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18.02 Multivariable Calculus
Fall 2007

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18.02 Lecture 30. – Tue, Nov 27, 2007

Handouts: practice exams 4A and 4B.

Clarification from end of last lecture: we derived the diffusion equation from 2 inputs. u = concentration, \mathbf{F} = flow, satisfy:

- 1) from physics: $\mathbf{F} = -k\nabla u$,
- 2) from divergence theorem: $\partial u / \partial t = -\text{div } \mathbf{F}$.

Combining, we get the diffusion equation: $\partial u / \partial t = -\text{div } \mathbf{F} = +k \text{div } (\nabla u) = k \nabla^2 u$.

Line integrals in space.

Force field $\mathbf{F} = \langle P, Q, R \rangle$, curve C in space, $d\vec{r} = \langle dx, dy, dz \rangle$

$$\Rightarrow \text{Work} = \int_C \mathbf{F} \cdot d\vec{r} = \int_C P dx + Q dy + R dz.$$

Example: $\mathbf{F} = \langle yz, xz, xy \rangle$. C : $x = t^3$, $y = t^2$, $z = t$. $0 \leq t \leq 1$. Then $dx = 3t^2 dt$, $dy = 2t dt$, $dz = dt$ and substitute:

$$\int_C \mathbf{F} \cdot d\vec{r} = \int_C yz dx + xz dy + xy dz = \int_0^1 6t^5 dt = 1$$

(In general, express (x, y, z) in terms of a *single* parameter: 1 degree of freedom)

Same \mathbf{F} , curve C' = segments from $(0, 0, 0)$ to $(1, 0, 0)$ to $(1, 1, 0)$ to $(1, 1, 1)$. In the xy -plane, $z = 0 \Rightarrow \mathbf{F} = xy\hat{\mathbf{k}}$, so $\mathbf{F} \cdot d\vec{r} = 0$, no work. For the last segment, $x = y = 1$, $dx = dy = 0$, so $\mathbf{F} = \langle z, z, 1 \rangle$ and $d\vec{r} = \langle 0, 0, dz \rangle$. We get $\int_0^1 1 dz = 1$.

Both give the same answer because \mathbf{F} is conservative, in fact $\mathbf{F} = \nabla(xyz)$.

Recall the fundamental theorem of calculus for line integrals:

$$\int_{P_0}^{P_1} \nabla f \cdot d\vec{r} = f(P_1) - f(P_0).$$

Gradient fields.

$$\mathbf{F} = \langle P, Q, R \rangle = \langle f_x, f_y, f_z \rangle ?$$

Then $f_{xy} = f_{yx}$, $f_{xz} = f_{zx}$, $f_{yz} = f_{zy}$, so $P_y = Q_x$, $P_z = R_x$, $Q_z = R_y$.

Theorem: \mathbf{F} is a gradient field if and only if these equalities hold (assuming defined in whole space or simply connected region)

Example: for which a, b is $axy\hat{\mathbf{i}} + (x^2 + z^3)\hat{\mathbf{j}} + (byz^2 - 4z^3)\hat{\mathbf{k}}$ a gradient field?

$$P_y = ax = 2x = Q_x \text{ so } a = 2; P_z = 0 = 0 = R_x; Q_z = 3z^2 = bz^2 = R_y \text{ so } b = 3.$$

Systematic method to find a potential: (carried out on above example)

$$f_x = 2xy, f_y = x^2 + z^3, f_z = 3yz^2 - 4z^3:$$

$$f_x = 2xy \Rightarrow f(x, y, z) = x^2y + g(y, z).$$

$$f_y = x^2 + g_y = x^2 + z^3 \Rightarrow g_y = z^3 \Rightarrow g(y, z) = yz^3 + h(z), \text{ and } f = x^2y + yz^3 + h(z).$$

$$f_z = 3yz^2 + h'(z) = 3yz^2 - 4z^3 \Rightarrow h'(z) = -4z^3 \Rightarrow h(z) = -z^4 + c, \text{ and } f = x^2y + yz^3 - z^4 + c.$$

Other method: $f(x_1, y_1, z_1) = f(0, 0, 0) + \int_{P_0}^{P_1} \mathbf{F} \cdot d\vec{r}$ (use a curve that gives an easy computation, e.g. 3 segments parallel to axes).

Curl: encodes by how much \mathbf{F} fails to be conservative.

$$\text{curl } \langle P, Q, R \rangle = (R_y - Q_z)\hat{\mathbf{i}} + (P_z - R_x)\hat{\mathbf{j}} + (Q_x - P_y)\hat{\mathbf{k}}.$$

How to remember the formula? Use the del operator $\nabla = \langle \partial/\partial x, \partial/\partial y, \partial/\partial z \rangle$.

Recall from last week that $\nabla \cdot \mathbf{F} = \langle \partial/\partial x, \partial/\partial y, \partial/\partial z \rangle \cdot \langle P, Q, R \rangle = P_x + Q_y + R_z = \text{div } \mathbf{F}$.

$$\text{Now we have: } \nabla \times \mathbf{F} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ P & Q & R \end{vmatrix} = \text{curl } \mathbf{F}.$$

Interpretation of curl for velocity fields: curl measures *angular velocity*.

Example: rotation around z -axis at constant angular velocity ω (trajectories are circles centered on z -axis): $\mathbf{v} = \langle -\omega y, \omega x, 0 \rangle$.

Then $\nabla \times \mathbf{v} = \dots = 0\hat{\mathbf{i}} + 0\hat{\mathbf{j}} + (\omega + \omega)\hat{\mathbf{k}} = 2\omega\hat{\mathbf{k}}$. So length of curl = twice angular velocity, and direction = axis of rotation.

A general motion can be complicated, but decomposes into various effects.

- curl measures the *rotation* component of a complex motion.

18.02 Lecture 31. – Thu, Nov 29, 2007

Handouts: PS11 solutions, PS12.

Stokes' theorem is the 3D analogue of Green's theorem for work (in the same sense as the divergence theorem is the 3D analogue of Green for flux).

$$\text{Recall } \text{curl } \mathbf{F} = \nabla \times \mathbf{F} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ P & Q & R \end{vmatrix} = \langle R_y - Q_z, P_z - R_x, Q_x - P_y \rangle.$$

Stokes' theorem: if C is a *closed curve*, and S *any* surface bounded by C , then

$$\oint_C \mathbf{F} \cdot d\vec{r} = \iint_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dS.$$

Orientation: compatibility of an orientation of C with an orientation of S (changing orientation changes sign on both sides of Stokes).

Rule: if I walk along C in positive direction, with S to my left, then $\hat{\mathbf{n}}$ is pointing up. (Various examples shown.)

Another formulation (right-hand rule): if thumb points along C (1-D object), index finger towards S (2-D object), then middle finger points along $\hat{\mathbf{n}}$ (3-D object).

More examples shown.

Example: Stokes vs. Green. If S is a portion of xy -plane bounded by a curve C counterclockwise, then $\oint_C \mathbf{F} \cdot d\vec{r} = \int_C P dx + Q dy$, by Green this is equal to $\iint_S (Q_x - P_y) dx dy = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{k}} dx dy = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{n}} dS$, so Green and Stokes say the same thing in this example.

Remark. In Stokes' theorem we are free to choose any surface S bounded by C ! (e.g. if C = circle, S could be a disk, a hemisphere, a cone, ...)

“Proof” of Stokes.

- 1) if C and S are in the xy -plane then the statement follows from Green.
- 2) if C and S are in an arbitrary plane: this also reduces to Green in the given plane. Green/Stokes works in any plane because of *geometric invariance* of work, curl and flux under rotations of space. They can be defined in purely geometric terms so as not to depend on the coordinate system (x, y, z) ; equivalently, we can choose coordinates (u, v, w) adapted to the given plane, and work

with those coordinates, the expressions of work, curl, flux will be the familiar ones replacing x, y, z with u, v, w .

3) in general, we can decompose S into small pieces, each piece is nearly flat (slanted plane); on each piece we have approximately work = flux by Green's theorem. When adding pieces, the line integrals over the inner boundaries cancel each other and we get the line integral over C ; the flux integrals add up to flux through S .

Example: verify Stokes for $\mathbf{F} = z\hat{\mathbf{i}} + x\hat{\mathbf{j}} + y\hat{\mathbf{k}}$, $C =$ unit circle in xy -plane (counterclockwise), $S =$ piece of paraboloid $z = 1 - x^2 - y^2$.

Direct calculation: $x = \cos \theta$, $y = \sin \theta$, $z = 0$, so $\oint_C \mathbf{F} \cdot d\vec{r} = \int_C z dx + x dy + y dz = \oint_C x dy = \int_0^{2\pi} \cos^2 \theta d\theta = \pi$.

By Stokes: $\text{curl } \mathbf{F} = \langle 1, 1, 1 \rangle$, and $\hat{\mathbf{n}} dS = \langle -f_x, -f_y, 1 \rangle dx dy = \langle 2x, 2y, 1 \rangle dx dy$.

$$\iint_S (\nabla \times \mathbf{F}) \cdot d\vec{S} = \iint_S \langle (2x + 2y + 1) \rangle dx dy = \iint_S 1 dx dy = \text{area}(\text{disk}) = \pi.$$

($\iint x dx dy = 0$ by symmetry and similarly for y).

18.02 Lecture 32. – Fri, Nov 30, 2007

Stokes and path independence.

Definition: a region is simply connected if every closed loop C inside it bounds some surface S inside it.

Example: the complement of the z -axis is not simply connected (shown by considering a loop encircling the z -axis); the complement of the origin is simply connected.

Topology: uses these considerations to classify for example surfaces in space: e.g., the mathematical proof that a sphere and a torus are “different” surfaces is that the sphere is simply connected, the torus isn't (in fact it has two “independent” loops that don't bound).

Recall: if \mathbf{F} is a gradient field then $\text{curl}(\mathbf{F}) = 0$.

Conversely, Theorem: if $\nabla \times \mathbf{F} = 0$ in a *simply connected* region then \mathbf{F} is conservative (so $\int \mathbf{F} \cdot d\vec{r}$ is path-independent and we can find a potential).

Proof: Assume R simply connected, $\nabla \times \mathbf{F} = 0$, and consider two curves C_1 and C_2 with same end points. Then $C = C_1 - C_2$ is a closed curve so bounds some S ; $\int_{C_1} \mathbf{F} \cdot d\vec{r} - \int_{C_2} \mathbf{F} \cdot d\vec{r} = \oint_C \mathbf{F} \cdot d\vec{r} = \iint_S \nabla \times \mathbf{F} \cdot \hat{\mathbf{n}} dS = 0$.

Orientability. We can apply Stokes' theorem to any surface S bounded by C ... provided that it has a well-defined normal vector! Counterexample shown: the Möbius strip. It's a one-sided surface, so we can't compute flux through it (no possible consistent choice of orientation of $\hat{\mathbf{n}}$). Instead, if we want to apply Stokes to the boundary curve C , we must find a two-sided surface with boundary C . (pictures shown).

Stokes and surface independence.

In Stokes we can choose any S bounded by C : so if a same C bounds two surfaces S_1, S_2 , then $\oint_C \mathbf{F} \cdot d\vec{r} = \iint_{S_1} \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iint_{S_2} \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS$? Can we prove directly that the two flux integrals are equal?

Answer: change orientation of S_2 , then $S = S_1 - S_2$ is a closed surface with $\hat{\mathbf{n}}$ outwards; so we can apply the divergence theorem: $\iint_S \text{curl } \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iiint_D \text{div}(\text{curl } \mathbf{F}) dV$. But $\text{div}(\text{curl } \mathbf{F}) = 0$,

always. (Checked by calculating in terms of components of \mathbf{F} ; also, symbolically: $\nabla \cdot (\nabla \times \mathbf{F}) = 0$, much like $u \cdot (u \times v) = 0$ for genuine vectors).

Review for Exam 4.

We've seen three types of integrals, with different ways of evaluating:

1) $\iiint f dV$ in rect., cyl., spherical coordinates (I re-explained the general setup and the formulas for dV); applications: center of mass, moment of inertia, gravitational attraction.

2) surface integrals, flux. Setting up $\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} dS$, by knowing formulas for $\hat{\mathbf{n}} dS$.

We have seen: planes parallel to coordinate planes (e.g. yz -plane: $\hat{\mathbf{n}} = \pm \hat{\mathbf{i}}$, $dS = dy dz$); spheres and cylinders ($\hat{\mathbf{n}}$ = straight out/in from center or axis; $dS = a dz d\theta$ for cylinders, $a^2 \sin \phi d\phi d\theta$ for spheres); if we can express $z = f(x, y)$, $\hat{\mathbf{n}} dS = \pm \langle -f_x, -f_y, 1 \rangle dx dy$ (recall $\langle \dots \rangle$ is not $\hat{\mathbf{n}}$ and $dx dy$ is not dS); if S has a given normal vector \vec{N} (e.g. if S is given by $g(x, y, z) = 0$), $\hat{\mathbf{n}} dS = \pm \vec{N} / (\vec{N} \cdot \hat{\mathbf{k}}) dx dy$.

3) line integrals $\int_C \mathbf{F} \cdot d\vec{r}$ ($= \int_C P dx + Q dy + R dz$), evaluate by parameterizing C and expressing in terms of a single variable.

While these various types of integrals are completely different in terms of interpretation and method of evaluation, we've seen some theorems that establish bridges between them:

a) (\iint vs \iiint) divergence theorem: S closed surface, $\hat{\mathbf{n}}$ outwards, then $\iint_S \mathbf{F} \cdot \hat{\mathbf{n}} dS = \iiint_D (\operatorname{div} \mathbf{F}) dV$.

b) (\int vs \iint) Stokes' theorem: C closed curve bounding S compatibly oriented, then $\int_C \mathbf{F} \cdot d\vec{r} = \iint_S (\nabla \times \mathbf{F}) \cdot \hat{\mathbf{n}} dS$.

Both sides of these theorems are integrals of the types discussed above, and are evaluated by the usual methods! (even if the integrand happens to be a div or a curl).

In fact, another conceptually similar bridge exists between no integral at all and line integral: the fundamental theorem of calculus, $f(P_1) - f(P_0) = \int_C \nabla f \cdot d\vec{r}$.

One more topic: given \mathbf{F} with $\operatorname{curl} \mathbf{F} = 0$, finding a potential function.