

The concept of optimality depends on what we wish to accomplish by decomposing \bar{g} into two components.

→ We can think of \bar{e} being the error in approximating \bar{g} by $c\bar{x}$. Hence, we can now define mathematically the component (or proj.) of a vector \bar{g} along vector \bar{x} to be $c\bar{x}$, where c is chosen to minimize the length of error vector $\bar{e} = \bar{g} - c\bar{x}$ (mag.)

→ You can see/observe that the perpendicular has the smallest mag. or norm

& it is the **unique one**.

Let us see that

$$c \|\bar{x}\| = \|\bar{g}\| \cos \theta \quad \text{--- (3)}$$

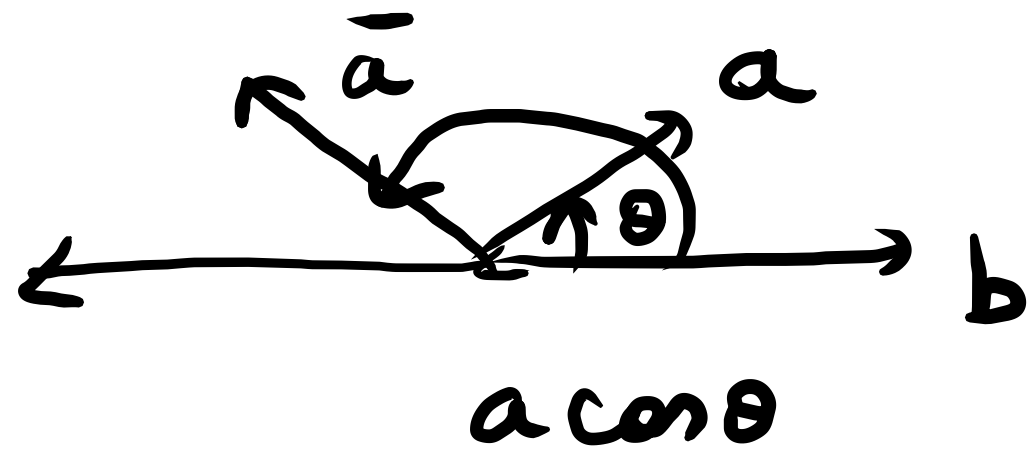
multiply by $\|\bar{x}\|$ both sides,
we get

$$c \|\bar{x}\|^2 = \|\bar{g}\| \|\bar{x}\| \cos \theta \\ = \langle \bar{g}, \bar{x} \rangle$$

$$\Rightarrow c = \frac{\langle \bar{g}, \bar{x} \rangle}{\langle \bar{x}, \bar{x} \rangle}$$

It is apparent that when \bar{g} & \bar{x} are perpendicular or orthogonal then \bar{g} has a zero

\therefore mag. of the component of \bar{g} along \bar{x} is $\|\bar{g}\| \cos \theta$
 \rightarrow this is for eq. (3) only.



Component along \bar{x} , consequently, $c=0$. Hence, \bar{g} & \bar{x} are said to be orthogonal if the inner (scalar or dot) prod. of two vectors is zero i.e.

$$\text{if } \langle \bar{g}, \bar{x} \rangle = 0$$

Gram-Schmidt procedure:-

Given M energy signals denoted as

$s_1(t), s_2(t), \dots, s_M(t)$, how do you create a complete orthonormal set of basis functions?

1. Start with $s_1(t)$, chosen from the set **arbitrarily**, the first basis function is defined by.

$$\Phi_1(t) = \frac{s_1(t)}{\sqrt{E_1}} \quad , \text{ where } E_1 \rightarrow \text{energy of}$$

Then, clearly $s_1(t) = \sqrt{E_1} \Phi_1(t)$

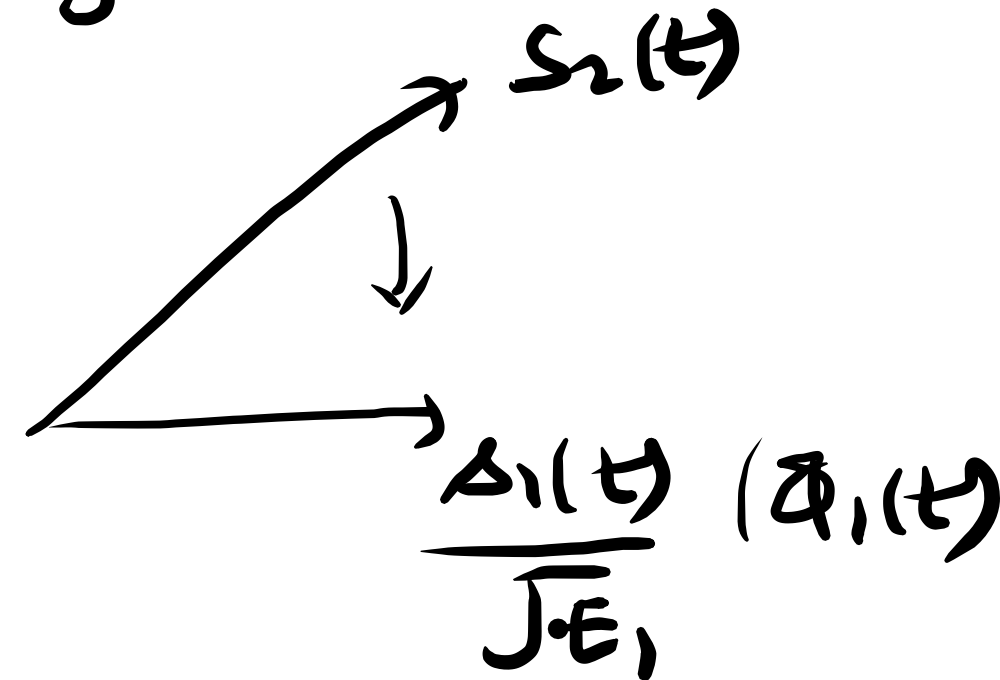
$$\text{or } s_{11} \Phi_1(t)$$

$$2. \text{ Define } s_2 = \frac{\int_0^T s_2(t) \Phi_1(t) dt}{\int_0^T \Phi_1^2(t) dt} \rightarrow 1$$

$$\text{let } g_2(t) = s_2(t) - s_{21} \Phi_1(t)$$

↳ this is orthogonal to $\Phi_1(t)$ over the interval $0 \leq t \leq T$.

$$\int_0^T s_1^2(t) dt$$



second basis function $\Phi_2(t) = \frac{s_2(t) - s_{21}\Phi_1(t)}{\sqrt{E_2 - s_{21}^2}}$

$$E_{num} = \int_0^T [s_2(t) - s_{21}\Phi_1(t)]^2 dt$$

$$E_2 = \int_0^T s_2^2(t) dt$$

$$= E_2 + s_{21}^2 - 2 \int_0^T s_2(t) s_{21} \Phi_1(t) dt$$

$$= E_2 + s_{21}^2 - 2 s_{21} \times s_{21} \quad (1, -1, 2)$$

$$= E_2 - s_{21}^2$$

3. Continuing in this fashion, we have

$$g_i(t) = s_i(t) - \sum_{j=1}^{i-1} s_{ij} \Phi_j(t) \quad \text{--- (4)}$$

where, $s_{ij} = \int_0^T s_i(t) \Phi_j(t) dt, j=1, 2, \dots, i-1$

We have discussed special cases of eq. (4) with $i=2$ & $i=1$ earlier.

Given the $g_i(t)$, we may define the set of basis functions

$$\phi_i(t) = \frac{g_i(t)}{\sqrt{\int_0^T g_i^2(t) dt}}, \quad i=1, 2, \dots, N$$

which form an orthonormal set.

$N \leq M$:-

1. The signals $s_1(t), s_2(t), \dots, s_M(t)$ form a linearly independent set, in which case $N=M$.

2. Signals $s_1(t)$, $s_2(t)$, \dots , $s_M(t)$ are not L.I., in which case $N < M$ & function $g_i(t)$ is zero for $i > N$

H.w. Why not $N > M$ possible?

Two examples: - Conventional Fourier series of a periodic signal

&

B.L. Signal's expansion in terms of its samples taken at the Nyquist rate.

Two distinctions: -

1. Form of basis functions $\phi_1(t)$, $\phi_2(t)$, \dots , $\phi_N(t)$

has not been specified.

i.e., $\left\{ \begin{array}{l} \text{not sinusoidal functions} \\ \text{(as in FS) or series " (as in} \\ \text{reconstruction formula)} \end{array} \right.$

(G5)
2. Here using finite no. of terms is not an approximation, wherein only the first N terms are significant rather an exact exp. where N & only N terms are significant.

H.W. Do EX 7.1.1 from comm. system engg.
by J.G. Proakis, 2nd edition