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## On the universality of almost every quantum logic gate

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Lloyd [Phys. Rev. Lett. **75**, 346 (1995)] showed that almost every quantum logic gate is universal in the sense that it can be used to approximate any unitary transformation. The argument relied on a more general fact whose proof was not given in detail. We give a complete proof of this more general fact. © *2000 American Institute of Physics*. [S0022-2488(00)02201-5]

In Ref. 1 Lloyd announced the following result. Let A and B be Hermitian matrices of dimension at least 2 and let  $\mathcal{L}$  be the Lie algebra they generate though commutation—that is,  $\mathcal{L}$  is the real linear span of the matrices

then for any  $L \in \mathcal{L}$ , the unitary matrix  $U = e^{iL}$  can be expressed in the form

$$U = e^{iAt_1}e^{iBt_2}e^{iAt_3}e^{iBt_4}\cdots. \tag{*}$$

This implies that almost every A and B are universal in the sense that any unitary matrix U can be realized by the expression (\*).

Informally, one thinks of A as the intrinsic Hamiltonian of a quantum system and takes B to be a different Hamiltonian resulting from some external influence which can be applied at will. By turning B on and off for successive time intervals of various lengths, one can achieve any time evolution of the form (\*), and the claim is that for almost every A and B this suffices to produce any desired unitary evolution. It follows that almost any quantum logic gate with two inputs is universal. This verifies a conjecture of Deutsch<sup>2</sup> and generalizes a result of Deutsch, Barenco, and Ekert.<sup>3</sup>

The proof was only sketched in Ref. 1, and actually the claim is not exactly true. The problem involves the use of negative values for  $t_j$ , clearly a practical impossibility, first explicitly in the expression

$$(e^{-iB\sqrt{t/n}}e^{-iA\sqrt{t/n}}e^{-iB\sqrt{t/n}}e^{iA\sqrt{t/n}})^n$$

and then implicitly in the assertion that the unitaries given by (\*) form a manifold. However, this problem is irrelevant to the main issue of universality of quantum logic gates, since, as Lloyd notes later in his paper,  $e^{iAt}$  can be approximated with arbitrary accuracy by the powers of some fixed  $e^{iAt_A}$ . This is true for both positive and negative values of t, so the difficulty does not invalidate the main line of argument.

Nonetheless, this objection is genuine, as the following example shows. Let A and B be the  $2\times 2$  matrices,

$$A = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix}, \quad B = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

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(a,b real). Then any expression of the form (\*) with non-negative  $t_j$  can be simplified to  $e^{iAt}$  for some  $t \ge 0$ . But if a/b is irrational, then  $U = e^{-iA}$  cannot be expressed in such a form. For then we would have

$$e^{-iA} = e^{iAt}$$
,

hence

$$\begin{pmatrix} e^{-ia} & 0 \\ 0 & e^{-ib} \end{pmatrix} = \begin{pmatrix} e^{iat} & 0 \\ 0 & e^{ibt} \end{pmatrix}$$

and

$$e^{i(at+a)} = 1 = e^{i(bt+b)}$$
.

so that at+a and bt+b are both integer multiples of  $2\pi$ . Since  $t \neq -1$  this implies that a/b = (at+a)/(bt+b) is rational.

In the previous example the operator  $e^{-iA}$  can be *approximated* by  $e^{iAt}$  for positive values of t. Indeed, in general if A is a Hermitian matrix of any finite dimension then given any s>0 we may find t>s such that  $\lambda t$  is approximately an integer multiple of  $2\pi$ , simultaneously for every eigenvalue  $\lambda$  of A. Then  $e^{iAt}$  approximates the identity matrix, so  $e^{iA(t-s)}$  approximates  $e^{-iAs}$ . This is essentially the same as the observation following Eq. (3) in Ref. 1.

Thus, the reasoning in Ref. 1 does imply that for every  $L \in \mathcal{L}$  the unitary matrix  $e^{iL}$  can be approximated by operators of the form (\*). This does not settle the question of exact representation of unitary matrices. However, the approximate result can be used to prove an exact result.

We review the argument that verifies this approximate result, which states: Let A and B be  $n \times n$  Hermitian matrices ( $n \ge 2$ ) and let  $\mathcal{L}$  be the Lie algebra they generate through commutation; then for any  $L \in \mathcal{L}$ , the unitary matrix  $U = e^{iL}$  can be approximated by finite products of the form (\*), with each  $t_j$  positive. The proof proceeds on the complexity of L. If L = A or L = B the conclusion is immediate. For L = i[A, B] we have

$$\lim_{k\to\infty} (e^{iA/\sqrt{k}}e^{iB/\sqrt{k}}e^{-iA/\sqrt{k}}e^{-iB/\sqrt{k}})^k = e^{iL}.$$

Thus  $e^{iL}$  can be approximated by finite products of the form (\*), if the  $t_j$  are allowed to take negative values. But as we previously noted, the negative exponents  $e^{-iA/\sqrt{k}}$  and  $e^{-iB/\sqrt{k}}$  can be approximated by positive exponents; so in fact  $e^{iL}$  can be approximated by finite products of the form (\*) with positive  $t_j$ . The conclusion can be inductively extended to more complex commutators.

Using this approximate result, we can now prove the desired theorem: For almost all Hermitian  $n \times n$  matrices A and B  $(n \ge 2)$ , every unitary  $n \times n$  matrix can be exactly represented in the form (\*) with each  $t_i \ge 0$ .

Observe first that it will suffice to consider only unitary matrices which are close to the identity matrix I, because any unitary matrix is a power of one close to the identity. That is, if  $U = e^{iL}$  then  $U = (e^{iL/k})^k$ , so we need only to represent  $e^{iL/k}$  in the form (\*) for large k.

Indeed, it will suffice to represent all unitaries close to any given unitary  $U_0$ . For any unitary close to I may be expressed as a power of  $U_0$  times a unitary close to  $U_0$ . In other words, if we can represent every V within some neighborhood of  $U_0$  then in particular we can represent  $U_0$  itself, and hence also anything of the form  $U_0^k V$ . Choosing a power of  $U_0$  which is sufficiently close to  $U_0^{-1}$ , this left multiplication by  $U_0^k$  will take the neighborhood of  $U_0$  onto a neighborhood of I.

Let  $\mathcal{L}$  be the Lie algebra generated by A and B through commutation. As noted in Ref. 1, for almost every A and B this algebra contains every Hermitian matrix. Therefore we assume this is the case.

Find unitaries  $V_1,\ldots,V_{n^2}$  such that  $\{V_1AV_1^{-1},\ldots,V_{n^2}AV_{n^2}^{-1}\}$  span the real vector space of all  $n\times n$  Hermitian matrices. This can be done provided  $\operatorname{tr}(A)\neq 0$  and A is not a scalar multiple of I, and such matrices A constitute a set of full measure; to avoid interrupting the main line of argument we postpone verification of this claim to the end of the proof. We may also take  $V_1=I$ . Now by the approximate representation result, find unitaries  $U_1,\ldots,U_{n^2-1}$  each in the form (\*) with positive  $t_j$ , such that  $U_j$  approximates  $e^{-iA}V_j^{-1}V_{j+1}$ . Then let  $U_{n^2}=e^{-iA}$  and  $U_0=e^{iA}U_1e^{iA}U_2\ldots e^{iA}U_{n^2}$ .

Let *M* be the manifold of all  $n \times n$  unitary matrices and consider the map  $\Phi: \mathbb{R}^{n^2} \to M$  defined by

$$\Phi(s_1,...,s_{n^2}) = e^{iAs_1}U_1e^{iAs_2}U_2\cdots e^{iAs_{n^2}}U_{n^2}.$$

If the  $s_j$  are all positive then  $\Phi(s_1,...,s_{n^2})$  is evidently representable in the form (\*) with positive  $t_j$ . We will show that  $\Phi$  maps neighborhoods of the point (1, 1,...,1) onto neighborhoods of  $U_0$ , and this will complete the proof.

By the implicit function theorem (see, e.g., Ref. 4), this will follow if we can show that the Jacobian of  $\Phi$  at (1, 1,...,1) is nonzero. Equivalently we may consider the map  $\Phi_0$  defined by  $\Phi_0(s_1,...,s_{n^2}) = \Phi(s_1,...,s_{n^2}) \cdot U_0^{-1}$ .

A simple calculation shows that

$$\left. \frac{\partial \Phi}{\partial s_j} \right|_{s_1 = \dots = s_{n^2} = 1} = i e^{iA} U_1 \dots e^{iA} U_{j-1} A e^{iA} U_j \dots e^{iA} U_{n^2}$$

and therefore

$$\left. \frac{\partial \Phi_0}{\partial s_j} \right|_{s_1 = \dots = s_n = 1} = i e^{iA} U_1 \cdots e^{iA} U_{j-1} A U_{j-1}^{-1} e^{-iA} \cdots U_1^{-1} e^{-iA}.$$

By the definition of the  $U_i$ 's, this simplifies to

$$\left. \frac{\partial \Phi_0}{\partial s_j} \right|_{s_1 = \dots = s_n^2 = 1} \approx i V_j A V_j^{-1},$$

which shows that partial derivatives of  $\Phi_0$  at the origin span the space of skew-Hermitian matrices, by our original choice of the  $V_j$ 's. Thus its Jacobian is nonzero at the point (1, 1, ..., 1), as desired.

This completes the proof, modulo the claim that if  $\operatorname{tr}(A) \neq 0$  and A is not a scalar multiple of I then the matrices  $VAV^{-1}$ , as V ranges over all unitaries, span the real vector space of Hermitian matrices. To verify this we use the fact that  $\langle A,B\rangle=\operatorname{tr}(AB)$  defines an inner product which makes the Hermitian matrices into a real Hilbert space. Thus if the matrices  $VAV^{-1}$  do not span this space then there must be a Hermitian matrix B such that  $\operatorname{tr}(VAV^{-1}B)=0$  for every unitary V. Taking V so that B and  $VAV^{-1}$  are simultaneously diagonalizable, we find that  $\sum a_ib_i=0$  where  $a_i$  are the eigenvalues of A—in any order—and  $b_i$  are the eigenvalues of B. Since A is not a scalar multiple of I, the  $a_i$  are not all identical, and since  $\operatorname{tr}(A)\neq 0$  neither are the  $b_i$ . Thus there must exist indices  $i_0$  and  $i_1$  such that  $a_{i_0}\neq a_{i_1}$  and  $b_{i_0}\neq b_{i_1}$ , and this implies that  $(a_{i_0}-a_{i_1})(b_{i_0}-b_{i_1})\neq 0$ , hence

$$a_{i_0}b_{i_0} + a_{i_1}b_{i_1} \neq a_{i_0}b_{i_1} + a_{i_1}b_{i_0}$$

Therefore we cannot also have  $\sum a_i b_i = 0$  for the rearrangement of the  $a_i$  which switches  $a_{i_0}$  and  $a_{i_1}$ . This contradiction shows that the matrices  $VAV^{-1}$  must span the space of Hermitian matrices.

An interesting feature of this solution is that one cannot achieve or even approximate unitary matrices close to the identity in arbitrarily short times. Because of the restriction to positive  $t_j$ , in short times one can only reach unitaries which are, so to speak, "on one side" of the identity. A more full explanation of this phenomenon in control-theoretic terms will be given elsewhere.

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<sup>&</sup>lt;sup>1</sup>S. Lloyd, Phys. Rev. Lett. **75**, 346 (1995).

<sup>&</sup>lt;sup>2</sup>D. Deutsch, Proc. R. Soc. London, Ser. A **425**, 73 (1989).

<sup>&</sup>lt;sup>3</sup>D. Deutsch, A. Barenco, and A. Ekert, Proc. R. Soc. London, Ser. A 449, 669 (1995).

<sup>&</sup>lt;sup>4</sup>J. E. Marsden, *Elementary Classical Analysis* (Freeman, New York, 1974).