- 1. Correction to text about Eqn. 5.1.48 (radial distribution function)
- 1.1. **What is the issue?** Fengyu Xi brought it to my attention that the text above Equation 5.1.48 is incorrect (even though the equation is correct). Below, I expand the derivation of 5.1.48 and explain why the original text above it was not correctly describing the steps that followed.
- 1.2. **Resolution.** The value of the radial distribution function at a distance r from a reference particle is equal to the angular average of $\rho(\mathbf{r})/\rho$:

$$g(r) = \frac{1}{\rho} \int d\mathbf{\hat{r}} \ \left\langle \rho(\mathbf{r}) \right\rangle_{N-1} = \frac{1}{\rho} \int d\mathbf{\hat{r}} \ \left\langle \sum_{j \neq i} \delta(\mathbf{r} - \mathbf{r}_j) \right\rangle_{N-1} \ , \label{eq:gradient}$$

where N is the total number of particles in the system, ρ denotes the average number density ($\rho \equiv (N/V)$) and \mathbf{r}_j is the distance of particle j from the origin, where particle i is located (Note: this means that \mathbf{r}_j is actually \mathbf{r}_{ij}). $\hat{\mathbf{r}}$ is the unit vector in the direction of \mathbf{r} . For simplicity we have written down the expression for g(r) for a given particle i, and hence the sum of $j \neq i$ is keeping i fixed, but in practice the expression is averaged over all equivalent particles i. The angular brackets denote the thermal average

(2)
$$\left\langle \cdots \right\rangle_{N-1} \equiv \frac{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^N)}(\cdots)}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^N)}},$$

where we integrate over N-1 coordinates, because particle i is held fixed. We can now write

(3)
$$\left(\frac{\partial g(\mathbf{r})}{\partial \mathbf{r}} \right) = \frac{1}{\rho} \frac{\partial}{\partial \mathbf{r}} \int d\hat{\mathbf{r}} \left\langle \sum_{j \neq i} \delta(\mathbf{r} - \mathbf{r}_j) \right\rangle$$

The only term that depends on r (the length of r) is the δ -function. We can therefore write

(4)
$$\left(\frac{\partial g(r)}{\partial r} \right) = \frac{1}{\rho} \int d\hat{\mathbf{r}} \left\langle \sum_{i \neq i} \hat{\mathbf{r}} \cdot \nabla_{\mathbf{r}} \delta(\mathbf{r} - \mathbf{r}_{j}) \right\rangle$$

The following sentence was incorrect: As the arguments of the δ -function is $\mathbf{r} - \mathbf{r}_j$, we can replace $\hat{\mathbf{r}} \cdot \nabla_{\mathbf{r}}$ by $-\hat{\mathbf{r}}_j \cdot \nabla_{\mathbf{r}_j}$ and perform a partial integration. This sentence is wrong, because

2 DF

the order is actually the following:

$$\begin{pmatrix}
\frac{\partial g(r)}{\partial r}
\end{pmatrix} = \frac{-1}{\rho} \frac{\int d\hat{\mathbf{r}} \int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})} \sum_{j\neq i} \hat{\mathbf{r}} \cdot \nabla_{\mathbf{r}} \delta(\mathbf{r} - \mathbf{r}_{j})}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})}}$$

$$= \frac{+1}{\rho} \frac{\int d\hat{\mathbf{r}} \int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})} \sum_{j\neq i} \hat{\mathbf{r}} \cdot \nabla_{\mathbf{r}_{j}} \delta(\mathbf{r} - \mathbf{r}_{j})}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})}}$$

$$= \frac{-\beta}{\rho} \frac{\int d\hat{\mathbf{r}} \int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})} \sum_{j\neq i} \delta(\mathbf{r} - \mathbf{r}_{j}) \hat{\mathbf{r}} \cdot \nabla_{\mathbf{r}_{j}} U(\mathbf{r}^{N})}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})}}$$

$$= \frac{\beta}{\rho} \frac{\int d\hat{\mathbf{r}} \int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})} \sum_{j\neq i} \delta(\mathbf{r} - \mathbf{r}_{j}) \hat{\mathbf{r}}_{j} \cdot \nabla_{\mathbf{r}_{j}} U(\mathbf{r}^{N})}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})}}$$

$$= \frac{-\beta}{\rho} \frac{\int d\hat{\mathbf{r}} \int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})} \sum_{j\neq i} \delta(\mathbf{r} - \mathbf{r}_{j}) \hat{\mathbf{r}}_{j} \cdot \nabla_{\mathbf{r}_{j}} U(\mathbf{r}^{N})}{\int d\mathbf{r}^{N-1} e^{-\beta U(\mathbf{r}^{N})}}$$

$$= \frac{\beta}{\rho} \int d\hat{\mathbf{r}} \left\langle \sum_{j\neq i} \delta(\mathbf{r} - \mathbf{r}_{j}) \hat{\mathbf{r}}_{j} \cdot \mathbf{F}_{j}(\mathbf{r}^{N}) \right\rangle_{N-1}$$
(5)

where $\hat{\mathbf{r}} \cdot \mathbf{F}_j \equiv F_j^{(r)}$ denotes the force on particle j in the radial direction. NOTE: In step 2, I replace differentiation of the delta function with respect to r by differentiation with respect to $-\mathbf{r}_j$. Step 3: partial integration. Here it is important to note that r does not depend in r_j . However, in step 3, we then use the fact that the delta function imposes $\mathbf{r} = \mathbf{r}_j$.

We can now integrate with respect to r

$$g(\mathbf{r}) = g(\mathbf{r} = 0) + \frac{\beta}{\rho} \int_{0}^{\mathbf{r}} d\mathbf{r}' \int d\mathbf{\hat{r}}' \left\langle \sum_{j \neq i} \delta(\mathbf{r}' - \mathbf{r}_{j}) F_{j}^{(\mathbf{r})}(\mathbf{r}^{N}) \right\rangle_{N-1}$$

$$= g(\mathbf{r} = 0) + \frac{\beta}{\rho} \int_{\mathbf{r}' < \mathbf{r}} d\mathbf{r}' \left\langle \frac{\sum_{j \neq i} \delta(\mathbf{r}' - \mathbf{r}_{j}) F_{j}^{(\mathbf{r})}(\mathbf{r}^{N})}{4\pi \mathbf{r}'^{2}} \right\rangle_{N-1}$$

$$= g(\mathbf{r} = 0) + \frac{\beta}{\rho} \sum_{j} \left\langle \theta(\mathbf{r} - \mathbf{r}_{j}) \frac{F_{j}^{(\mathbf{r})}(\mathbf{r}^{N})}{4\pi \mathbf{r}_{j}^{2}} \right\rangle_{N-1},$$
(6)

where θ denotes the Heaviside step function. To make a connection to the results of Borgis *et al.*, we note that in a homogeneous system, all particles i of the same species are equivalent. We can therefore write

$$g(r) = g(r=0) + \frac{\beta}{N\rho} \sum_{i=1}^{N} \sum_{j \neq i} \left\langle \theta(r-r_{ij}) \frac{F_{j}^{(r)}(\textbf{r}^{N})}{4\pi r_{ij}^{2}} \right\rangle_{N-1} \; . \label{eq:gradient}$$

g(r) 3

But i and j are just dummy indices. So we obtain the same expression for g(r) by permuting i and j, except that if $\hat{\mathbf{r}} = \hat{\mathbf{r}}_{ij}$, then $\hat{\mathbf{r}} = -\hat{\mathbf{r}}_{ji}$. Adding the two equivalent expressions for g(r) and dividing by two, we get

(7)
$$g(r) = g(r = 0) + \frac{\beta}{2N\rho} \sum_{i=1}^{N} \sum_{j \neq i} \left\langle \theta(r - r_{ij}) \frac{F_{j}^{(r)}(\textbf{r}^{N}) - F_{i}^{(r)}(\textbf{r}^{N})}{4\pi r_{ij}^{2}} \right\rangle_{N-1}$$

equation (7) is equivalent to the result of Borgis *et al.*.

The remarkable feature of equation (7) is that g(r) depends not just on the number of pairs at distance r, but on all pair distances less than r. We stress that we have not assumed that the interactions in the system are pairwise additive: $F_i - F_j$ is not a pair force.