

Modeling of Measurement-based Quantum Computing on IBM Q Experience Devices

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One-way quantum computing

One-way quantum computing is one of many frameworks of measurement-based quantum computation (MBQC).

In such framework, a quantum computation is not realized through an explicit unitary evolution of qubits mediated by unitary quantum gates as in the quantum circuit model.

Rather, a computation is realized through sequential single-qubit measurements on an initial resource state [2].

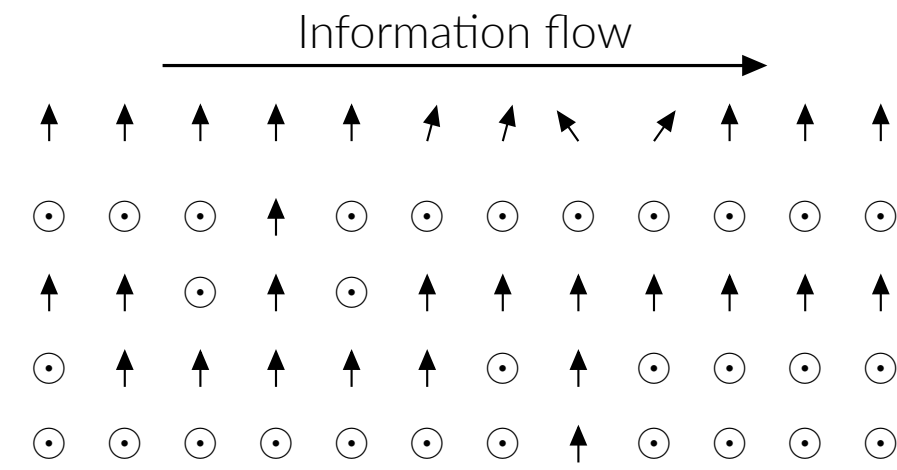


Figure 1. Propagation of information through an entangled state $|\phi\rangle_G$ via measurements. Tilted arrows represented $x - y$ plane measurements, vertical arrows represent measurements of $\hat{\sigma}_x$ and the circles denote measurement of $\hat{\sigma}_z$.

Graph states for quantum computation

The initial resource states are provided in the form of highly entangled multi-qubit states called **graph states**

$$|G\rangle = \prod_{\{k,l\} \in E} CZ^{kl} |+\rangle^{\otimes |V|} \quad (1)$$

The two-qubit gate CZ^{kl} (controlled $\hat{\sigma}_z$) is applied between pairs of vertices according to the edge set E of the underlying graph $G = (V, E)$, describing the topology of the state.

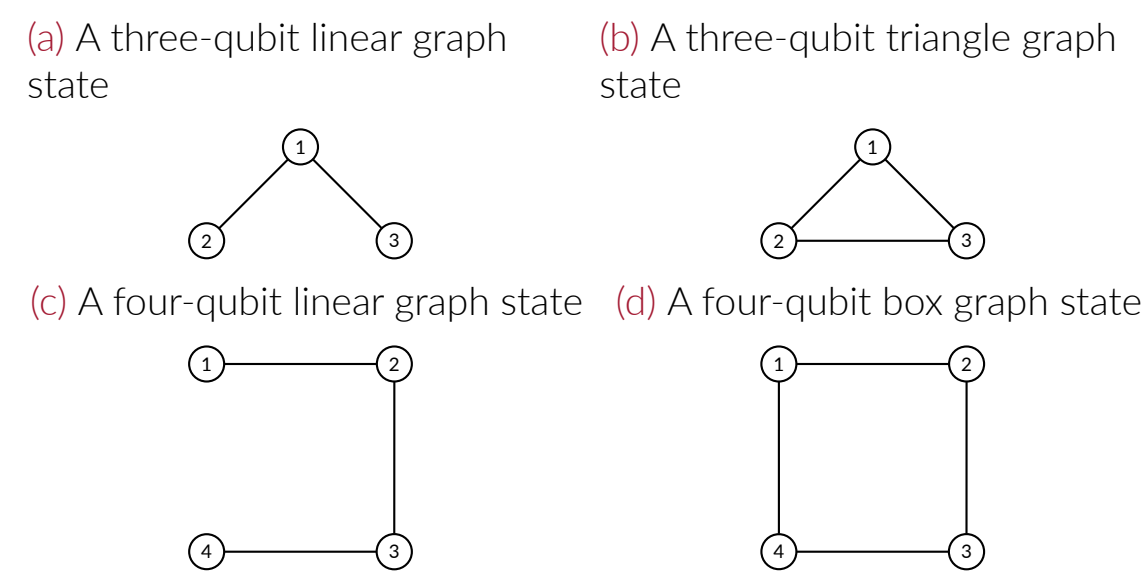


Figure 2. A few representatives graph states.

Local unitaries

Two graph state vectors $|G\rangle$ and $|G'\rangle$ and their underlying graphs G and G' are said to be equivalent to one another under a local unitary \hat{U}

$$|G'\rangle = \hat{U} |G\rangle \quad (2)$$

The graph states $|G\rangle$ and $|G'\rangle$ are said to belong to some equivalence class and within such a class, all graph states are equivalent modulo local unitaries (LU).

Edge local complementation and equivalence classes

From an underlying graph $G = (V, E)$ of a graph state $|G\rangle$, a new graph $G' = (V, E')$ and an associated graph state $|G'\rangle$ can be realized via the complement of a subgraph induced by the neighborhood of some vertex $k \in V$ of G . Then the two graph states $|G\rangle$ and $|G'\rangle$ are said to belong to the same **LU-equivalence** class [1].

Equivalence classes of graph states

Unitary for edge complementation

The mapping $\tau^k : G \mapsto G'$ under edge local complementation is mediated by a unitary operation \hat{U}^k [1].

$$|\tau^k(G)\rangle = \hat{U}^k(G) |G\rangle \quad (3)$$

where \hat{U}^k is of the form

$$\hat{U}^k(G) = (-i\hat{\sigma}_x^k)^{1/2} \prod_{l \in \eta_k} (i\hat{\sigma}_z^l)^{1/2} \quad (4)$$

Four-qubit linear graph state $\stackrel{w}{\equiv}$ Four-qubit box graph state

Starting from the four-qubit linear graph state as in Fig. 2c, a sequential application of edge local complementation on the vertices 3, 2 and 3 produces a four-qubit box graph state similar to the one in Fig. 2d

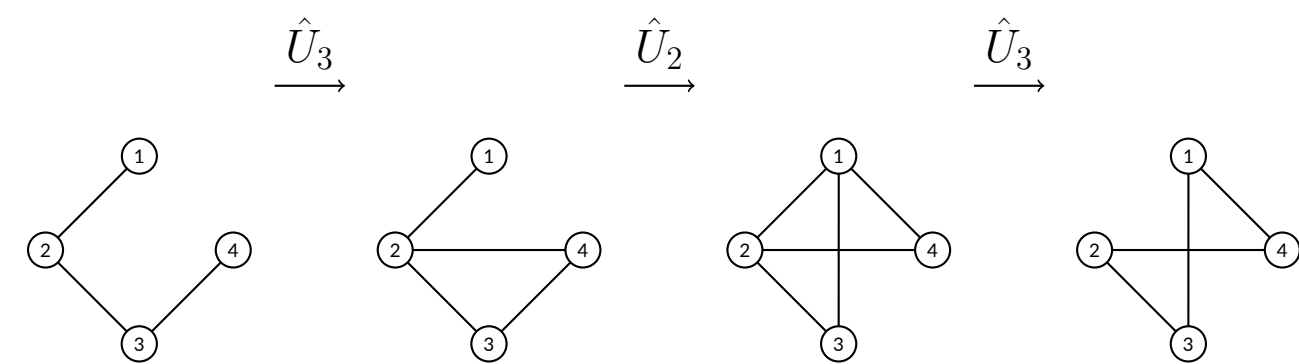


Figure 3. Sequentially applying the edge local complementation rule on a four-qubit linear graph state showing that such a state is LU-equivalent to the four-qubit box graph state under unitary transformation $\hat{U}^3 \hat{U}^2 \hat{U}^3 = \hat{\sigma}_z^1 \otimes \hat{H}^2 \otimes \hat{H}^3 \otimes \hat{\sigma}_z^4$

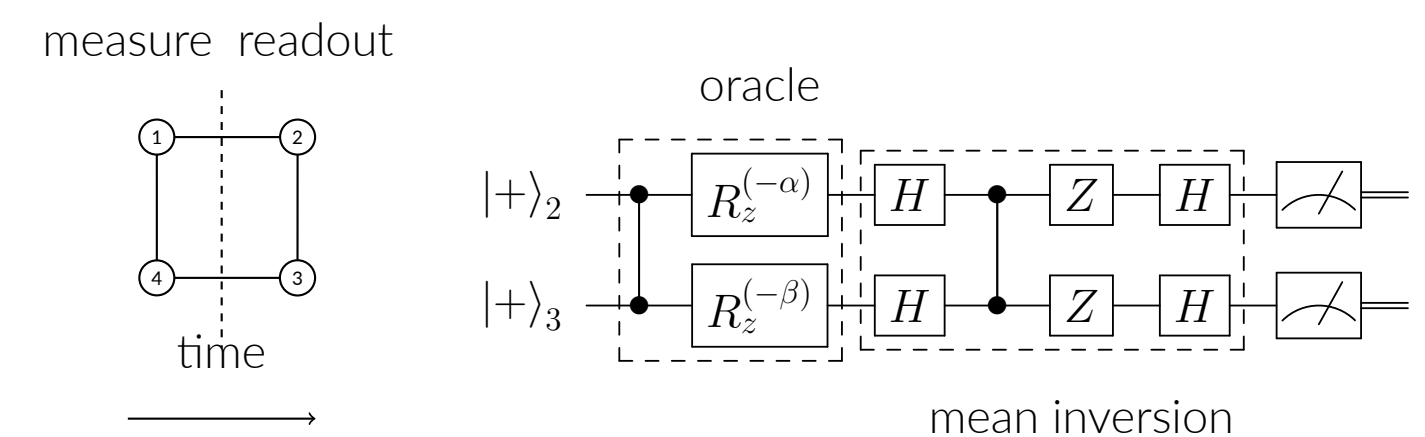
Under the vertex-mapping bijection that swaps vertices 1 and 4, the graphs in the last step of Fig. 3 and Fig. 2d are isomorphic to one another.

Case Study 00 : Grover's quantum search algorithm

Two-qubit quantum search on four-qubit box graph state

A two-qubit measurement-based version of Grover's quantum search algorithm can be realized on the four-qubit box graph state [4].

Through measurements of qubits 1 and 4 of Fig. 2d in the basis $B_j(\alpha) = \{|+\alpha\rangle_j, |-\alpha\rangle_j\}$ where $|\pm\alpha\rangle_j = \frac{1}{\sqrt{2}}(|0\rangle_j \pm e^{i\alpha}|1\rangle_j)$ ($\alpha \in \mathbb{R}$). Reading out qubits 2 and 3 in the basis $B_j(\pi) = \frac{1}{\sqrt{2}}(|0\rangle_j \mp |1\rangle_j)$.



Each of the four cases, where the oracle tags the elements 00, 01, 10 and 11 are realized by the measurement settings $\alpha\beta = \pi\pi, \pi 0, 0\pi$ and 00 respectively.

Caveats

The outcome of a measurement in the basis $B_i(\alpha)$ on qubit i is α_i , which is $O(1)$ for a measurement of $|+\alpha\rangle$ ($|-\alpha\rangle$).

If the measurement on qubits 1 and 4 are anything other than $\alpha_1 = 0$ and $\alpha_4 = 0$, Pauli errors are introduced and there is a need to reinterpret the outcomes on qubits 2 and 3. In such a scenario, a feedforward of the outcomes must be applied $\{\alpha_2, \alpha_3\} \mapsto \{\alpha_2 \oplus \alpha_4, \alpha_3 \oplus \alpha_1\}$.

Case study 01 : Improved quantum search

The LU-equivalence and isomorphism of the four-qubit linear and box graph state offers a way to improve the performance of the former implementation.

The resultant LU-equivalent four-qubit linear graph state has one less CZ gate than the four-qubit box state, which should have an improvement upon the performance of the former implementation since generally NISQ two-qubit gates are much more error-prone than single-qubit gates.

The measurement procedure is the same as in the former implementation, however the feedforward relation has modified to $\{\alpha_2, \alpha_3\} \mapsto \{\alpha_2 \oplus \alpha_1, \alpha_3 \oplus \alpha_4\}$ since qubits 1 and 4 have to be swapped.

Qiskit experiments

To model the algorithms described in the two case studies, the quantum circuits in Fig. 4 and Fig. 5 were executed on the IBM Q 5 vigo machine, with trials of 8100 shots.

Since the IBM Q 5 vigo machine doesn't have the necessary topology to realize a four-qubit box graph state, the mapping of logical qubits to physical qubits was chosen as to minimize the number of swap gates in the resulting transpiled circuit.

To further limit the number of gates, the feedforward operations which would be CNOT gates on the relevant qubits were done post-experiment. Similarly, the swap of qubits 2 and 3 in the algorithm in case study 01 was also done post-experiment.

The four-qubit linear graph state passed a multi-partite stabilizer-based entanglement witness test [3], giving an expectation value of $\langle \mathcal{W} \rangle = -0.45 \pm 0.261$ while the four-qubit box graph state narrowly failed the entanglement witness test, giving $\langle \mathcal{W} \rangle = 0.003 \pm 0.260$.

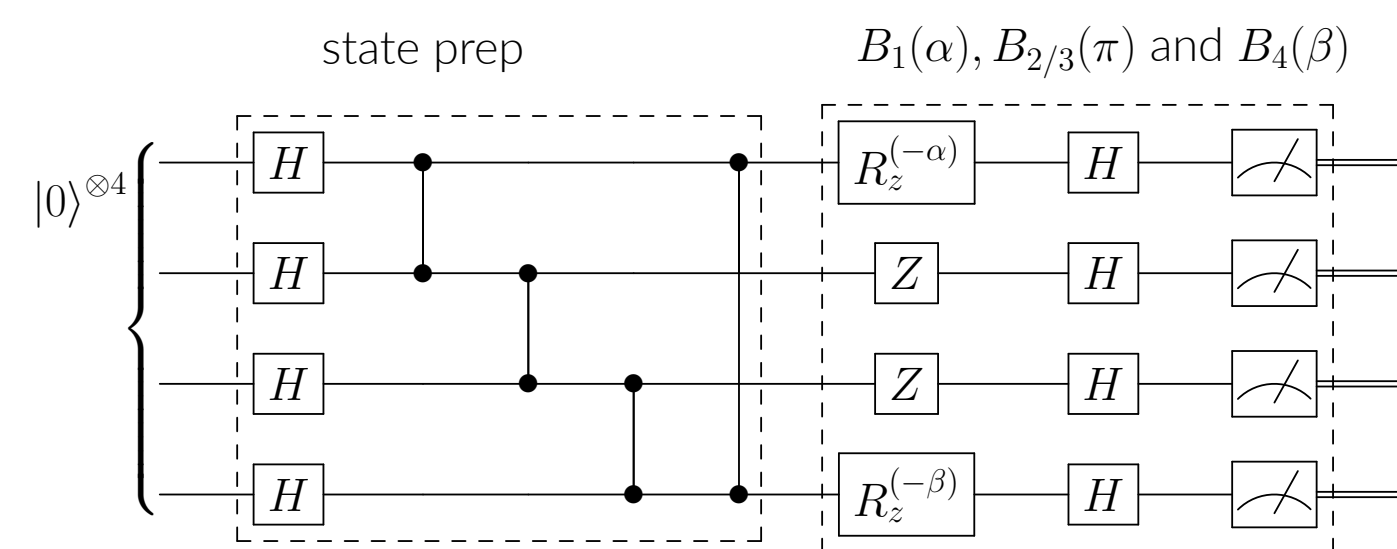


Figure 4. Measurement-based two-qubit Grover's algorithm on a four-qubit box graph state

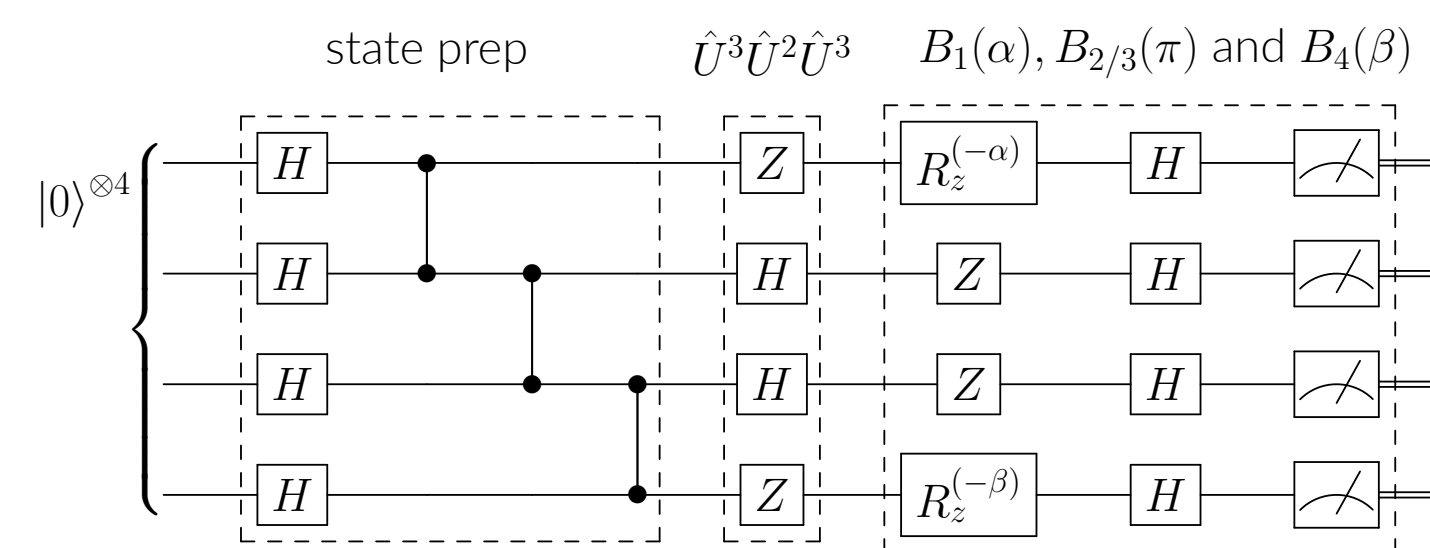


Figure 5. Measurement-based two-qubit Grover's algorithm on a four-qubit linear graph state

Results

Case study 00

The results below were obtained for the algorithm described in case study 00, showing the feedforward and no feedforward results. In all cases, the success probabilities are approximately 75%.

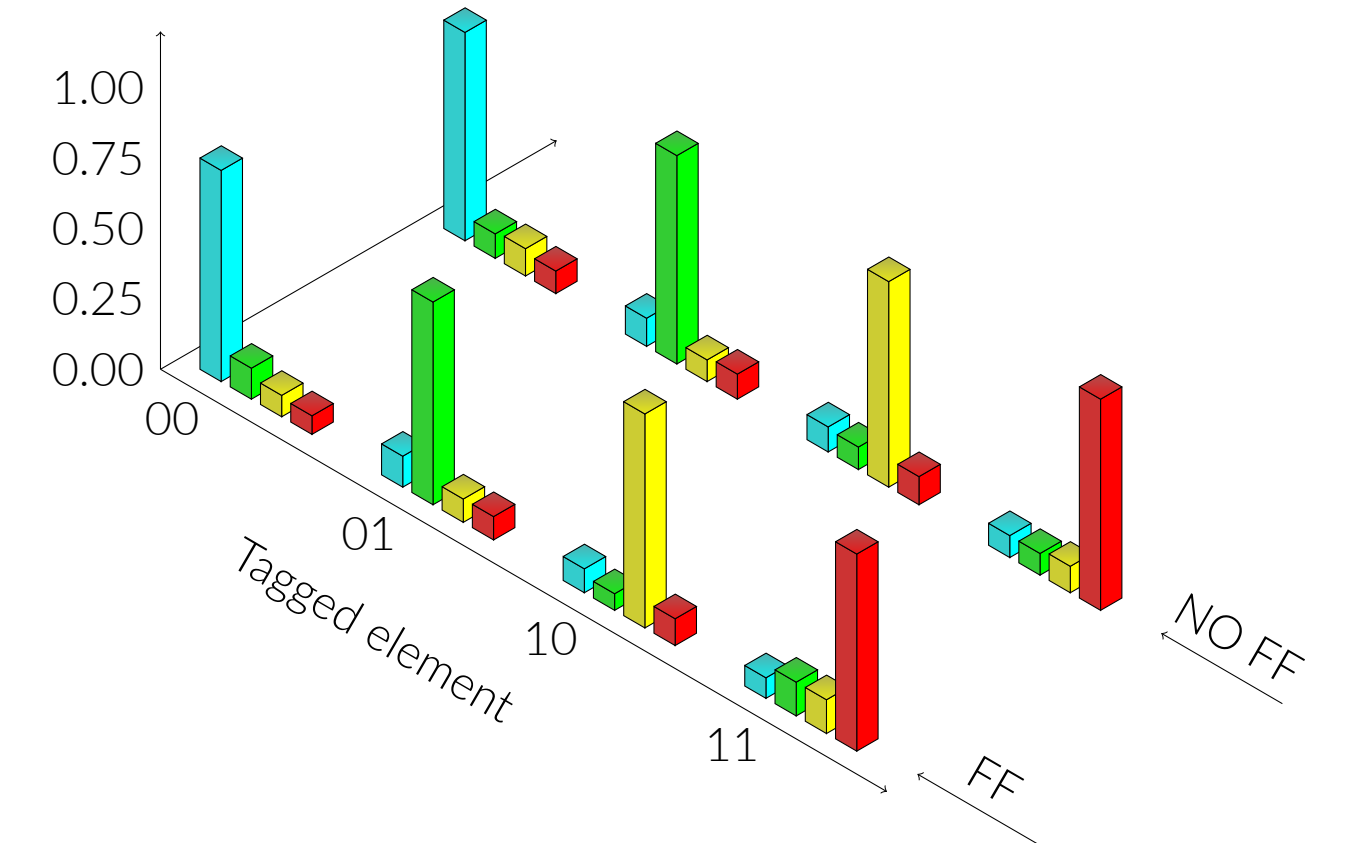


Figure 6. The measured outputs of Grover's algorithm on ibmq_vigo. The data labeled as 'NO FF' show the outputs in those cases the outputs $\{\alpha_1, \alpha_4\}$ were both zero. The data labeled 'FF' show the outputs to which the feedforward relation was applied

Case study 01

The results below were those obtained for the algorithm described in case study 01. Success probabilities are all well over 80%, making this computation almost as good as the computation in the quantum circuit model computation albeit the greater number of gates.

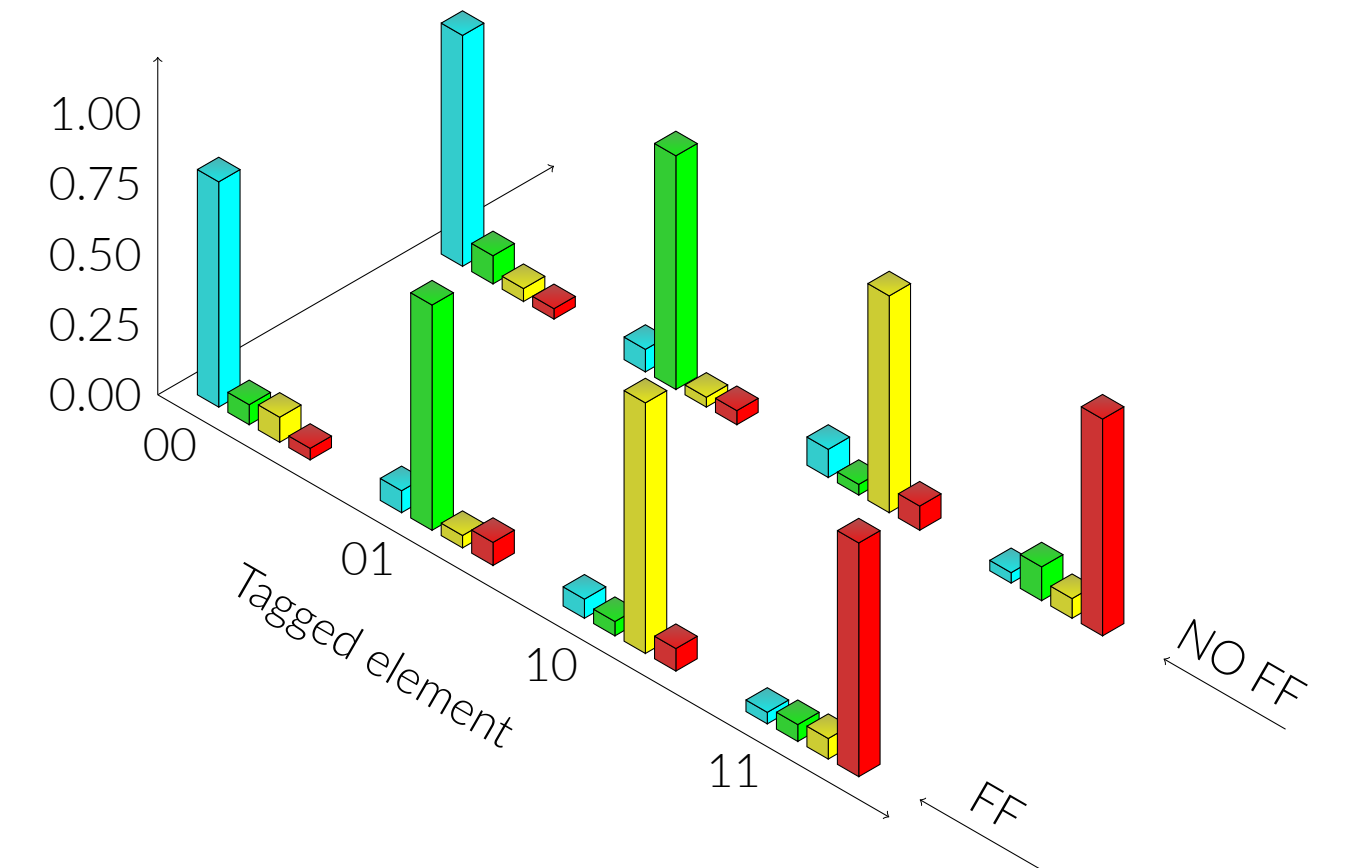


Figure 7. The measured outputs of Grover's algorithm on ibmq_vigo. The data labeled as 'NO FF' show the outputs in those cases the outputs $\{\alpha_1, \alpha_4\}$ were both zero. The data labeled 'FF' show the outputs to which the feedforward relation was applied

References

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