

**PROBLEM SET 8***Reading in Kittel & Kroemer:*

- Ch. 8, pp. 227-232 (Heat and Work)
- Ch. 7 (Fermi and Bose Gases)

1. *Heat engine with finite reservoirs.* Two identical bodies, each characterized by a heat capacity at constant pressure  $C_p$  which is independent of temperature, are used as heat reservoirs for a heat engine (Fig. 1). The bodies remain at constant pressure. Initially, their temperatures are  $\tau_1$  and  $\tau_2$ , respectively; finally, as a result of the operation of the heat engine, the bodies will attain a common final temperature  $\tau_f$ .
  - a. What is the total amount of work  $W$  done by the engine? Express the answer in terms of  $C_p$ ,  $\tau_1$ ,  $\tau_2$ , and  $\tau_f$ .  
*Hint:* By definition, the heat engine operates in a cycle, i.e., at the end of the process, the heat engine is in the same macrostate in which it began.
  - b. Use arguments based upon entropy considerations to derive an inequality relating  $\tau_f$  to the initial temperatures  $\tau_1$  and  $\tau_2$ .
  - c. For given initial temperatures  $\tau_1$  and  $\tau_2$ , what is the maximum amount of work obtainable from the engine?

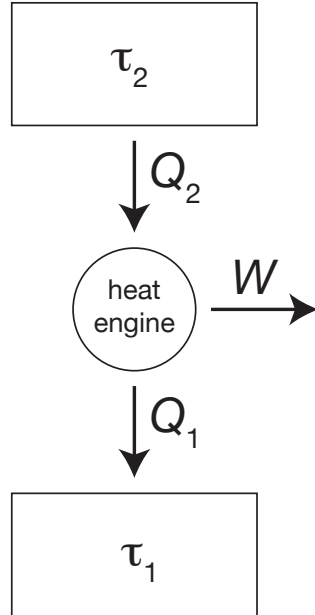


FIG. 1. Schematic of the heat engine and reservoirs in Problem 1.

2. *Pressure of a degenerate Fermi gas.* The equation of state of a strongly interacting Fermi gas, relevant to the study of neutron stars, is difficult to calculate from first principles but can be measured in experiments with ultracold atoms [1]. These experiments make use of the Gibbs-Duhem equation

$$dp = n d\mu + s d\tau \quad (1)$$

to determine the pressure  $p$  from measurements of the atomic density  $n \equiv N/V$  as a function of chemical potential at fixed temperature  $\tau$ ;  $s \equiv \sigma/V$  denotes the entropy density.

- a. Derive the Gibbs-Duhem equation (Eq. 1).
- b. Show that a *non-interacting* Fermi gas in the ground state exerts a pressure

$$p = \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \left( \frac{N}{V} \right)^{5/3}. \quad (2)$$

3. *Mass-radius relationship for white dwarfs* (K&K 7.6). Consider a white dwarf star of mass  $M$  and radius  $R$ . Let the electrons be degenerate but non-relativistic; the protons are nondegenerate.

- a. Show that the order of magnitude of the gravitational self-energy is  $-GM^2/R$ , where  $G$  is the gravitational constant. (If the mass density is constant within the sphere of radius  $R$ , the exact potential energy is  $-3GM^2/5R$ .)
- b. Show that the order of magnitude of the kinetic energy of the electrons in the ground state is

$$\frac{\hbar^2 N^{5/3}}{mR^2} \approx \frac{\hbar^2 M^{5/3}}{mM_H^{5/3} R^2}, \quad (3)$$

where  $m$  is the mass of an electron and  $M_H$  is the mass of a proton.

- c. Show that if the gravitational and kinetic energies are of the same order of magnitude (as required by the virial theorem of mechanics),  $M^{1/3}R \approx 10^{20} \text{g}^{1/3} \text{cm}$ .
  - d. If the mass is equal to that of the Sun ( $2 \times 10^{33} \text{g}$ ), what is the density of the white dwarf?
  - e. It is believed that pulsars are stars composed of a cold degenerate gas of neutrons. Show that for a neutron star  $M^{1/3}R \approx 10^{17} \text{g}^{1/3} \text{cm}$ . What is the value of the radius for a neutron star with a mass equal to that of the Sun? Express the result in km.
4. *Boson gas in one dimension* (from K&K 7.9). Calculate the integral representing the number  $N_e$  of atoms in excited states for a one-dimensional gas of non-interacting bosons. Show that the integral does not converge. This result suggests that an ideal boson gas in one dimension does not form a condensate. *Note:* take  $\mu = 0$  for the calculation, and give a justification for this approximation.
  5. *Energy, heat capacity, and entropy of a 3D Bose gas* (from K&K 7.8). Find expressions as a function of temperature in the region  $\tau < \tau_E$  for the energy, heat capacity, and entropy of a gas of  $N$  non-interacting bosons of spin zero confined to a volume  $V$ . Put the definite integral in dimensionless form; it need not be evaluated.
  6. *Fluctuations of Fermi and Bose gases* (from K&K 7.11-7.12).

- a. Show for a single orbital of a non-interacting fermion system that

$$\langle (\Delta N)^2 \rangle = \langle N \rangle (1 - \langle N \rangle), \quad (4)$$

where  $\langle N \rangle$  is the average number of fermions in that orbital. (By definition,  $\Delta N = N - \langle N \rangle$ .)

- b. Show that for a single orbital of a non-interacting boson system,

$$\langle (\Delta N)^2 \rangle = \langle N \rangle (1 + \langle N \rangle). \quad (5)$$

- c. Sketch the variance  $\langle (\Delta N)^2 \rangle$  vs.  $\langle N \rangle$  for fermions and bosons on a single plot, labeling each. Notice that for fermions, the fluctuation vanishes for orbitals with energies deep enough below the Fermi energy so that  $\langle N \rangle = 1$ . By contrast, for bosons, if the occupancy is large ( $\langle N \rangle \gg 1$ ), the fractional fluctuations are of the order of unity:  $\langle N^2 \rangle / \langle N \rangle^2 \approx 1$ , so that the actual fluctuations can be enormous. It has been said that “bosons travel in flocks.”

7. *Optional review problems (more will soon be posted separately):*

- K&K 7.2
- K&K 7.10

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[1] Immanuel Bloch, Jean Dalibard, and Sylvain Nascimbene, “Quantum simulations with ultracold quantum gases,” [Nature Physics](#) **8**, 267–276 (2012).