

1 Atomic Units

a. Speed of light

Solution: To find the speed of light in atomic units it would be useful to remember the *fine structure constant* defined as

$$\frac{e^2}{4\pi\epsilon_0\hbar c} = \frac{1}{137}, \quad (1)$$

but since $e = \hbar = 1/4\pi\epsilon_0 = 1$ eq. 1 turns into

$$\begin{aligned} \frac{1}{c} &= \frac{1}{137}, \\ \Rightarrow c &= 137. \end{aligned} \quad (2)$$

b. Bohr's Magnetron

Solution: Bohr magneton is defined as

$$\mu_B = \frac{e\hbar}{2m_e}, \quad (3)$$

now, from the fact that $e = \hbar = m_e = 1$ we can deduce the numerical value

$$\mu_B = \frac{e\hbar}{2m_e} = \frac{1}{2}. \quad (4)$$

In both questions we went from SI units to atomic ones.

2 Magnetic Moment

a. Calculate the magnetic moment

Solution: For this problem first it might be useful to remember the definition of the \hat{L} operator

$$\hat{L} = \frac{\hbar}{i} \vec{r} \times \nabla \quad (5)$$

Now let's begin our calculations

$$\begin{aligned} \vec{m} &= \frac{1}{2} \int \vec{r} \times \vec{j}_e d^3r = \frac{1}{2} \int \vec{r} \times \left(\frac{-e\hbar}{2im_e} \right) (\psi_{n,l,m}^* \nabla \psi_{n,l,m} - \psi_{n,l,m} \nabla \psi_{n,l,m}^*) d^3r, \\ &= \frac{1}{2} \int \left(\frac{-e\hbar}{2im_e} \right) \psi_{n,l,m}^* \vec{r} \times \nabla \psi_{n,l,m} - \left(\frac{-e\hbar}{2im_e} \right) \psi_{n,l,m} \vec{r} \times \nabla \psi_{n,l,m}^* d^3r, \\ &= \frac{1}{2} \left(\frac{-e}{2m_e} \right) \int \psi_{n,l,m}^* \hat{L} \psi_{n,l,m} - \psi_{n,l,m} \hat{L} \psi_{n,l,m}^* d^3r, \quad \text{using eq. 5,} \\ &= \frac{1}{2} \left(\frac{-e}{2m_e} \right) 2 \langle \psi_{n,l,m} | \hat{L} | \psi_{n,l,m} \rangle, \\ &= \frac{-e}{2m_e} \langle \psi_{n,l,m} | \hat{L} | \psi_{n,l,m} \rangle, \end{aligned}$$

$$= \frac{-e\hbar}{2\hbar m_e} \langle \psi_{n,l,m} | \hat{L} | \psi_{n,l,m} \rangle,$$

and finally

$$= \frac{-\mu_B}{\hbar} \langle \psi_{n,l,m} | \hat{L} | \psi_{n,l,m} \rangle, \quad \text{using eq. 7} \quad (6)$$

where in the last step we used the definition of Bohr magneton

$$\mu_B = \frac{e\hbar}{2m_e}. \quad (7)$$

So as we can see there's a direct connection between \vec{m} and the expectation value of the angular momentum operator.

b. z-components

Solution: If we are only interested in the z component of the magnetic moment then from eq. 6 we can isolate the z term of \hat{L} , and the resulting equation will be

$$m_z = -\frac{\mu_B}{\hbar} \langle \psi_{n,l,m} | \hat{L}_z | \psi_{n,l,m} \rangle, \quad (8)$$

if we remember, the eigenvalue equation for \hat{L}_z is

$$\hat{L}_z | \psi_{n,l,m} \rangle = \hbar m | \psi_{n,l,m} \rangle. \quad (9)$$

Plugging eq. 9 into eq. 8 we get

$$\begin{aligned} m_z &= -\frac{\mu_B}{\hbar} \langle \psi_{n,l,m} | \hat{L}_z | \psi_{n,l,m} \rangle, \\ &= -\frac{\mu_B}{\hbar} \langle \psi_{n,l,m} | \hbar m | \psi_{n,l,m} \rangle, \\ &= -\mu_B m \langle \psi_{n,l,m} | \psi_{n,l,m} \rangle, \\ &= -\mu_B m. \end{aligned}$$

Therefore

$$[m_{n,l,m}]_z = -\mu_B m, \quad (10)$$

and we can use this to calculate what we need.

- $[m_{100}]_z = -\mu_B(m) = -\mu_B(0) = 0$
- $[m_{200}]_z = 0$, for the same reason as above
- $[m_{21-1}]_z = -\mu_B(-1) = \mu_B$
- $[m_{210}]_z = -\mu_B(0) = 0$
- $[m_{211}]_z = -\mu_B(1) = -\mu_B$
- $[m_{532}]_z = -\mu_B(2) = -2\mu_B$

3 Charge States

a. Formal Charge

Solution: In the case of the first compound we have KCrF_3 , the charges we know are

- $\text{K} = 1$.
- $\text{F}_3 = 3(-1) = -3$.

So in order to neutralize the charge we need

- $\text{Cr} = -(1 - 3) = -(-2) = 2$.

The next compound is SrMnO_3 , the charges we know are

- $\text{Sr} = 2$.
- $\text{O}_3 = 3(-2) = -6$

Hence, the charge we need is

- $\text{Mn} = -(2 - 6) = -(-4) = 4$.

4 Atomic Radii

a. Read article

Solution: I did not read it but it's on my to-do list now ;)