

TW364: Applied Fourier Analysis

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Lecture 5

Periodic driving force

The Fourier series can be useful in finding a **solution** of a differential equation (DE) describing a system with **periodic driving force**.

Undamped **oscillator** with external (driving) force $f(t)$, mass m , and spring constant k is described by the DE:

$$m \frac{d^2 x}{dt^2} + kx = f(t)$$



Suppose $m = \frac{1}{16}$, $k = 4$, and $f(t)$ is the periodic extension of

$$f(t) = \pi t, \quad -1 < t < 1$$

Note: in a practical sense we only consider $f(t)$ for $t \geq 0$, but when applying the Fourier series we'll also regard $t < 0$.

$f(t)$ is the **odd** extension of πt , $0 < t < 1$, so let's get the **sine** series:



$$f(t) = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \sin(n\pi t)$$

The DE becomes

$$\frac{1}{16} \frac{d^2 x}{dt^2} + 4x = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \sin(n\pi t)$$

This DE is linear and non-homogeneous. By the method of **undetermined coefficients**, we assume a particular solution similar in form to the RHS and its derivatives, i.e.

$$x_p = \sum_{n=1}^{\infty} [A_n \cos(n\pi t) + B_n \sin(n\pi t)]$$



Plug x_p into the LHS of the DE:

$$\sum_{n=1}^{\infty} \left[\left(\frac{-n^2\pi^2}{16} + 4 \right) A_n \cos(n\pi t) + \left(\frac{-n^2\pi^2}{16} + 4 \right) B_n \sin(n\pi t) \right] = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \sin(n\pi t)$$



Solve for A_n and B_n :

$$A_n = 0, \quad B_n = \frac{32(-1)^{n+1}}{n(64 - n^2\pi^2)}$$

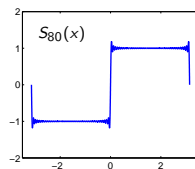
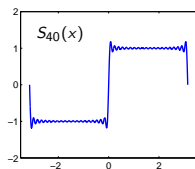
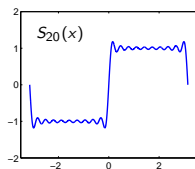
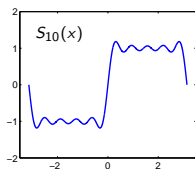
$$\therefore x_p = \sum_{n=1}^{\infty} \frac{32(-1)^{n+1}}{n(64 - n^2\pi^2)} \sin(n\pi t)$$

The Gibbs phenomenon

In Lecture 4 we found the Fourier series of $f(x) = \begin{cases} -1, & -\pi \leq x < 0 \\ 1, & 0 \leq x \leq \pi \end{cases}$

$$f(x) = \sum_{n=1}^{\infty} \frac{2-2(-1)^n}{n\pi} \sin(nx)$$

Partial sum approximations:



The **oscillations** exhibited by these approximations around **function discontinuities** (also in the periodic extension of f) are referred to as the **Gibbs phenomenon**.

The magnitude of the “overshoot” in $S_N(x)$ does **not** disappear as $N \rightarrow \infty$, but approaches a finite limit.

It turns out that this limit is about **8.95%** of the magnitude of the particular jump discontinuity.

MATLAB DEMO

Observing this percentage for the previous example...

If the Gibbs overshoot does not vanish in the limit, does that not **contradict** the **convergence** theorem of the Fourier series?

No! The theorem describes **point-wise convergence**.



Since the region where oscillations occur get narrower as N increases, the oscillations will disappear in the neighbourhood of any x where $f(x)$ is continuous, even for x arbitrarily close to the discontinuity.