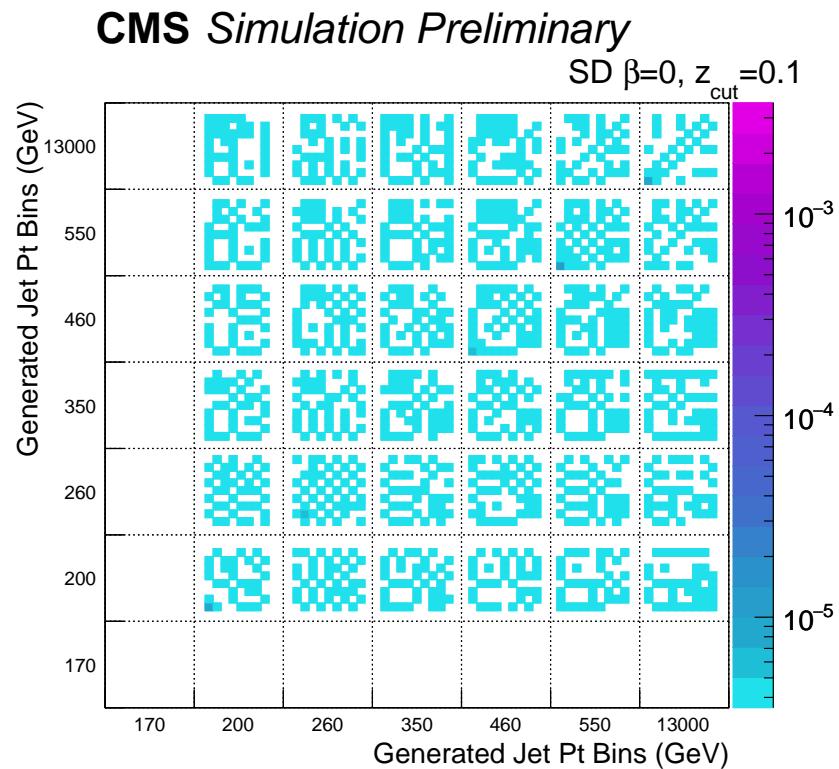


Figure 125: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...0 GeV).

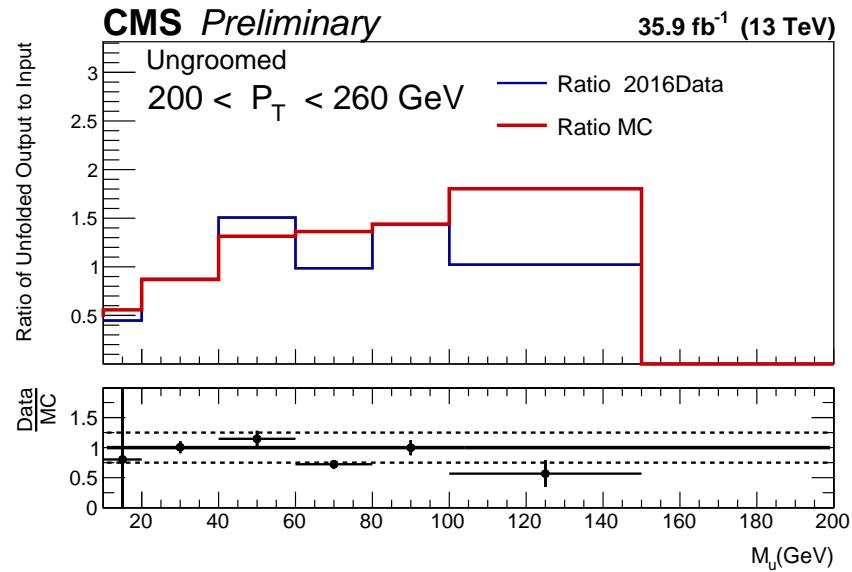


Figure 126: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 200-260 GeV.

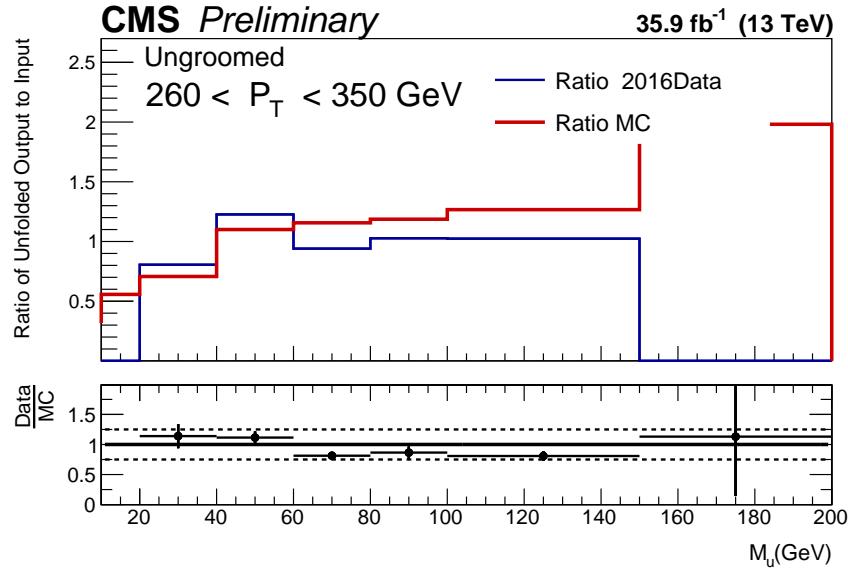


Figure 127: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 260-350 GeV.

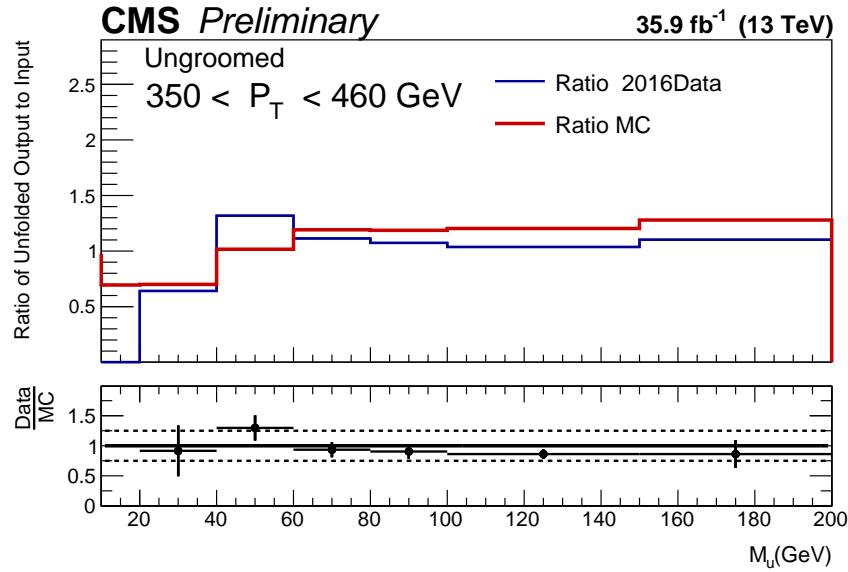


Figure 128: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 350-460 GeV.

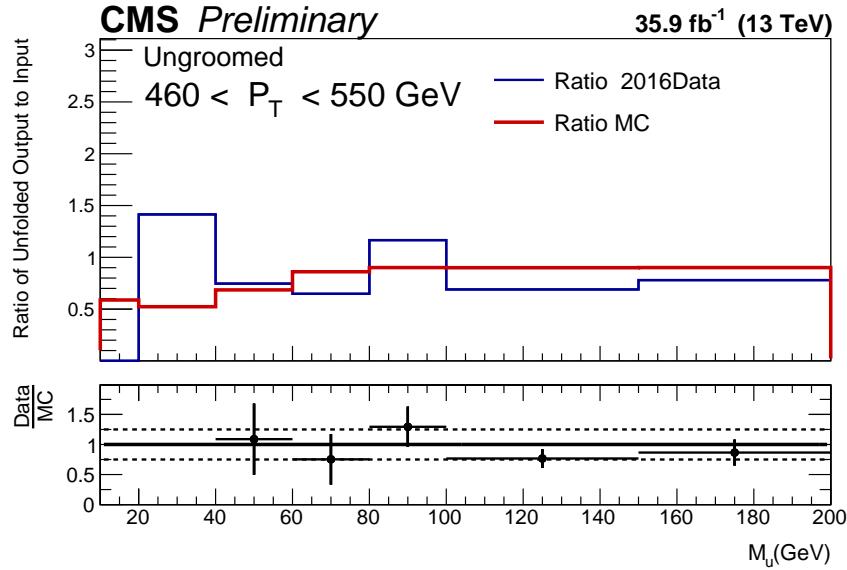


Figure 129: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 460-550 GeV.

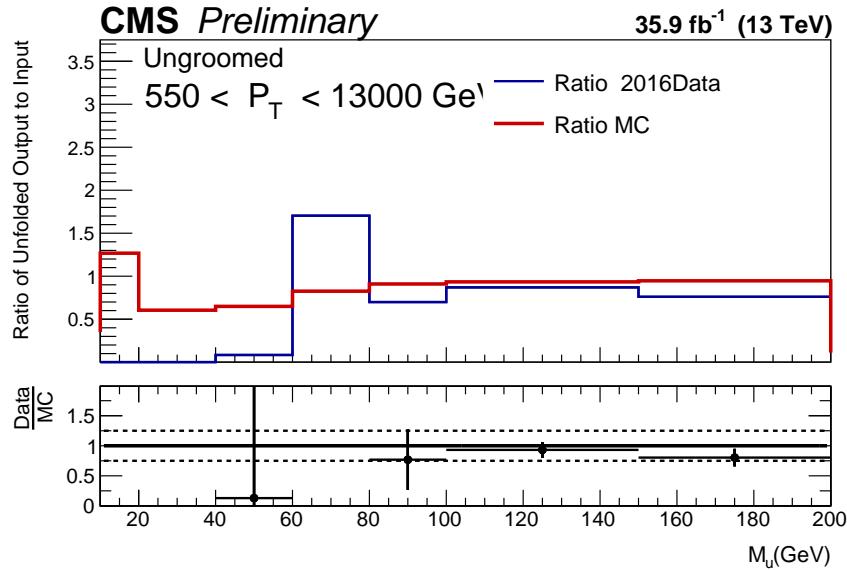


Figure 130: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 550-13000 GeV.

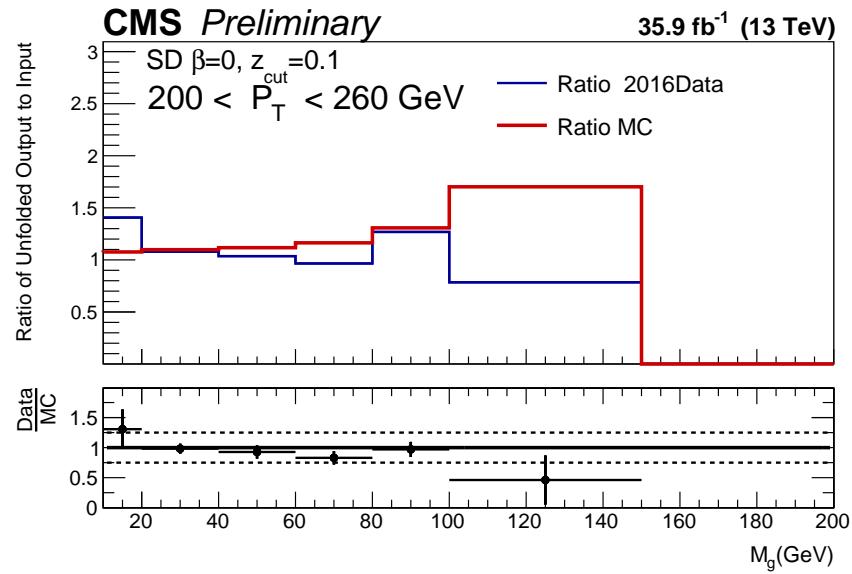


Figure 131: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 200-260 GeV.

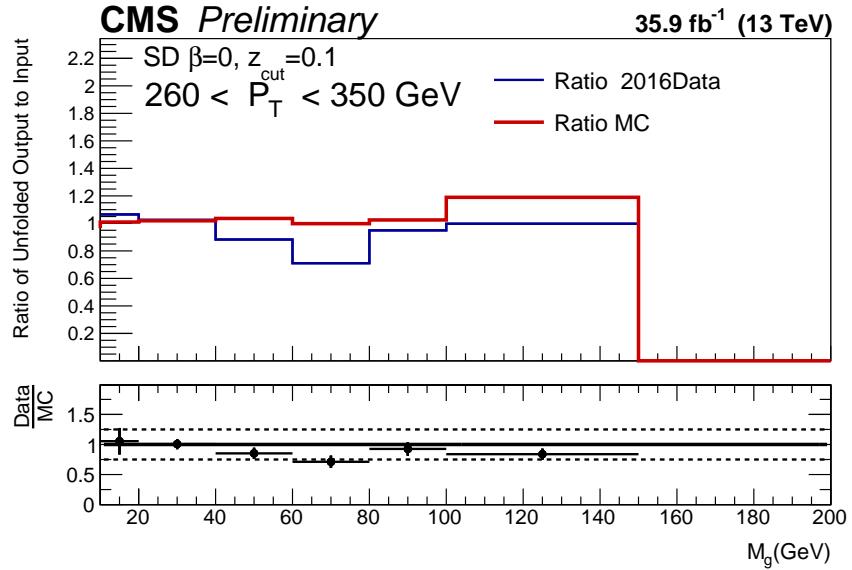


Figure 132: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 260-350 GeV.

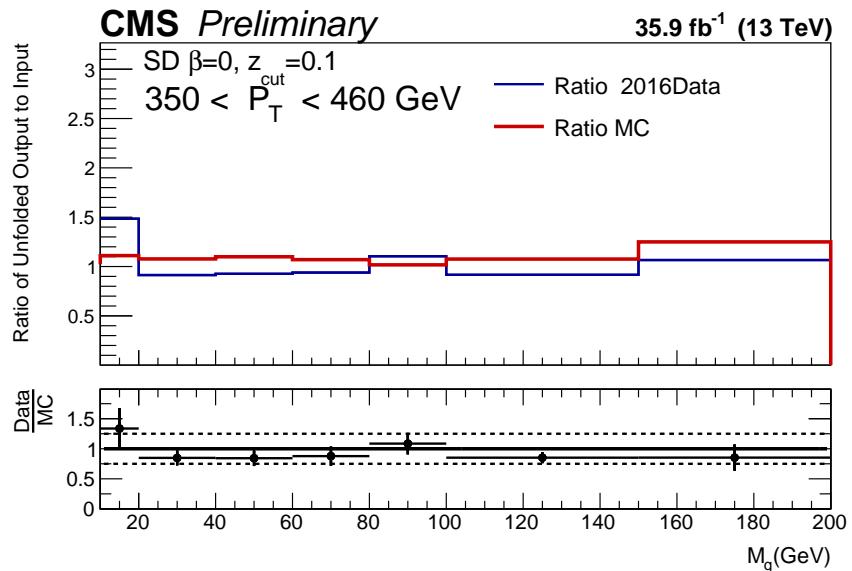


Figure 133: Ratio of unfolded over raw data and MC for groomed jets, p_{T} 350-460 GeV.

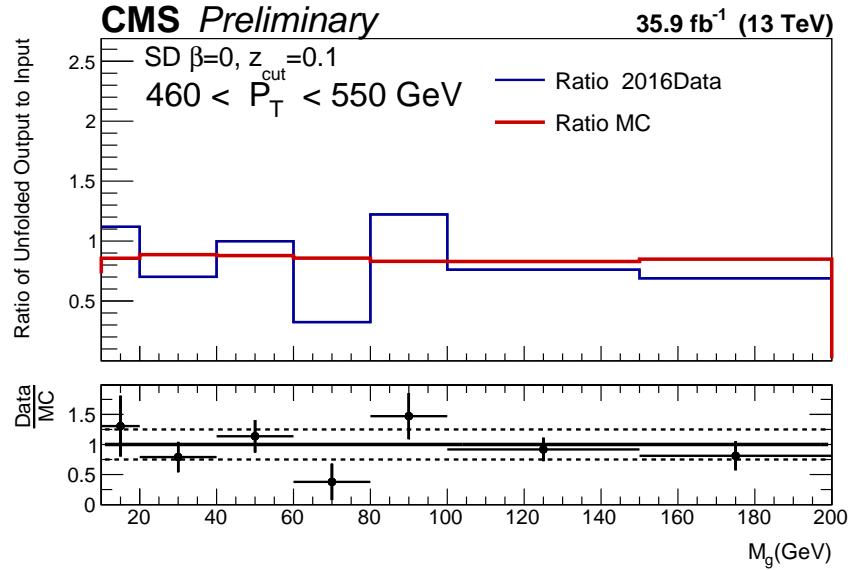


Figure 134: Ratio of unfolded over raw data and MC for groomed jets, p_T 460-550 GeV.

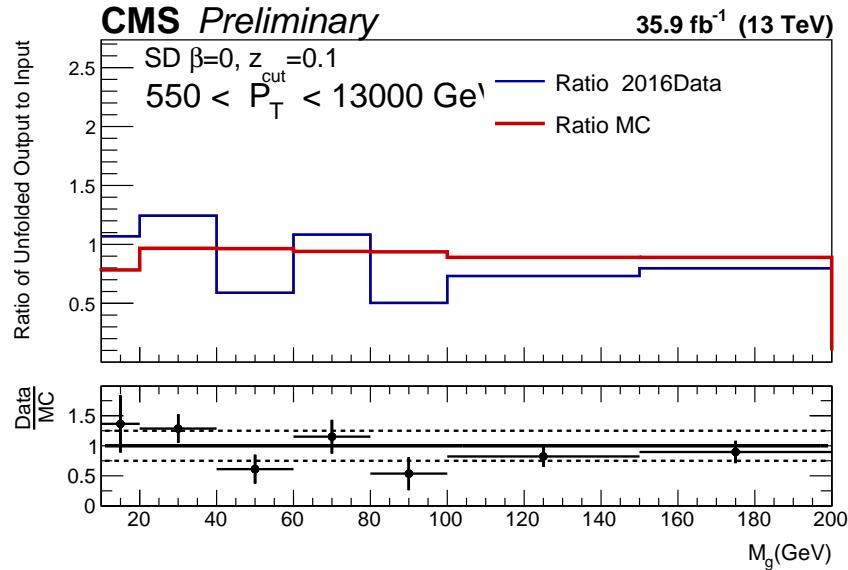


Figure 135: Ratio of unfolded over raw data and MC for groomed jets, p_T 550-13000 GeV.

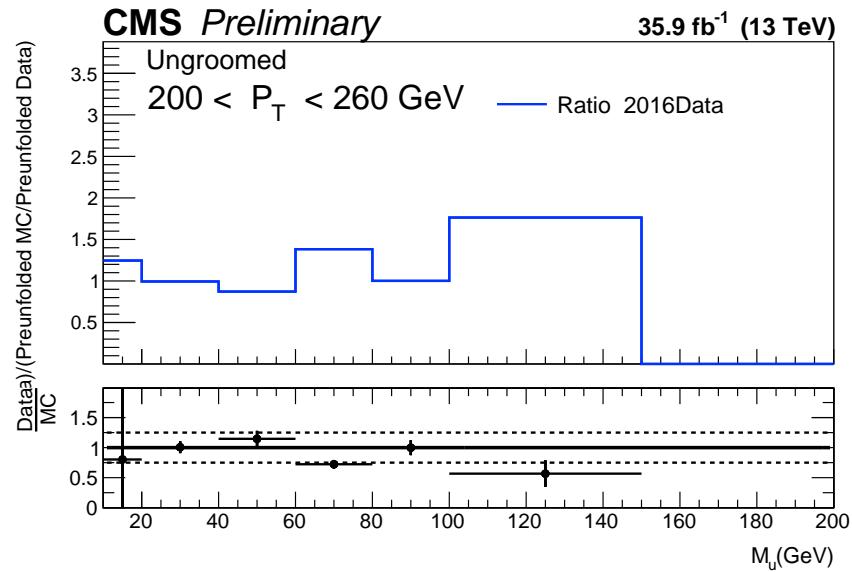


Figure 136: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 200-260 GeV.

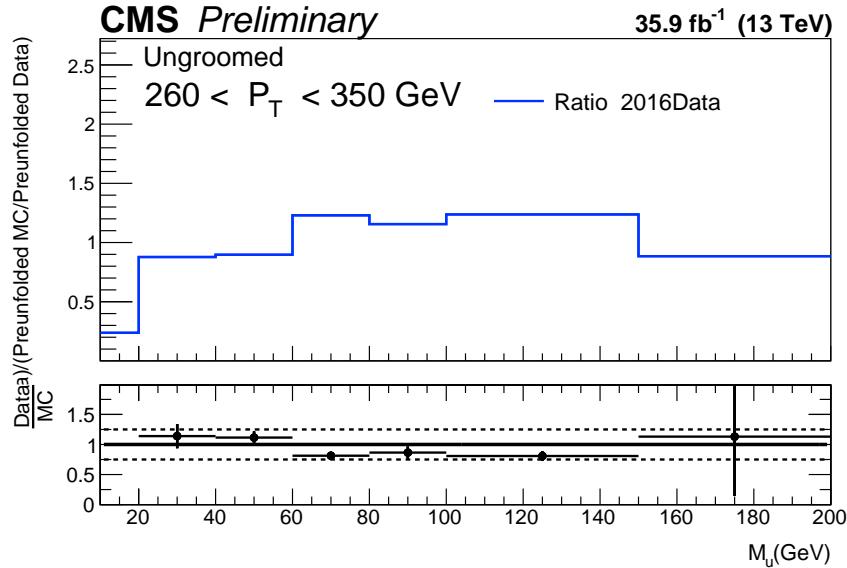


Figure 137: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 260-350 GeV.

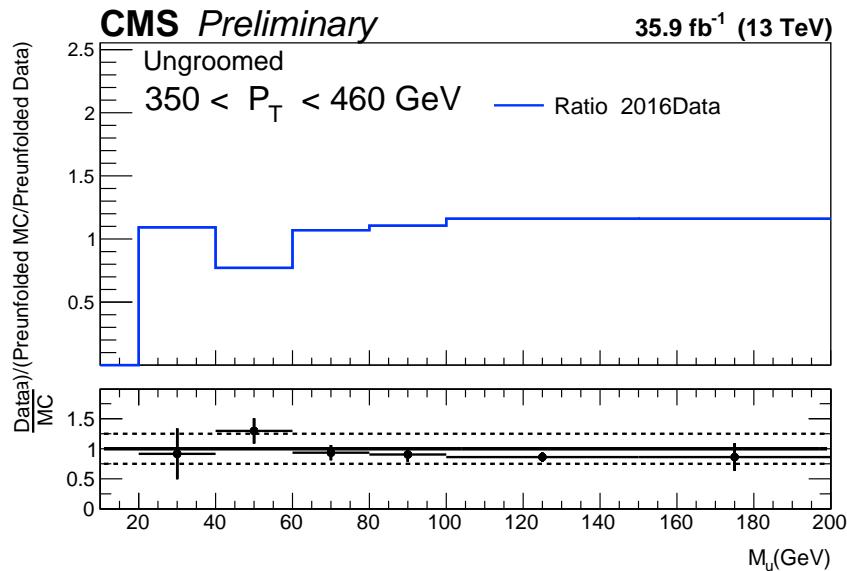


Figure 138: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 350-460 GeV.

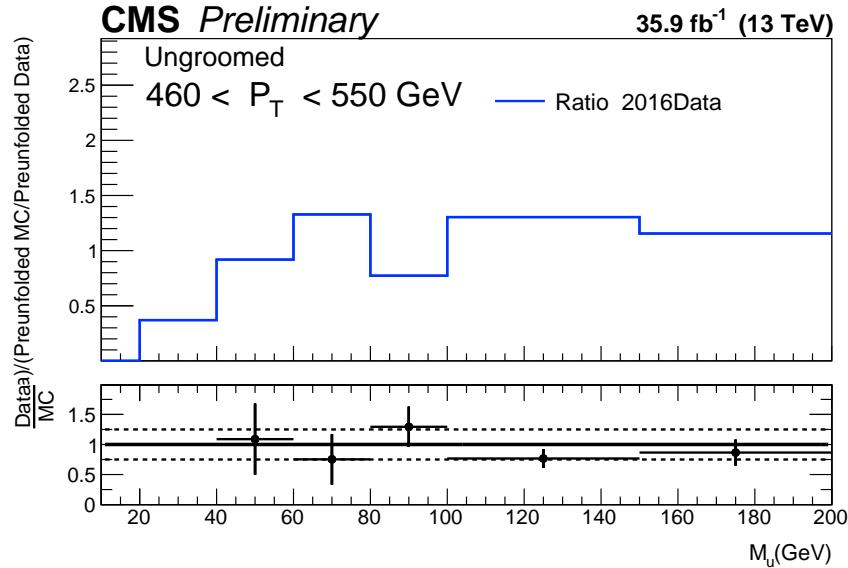


Figure 139: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 460-550 GeV.

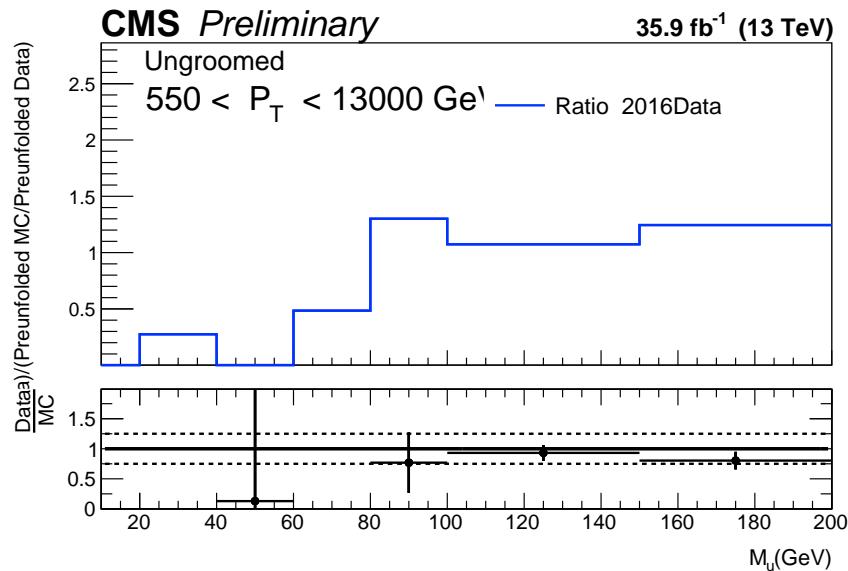


Figure 140: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 550-13000 GeV.

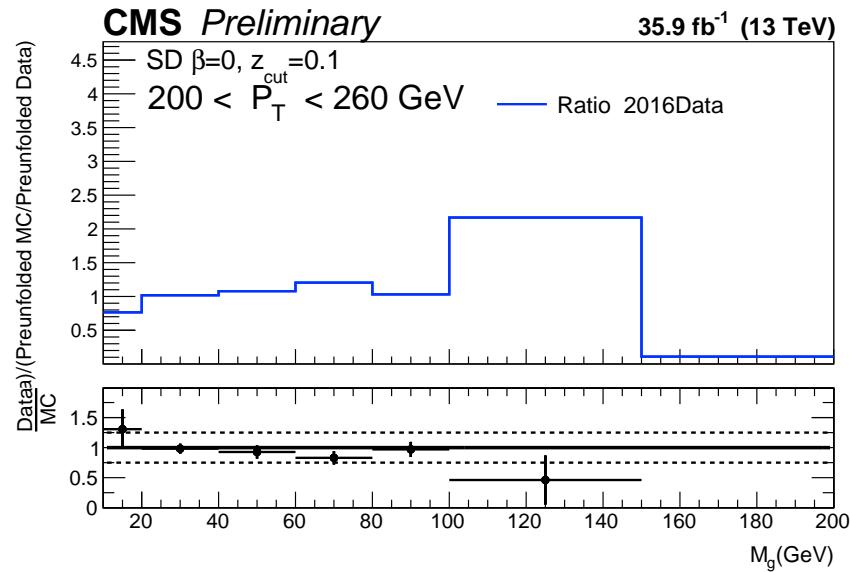


Figure 141: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_{T} 200-260 GeV.

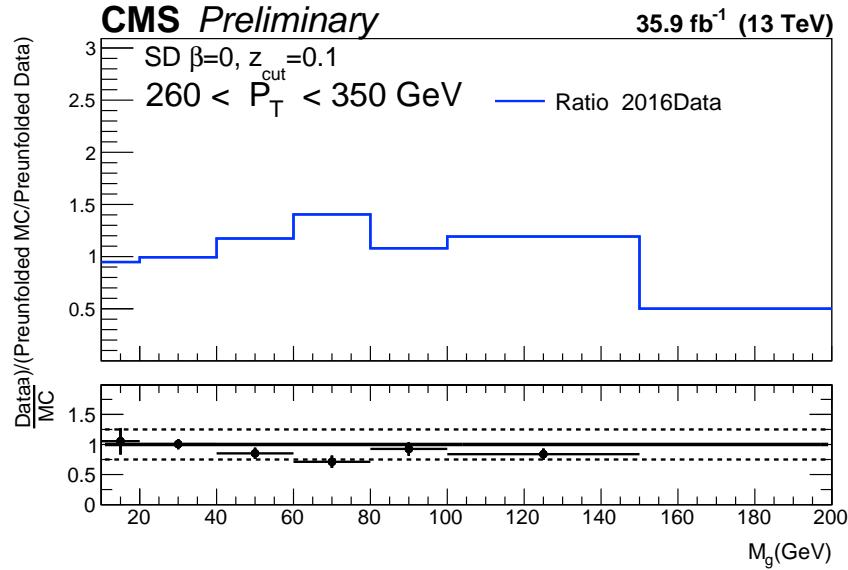


Figure 142: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 260-350 GeV.

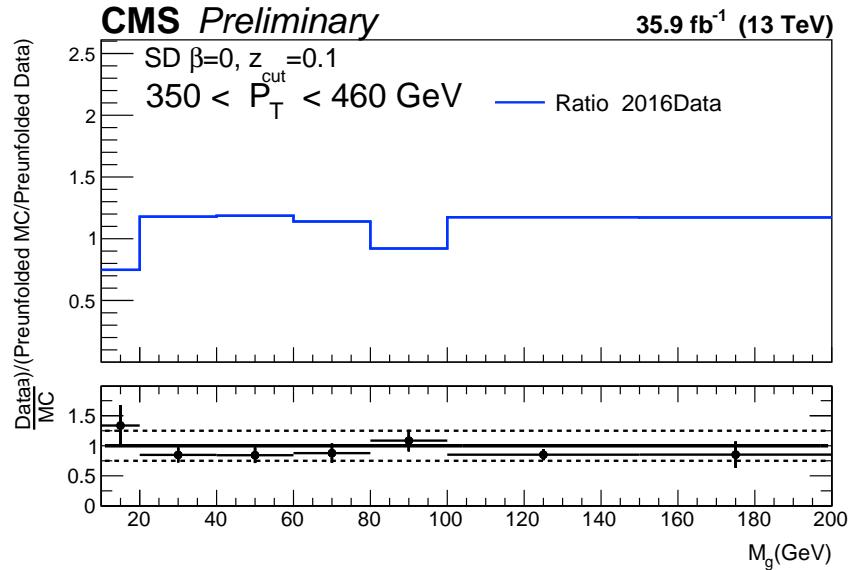


Figure 143: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 350-460 GeV.

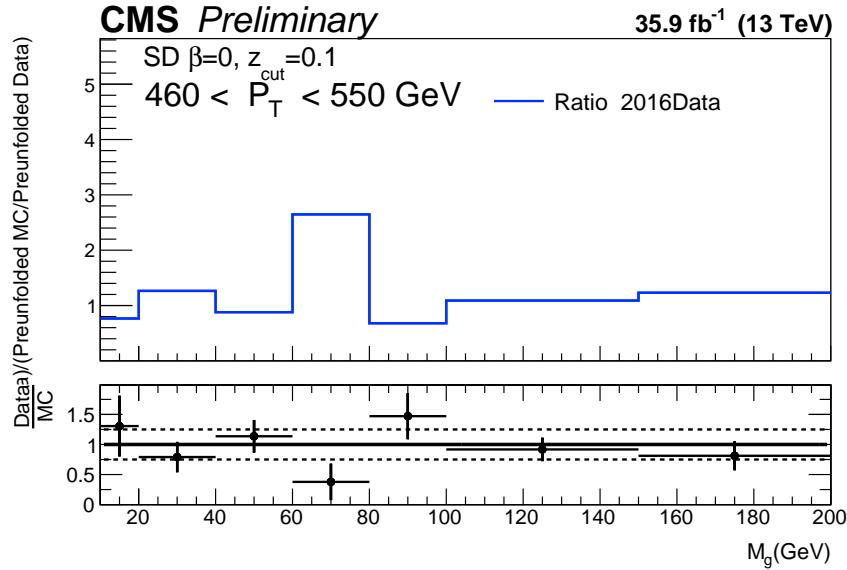


Figure 144: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 460-550 GeV.

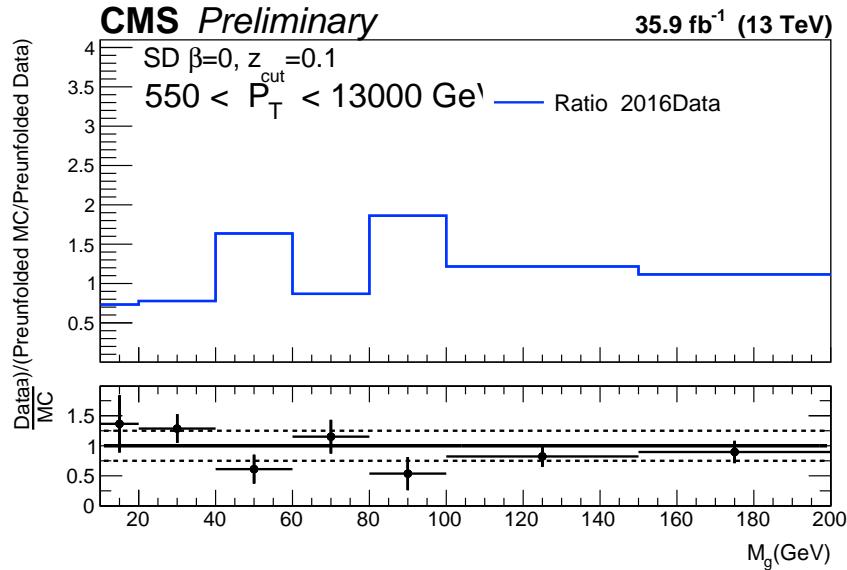


Figure 145: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 550-13000 GeV.

.2 2017 data results

This Appendix shows the distributions from the "Detector Response" through the "Results" sections of the main analysis note with only the 2017 data rather than the full Run 2 statistics seen in the main body of the note.

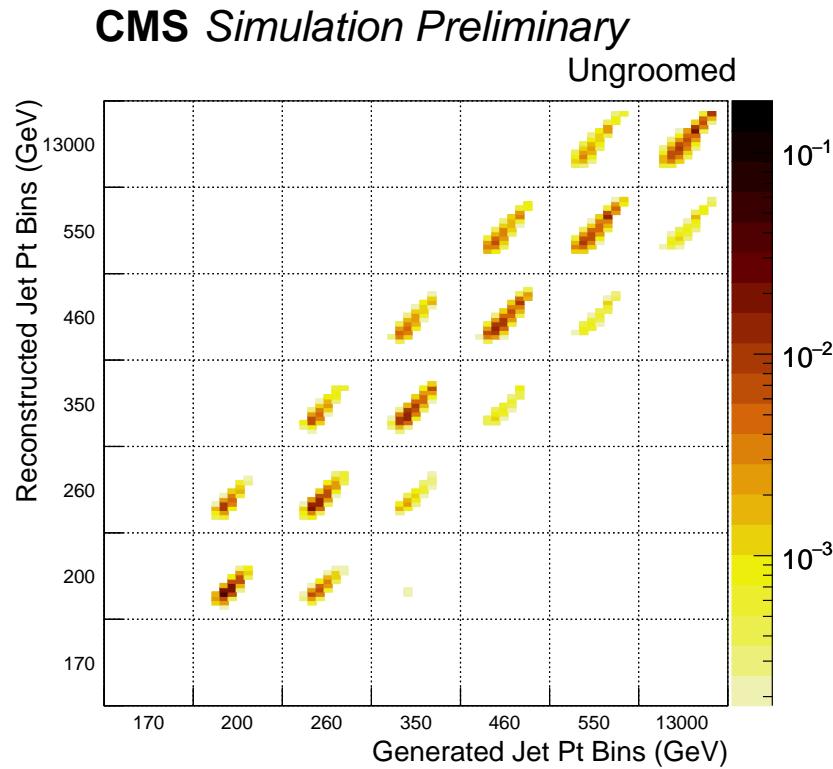


Figure 146: Two-dimensional response matrix for ungroomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). While the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

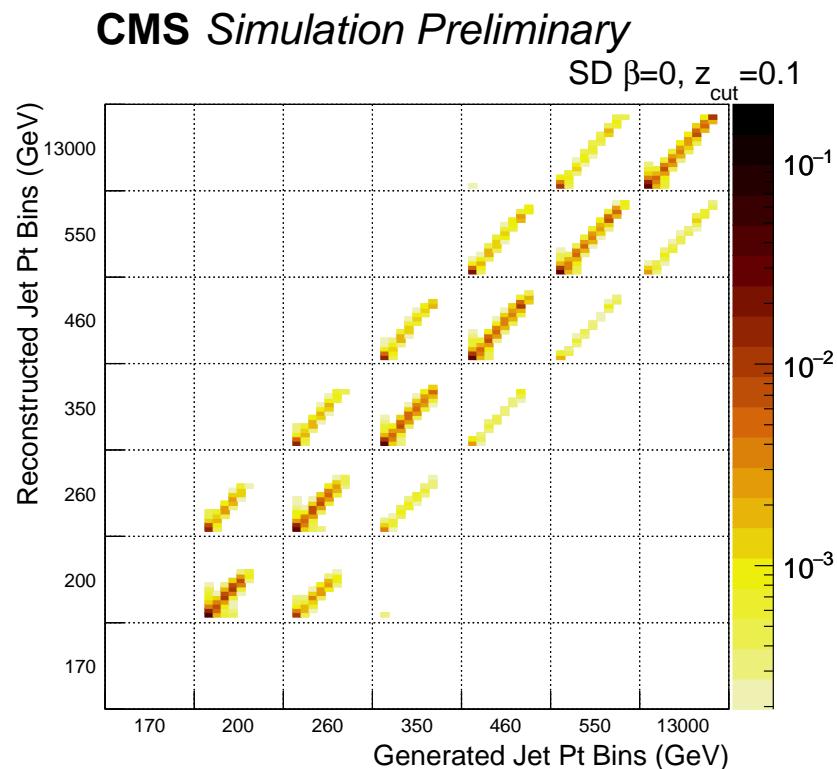


Figure 147: Two-dimensional response matrix for groomed jets $\beta = 0, z_{cut} = 0.1$. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV) while the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

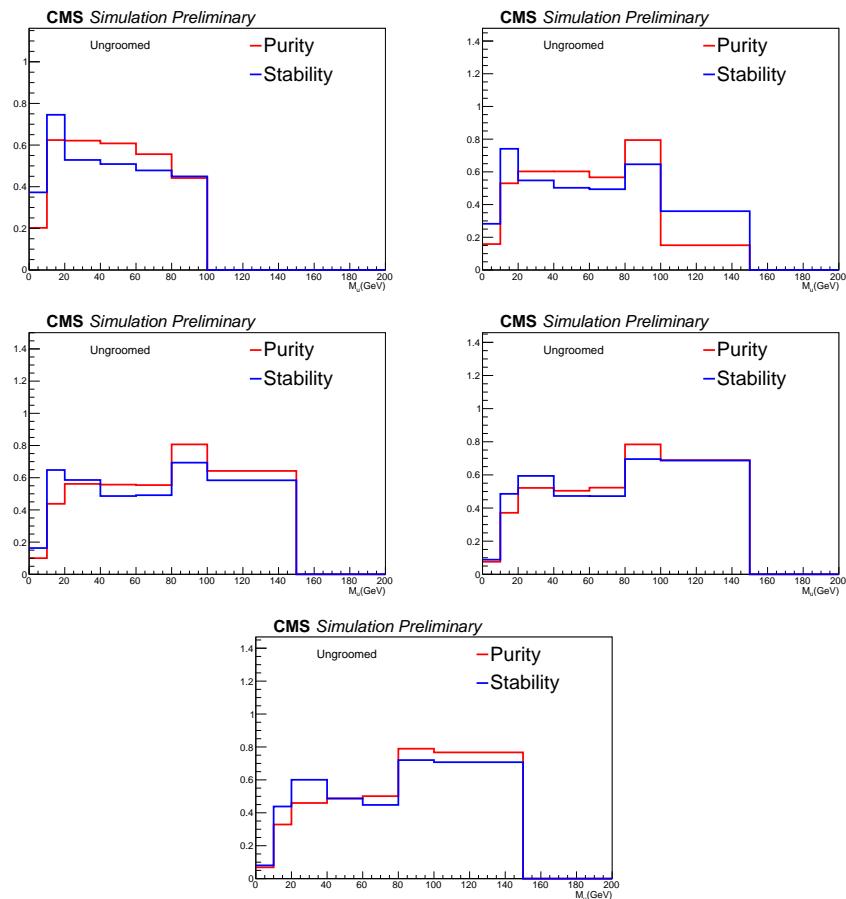


Figure 148: Purity and stability for ungroomed jets in various p_T bins.

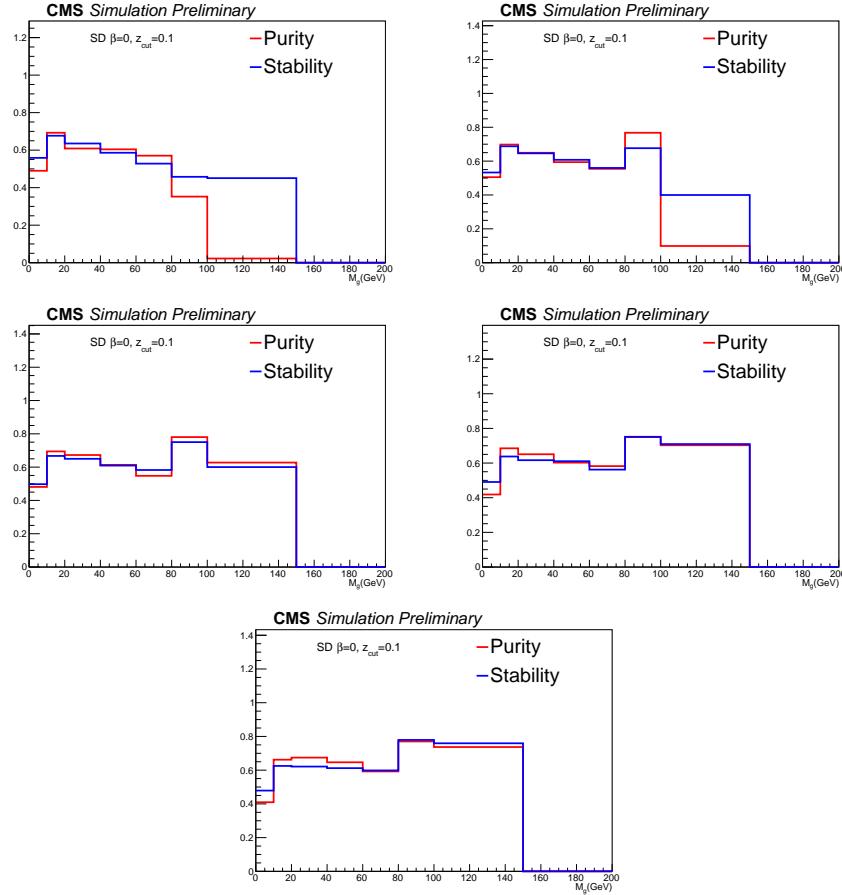


Figure 149: Purity and stability for groomed jets, where the soft-drop criterion was applied with $\beta = 0$ and $z_{cut} = 0.1$, in various p_T bins.

Figure-150 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section. Figure-151 shows the same distributions for the raw dimensionless mass.

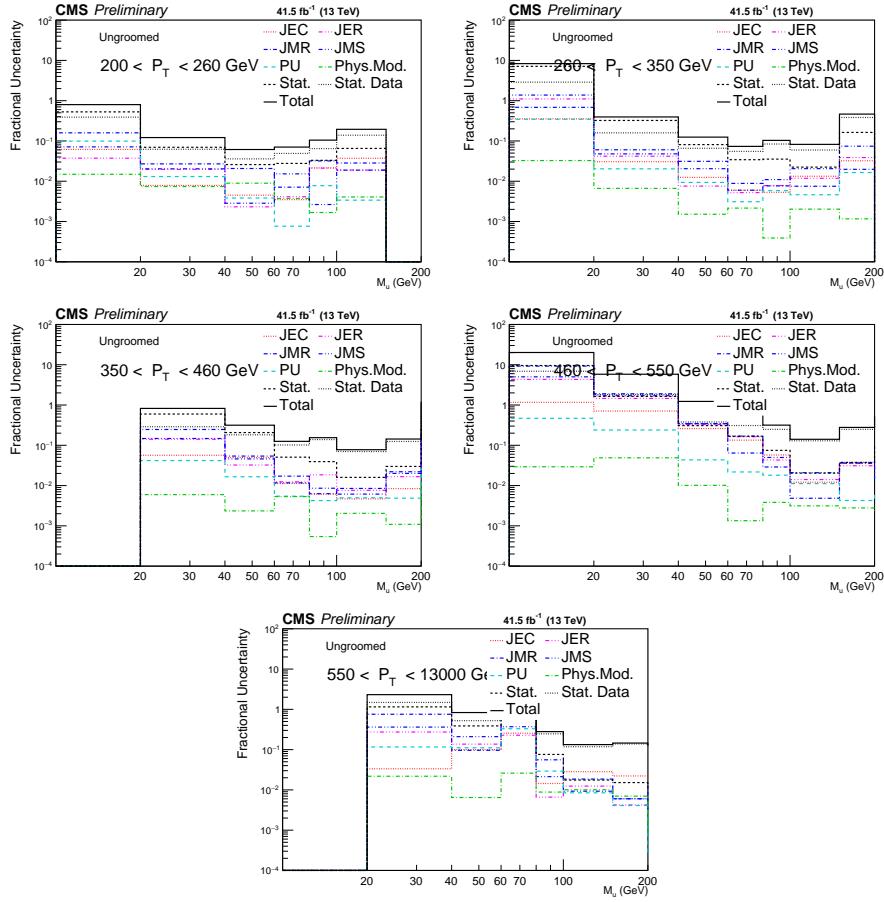


Figure 150: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw jet mass.

Figure-152 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section. Figure-153 shows the same distributions for the groomed dimensionless mass.

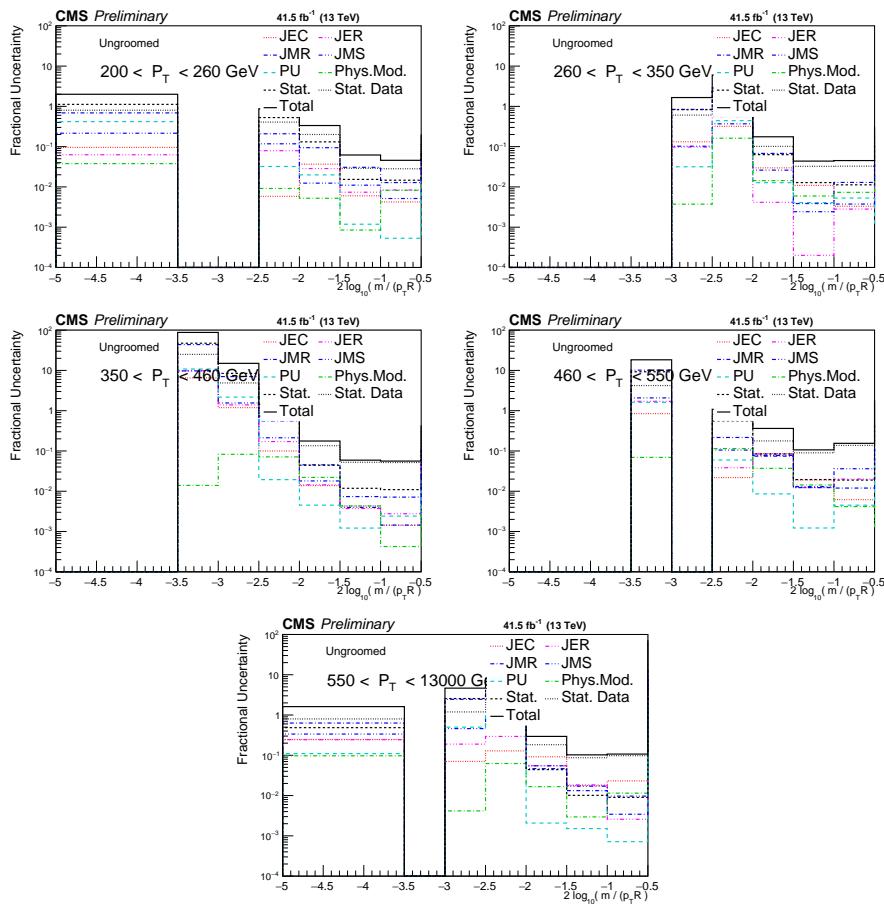


Figure 151: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw dimensionless jet mass.

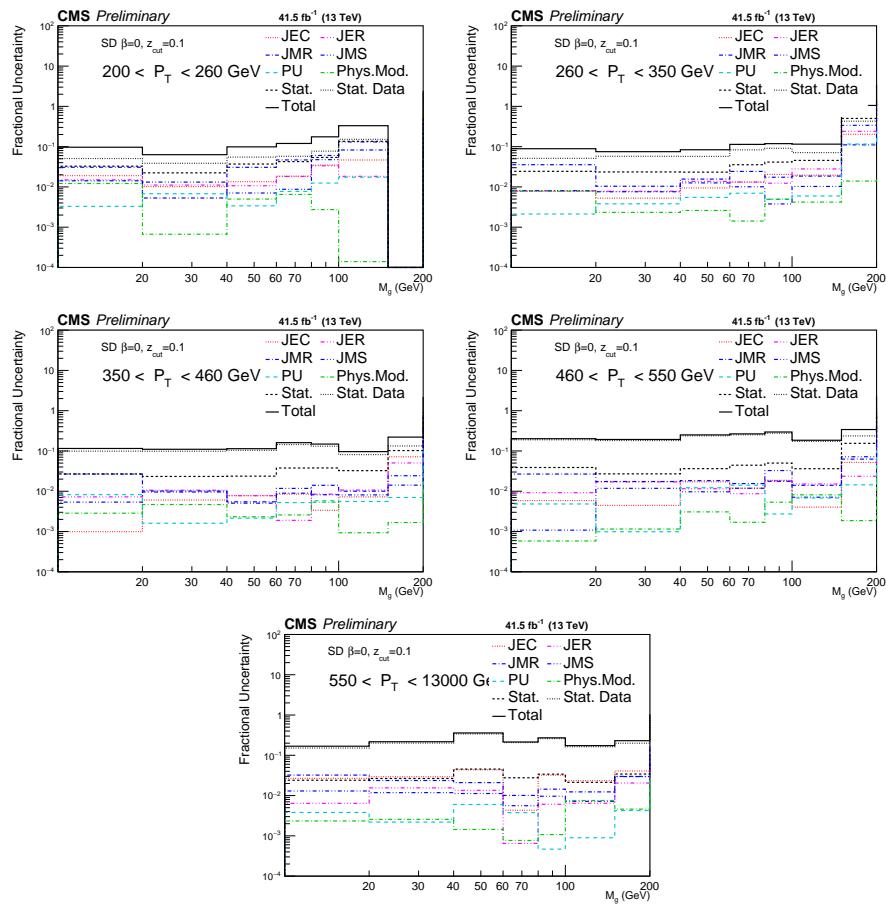


Figure 152: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

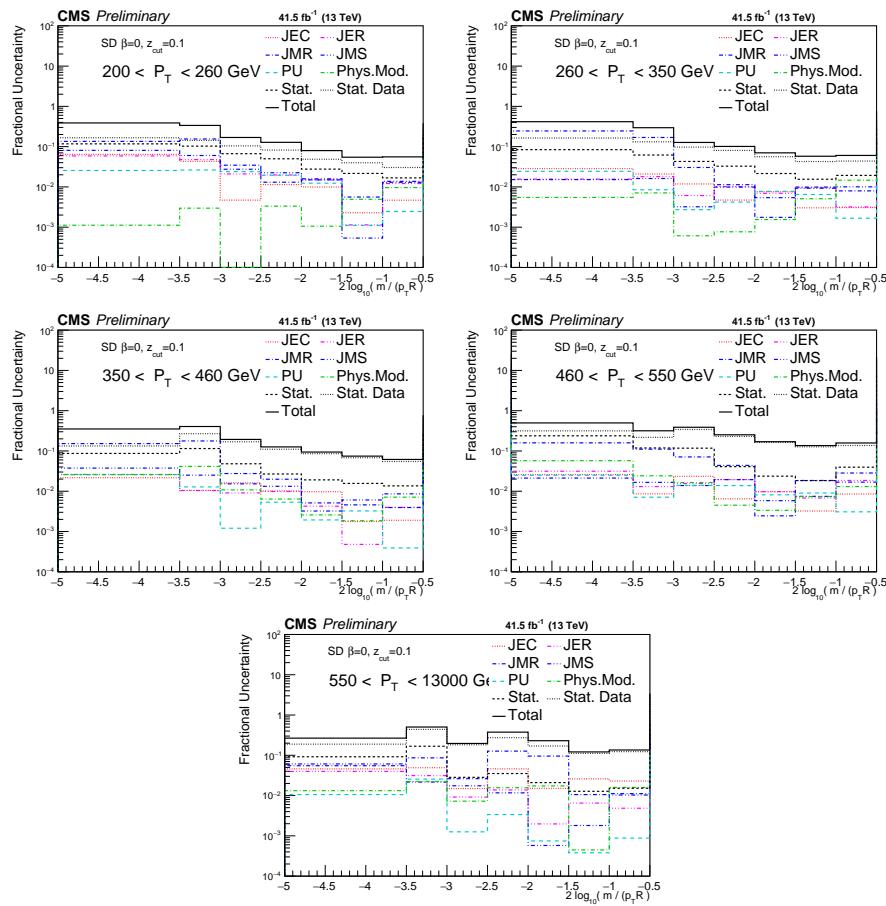


Figure 153: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

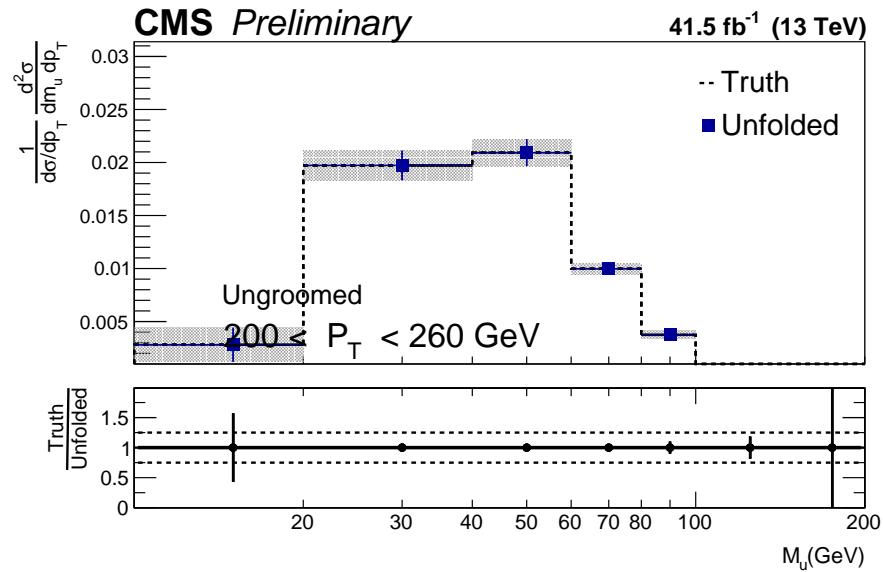


Figure 154: Closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

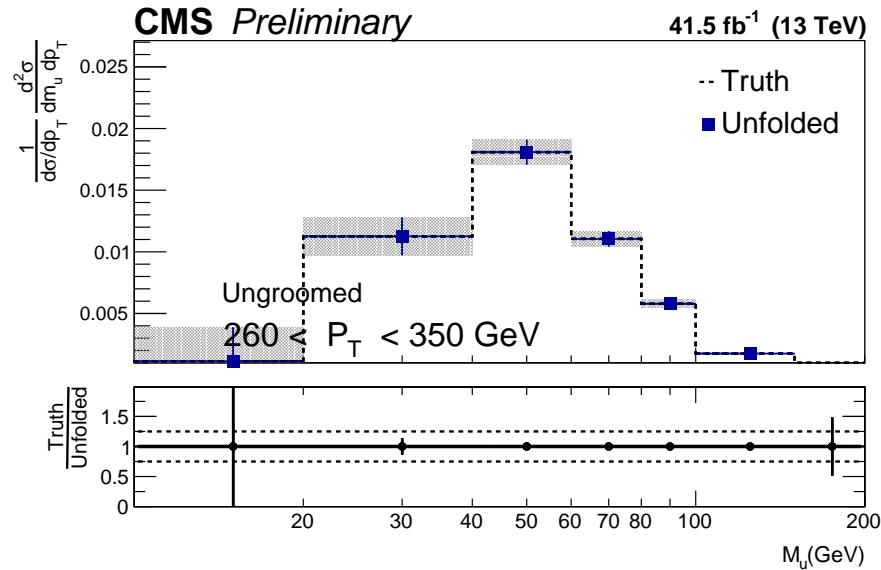


Figure 155: Closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

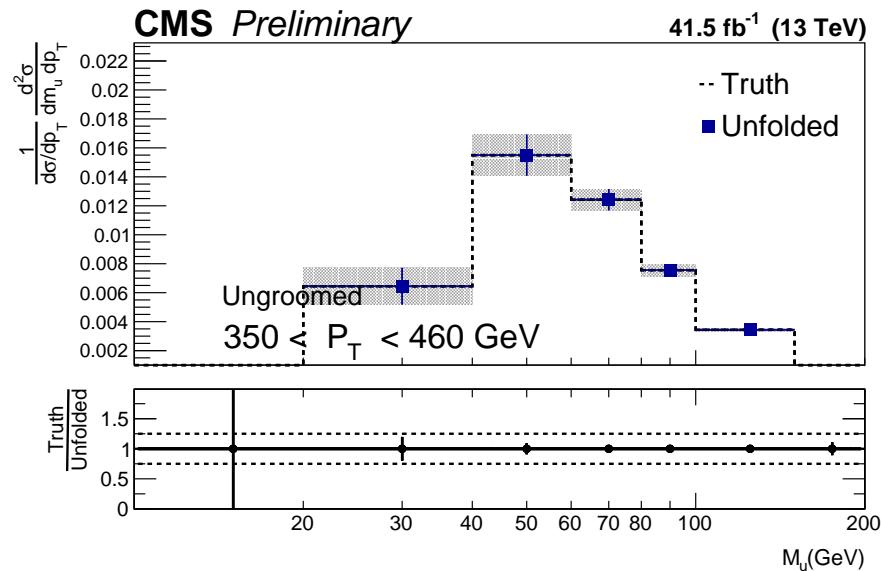


Figure 156: Closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

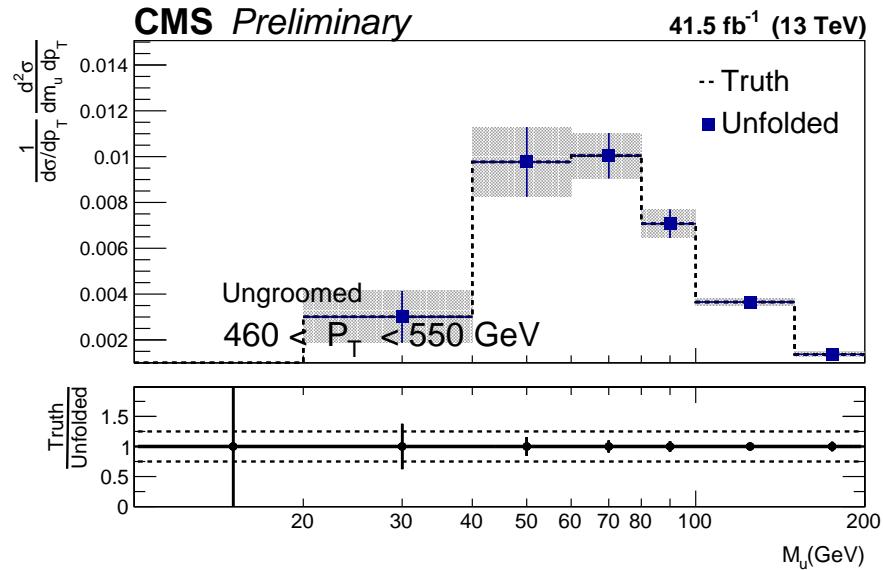


Figure 157: Closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

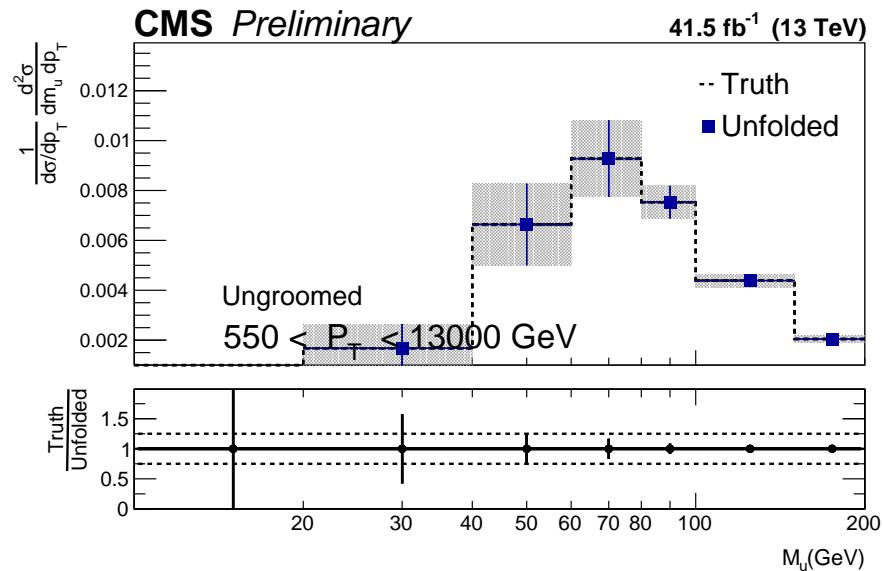


Figure 158: Closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

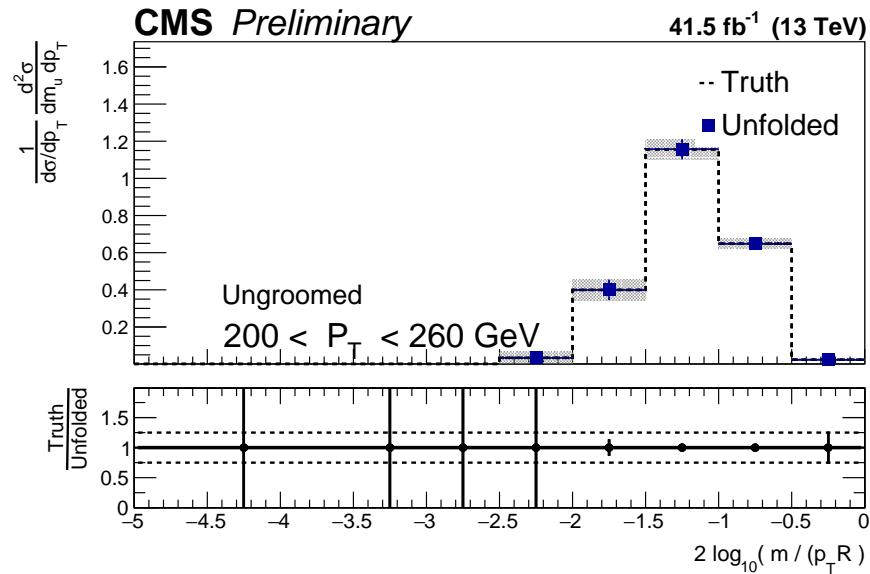


Figure 159: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

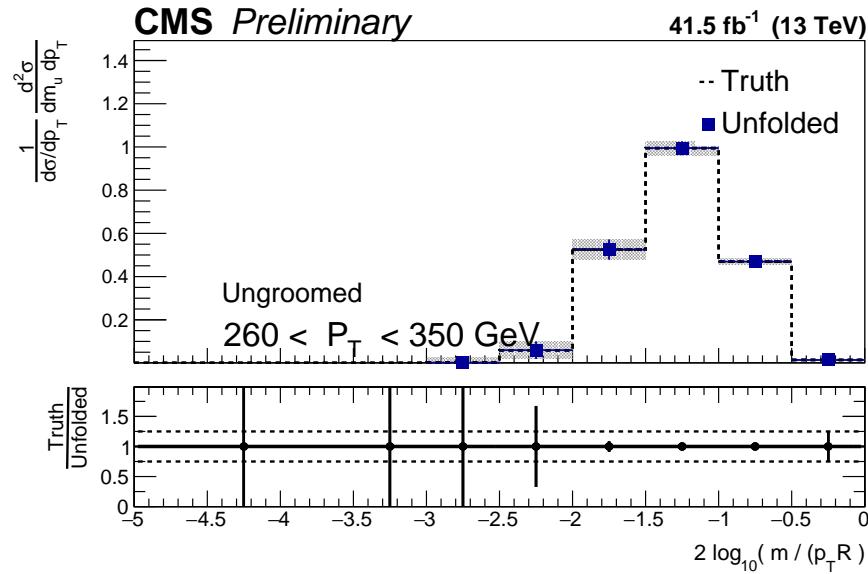


Figure 160: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

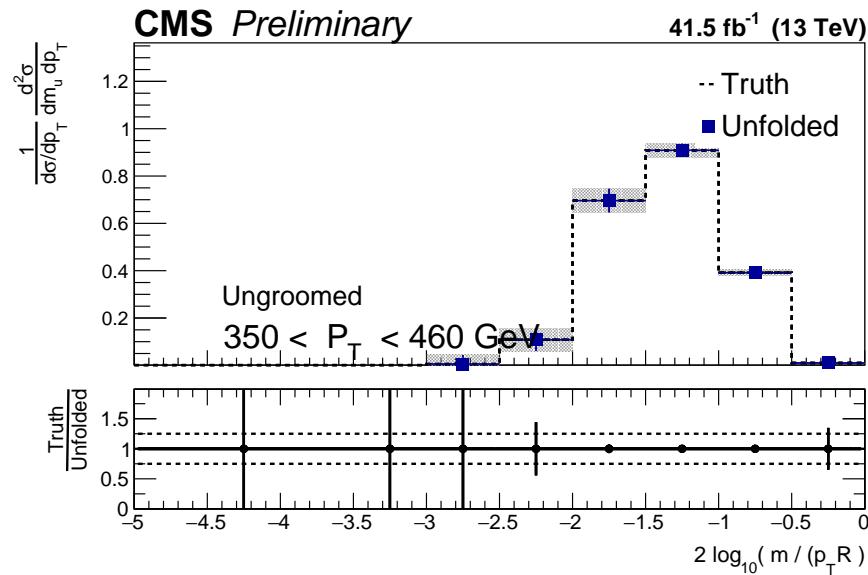


Figure 161: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

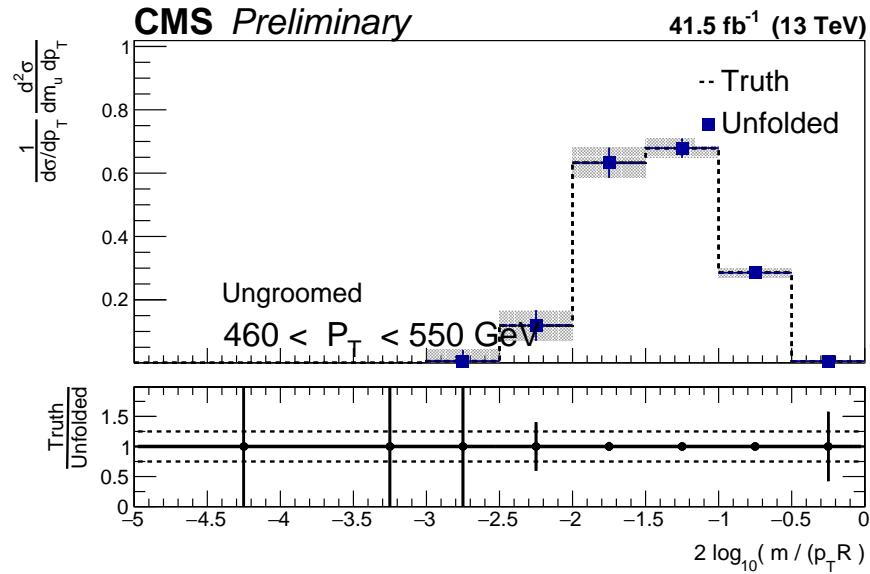


Figure 162: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

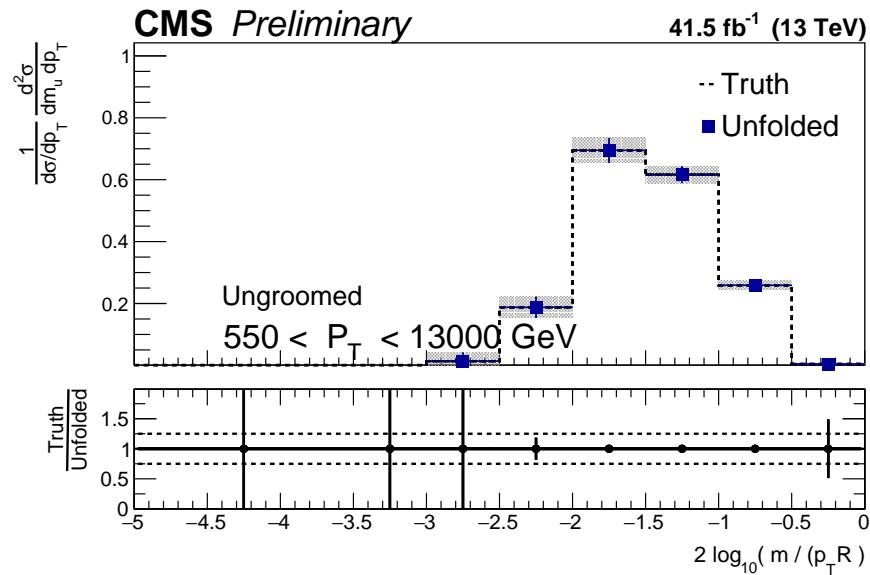


Figure 163: Dimensionless mass closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

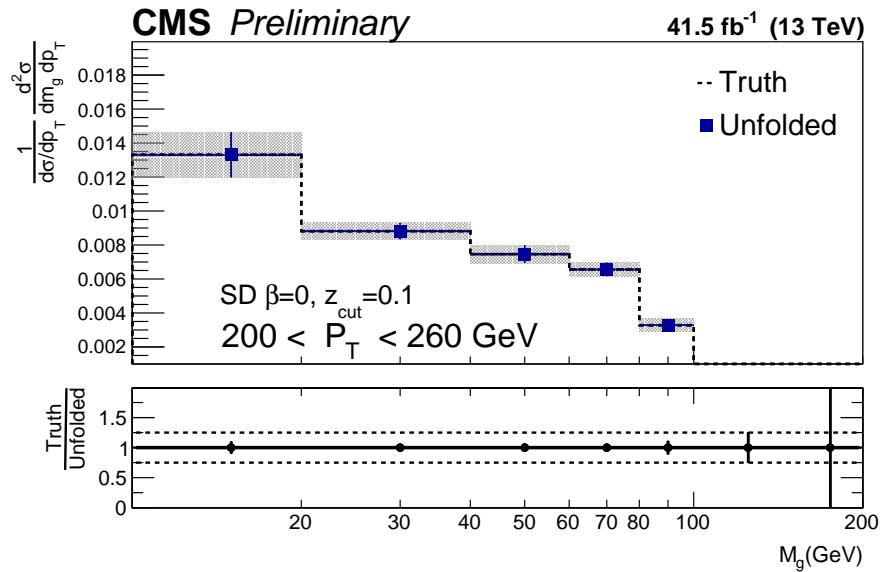


Figure 164: Closure test of groomed reconstructed Monte Carlo, p_{T} 200-260 GeV.

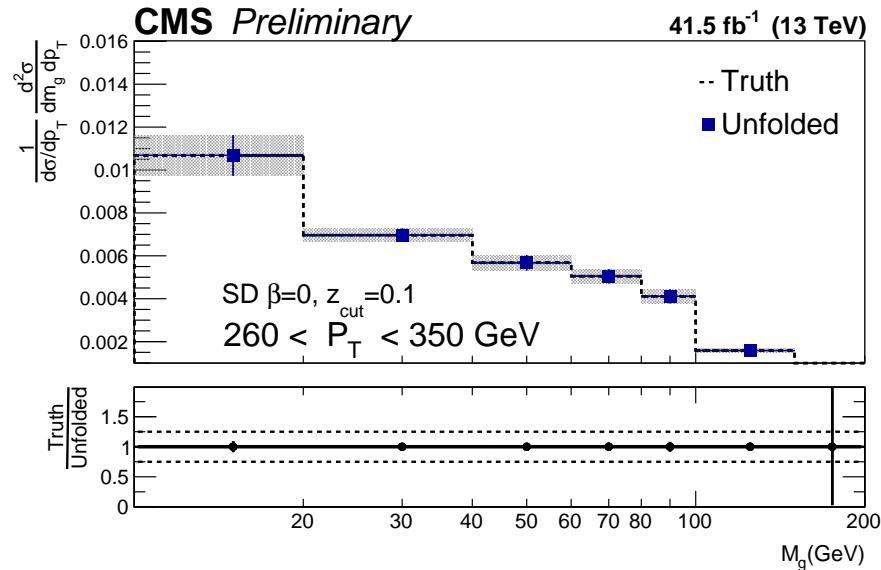


Figure 165: Closure test of groomed reconstructed Monte Carlo, p_T 260–350 GeV.

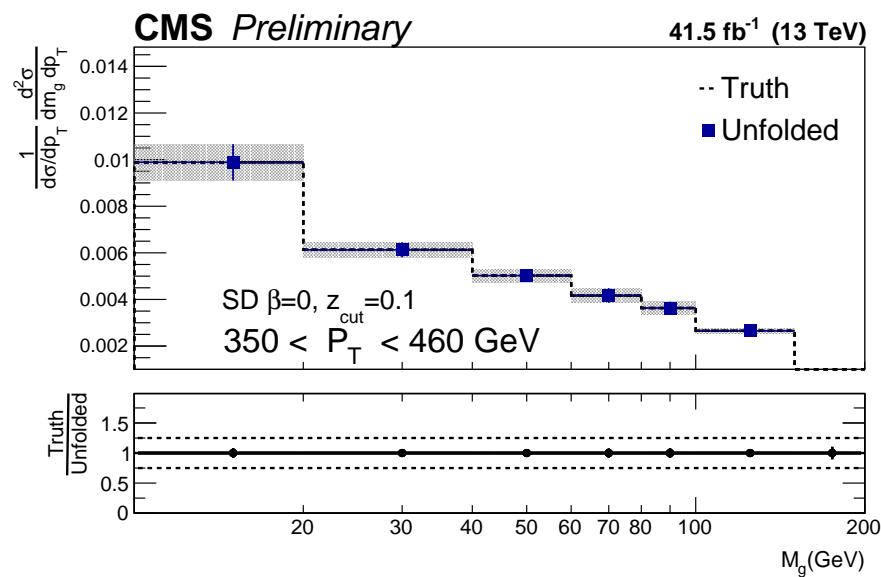


Figure 166: Closure test of groomed reconstructed Monte Carlo, p_T 350–460 GeV.

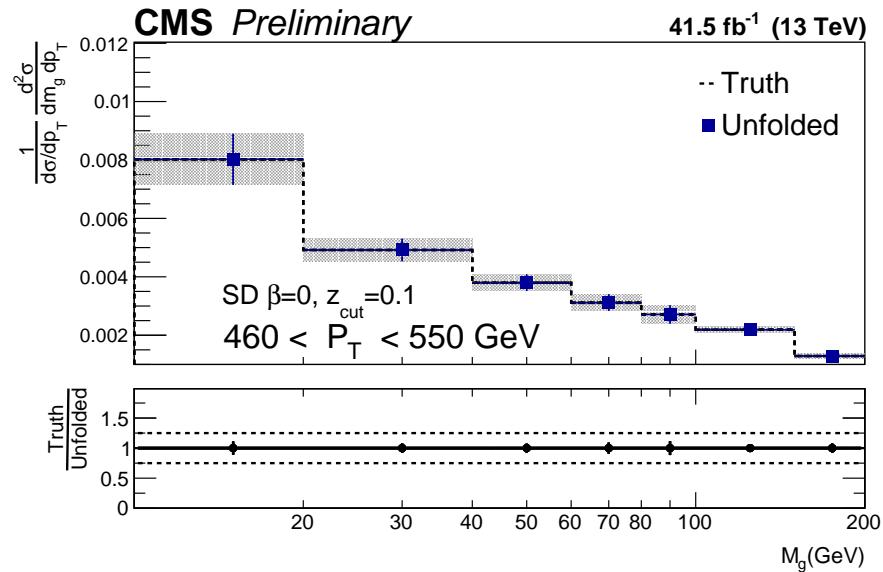


Figure 167: Closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

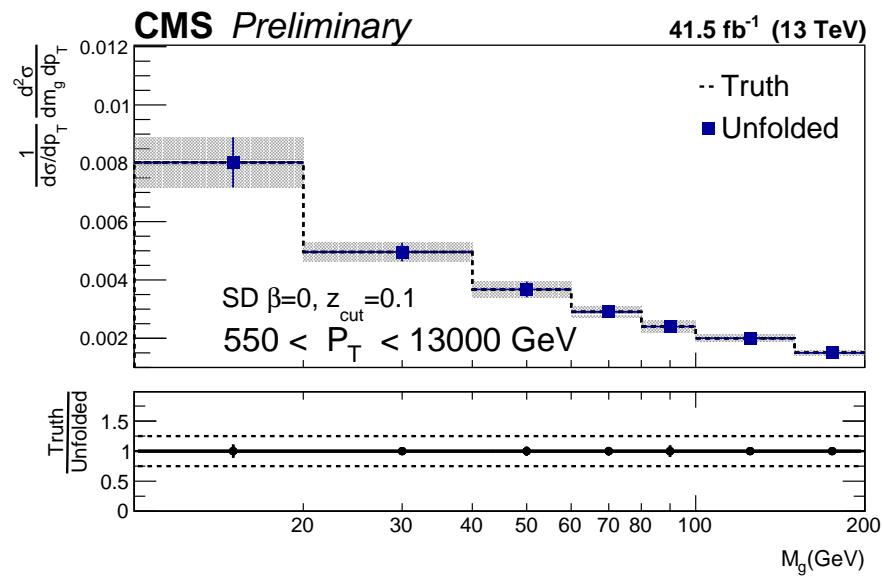


Figure 168: Closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

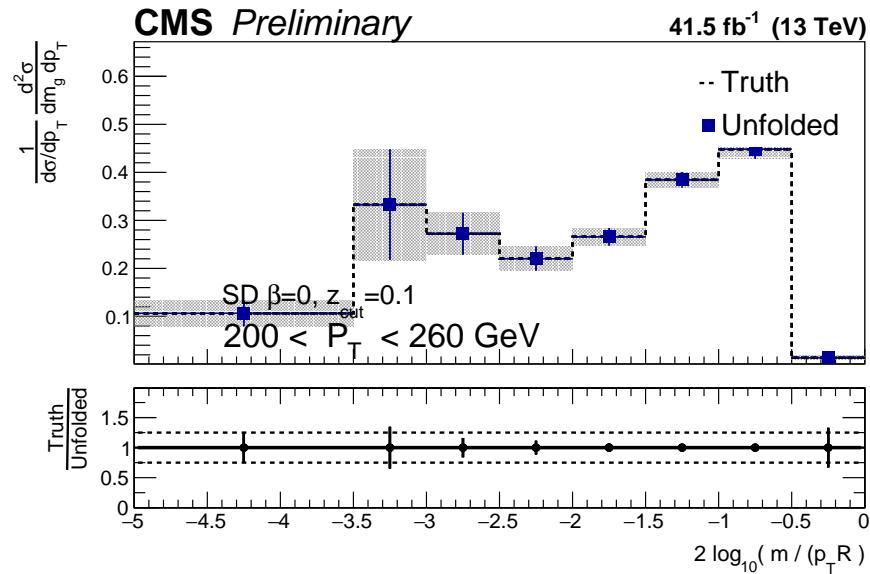


Figure 169: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

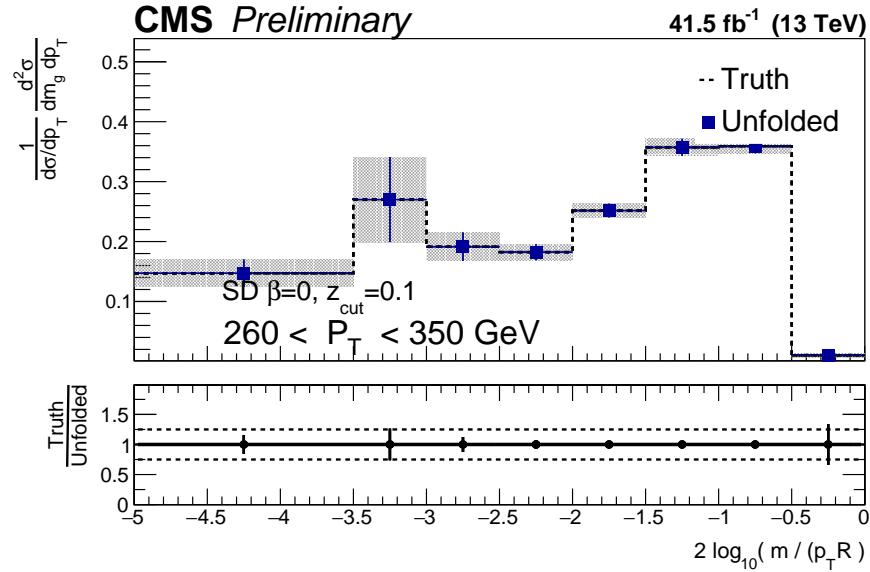


Figure 170: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

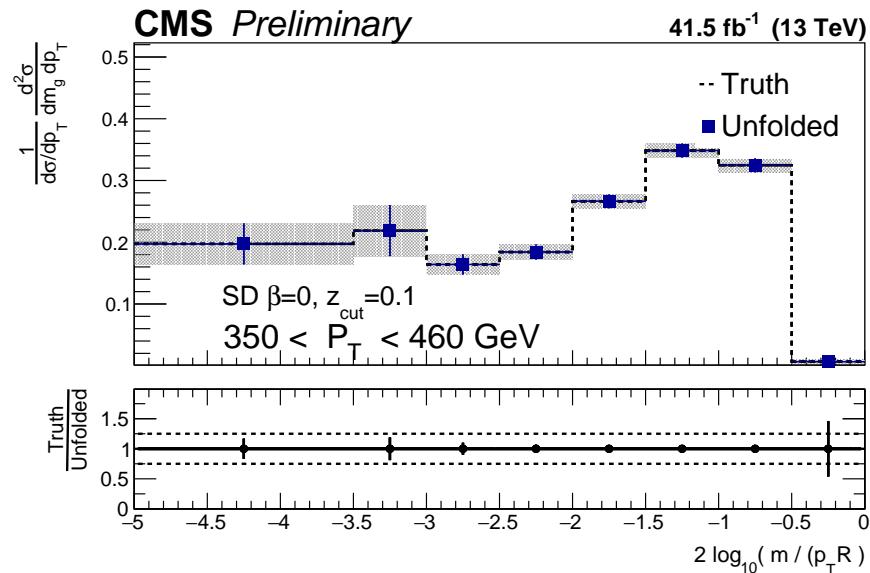


Figure 171: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

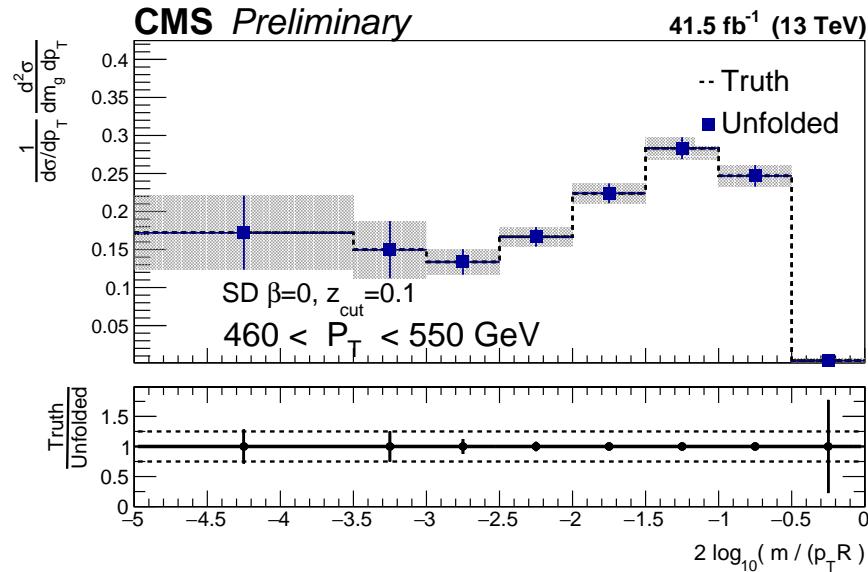


Figure 172: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

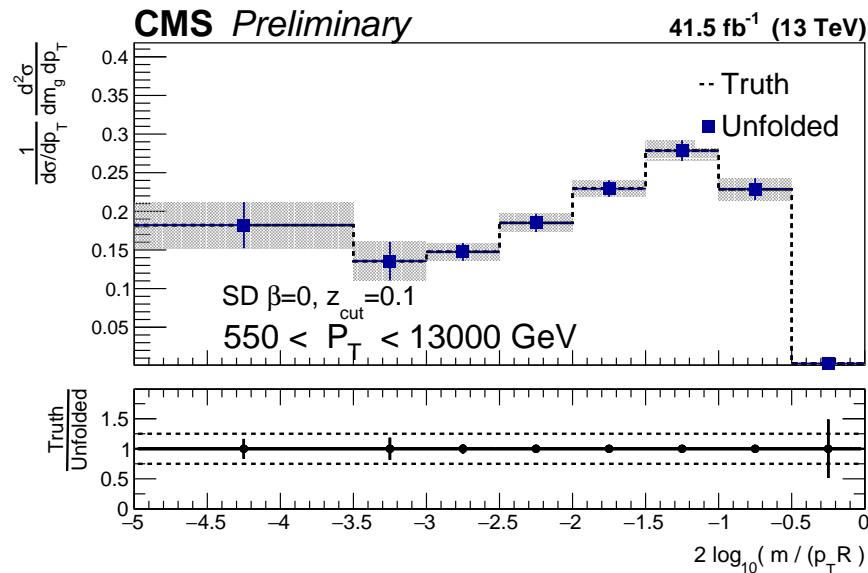


Figure 173: Dimensionless mass closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

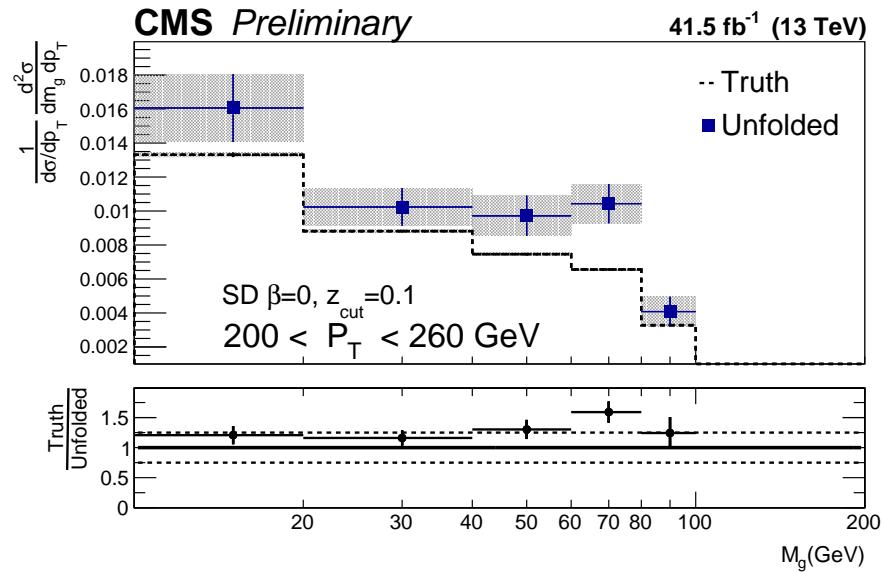


Figure 174: Physical bias test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

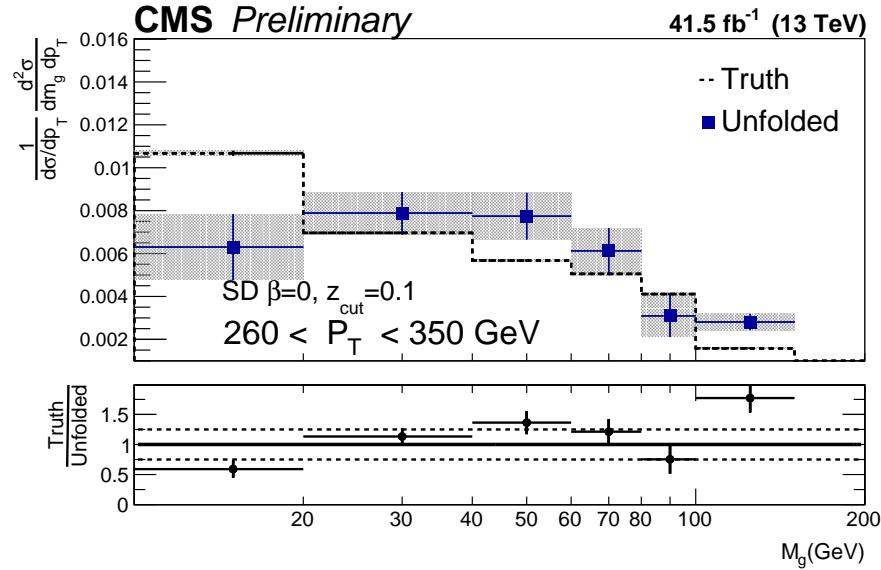


Figure 175: Physical bias test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

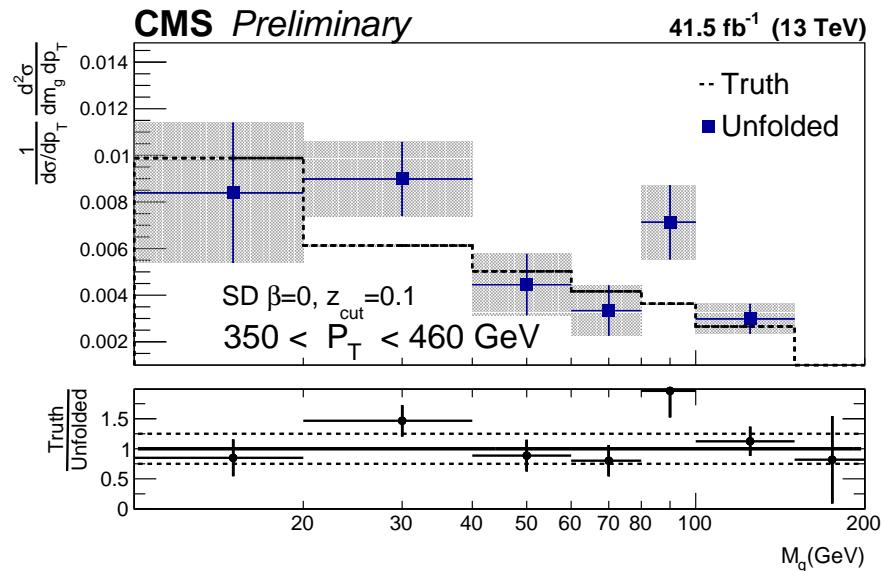


Figure 176: Physical bias test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

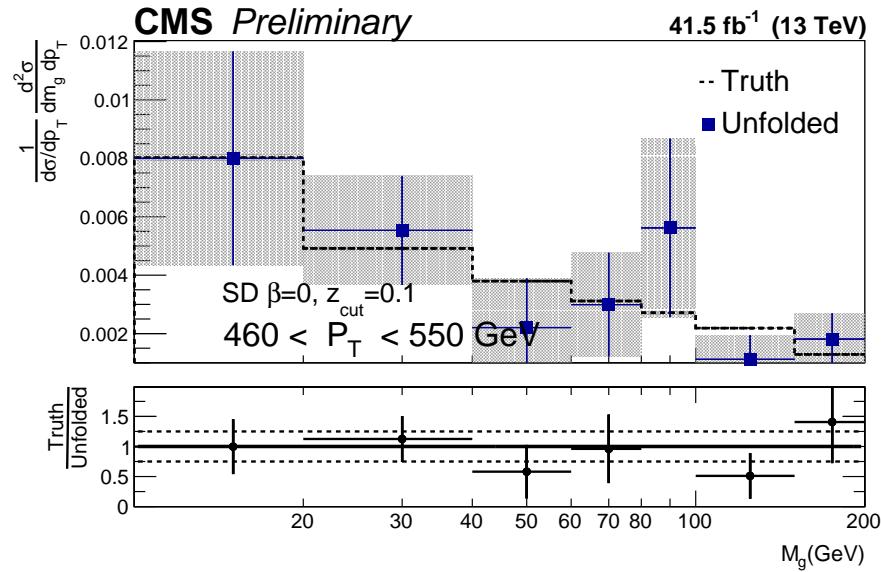


Figure 177: Physical bias test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

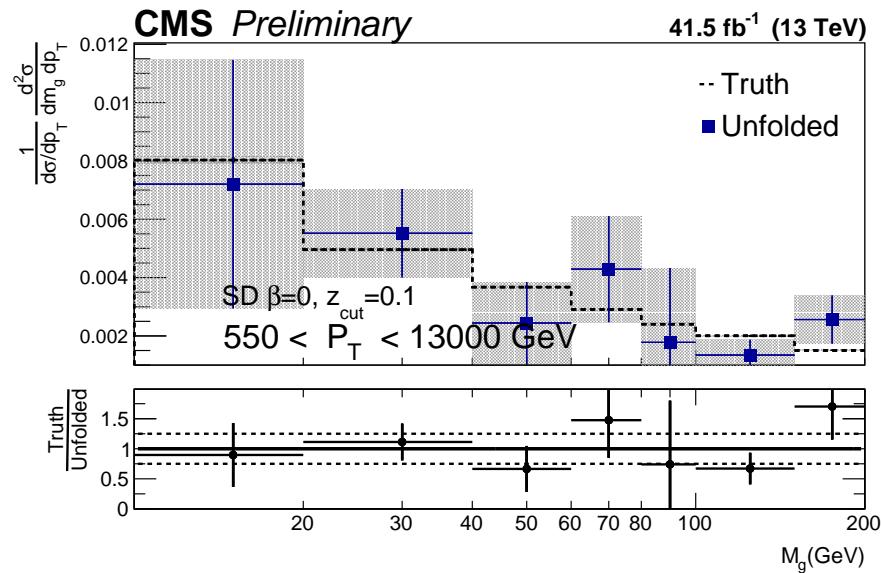


Figure 178: Physical bias test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

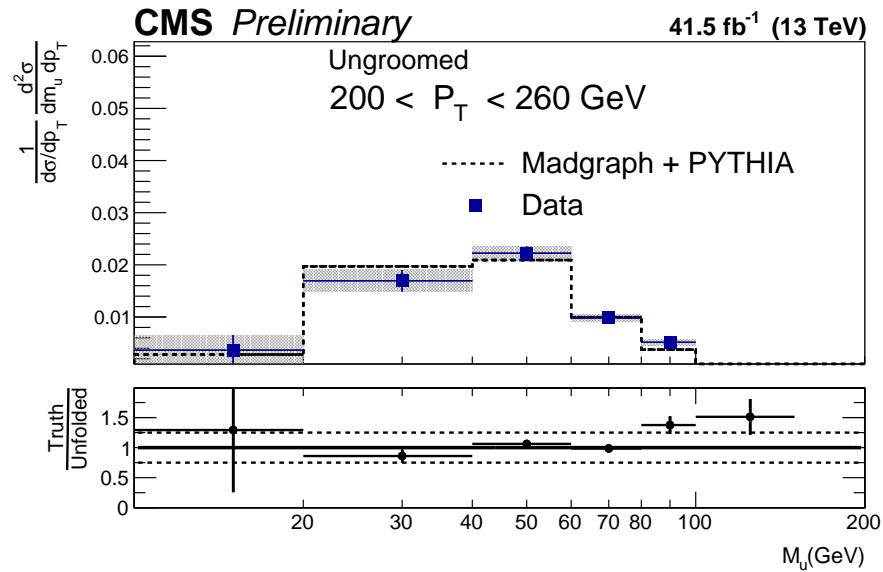


Figure 179: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 200-260 GeV.

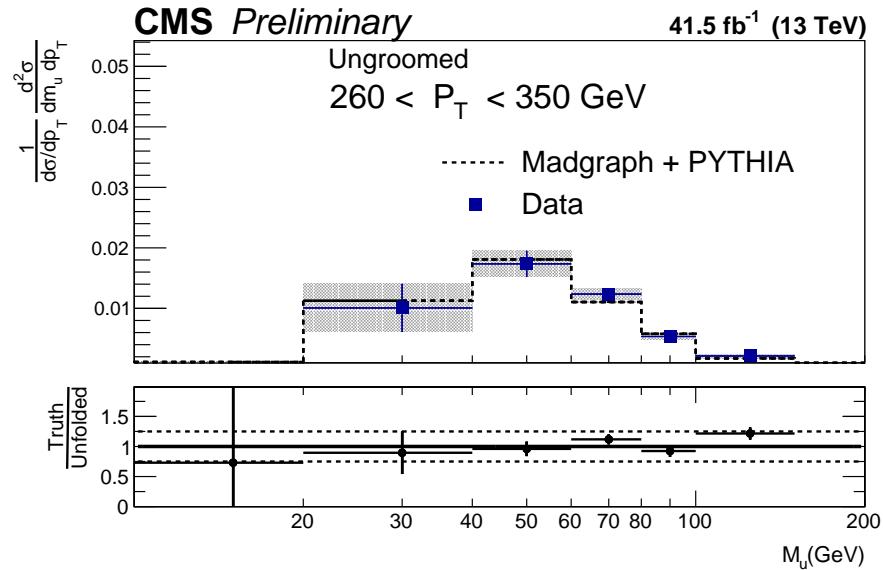


Figure 180: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 260–350 GeV.

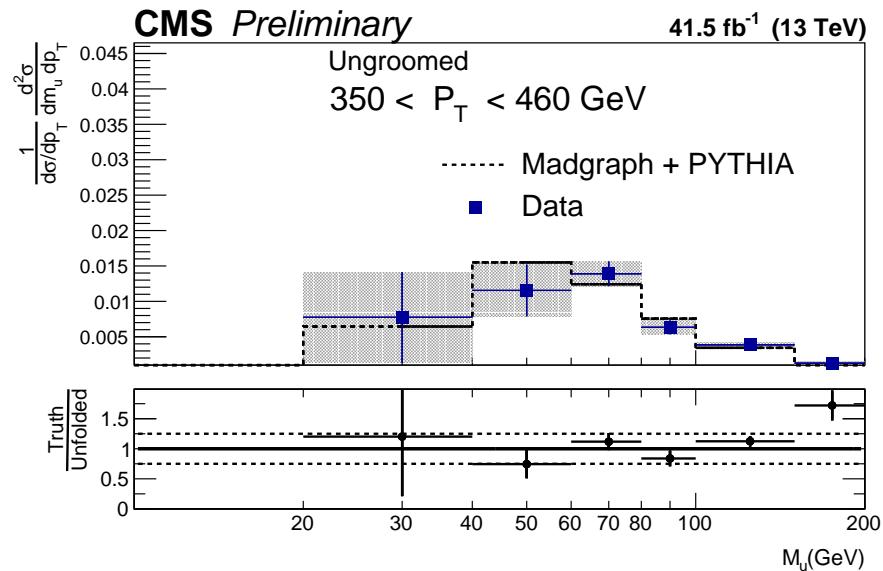


Figure 181: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 350–460 GeV.

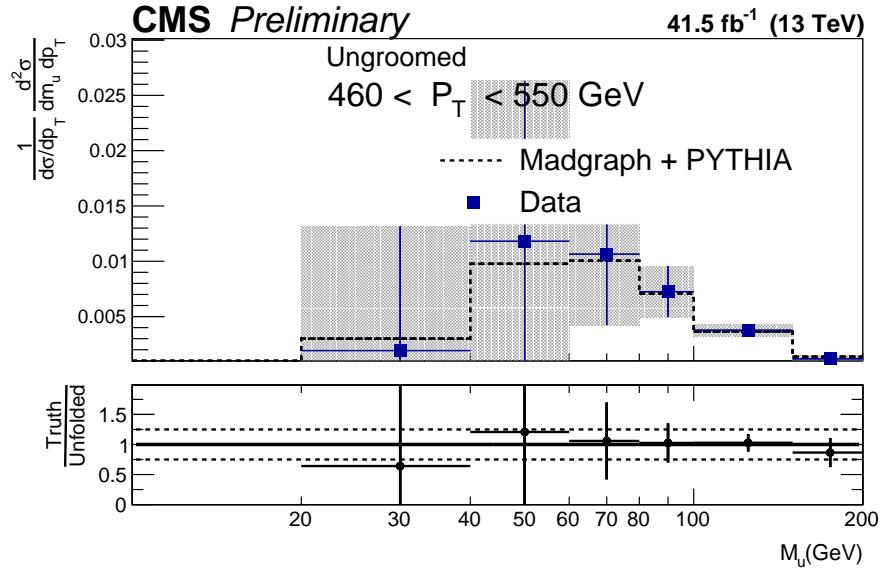


Figure 182: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 460-550 GeV.

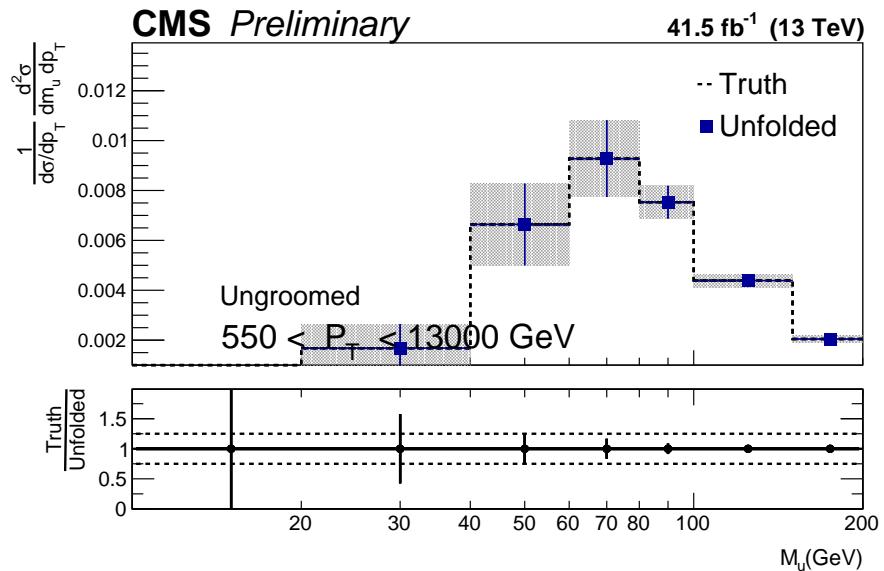


Figure 183: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 550-13000 GeV.

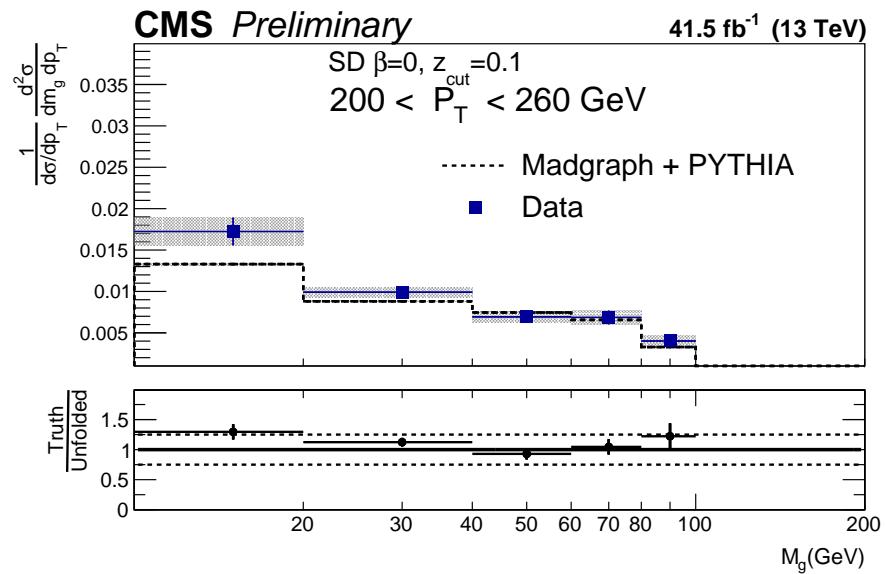


Figure 184: Normalized cross section results with respect to jet mass for groomed jets, p_{T} 200-260 GeV.

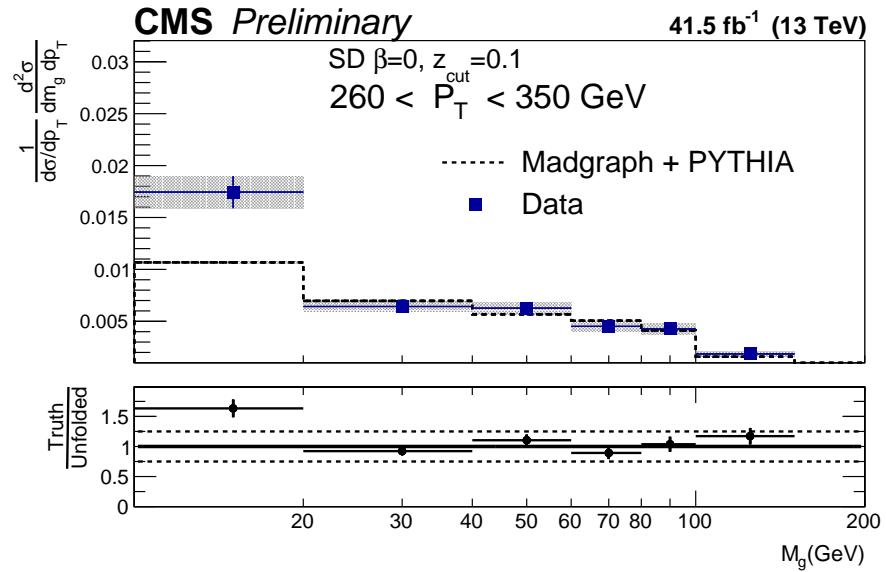


Figure 185: Normalized cross section results with respect to jet mass for groomed jets, p_T 260-350 GeV.

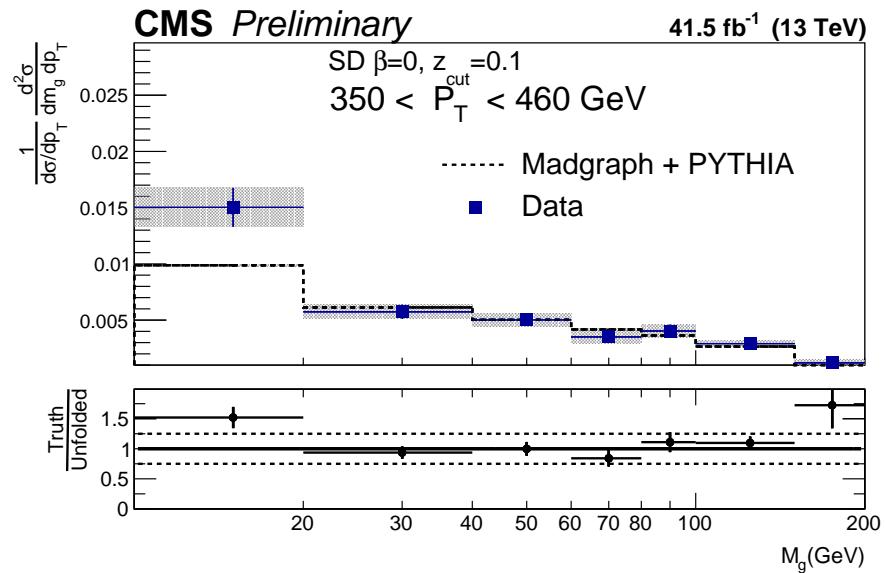


Figure 186: Normalized cross section results with respect to jet mass for groomed jets, p_T 350-460 GeV.

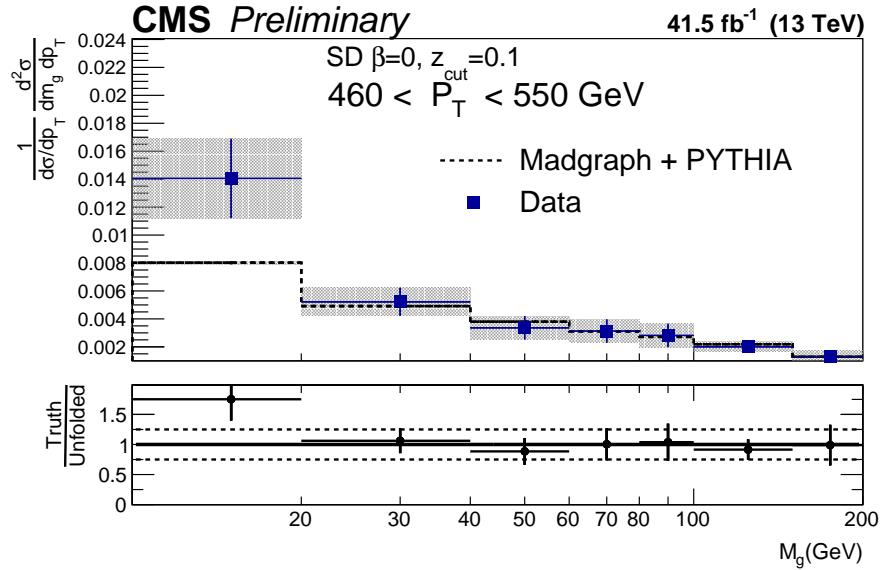


Figure 187: Normalized cross section results with respect to jet mass for groomed jets, p_T 460-550 GeV.

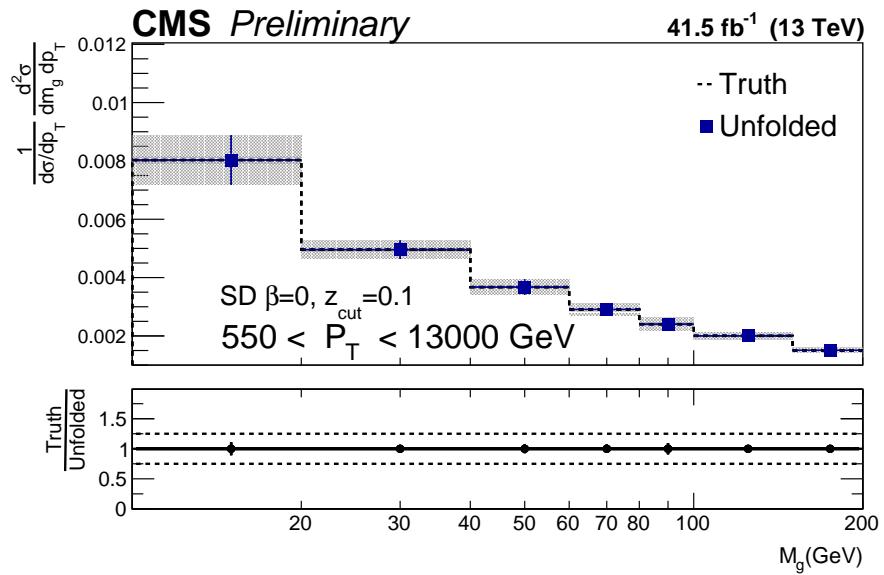


Figure 188: Normalized cross section results with respect to jet mass for groomed jets, p_T 550-13000 GeV.

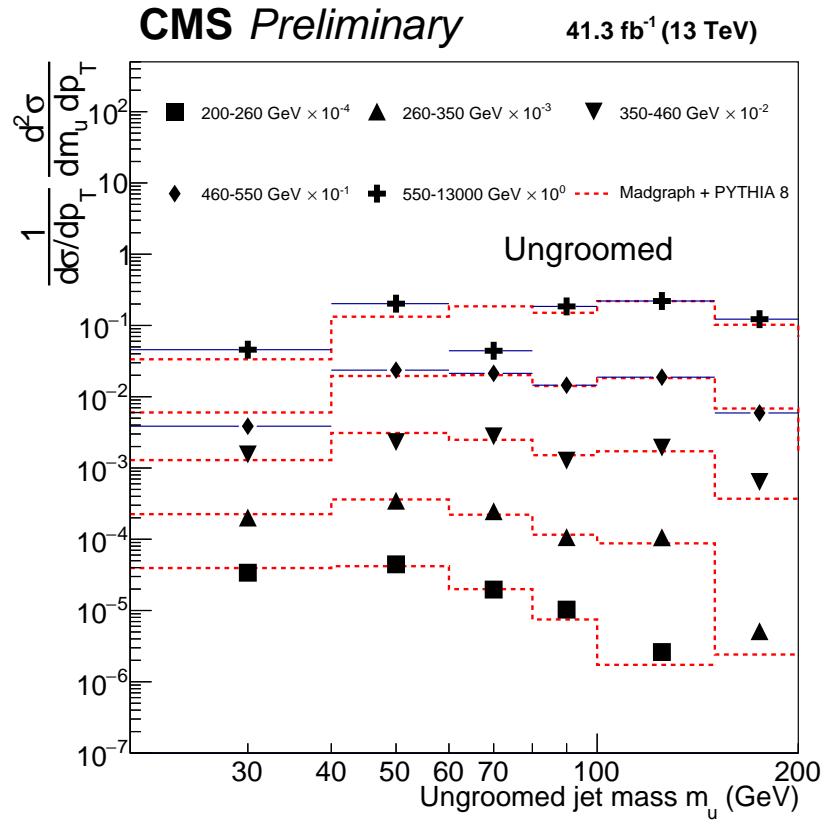


Figure 189: Results for ungroomed reconstructed unfolding with jet mass.

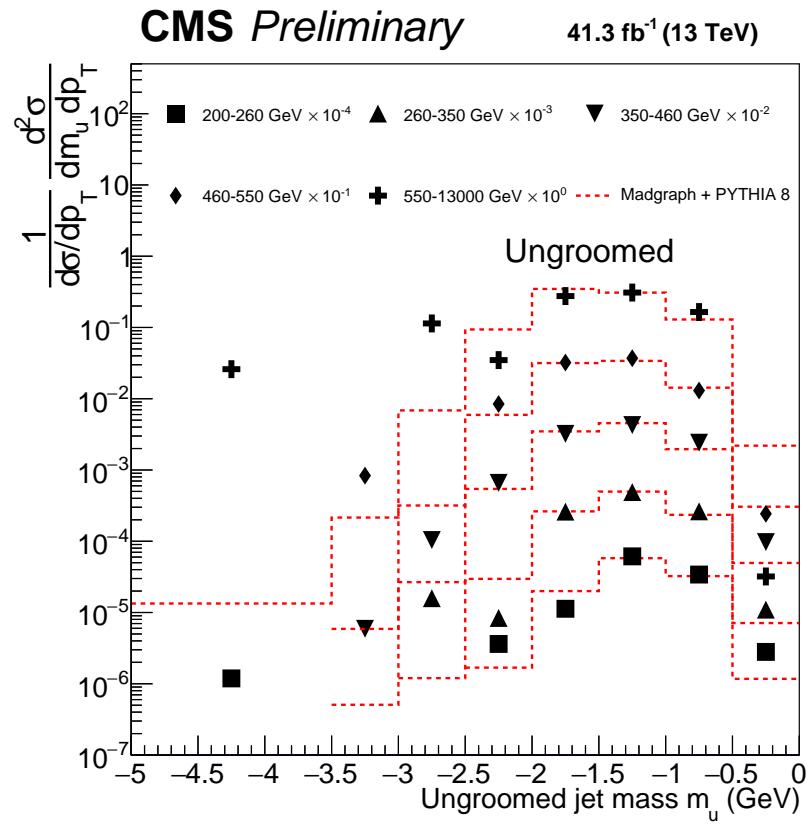


Figure 190: Results for ungroomed reconstructed unfolding with jet dimensionless mass.

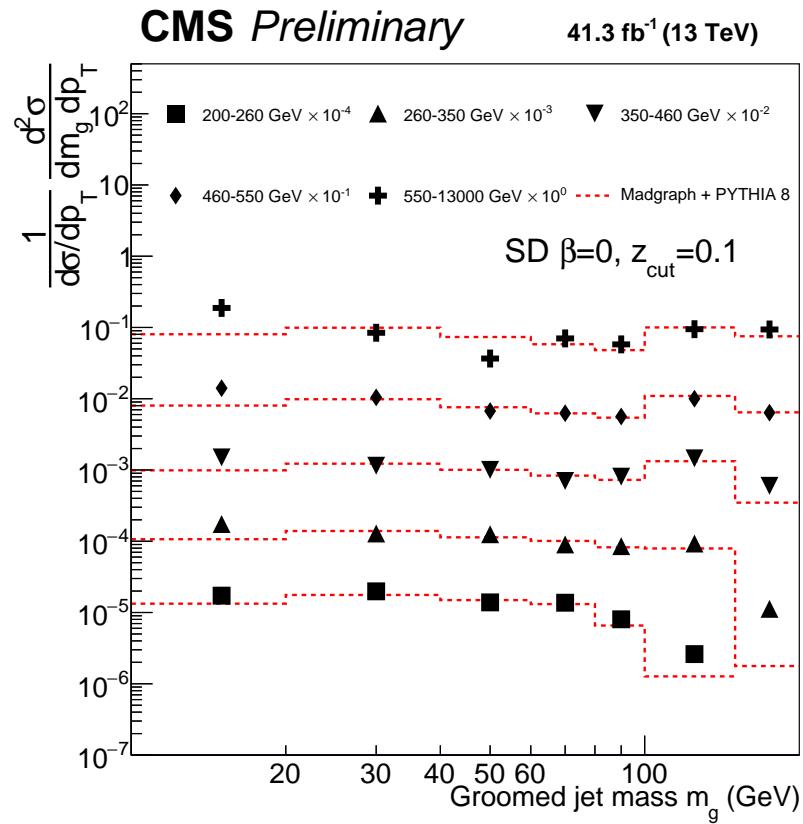


Figure 191: Results for groomed reconstructed unfolding with jet mass.

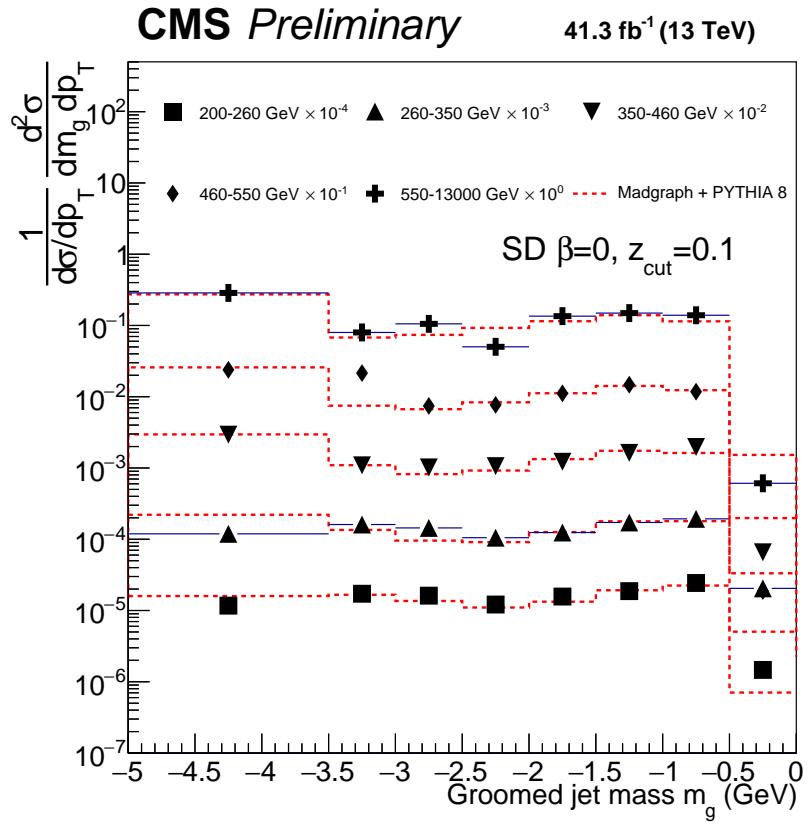


Figure 192: Results for groomed reconstructed unfolding with jet dimensionless mass.

The resulting correlation matrix without systematic uncertainties for the ungroomed jets is shown in Fig. 193, and for the groomed jets is shown in Fig. 195. The same figures with systematic uncertainties are shown for the ungroomed jets in Fig. 197, and for the groomed jets is shown in Fig. 199.

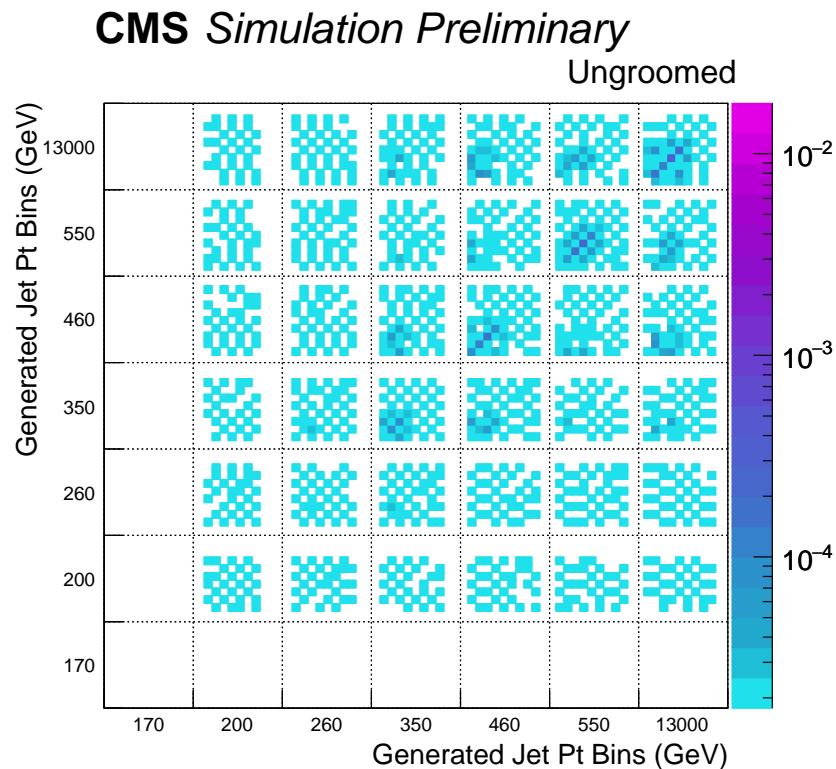


Figure 193: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 10, 20, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 350, \dots, 13000$ GeV).

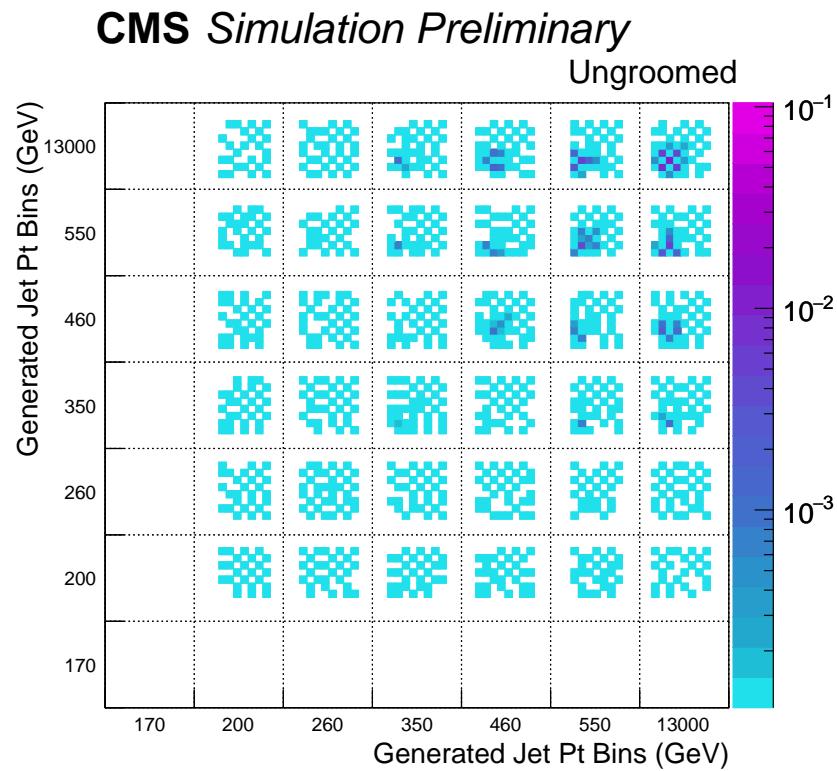


Figure 194: Correlation matrix for ungroomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 350, ... 13000 GeV).

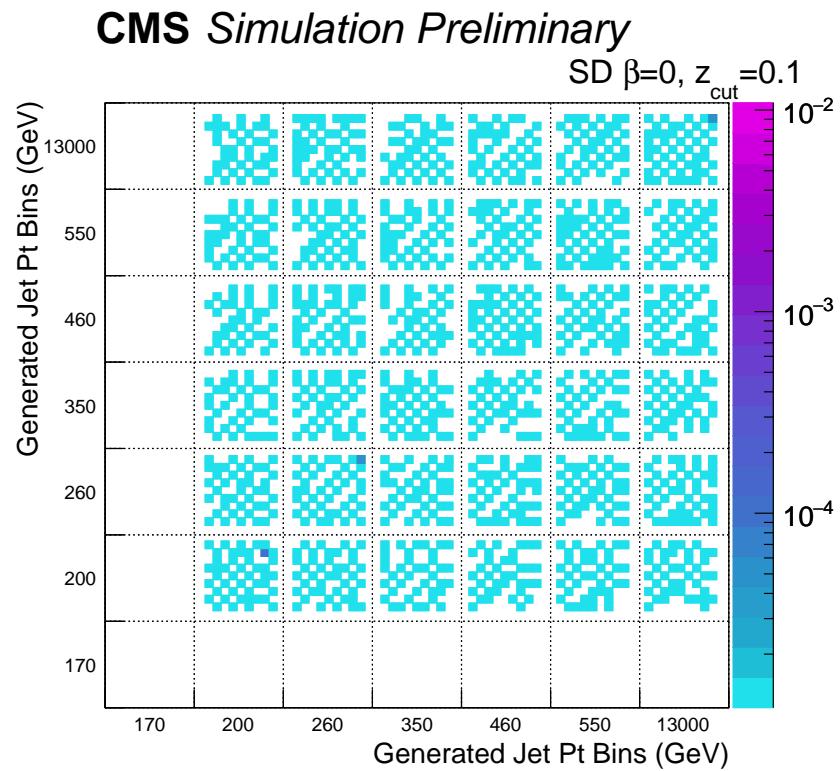


Figure 195: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...760 GeV).

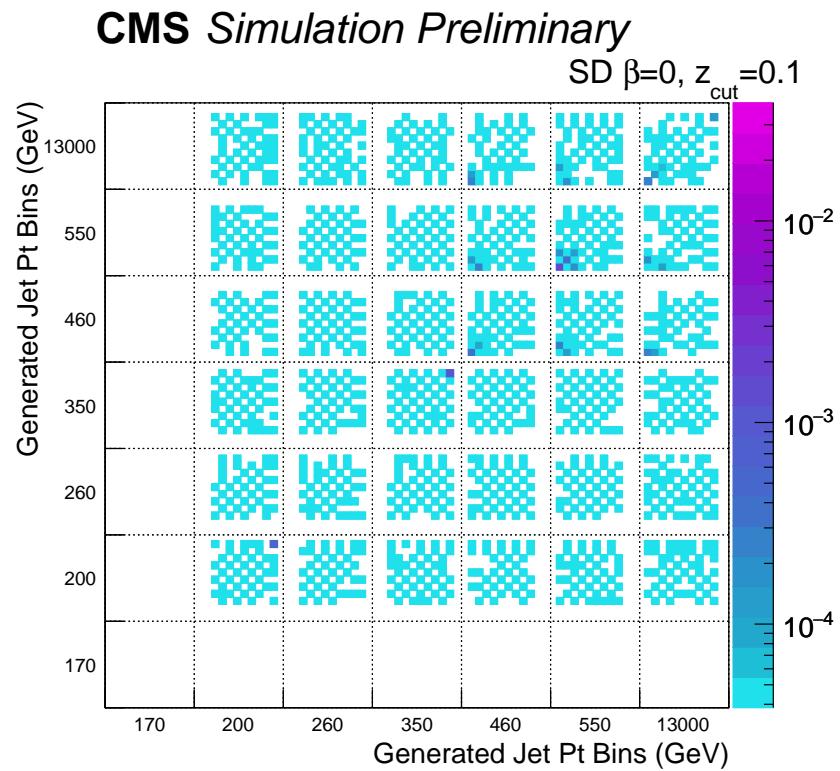


Figure 196: Correlation matrix for groomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 340, ... 760 GeV).

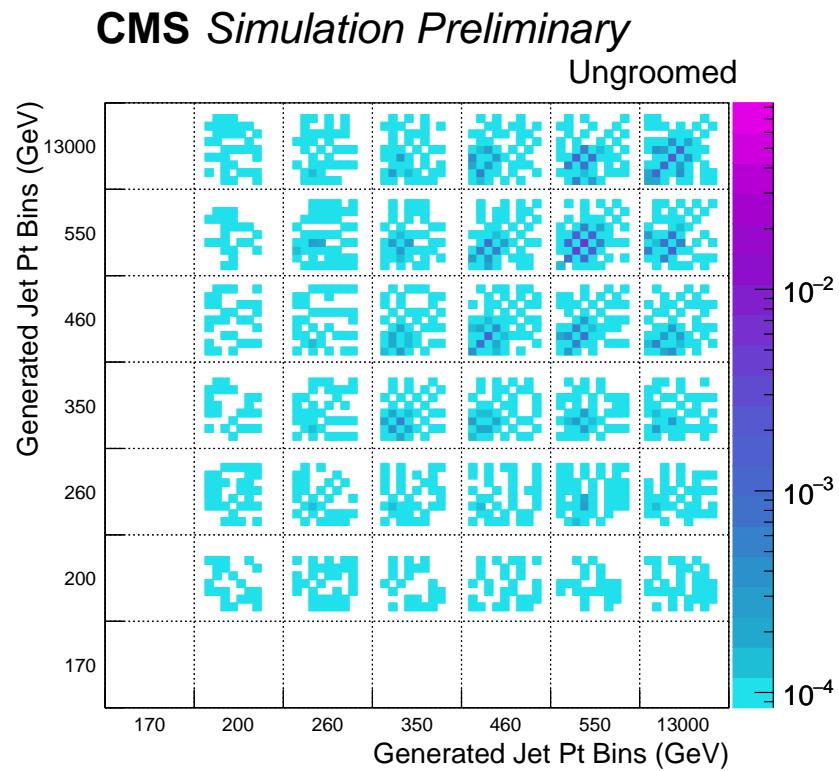


Figure 197: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,260,...13000 GeV).

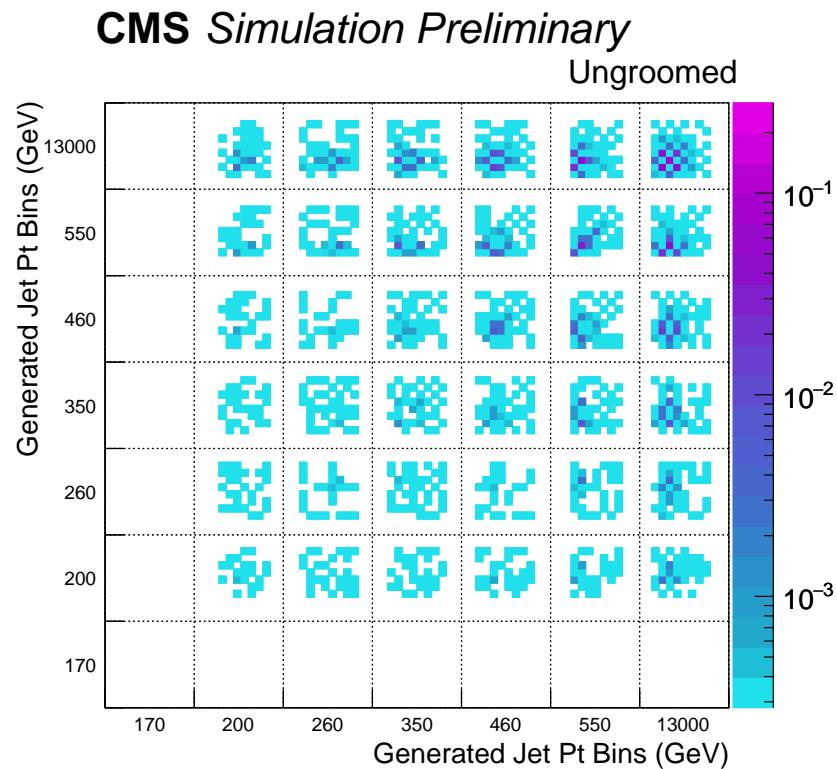


Figure 198: Correlation matrix for ungroomed jets. The bins are set such that the jet dimensionless mass is indexed by the minor blocks (-5, -3.5, -3, ... 0), while the p_T is indexed by the major blocks (200, 260, ... 13000 GeV).

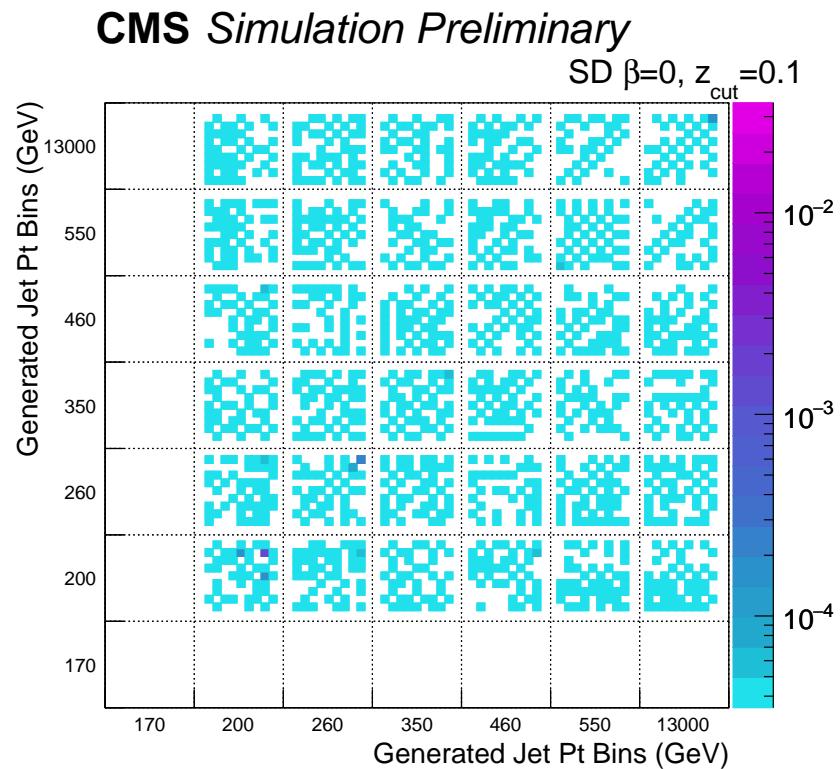


Figure 199: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...0 GeV).

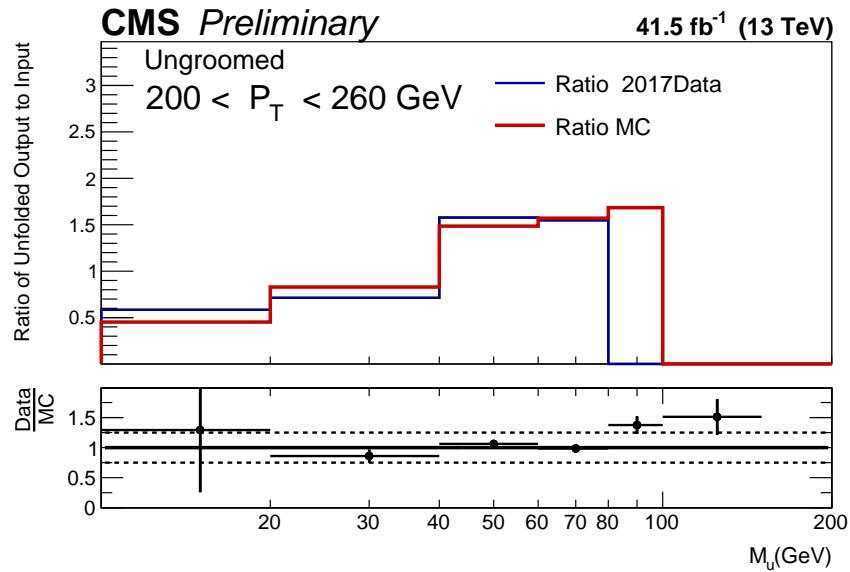


Figure 200: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 200-260 GeV.

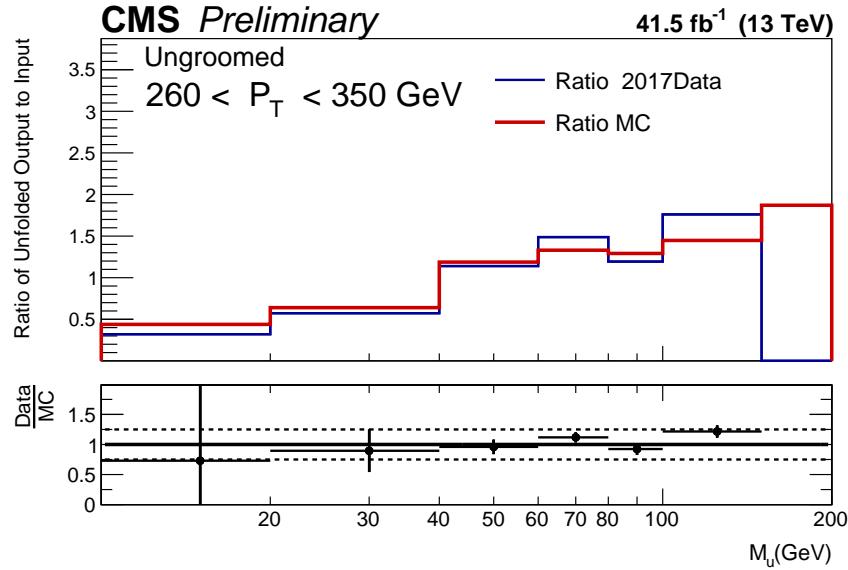


Figure 201: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 260-350 GeV.

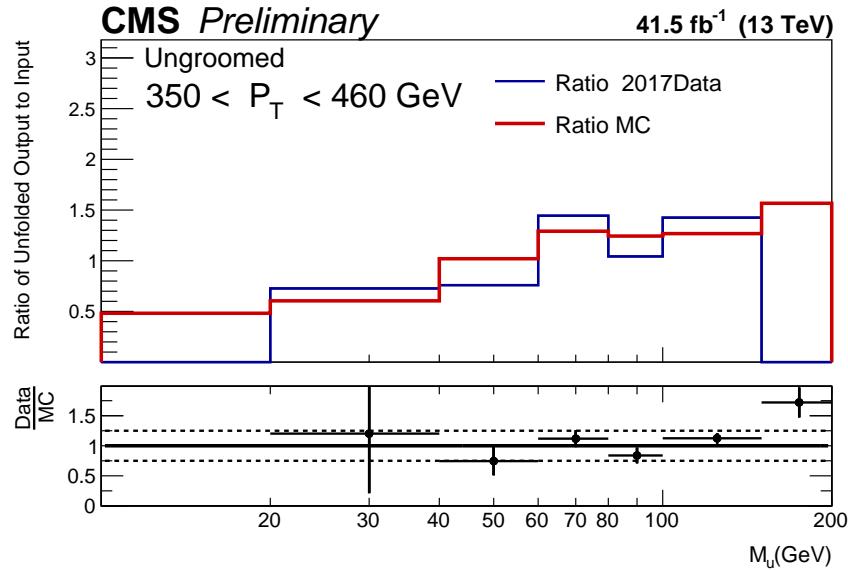


Figure 202: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 350-460 GeV.

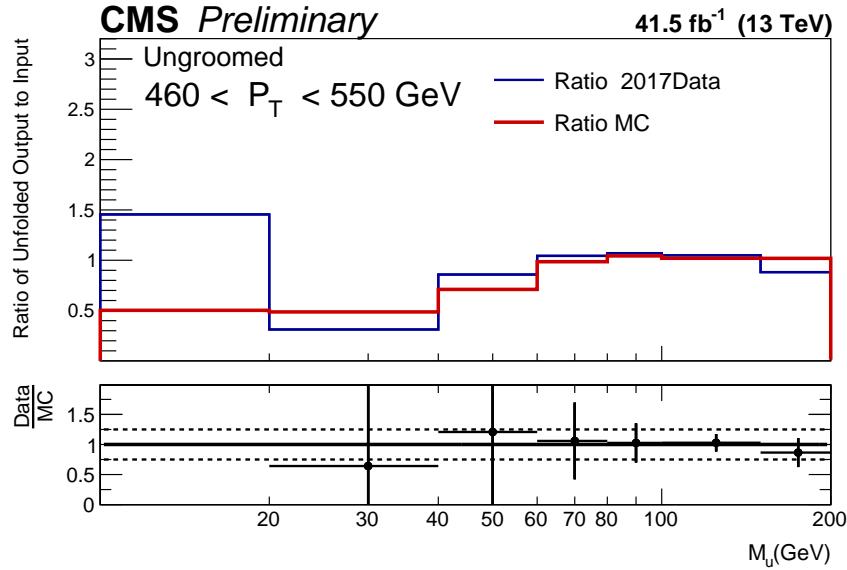


Figure 203: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 460-550 GeV.

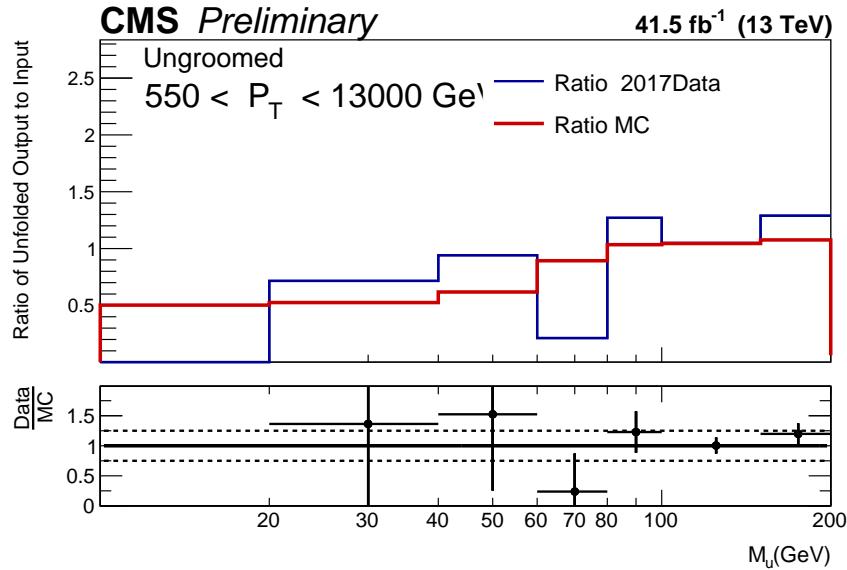


Figure 204: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 550-13000 GeV.

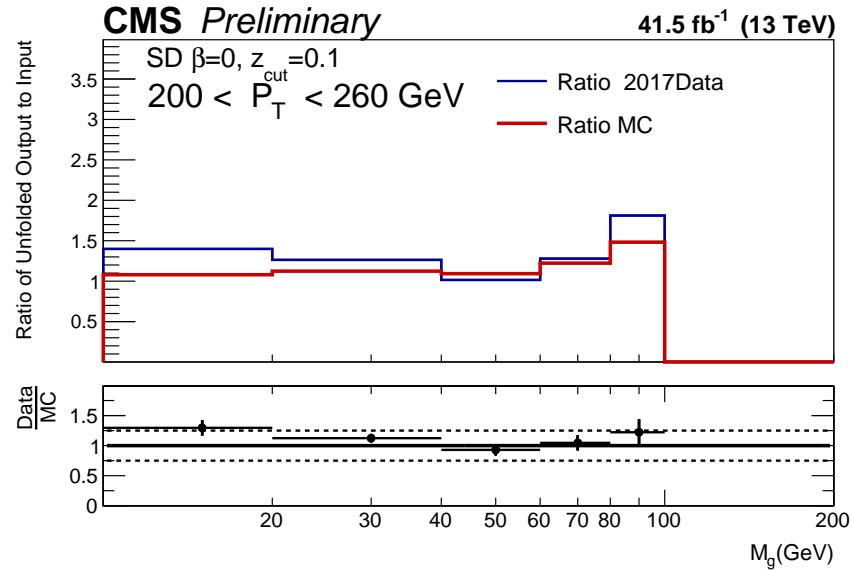


Figure 205: Ratio of unfolded over raw data and MC for groomed jets, p_T 200-260 GeV.

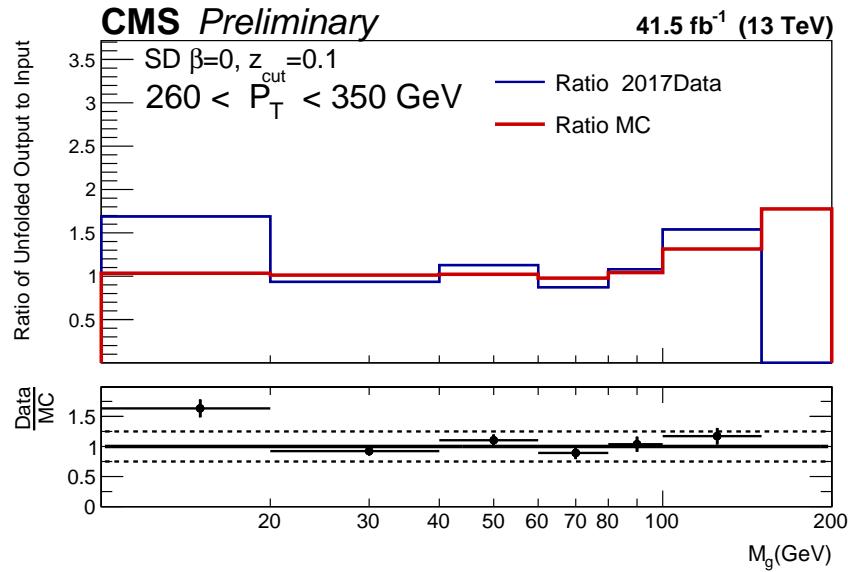


Figure 206: Ratio of unfolded over raw data and MC for groomed jets, p_T 260-350 GeV.

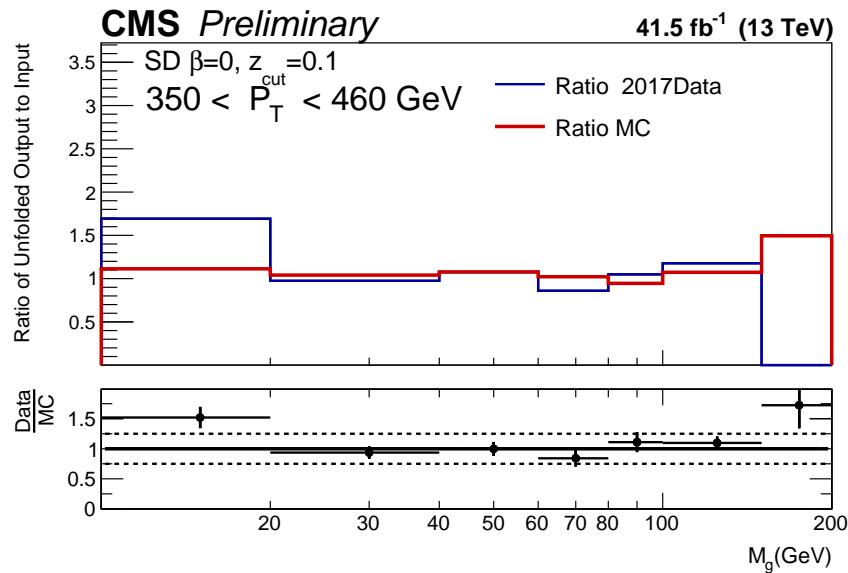


Figure 207: Ratio of unfolded over raw data and MC for groomed jets, p_T 350-460 GeV.

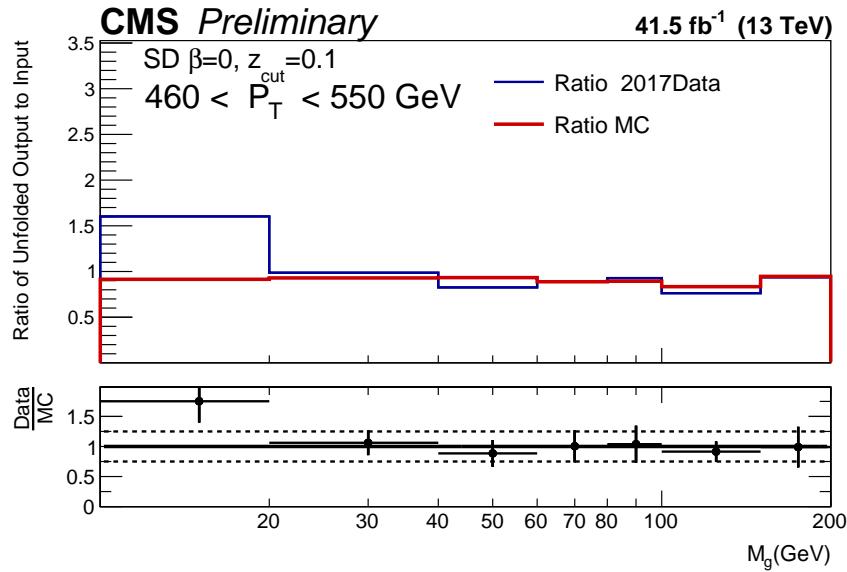


Figure 208: Ratio of unfolded over raw data and MC for groomed jets, p_T 460-550 GeV.

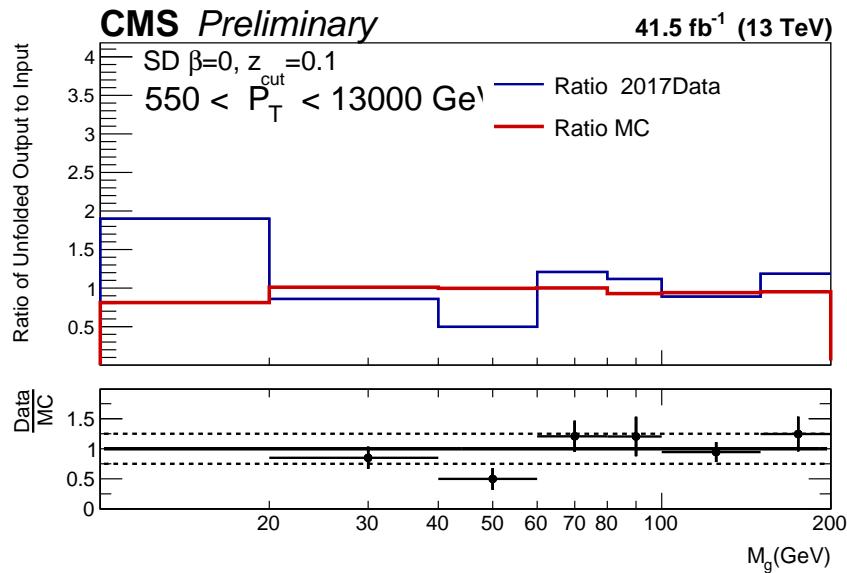


Figure 209: Ratio of unfolded over raw data and MC for groomed jets, p_T 550-13000 GeV.

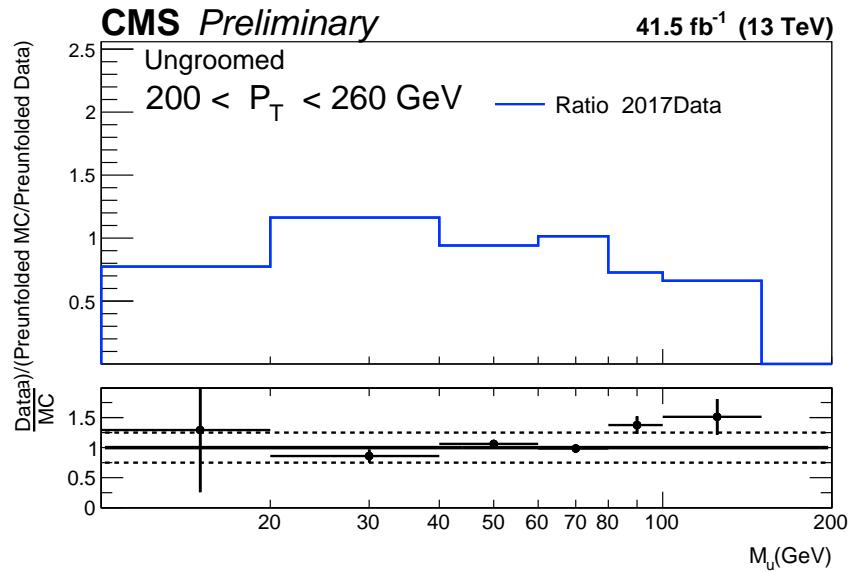


Figure 210: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 200-260 GeV.

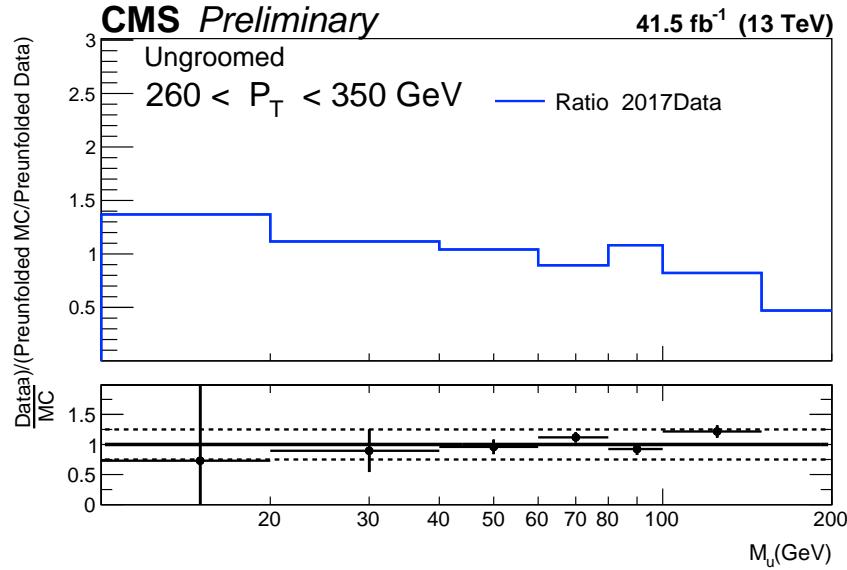


Figure 211: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 260-350 GeV.

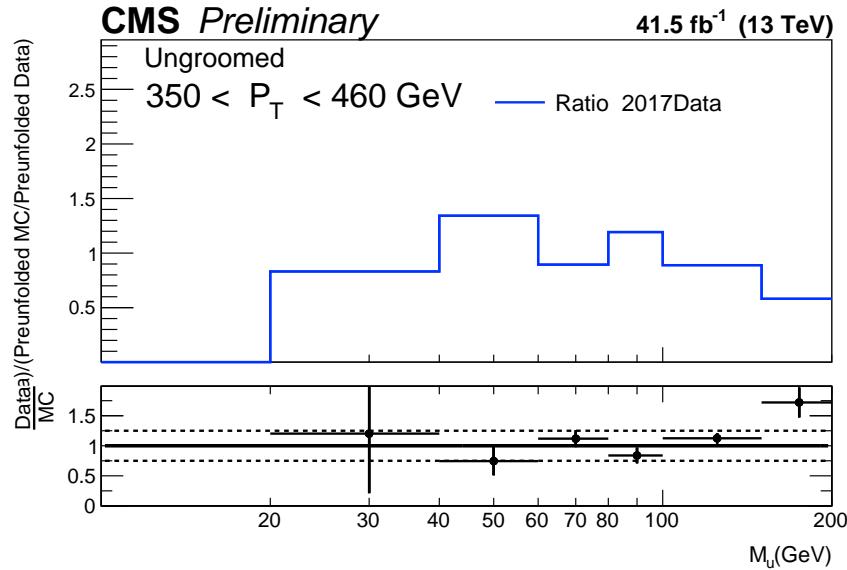


Figure 212: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 350-460 GeV.

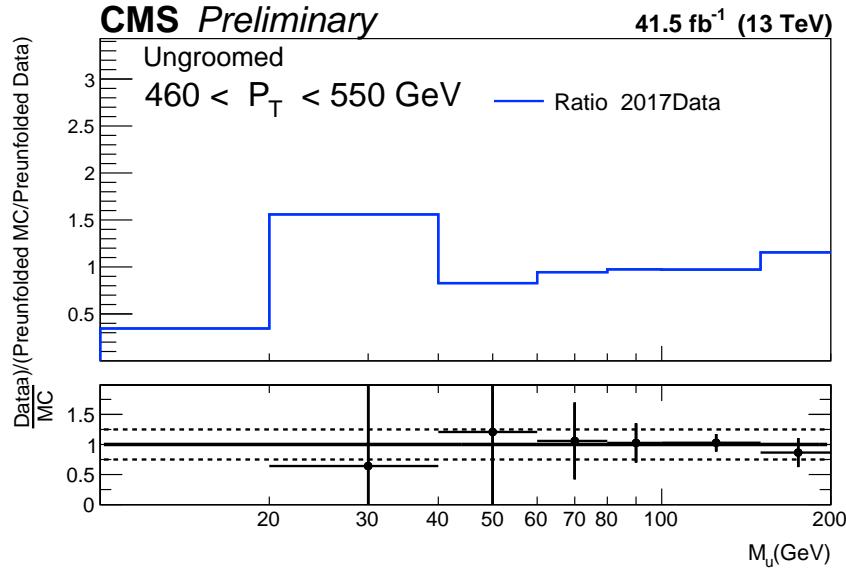


Figure 213: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 460-550 GeV.

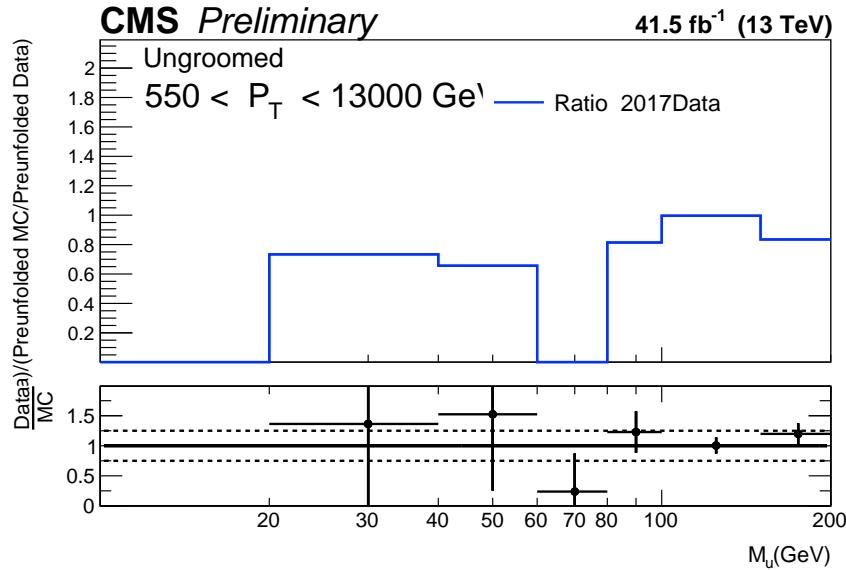


Figure 214: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 550-13000 GeV.

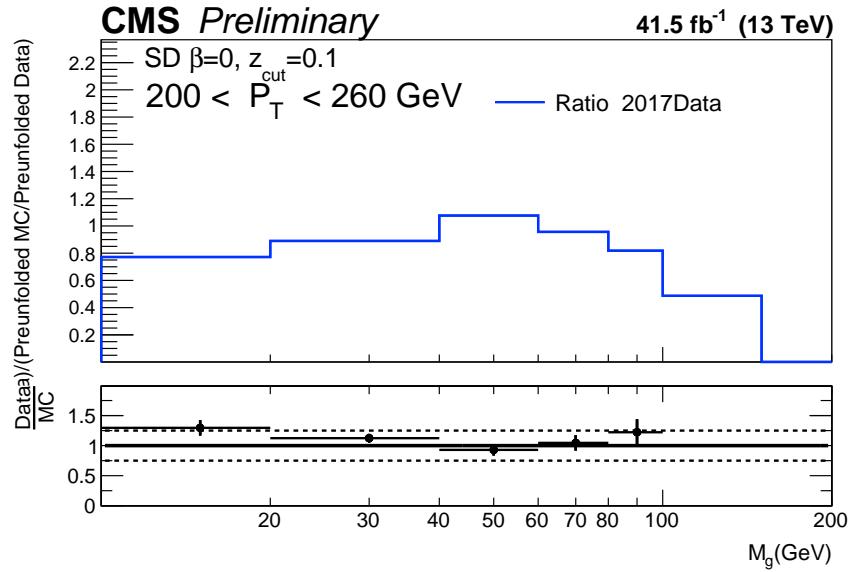


Figure 215: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 200-260 GeV.

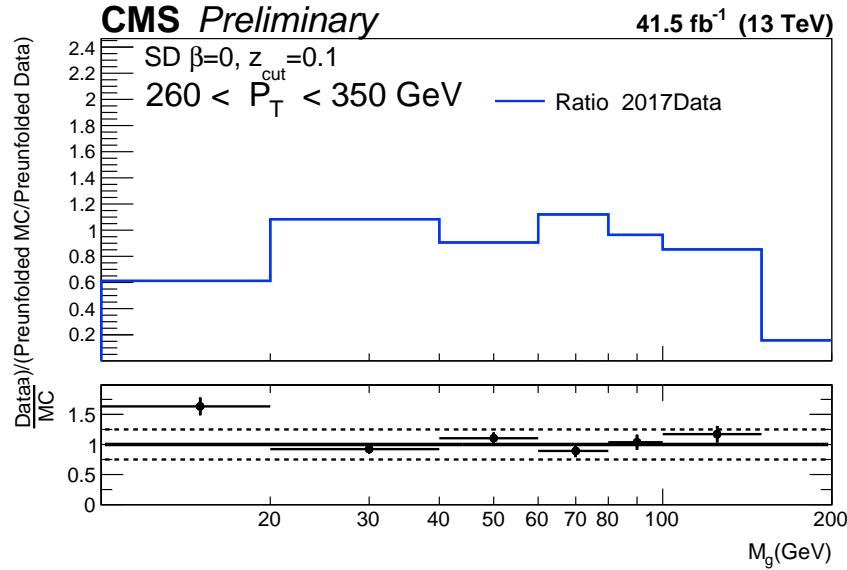


Figure 216: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 260-350 GeV.

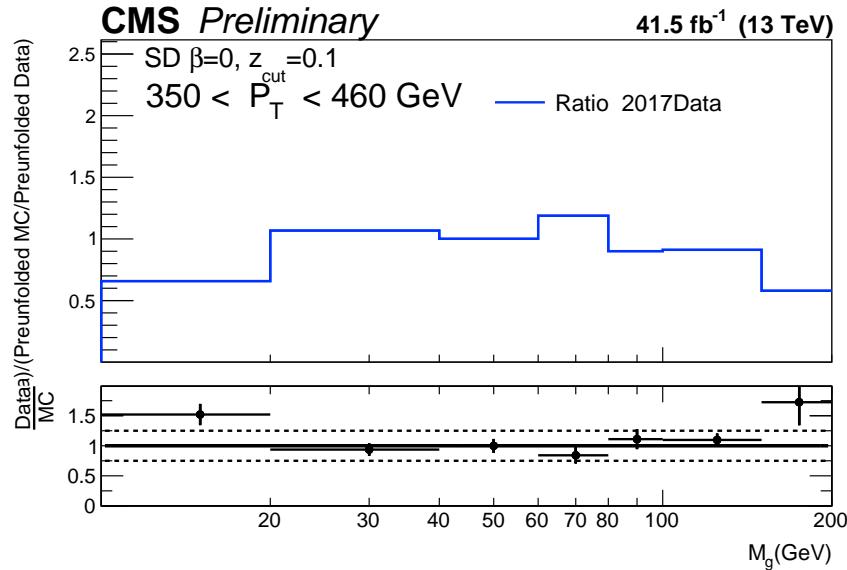


Figure 217: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 350-460 GeV.

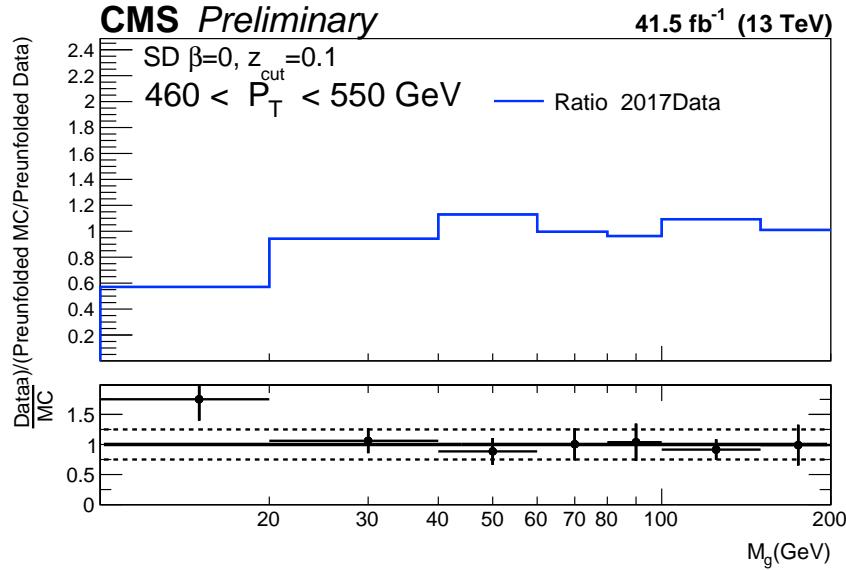


Figure 218: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 460-550 GeV.

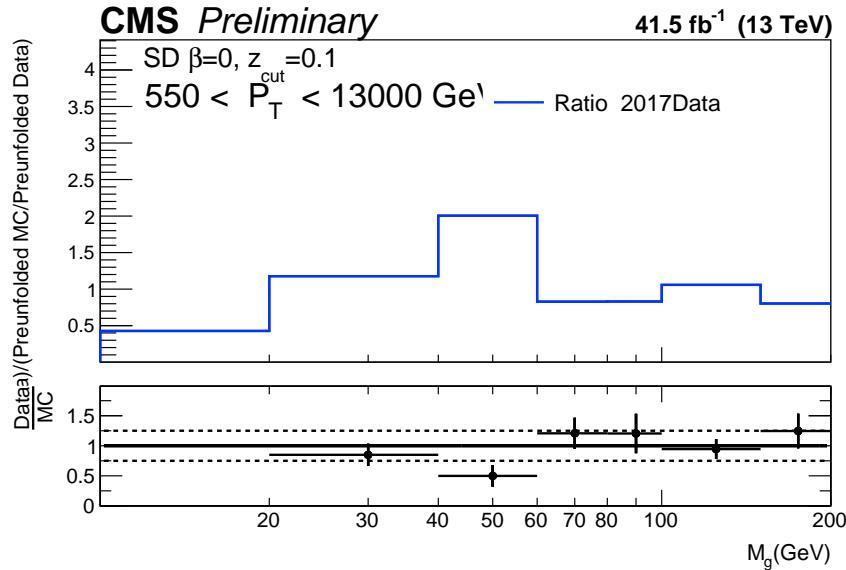


Figure 219: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 550-13000 GeV.

.3 2018 data results

This Appendix shows the distributions from the "Detector Response" through the "Results" sections of the main analysis note with only the 2018 data rather than the full Run 2 statistics seen in the main body of the note.

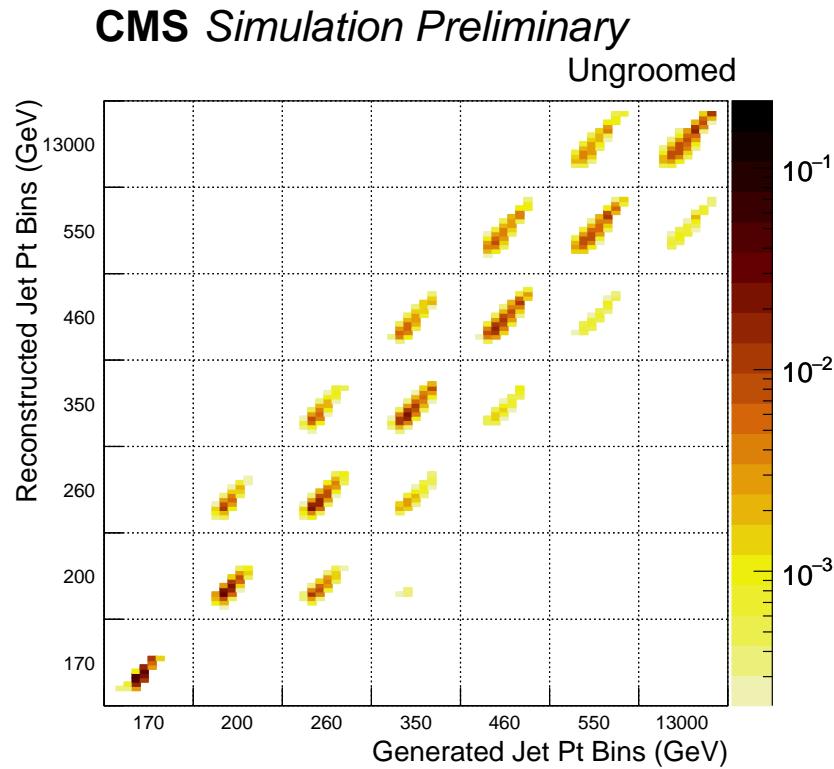


Figure 220: Two-dimensional response matrix for ungroomed jets. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV). While the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

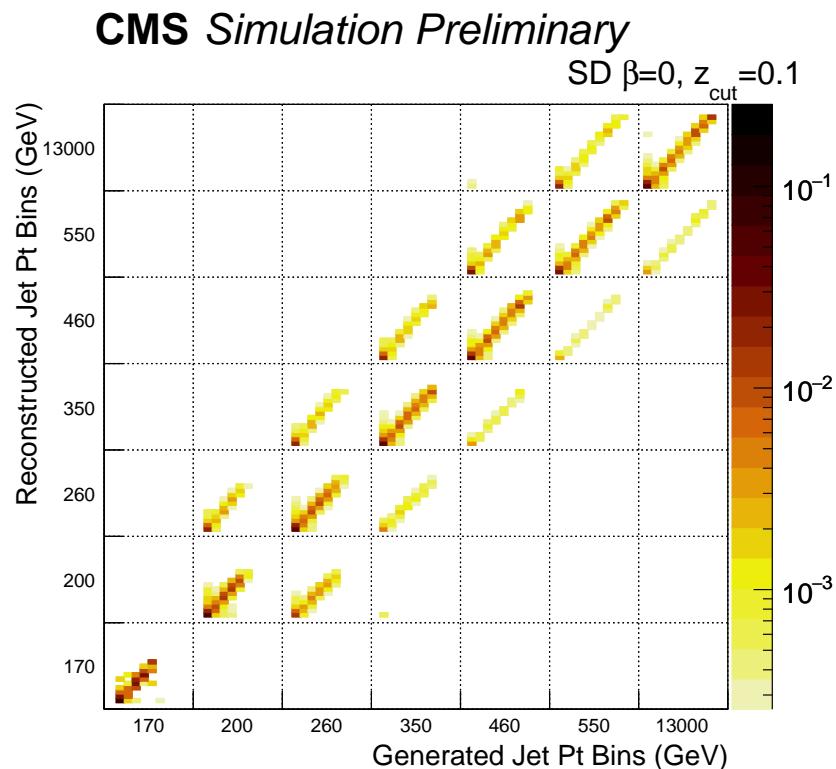


Figure 221: Two-dimensional response matrix for groomed jets $\beta = 0, z_{cut} = 0.1$. The bins are set such that the p_T is indexed by the major blocks (200, 260, 350, 460, 550, 13000 GeV) while the jet mass is indexed by the minor blocks on the x (y) axis 0, 10, 20, 40, 60, 80, 100, 150, 200, 13000 (0, 5, 10, 15, 20, 30, 40, 50, 60, 70, 80, 90, 100, 125, 150, 200, 1000, 13000) GeV.

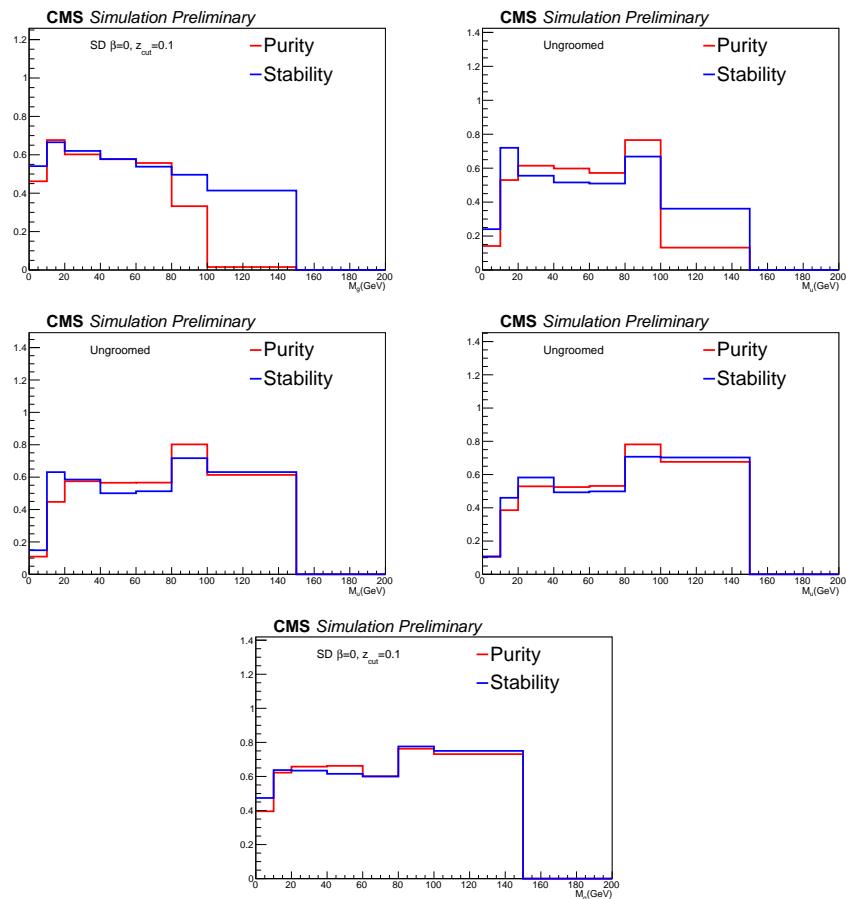


Figure 222: Purity and stability for ungroomed jets in various p_T bins.

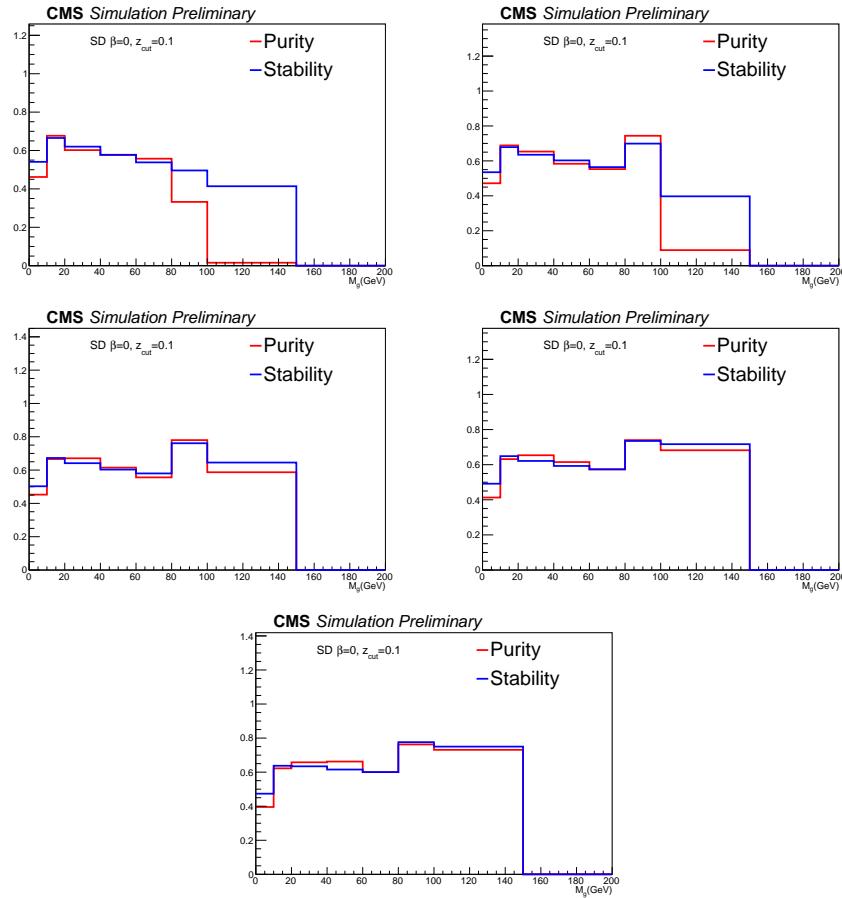


Figure 223: Purity and stability for groomed jets, where the soft-drop criterion was applied with $\beta = 0$ and $z_{cut} = 0.1$, in various p_T bins.

Figure-224 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

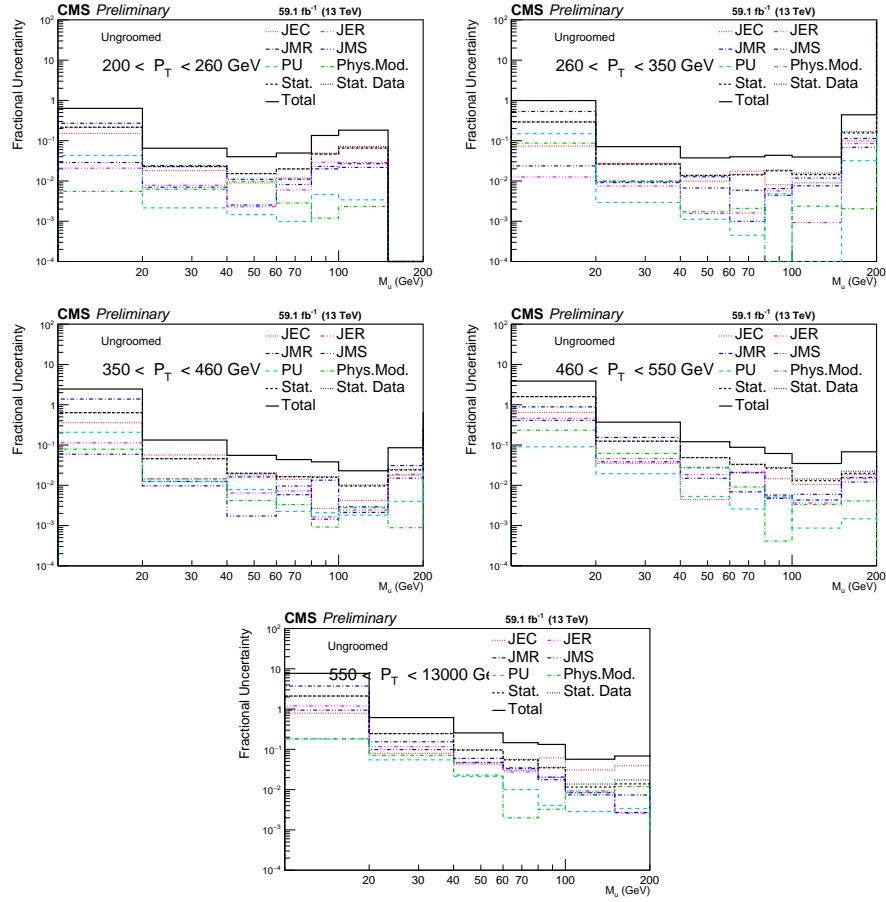


Figure 224: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the raw jet mass.

Figure-225 shows the uncertainty bands for all of the various sources listed above for each jet mass and p_T bin for the normalized cross section.

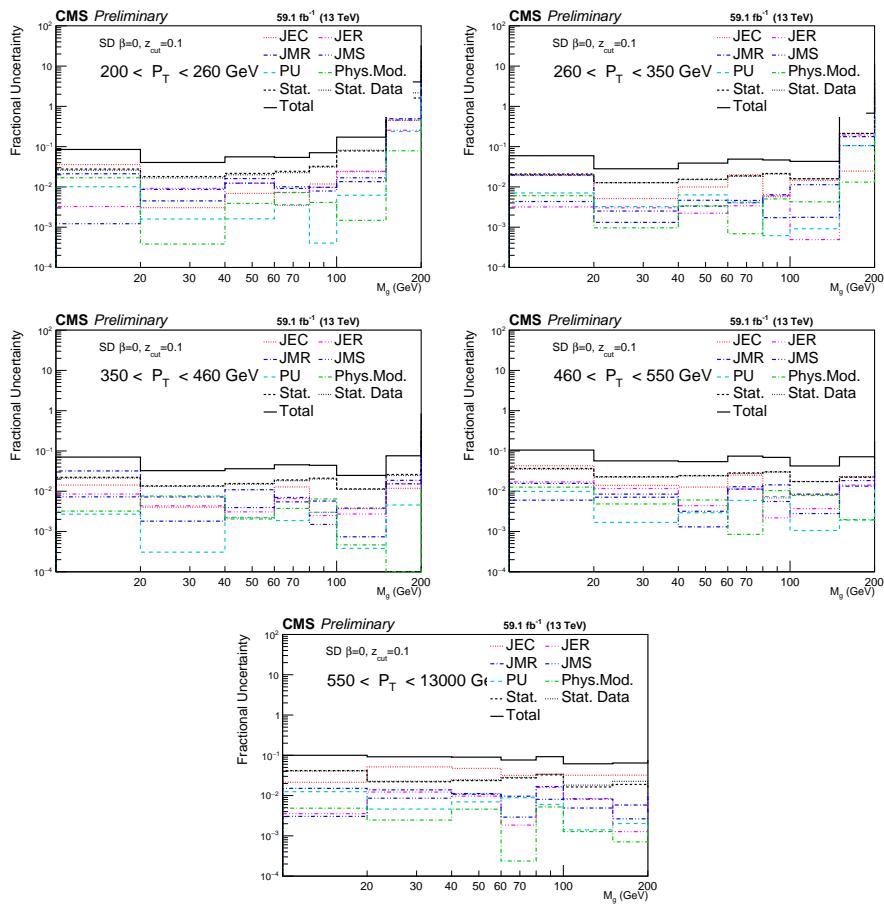


Figure 225: Systematic uncertainties are significantly reduced in the normalized cross section after unfolding for the various p_T bins for the jet mass after grooming with default Soft-Drop.

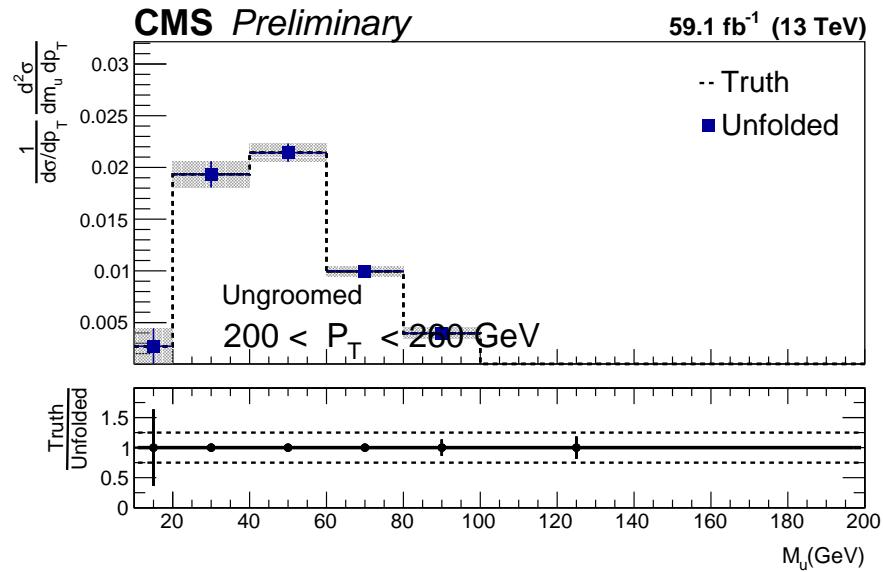


Figure 226: Closure test of ungroomed reconstructed Monte Carlo, p_T 200-260 GeV.

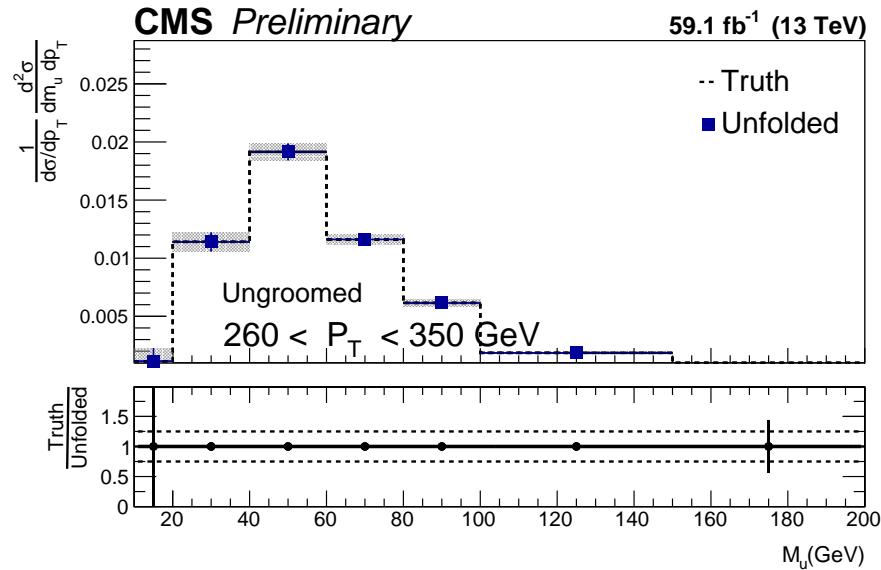


Figure 227: Closure test of ungroomed reconstructed Monte Carlo, p_T 260-350 GeV.

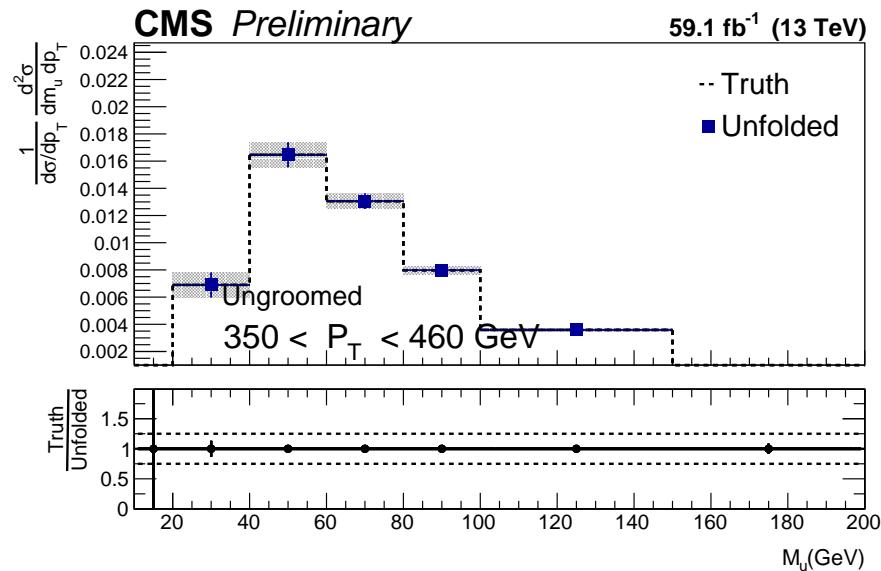


Figure 228: Closure test of ungroomed reconstructed Monte Carlo, p_T 350-460 GeV.

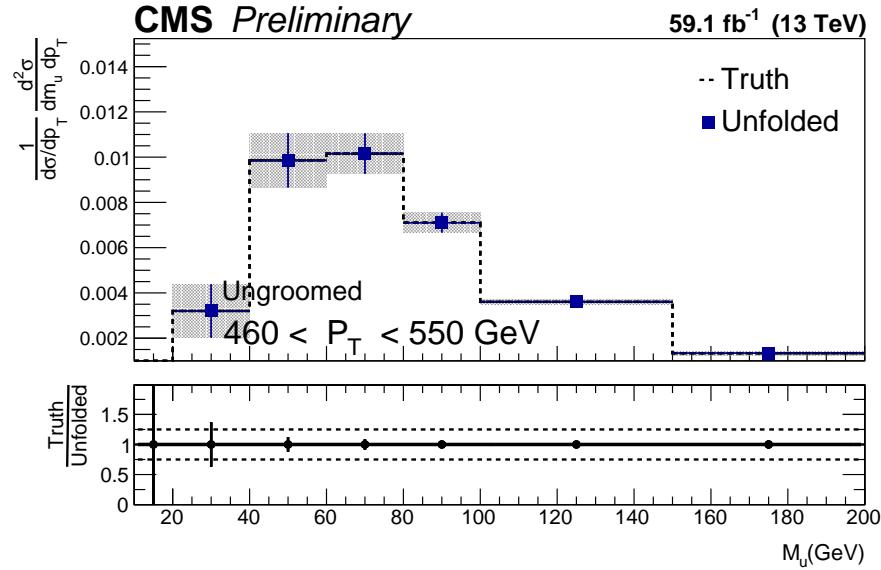


Figure 229: Closure test of ungroomed reconstructed Monte Carlo, p_T 460-550 GeV.

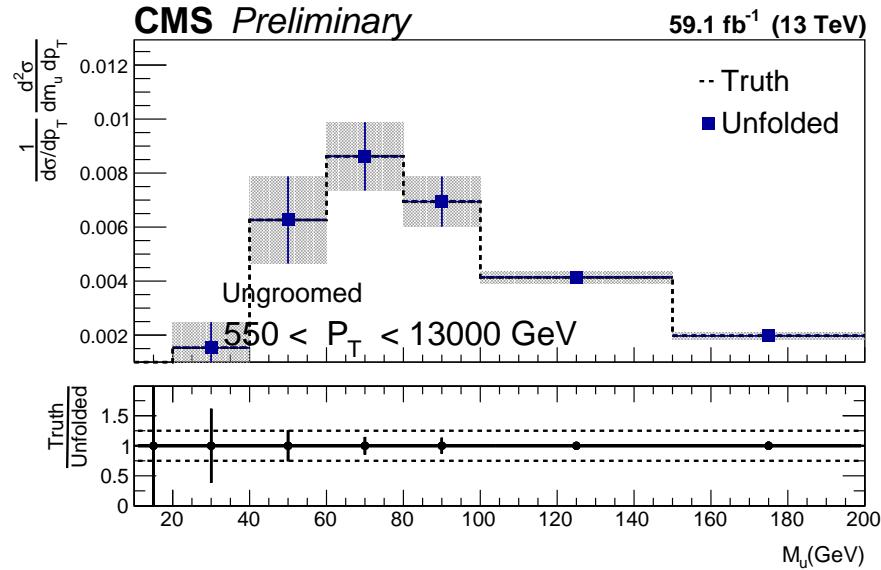


Figure 230: Closure test of ungroomed reconstructed Monte Carlo, p_T 550-13000 GeV.

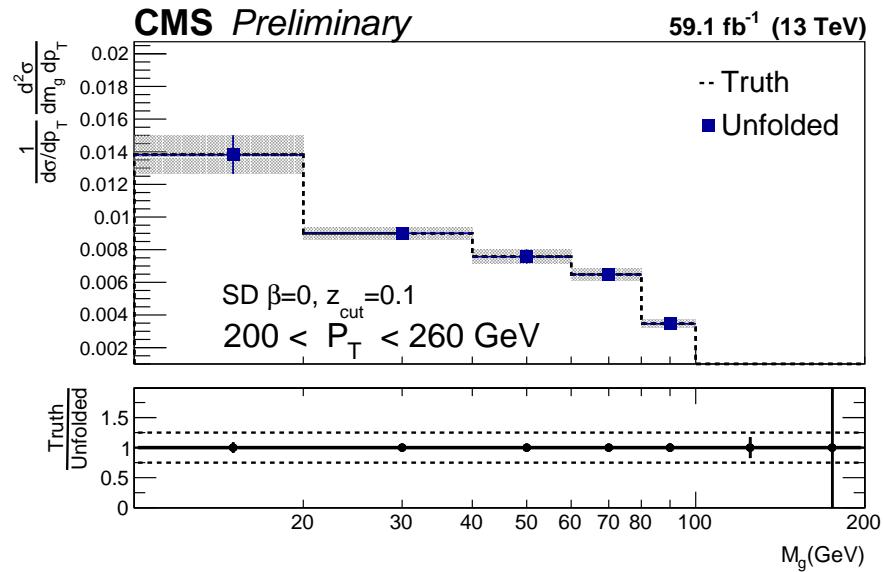


Figure 231: Closure test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

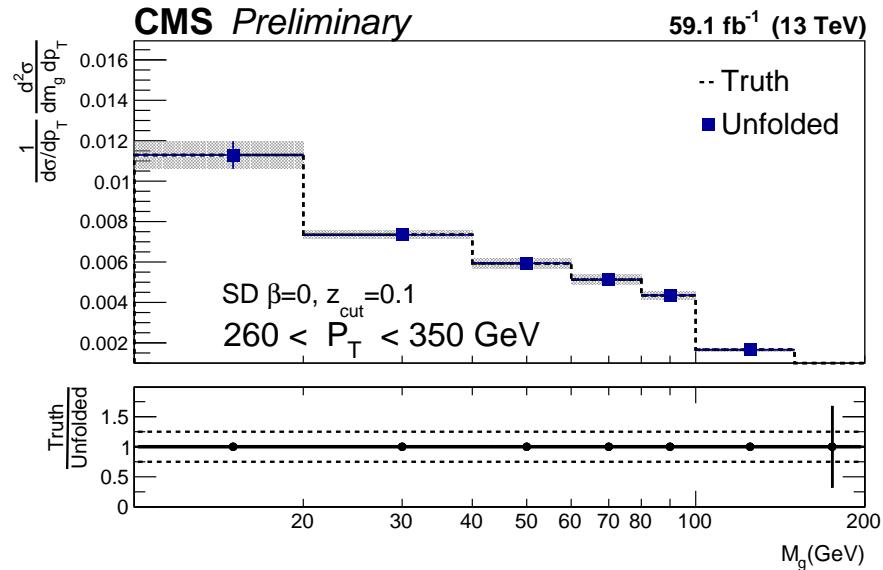


Figure 232: Closure test of groomed reconstructed Monte Carlo, p_T 260–350 GeV.

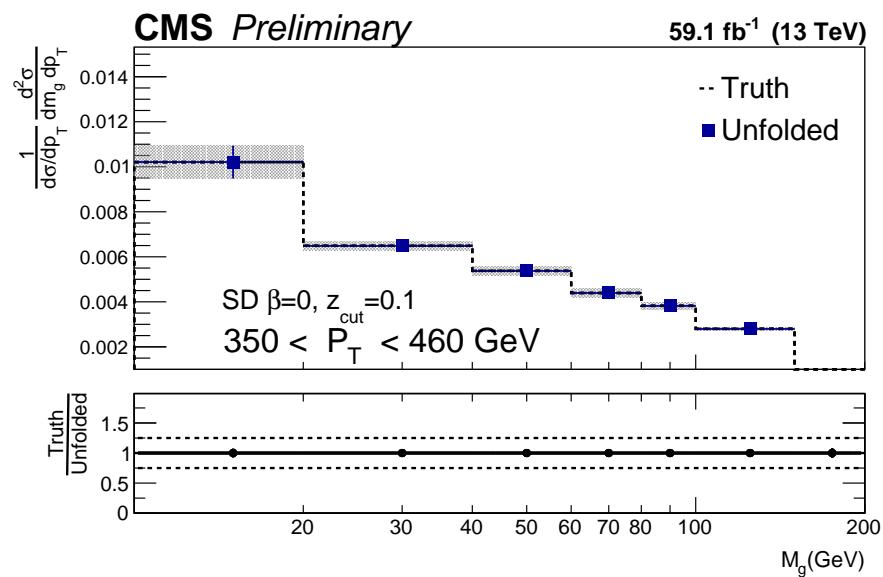


Figure 233: Closure test of groomed reconstructed Monte Carlo, p_T 350–460 GeV.

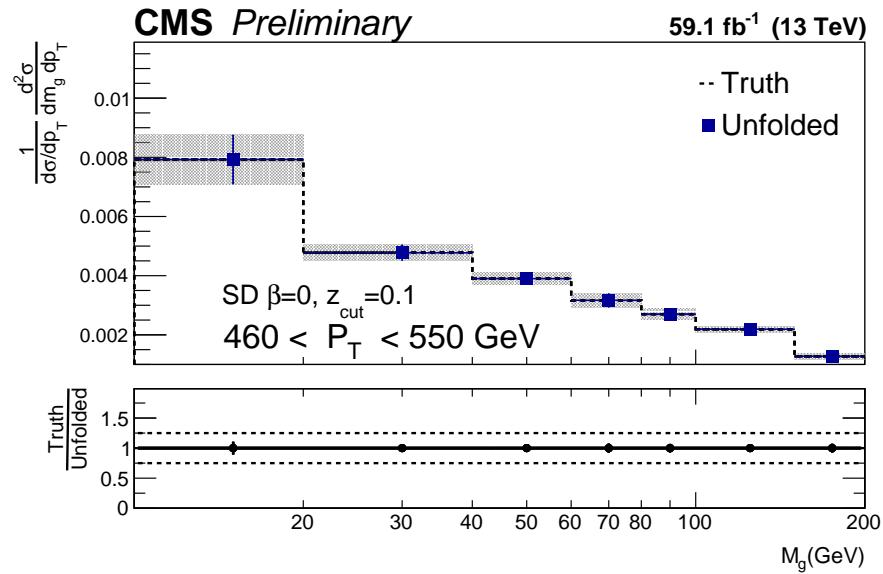


Figure 234: Closure test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

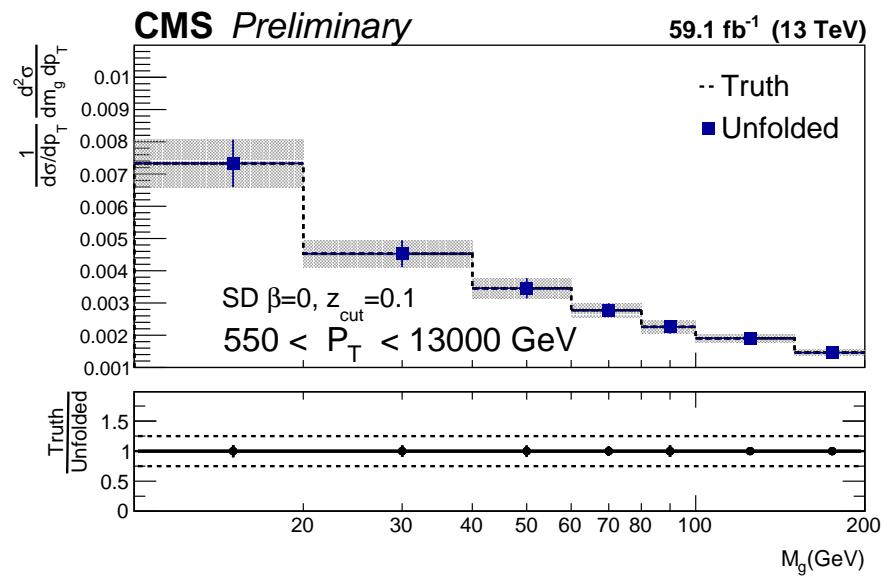


Figure 235: Closure test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

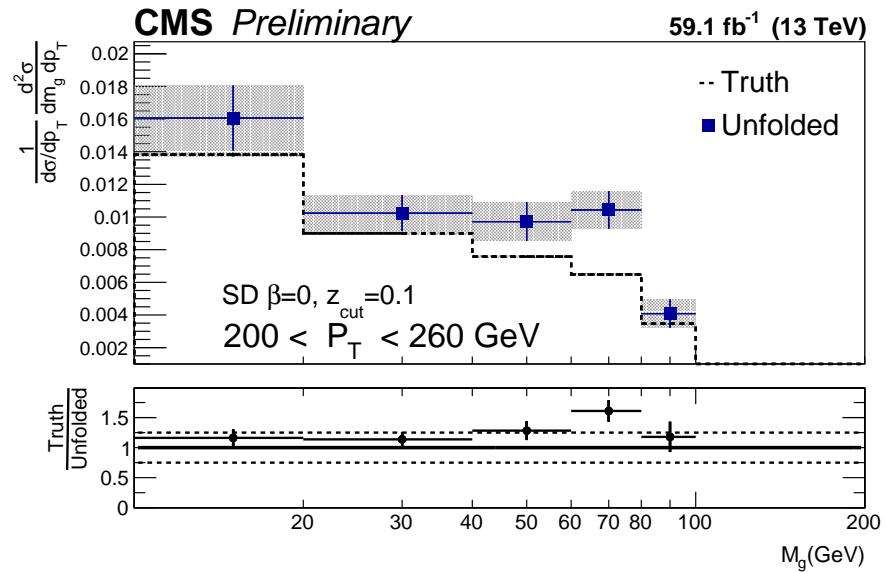


Figure 236: Physical bias test of groomed reconstructed Monte Carlo, p_T 200-260 GeV.

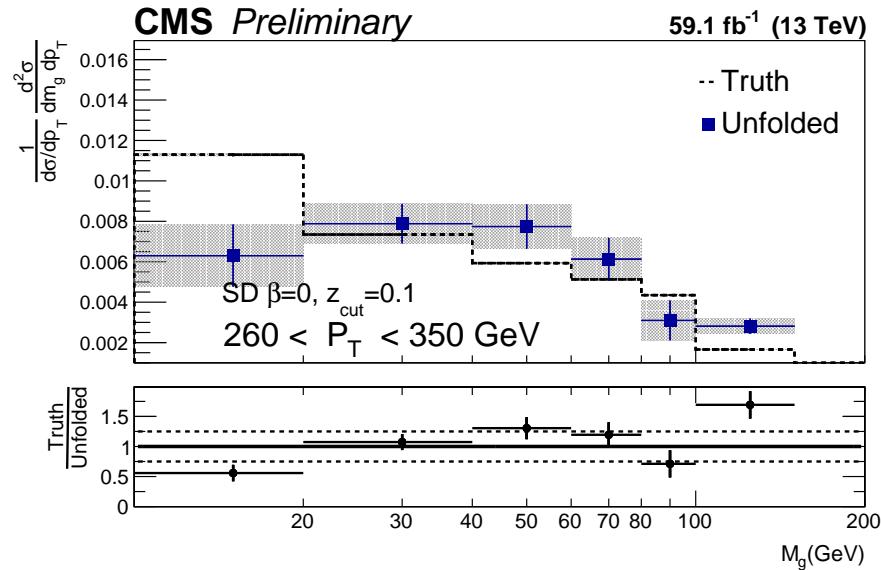


Figure 237: Physical bias test of groomed reconstructed Monte Carlo, p_T 260-350 GeV.

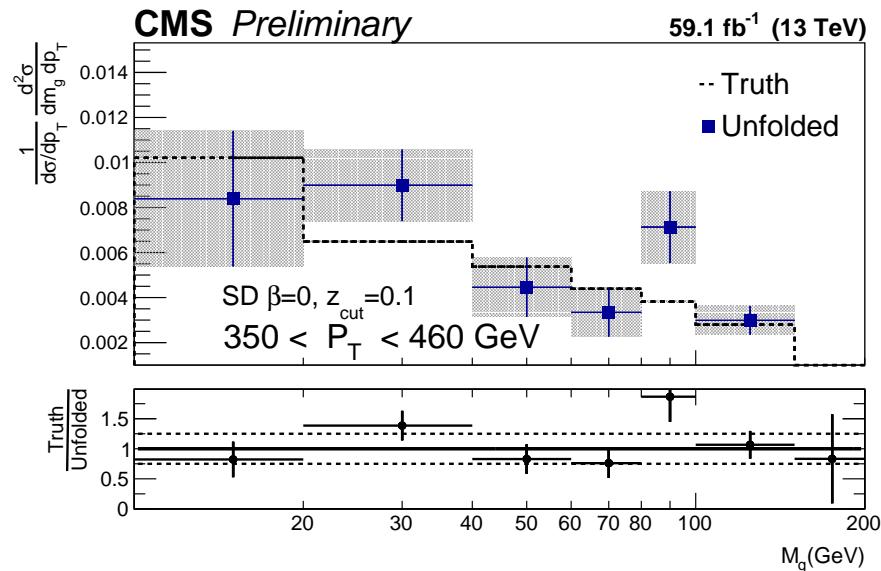


Figure 238: Physical bias test of groomed reconstructed Monte Carlo, p_T 350-460 GeV.

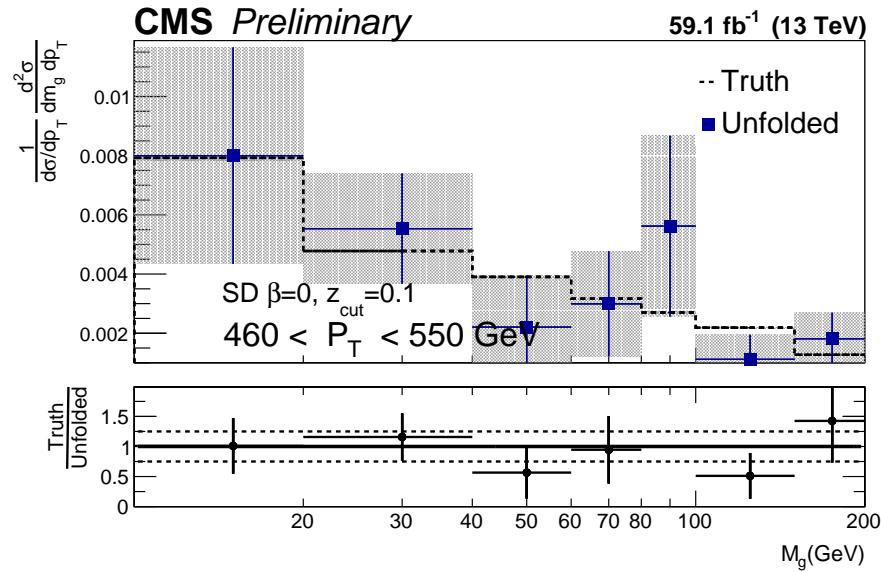


Figure 239: Physical bias test of groomed reconstructed Monte Carlo, p_T 460-550 GeV.

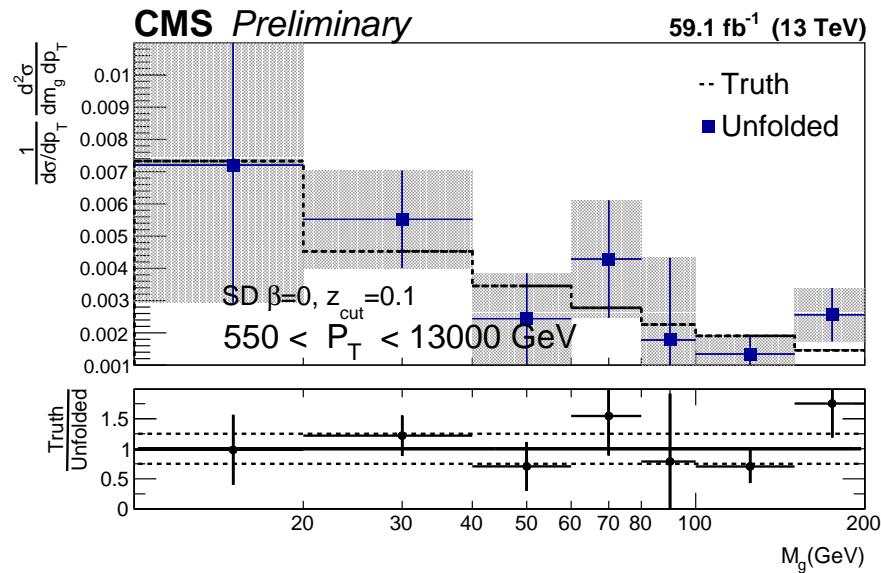


Figure 240: Physical bias test of groomed reconstructed Monte Carlo, p_T 550-13000 GeV.

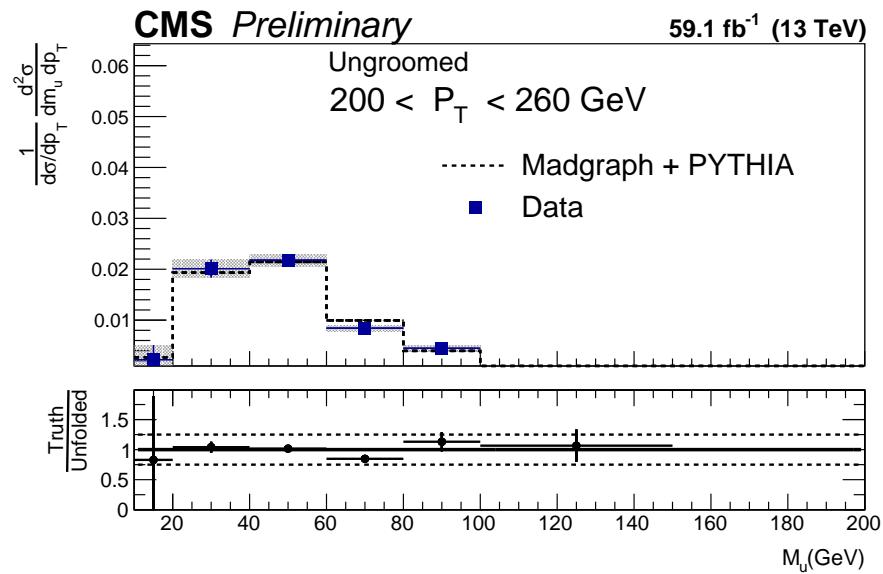


Figure 241: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 200-260 GeV.

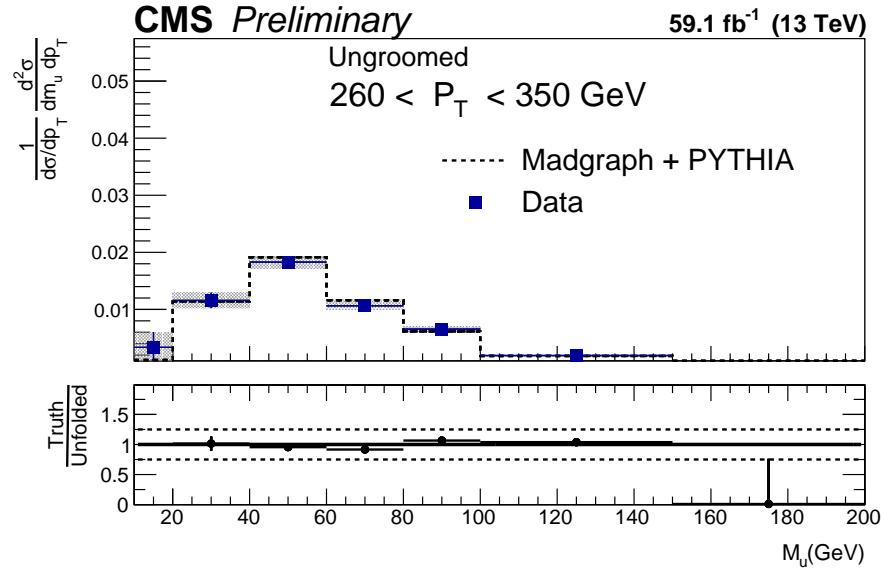


Figure 242: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 260–350 GeV.

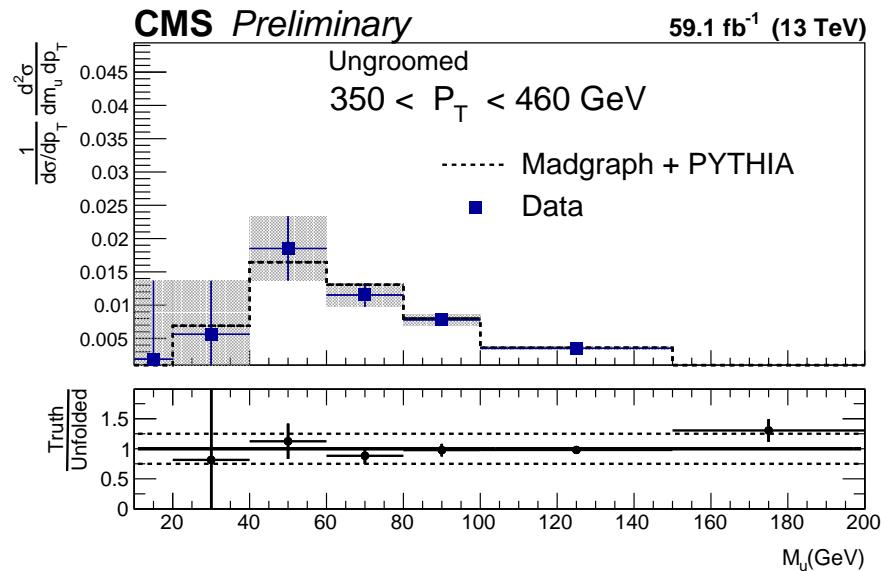


Figure 243: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 350–460 GeV.

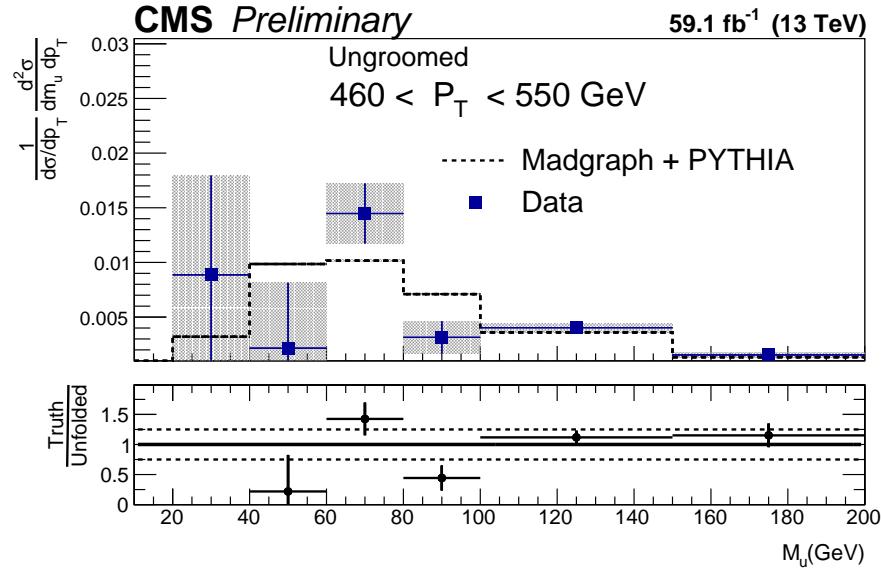


Figure 244: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 460–550 GeV.

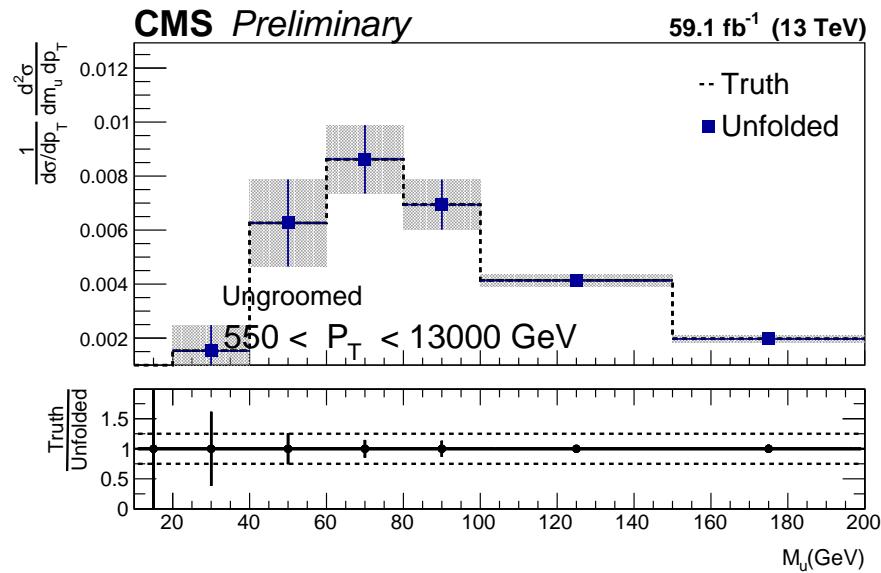


Figure 245: Normalized cross section results with respect to jet mass for ungroomed jets, p_T 550–13000 GeV.

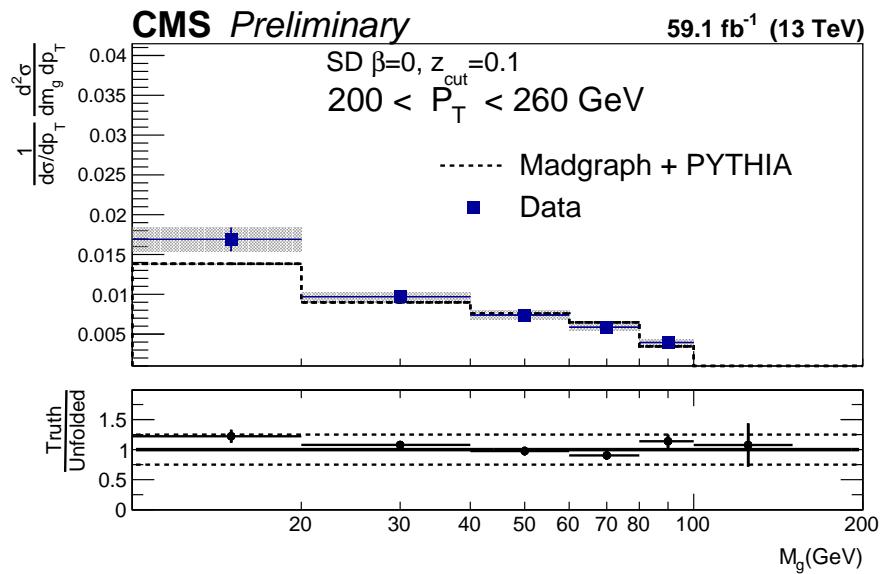


Figure 246: Normalized cross section results with respect to jet mass for groomed jets, p_{T} 200-260 GeV.

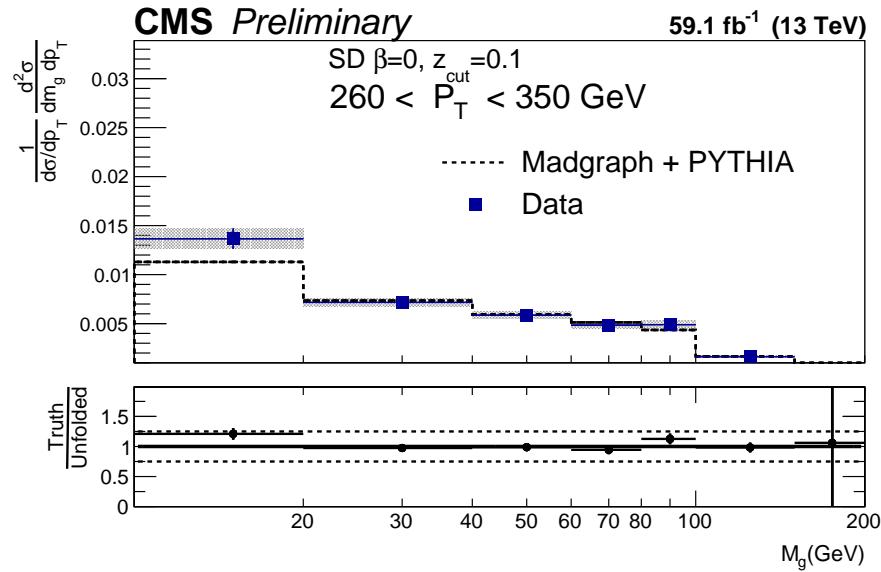


Figure 247: Normalized cross section results with respect to jet mass for groomed jets, p_T 260-350 GeV.

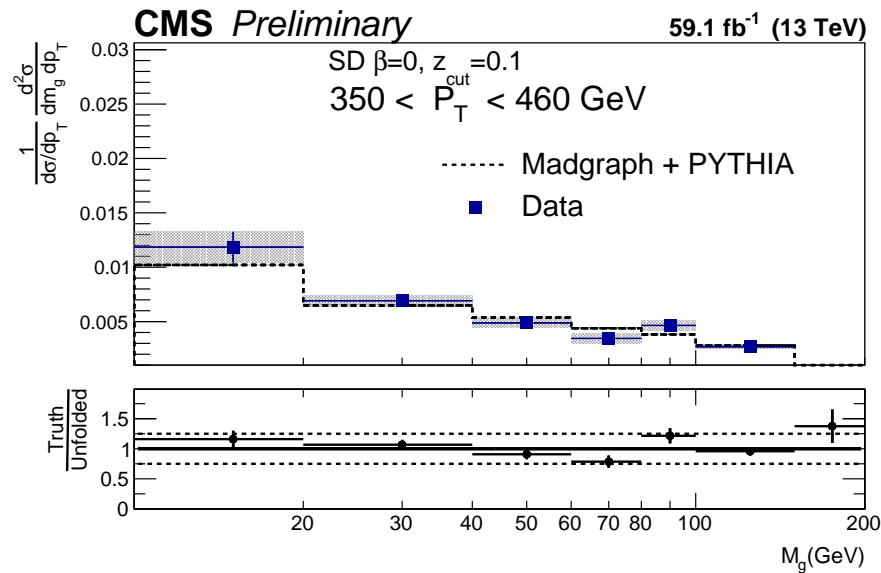


Figure 248: Normalized cross section results with respect to jet mass for groomed jets, p_T 350-460 GeV.

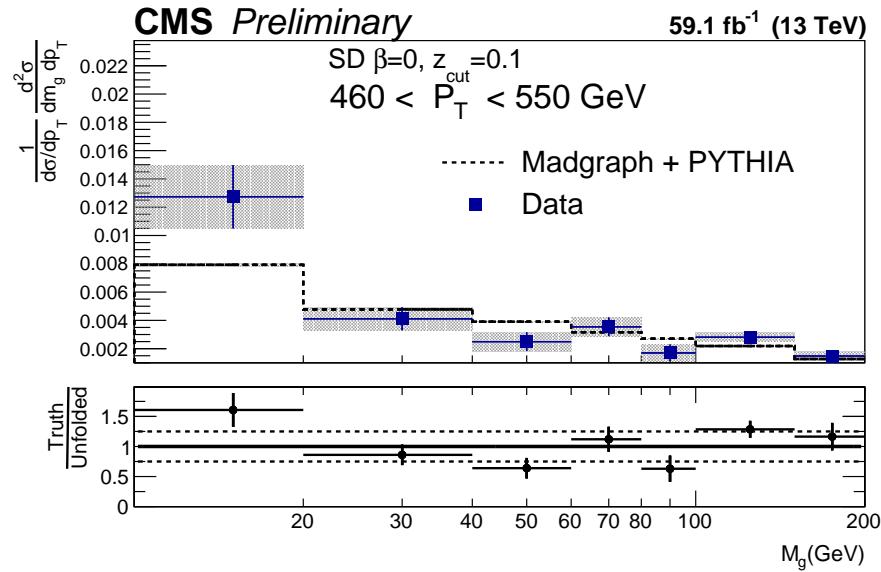


Figure 249: Normalized cross section results with respect to jet mass for groomed jets, p_T 460-550 GeV.

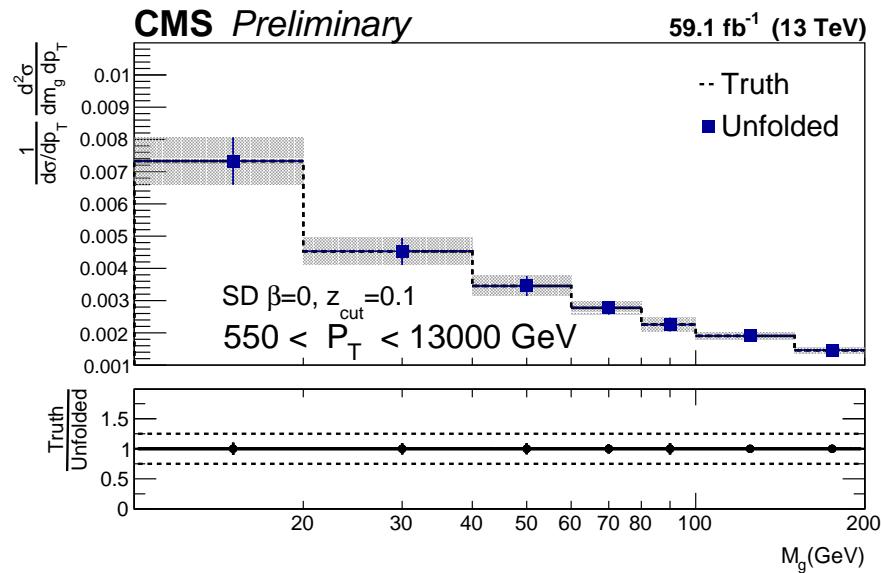


Figure 250: Normalized cross section results with respect to jet mass for groomed jets, p_T 550-13000 GeV.

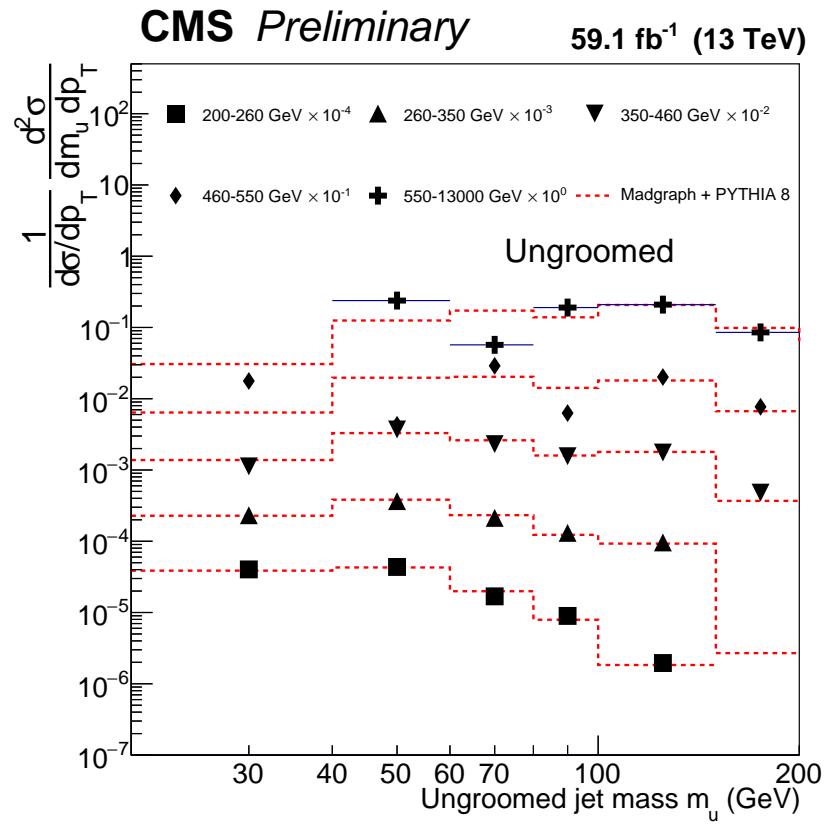


Figure 251: Results for ungroomed reconstructed unfolding with jet mass.

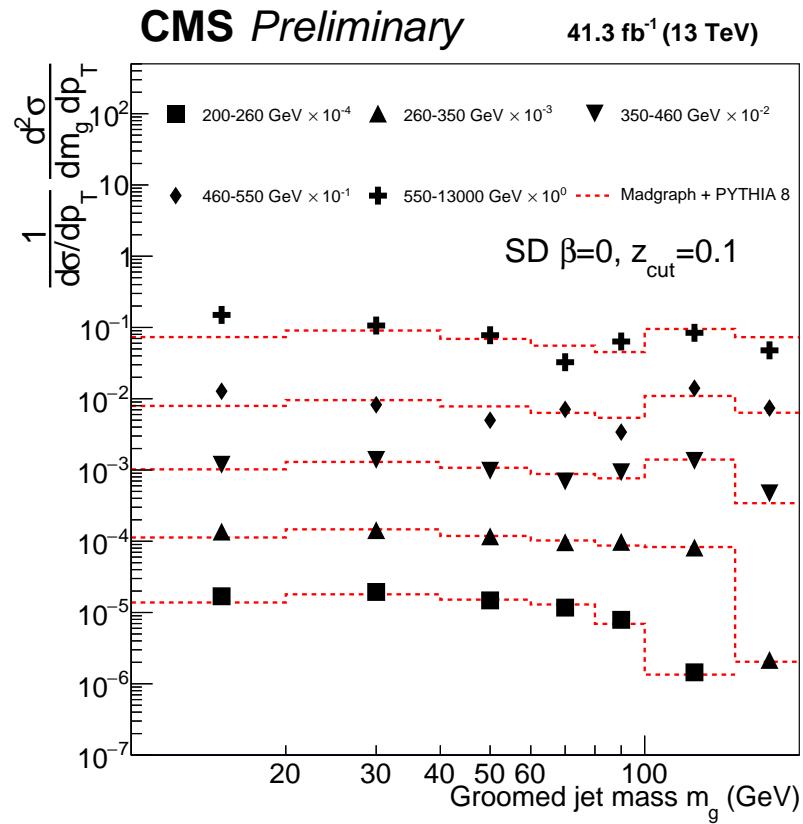


Figure 252: Results for groomed reconstructed unfolding with jet mass.

The resulting correlation matrix without systematic uncertainties for the ungroomed jets is shown in Fig. 253, and for the groomed jets is shown in Fig. 254. The same figures with systematic uncertainties are shown for the ungroomed jets in Fig. 255, and for the groomed jets is shown in Fig. 256.

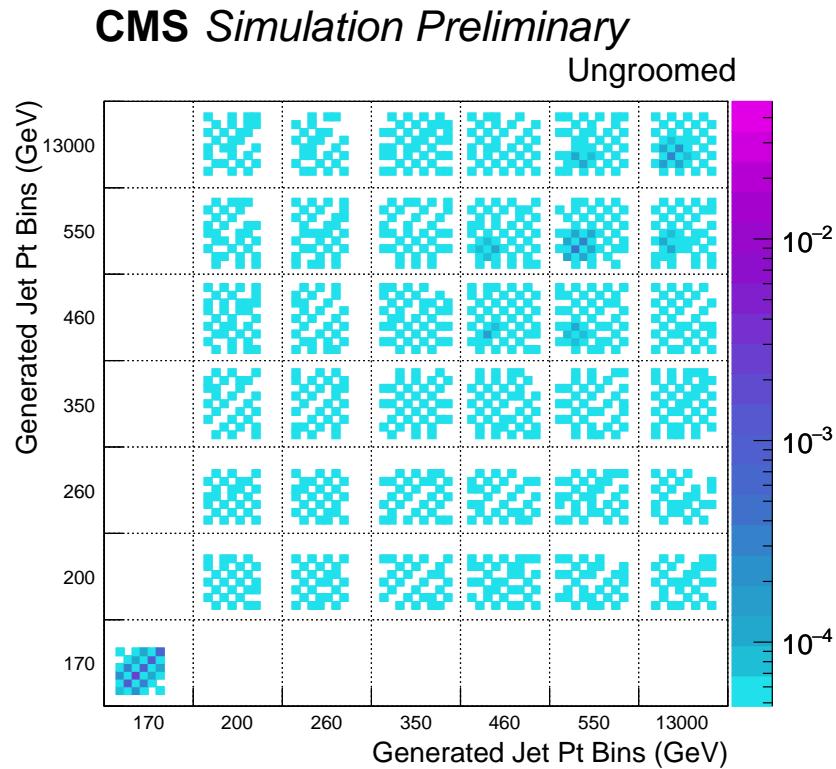


Figure 253: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 10, 20, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 350, \dots, 13000$ GeV).

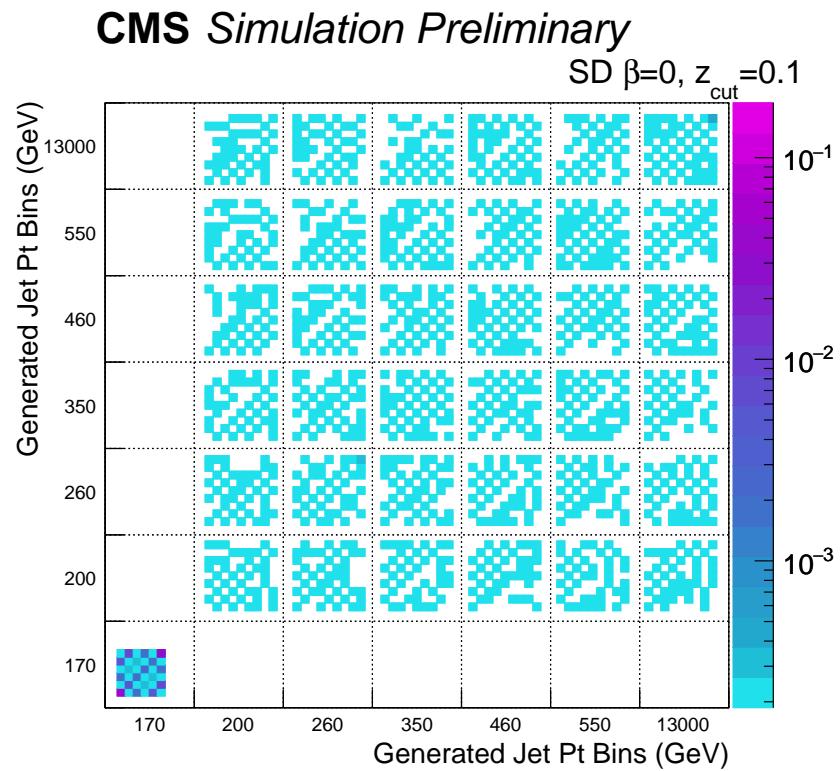


Figure 254: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks ($0, 20, 40, \dots, 200$ GeV), while the p_T is indexed by the major blocks ($200, 340, \dots, 760$ GeV).

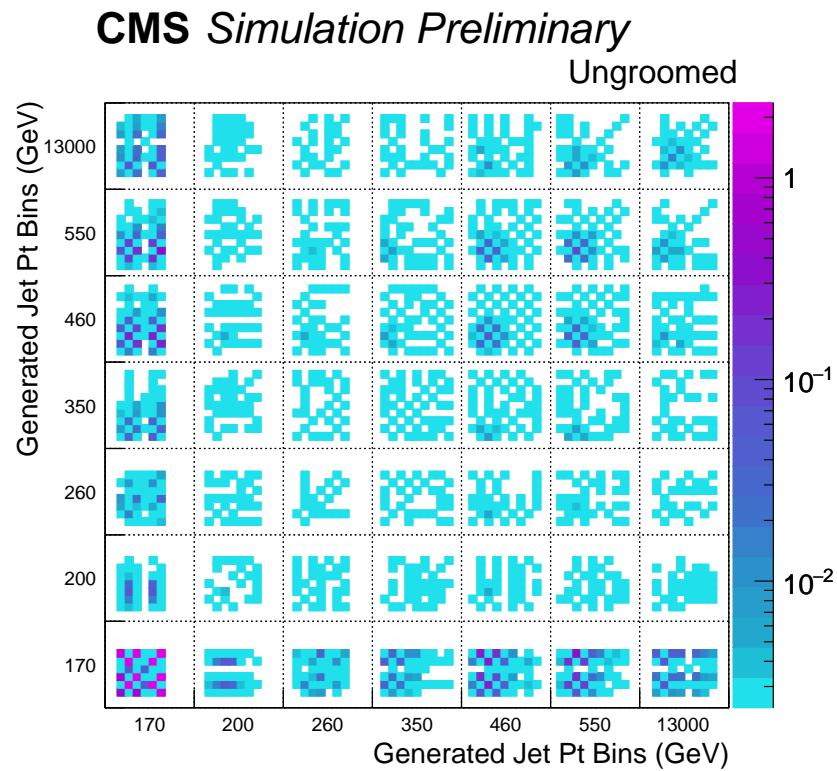


Figure 255: Correlation matrix for ungroomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,... 200 GeV), while the p_T is indexed by the major blocks (200,260,... 13000 GeV).

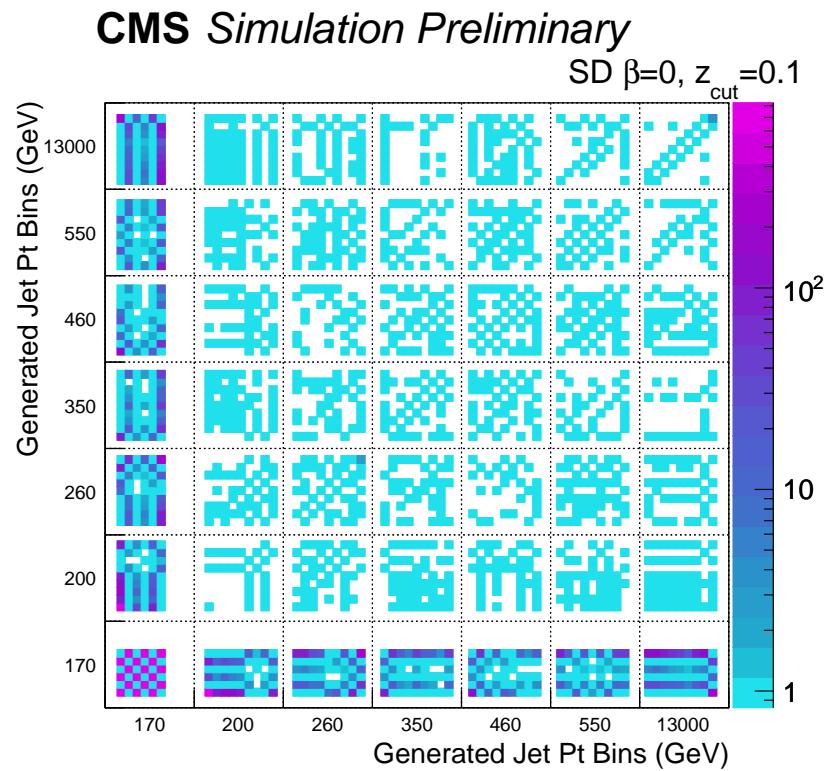


Figure 256: Correlation matrix for groomed jets. The bins are set such that the jet mass is indexed by the minor blocks (0,20,40,...200 GeV), while the p_T is indexed by the major blocks (200,340,...0 GeV).

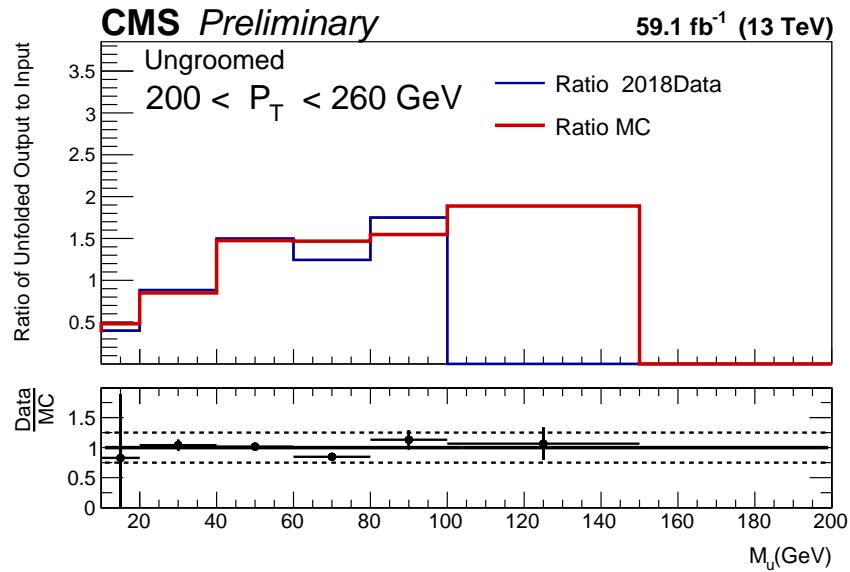


Figure 257: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 200-260 GeV.

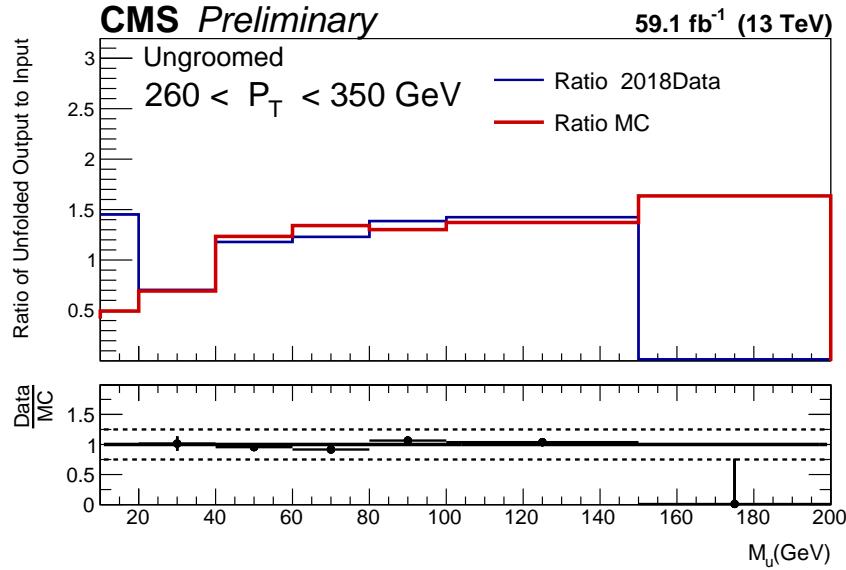


Figure 258: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 260-350 GeV.

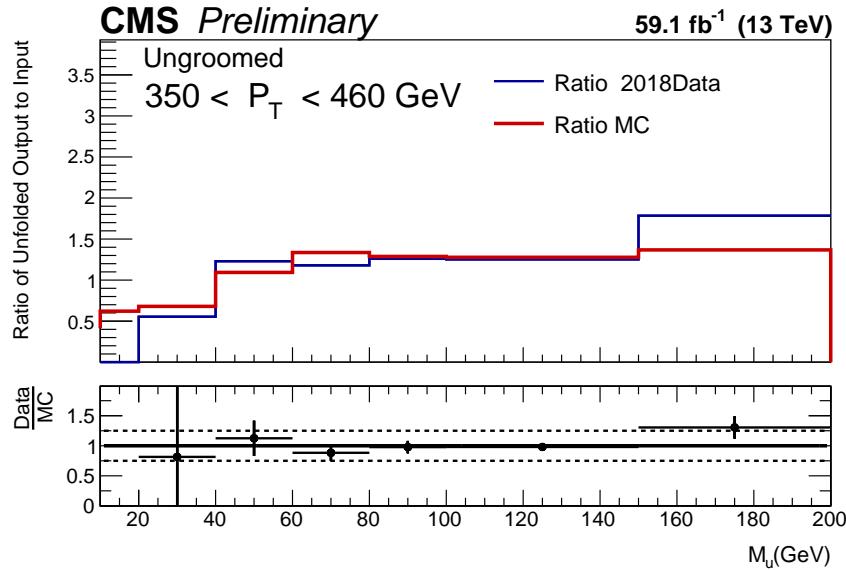


Figure 259: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 350-460 GeV.

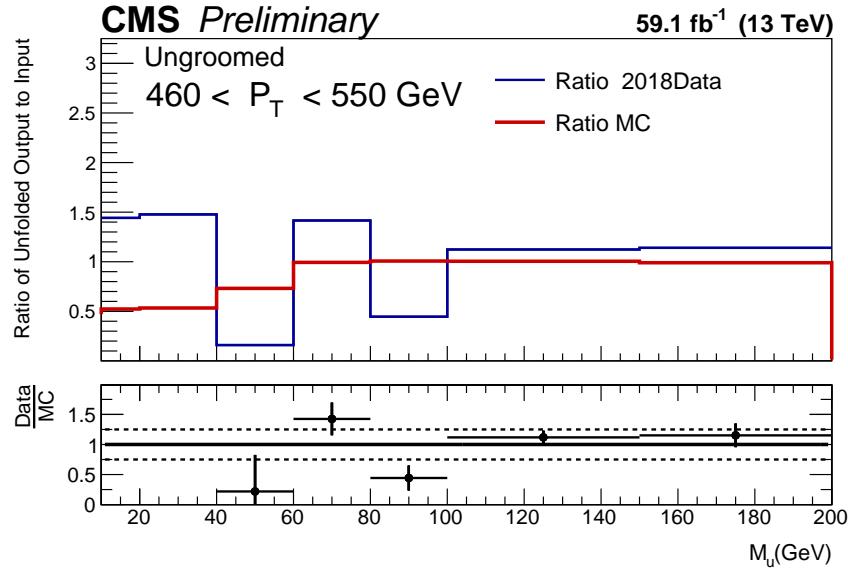


Figure 260: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 460-550 GeV.

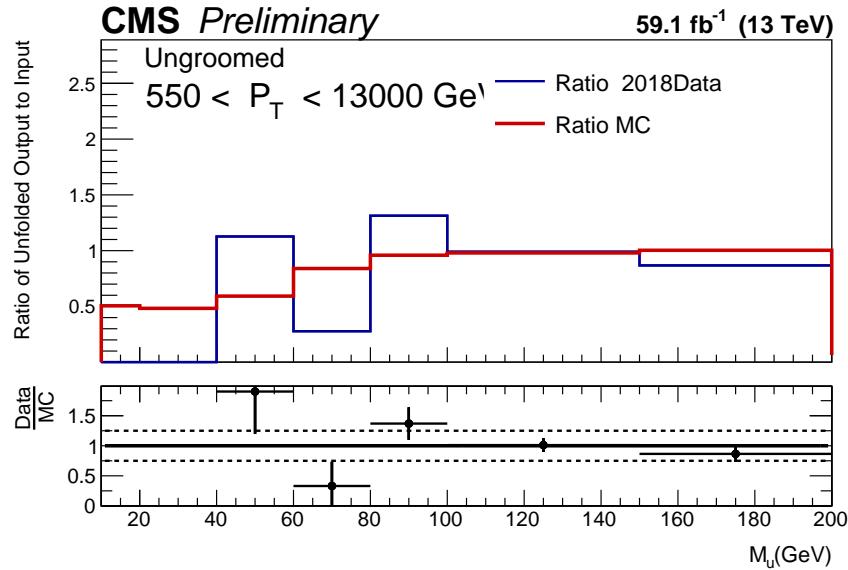


Figure 261: Ratio of unfolded over raw data and MC for ungroomed jets, p_T 550-13000 GeV.

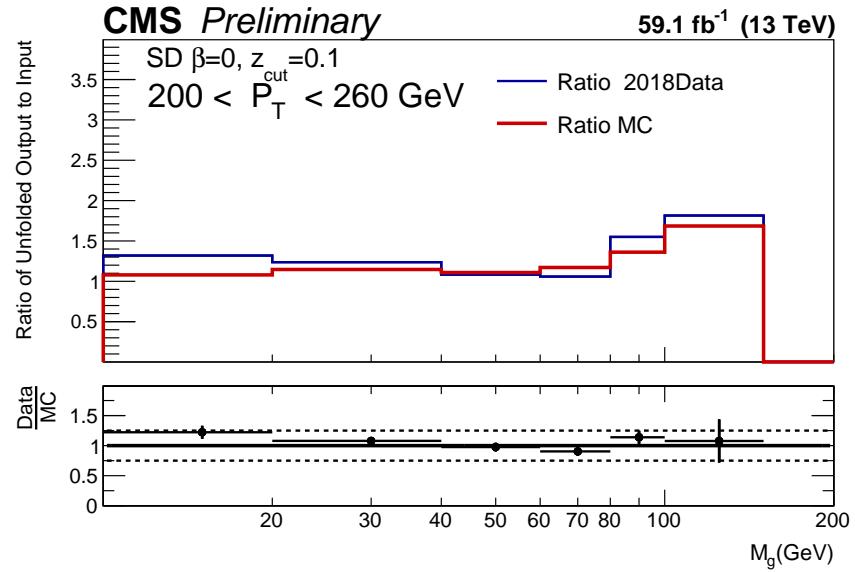


Figure 262: Ratio of unfolded over raw data and MC for groomed jets, p_T 200-260 GeV.

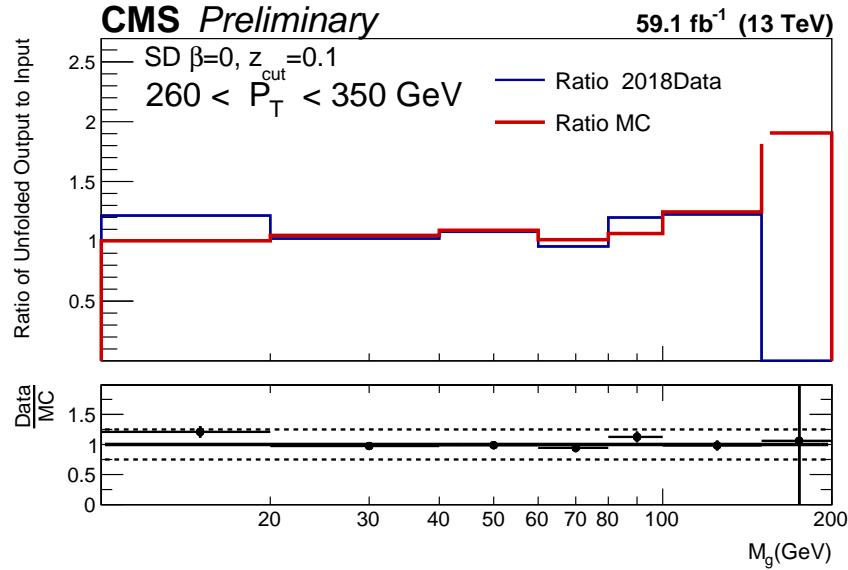


Figure 263: Ratio of unfolded over raw data and MC for groomed jets, p_T 260-350 GeV.

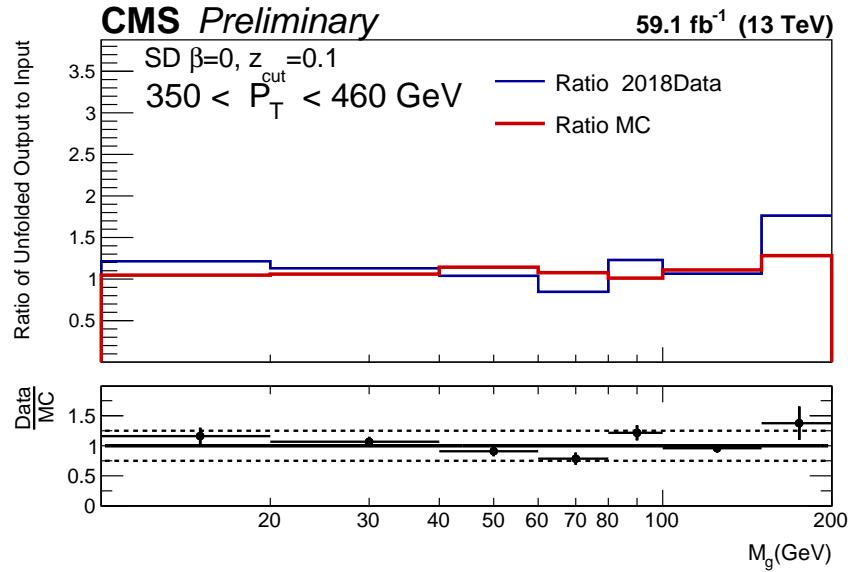


Figure 264: Ratio of unfolded over raw data and MC for groomed jets, p_T 350-460 GeV.

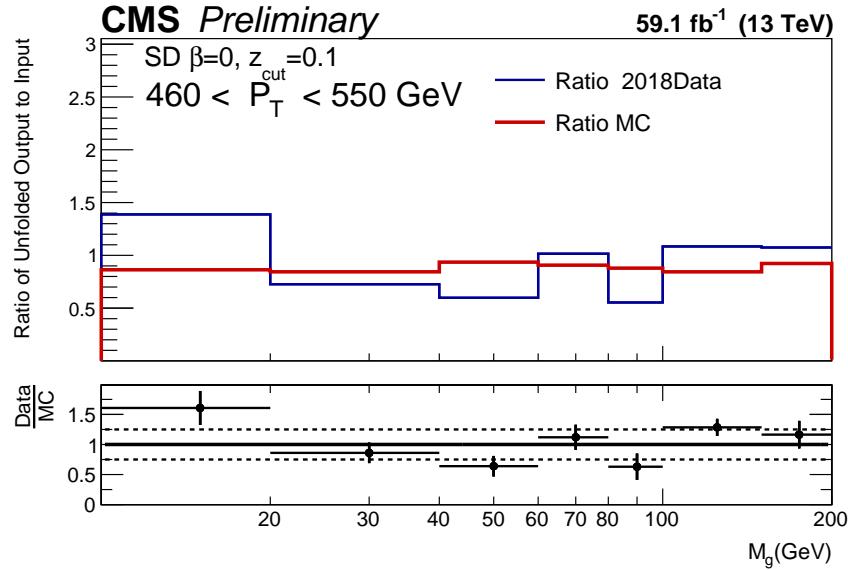


Figure 265: Ratio of unfolded over raw data and MC for groomed jets, p_T 460-550 GeV.

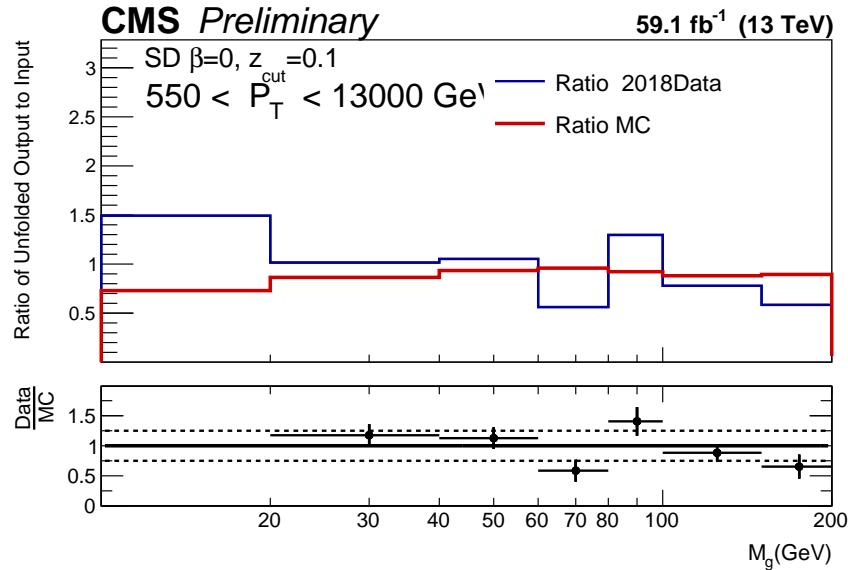


Figure 266: Ratio of unfolded over raw data and MC for groomed jets, p_T 550-13000 GeV.

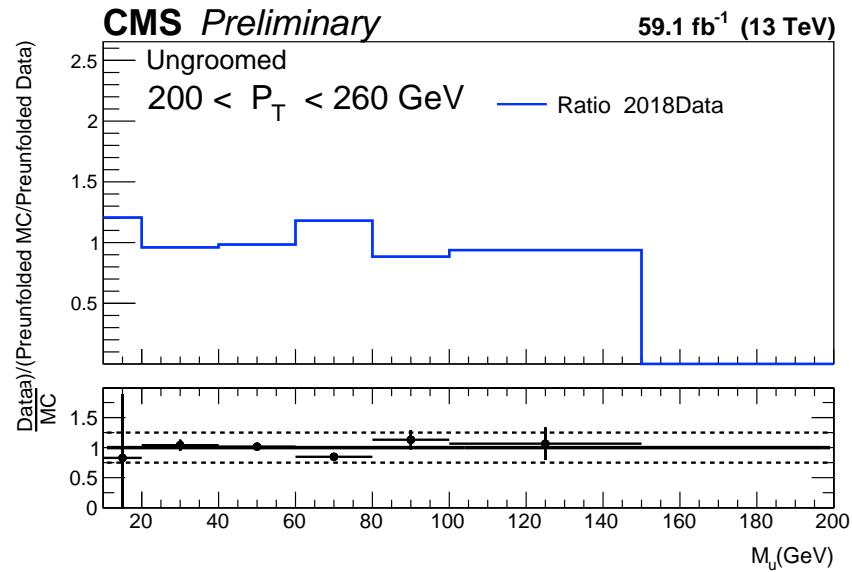


Figure 267: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 200-260 GeV.

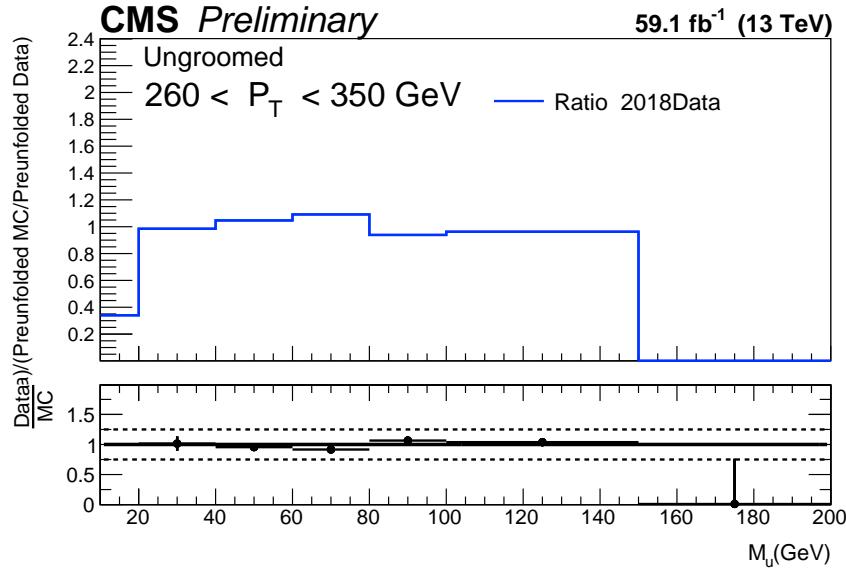


Figure 268: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 260-350 GeV.

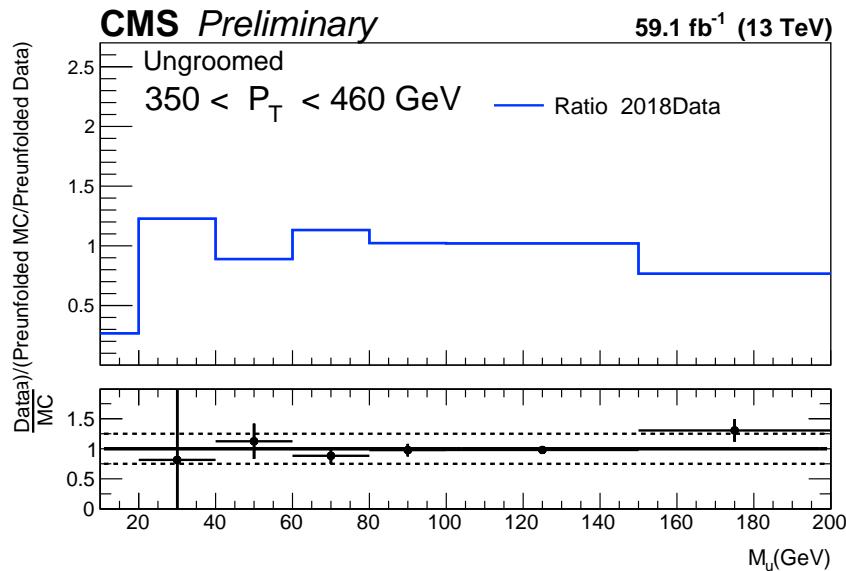


Figure 269: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 350-460 GeV.

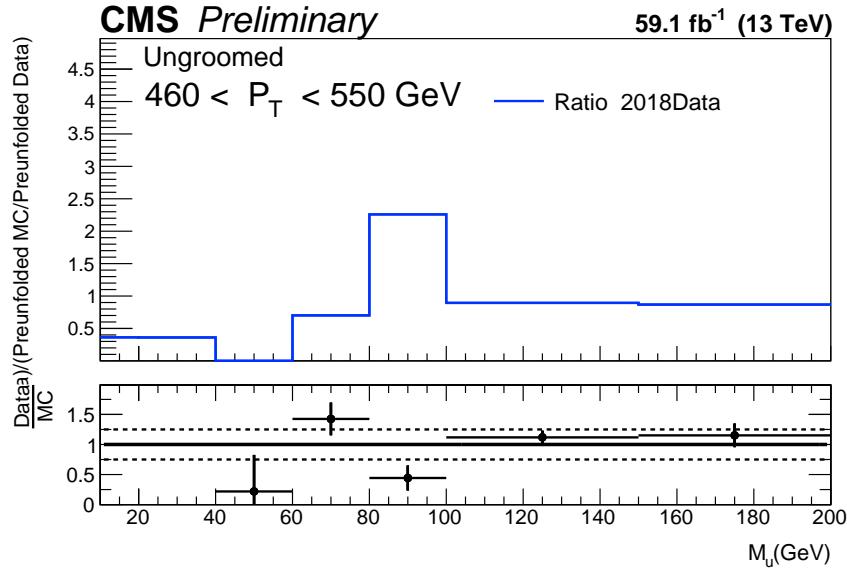


Figure 270: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 460-550 GeV.

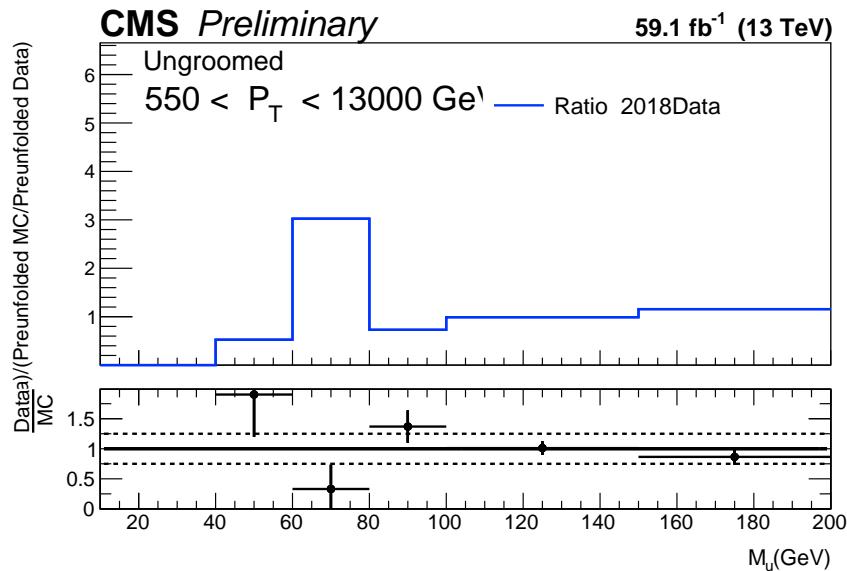


Figure 271: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for ungroomed jets, p_T 550-13000 GeV.

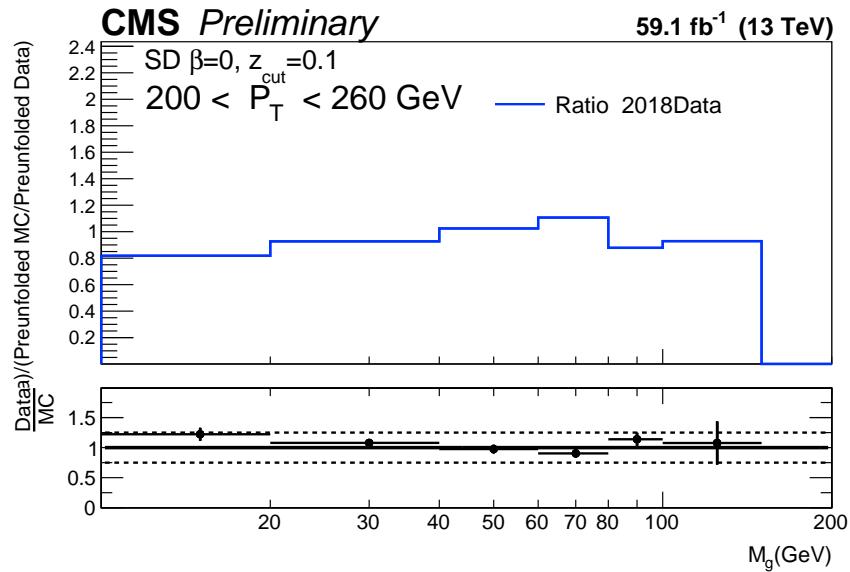


Figure 272: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 200-260 GeV.

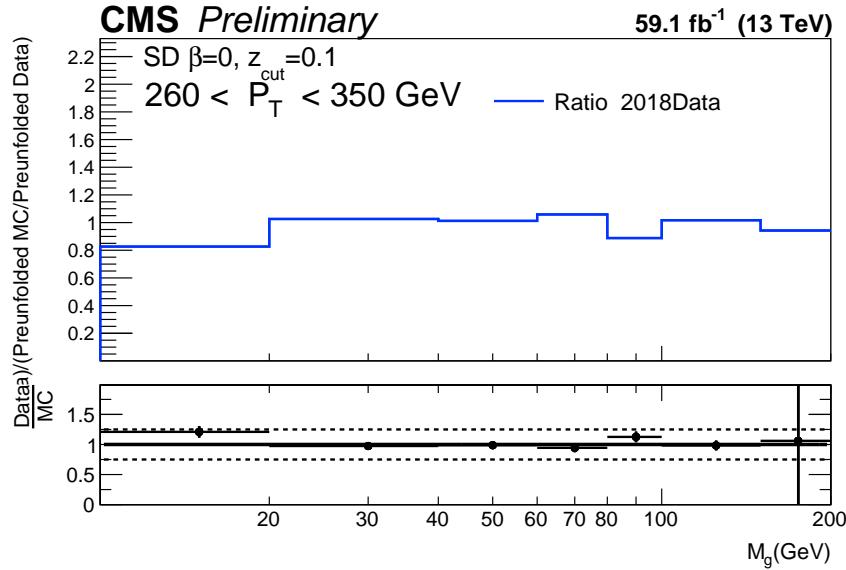


Figure 273: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 260-350 GeV.

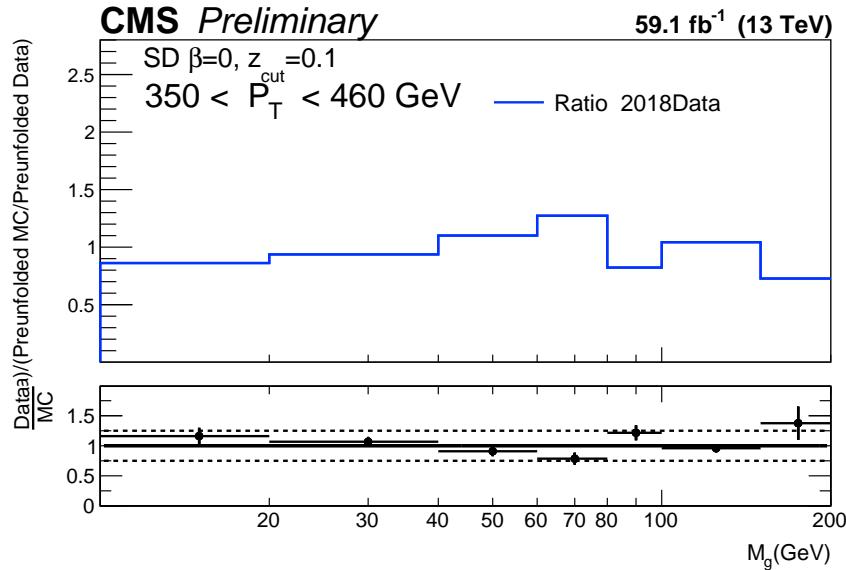


Figure 274: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 350-460 GeV.

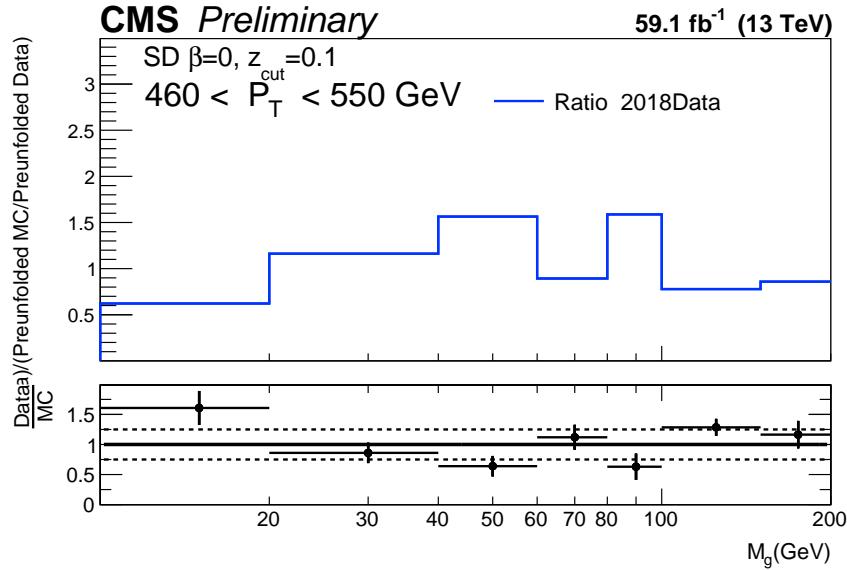


Figure 275: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 460-550 GeV.

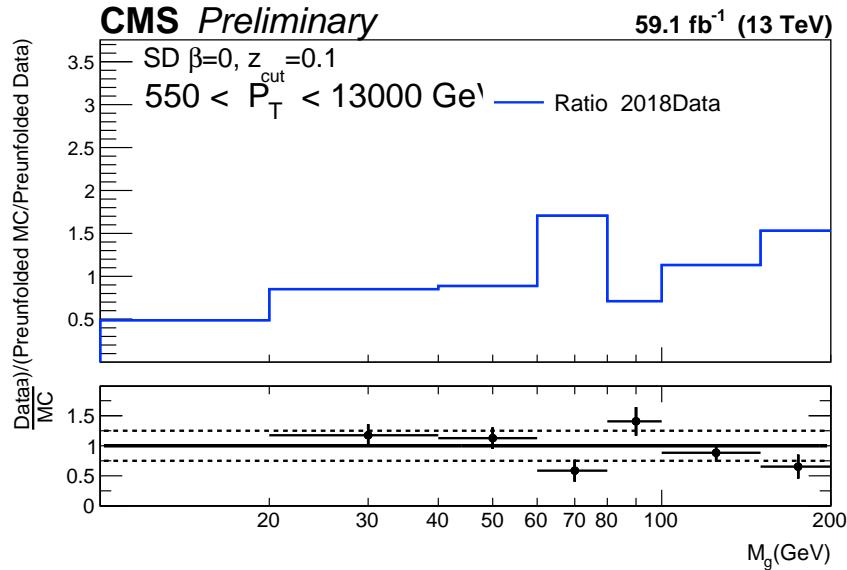


Figure 276: Ratio of generator truth over unfolded data, divided by the ratio of raw MC over raw data for groomed jets, p_T 550-13000 GeV.

Relativistic Kinematics

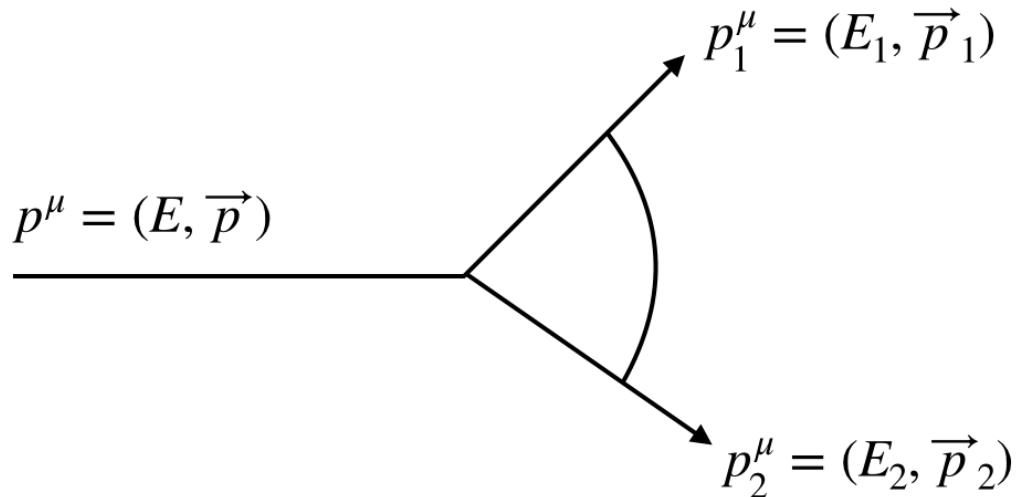


Figure A.1: A simple $1 \rightarrow 2$ decay described by relativistic kinematics.

Consider the LO process depicted in A.1, of a simple $1 \rightarrow 2$ decay. at LO there is a kinematic turn-off at $p_T R / 2$ (where R is the distance parameter for the jet clustering), from the relativistic kinematics of a $1 \rightarrow 2$ decay. However, for real jets the turn-off is closer to $p_T R / \sqrt{2}$ due to stochastic effects.

To see this, consider a particle of energy E and mass m decaying to two massless particles, each with an energy $E/2$ and separated by an angle θ . The mass must satisfy $m^2 < \frac{E^2}{2}(1 - \cos \theta)$. In the small angle limit, this would be $m^2 < E^2\theta^2/4$, or $m < E\theta/2$.

According to relativistic kinematics the interaction can be described by the following equation:

$$p^\mu p_\mu = (p_1 + p_2)^\mu (p_1 + p_2)_\mu$$

$$p^\mu p_\mu = (p_1 + p_2)^\mu (p_1 + p_2)_\mu$$

$$m^2 = (E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2) \cdot (\vec{p}_1 + \vec{p}_2)$$

$$m^2 \simeq 2E_1 E_2 (1 - \cos \theta)$$

If the energies of the splitting particles are equal then the equation simplifies since $E_1 = E_2 = \frac{E}{2}$.

$$m^2 < \frac{E^2}{2}(1 - \cos \theta)$$

$$\frac{2m^2}{E^2} < (1 - \cos \theta)$$

Using the small angle approximation this simplifies further.

$$(1 - (1 - \frac{\theta^2}{2})) \simeq \frac{\theta^2}{2}$$

Solving for mass, one finds :

$$m < E\theta/2$$

With more particle decays, the stochastic nature of the shower increases this to $m < E\theta/\sqrt{2}$. Thus, a leading-order ($1 \rightarrow 2$) decay will have a faster kinematic turn-off than an all-orders ($1 \rightarrow \text{many}$) decay.

Solving for theta:

$$\theta < \frac{2}{\gamma} \text{ where } \gamma \text{ is the lorentz factor } \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}.$$

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