1 Evolution of operators

a. Time evolution of fermionic annihilator

Solution: Our hamiltonian for this problem is

$$H_0 = \sum_{k\sigma} \varepsilon_k c_{k\sigma}^{\dagger} c_{k\sigma}. \tag{1}$$

We can prove trivially by using Taylor series for t that

$$[e^{\pm iH_0t}, H_0] = 0. (2)$$

If we take the following relation

$$[ABC, D] = AB[C, D] + A[B, D]C + [A, D]BC,$$
 (3)

we can easily calculate the next commutator

$$[e^{iH_0t}c_{k\sigma}e^{-iH_0t}, H_0] = e^{iH_0t}c_{k\sigma}[e^{-iH_0t}, H_0] + e^{iH_0t}[c_{k\sigma}, H_0]e^{-iH_0t} + [e^{iH_0t}, H_0]c_{k\sigma}e^{-iH_0t},$$

$$= e^{iH_0t}[c_{k\sigma}, H_0]e^{-iH_0t}.$$
(4)

Finally we can use the relation

$$[A, BC] = B[A, C] + [A, B]C,$$
 (5)

to calculate the following commutator

$$[c_{k\sigma}, c_{j\gamma}^{\dagger} c_{j\gamma}] = c_{j\gamma}^{\dagger} [c_{k\sigma}, c_{j\gamma}] + [c_{k\sigma}, c_{j\gamma}^{\dagger}] c_{j\gamma},$$

$$= c_{j\gamma}^{\dagger} (c_{k\sigma} c_{j\gamma} - c_{j\gamma} c_{k\sigma}) + (c_{k\sigma} c_{j\gamma}^{\dagger} - c_{j\gamma}^{\dagger} c_{k\sigma}) c_{j\gamma},$$

$$= c_{j\gamma}^{\dagger} c_{k\sigma} c_{j\gamma} + c_{k\sigma} c_{j\gamma}^{\dagger} c_{j\gamma} - (c_{j\gamma}^{\dagger} c_{j\gamma} c_{k\sigma} + c_{j\gamma}^{\dagger} c_{k\sigma} c_{j\gamma}),$$

$$= \{c_{j\gamma}^{\dagger} c_{k\sigma}\} c_{j\gamma} - c_{j\gamma}^{\dagger} \{c_{j\gamma} c_{k\sigma}\},$$

$$= \delta_{jk} \delta_{\sigma\gamma} c_{j\gamma}.$$
(6)

Where in the last step we used the anti-commutation rules of the fermionic creation and annihilation operators.

Now we can finally start. Let's consider the time evolution equation in the Heisenberg picture

$$i\frac{d}{dt}c_{k\sigma}(t) = [c_{k\sigma}(t), H_0],\tag{7}$$

taking

$$c_{k\sigma}(t) = e^{iH_0t}c_{k\sigma}e^{-iH_0t},\tag{8}$$

the time evolution equation takes the form

$$i\frac{d}{dt}c_{k\sigma}(t) = [e^{iH_0t}c_{k\sigma}e^{-iH_0t}, H_0].$$
 (9)

Luckily we can use eq. 4 to move further and obtain

$$i\frac{d}{dt}c_{k\sigma}(t) = e^{iH_0t}[c_{k\sigma}, H_0]e^{-iH_0t}.$$
 (10)

Now is time to plug the hamiltonian into this commutator

$$i\frac{d}{dt}c_{k\sigma}(t) = e^{iH_0t}[c_{k\sigma}, \sum_{k\sigma} \varepsilon_k c_{k\sigma}^{\dagger} c_{k\sigma}]e^{-iH_0t}, \tag{11}$$

since the commutator is linear and $\varepsilon_k \in \mathbb{R}$ we can transform the equation to this one

$$i\frac{d}{dt}c_{k\sigma}(t) = e^{iH_0t} \sum_{k\sigma} \varepsilon_k [c_{k\sigma}, c_{k\sigma}^{\dagger} c_{k\sigma}] e^{-iH_0t}, \qquad (12)$$

and at this point we are lucky again, because now we can use eq. 6 to continue our calculations

$$i\frac{d}{dt}c_{k\sigma}(t) = e^{iH_0t} \sum_{k\sigma} \varepsilon_k \delta_{jk} \delta_{\sigma\gamma} c_{j\gamma} e^{-iH_0t},$$

$$= e^{iH_0t} \varepsilon_k c_{k\sigma} e^{-iH_0t},$$

$$= \varepsilon_k e^{iH_0t} c_{k\sigma} e^{-iH_0t},$$

$$= \varepsilon_k c_{k\sigma}(t),$$
(13)

therefore

$$\frac{d}{dt}c_{k\sigma}(t) = -i\varepsilon_k c_{k\sigma}(t),\tag{14}$$

at this point the equation looks very familiar, and if it looks like a duck, swims like a duck, and quacks like a duck, then it probably is a duck. To make sure we suppose it is

$$c_{k\sigma}(t) = e^{\alpha t} c_{k\sigma}. \tag{15}$$

If we differentiate this operator with respect to t, we will have

$$\frac{d}{dt}c_{k\sigma}(t) = \frac{d}{dt}e^{\alpha t}c_{k\sigma} = \alpha e^{\alpha t}c_{k\sigma} = \alpha c_{k\sigma}(t). \tag{16}$$

Comparing eqs. 14 and 16 we finally get that

$$\alpha = -i\varepsilon_k,\tag{17}$$

so finally, our solution will be

$$c_{k\sigma}(t) = e^{-i\varepsilon_k t} c_{k\sigma}. \quad \blacksquare \tag{18}$$

There are two ways of getting the creation operator, the hard one would be to repeat similar calculations with $c_{k\sigma}^{\dagger}$, which we will not perform here. The easy one is just to take

$$c_{k\sigma}^{\dagger}(t) = (c_{k\sigma}(t))^{\dagger} = (e^{-i\varepsilon_k t} c_{k\sigma})^{\dagger} = e^{i\varepsilon_k t} c_{k\sigma}^{\dagger}. \tag{19}$$

2 Non-interacting Green function

a. Project exercise II

Solution:

3 Matsubara sums and chemical potential