

Ay 20 #2 - The Solar System: planets, moons, and minor planets

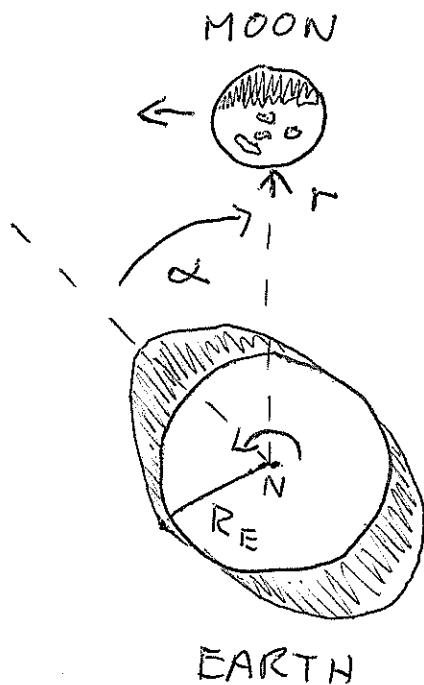
* Gravity (Newtonian)

$$\underline{F} = \frac{G m_1 m_2}{|\underline{r}|^3} \underline{r}, \quad U = - \frac{G m_1 m_2}{|\underline{r}|}$$

From this, and the standard laws of classical mechanics, most astrophysical dynamics can be solved. And more!

Consider two examples...

Tidal locking and spin-downs



The "tidal" force gradient across the Earth excites a deformation in the Earth, which leads the lunar orbit due to the Earth's rotation. The asymmetry causes the Earth to spin down.

A rough calculation:

We need to estimate the Torque

$$\tau = \frac{dL}{dt} = I_E \frac{d\Omega}{dt} \quad (I_E \approx \frac{2}{5} M_E R_E^2)$$

First, $\tau \approx \frac{2 R_E}{?} F_T \sin \alpha$ ↖ "force due to tides"
not "tidal force".

$$F_T \approx \frac{dF_G}{dr} \cdot \frac{2 R_E}{?} \cos \alpha, \quad \text{where}$$

$$\frac{dF_G}{dr} = - \frac{2 G M m_T}{r^3} = - \frac{2 G M}{r^3} \pi R_E^2 h \rho$$

h is the mean tidal height. Finally,

$$2\pi \frac{d\alpha}{dt} = - \frac{2 R_E \sin \alpha}{\frac{2}{5} M_E R_E^2} \cdot 2 R_E \cos \alpha \cdot \frac{2 G M}{r^3} \pi R_E^2 h \rho.$$

$$\frac{d\alpha}{dt} = - \frac{5 h \rho G M R_E^2 \sin 2\alpha}{M_E r^3}$$

Constants : $h \sim 10^2 \text{ cm}$, $\rho \sim 1 \text{ g cm}^{-3}$,
 $G = 6.7 \times 10^{-8} \text{ cm}^3 \text{ g}^{-1} \text{ s}^{-2}$, $M = 7 \times 10^{25} \text{ g}$,
 $R_E = 6.4 \times 10^8 \text{ cm}$, $\sin 2\alpha \sim 1/13$,
 $M_E = 6 \times 10^{27} \text{ g}$, $r \sim 4 \times 10^{10} \text{ cm}$.

Then, $\frac{d\gamma}{dt} \approx -2 \times 10^{-22} \text{ Hz s}^{-1}$

Now, because $\frac{dP}{dt} = -P^2 \frac{d\gamma}{dt} \approx 1.5 \times 10^{-12} \text{ s s}^{-1}$

or 1 s every $\sim 20,000$ years.

Note the (several) first-order approximations!

Gravitational radiation

The LIGO detectors are sensitive to
 "strain amplitudes", or distortions in the
 space-like part of the metric, of $\sim 10^{-21}$
 at 10^2 Hz .

Consider units where $G = c = 1$.

The gravitational wave (GW) strain is easily approximated by

$$h \sim \frac{M}{D}$$

The luminosity is then $L \sim h^2$.

In "real" terms,

$$h \sim \frac{G}{c^2} \frac{M}{D}, \quad L \sim \frac{c^3}{16\pi G} h^2 \cdot 4\pi D^2$$

Where do the pre-factors come from?

Binary neutron star: $M = 2 \times 2.8 M_{\odot}$.
 \uparrow
solar mass

Gives $h_{LIGO} = 10^{-21}$,

$$D = 4 \times 10^{26} \text{ cm} = 130 \text{ Mpc.}$$

$L = 2 \times 10^{49} \text{ erg s}^{-1}$, or $2 \times 10^{28} \text{ g}$ converted
to energy per s (3 Earths!).

* Gravitational orbits.

In polar coordinates, the equations of motions is

$$r \left(\frac{d\theta}{dt} \right)^2 - \frac{d^2 r}{dt^2} = \frac{GM_1}{r^2} \quad (\theta \text{ is "mean anomaly"})$$

At any time, the total energy of the system is

$$E = \frac{1}{2} \mu v^2 - \frac{GM\mu}{r}, \quad M = M_1 + M_2, \\ \mu = \frac{M_1 M_2}{M} \quad (\text{reduced mass})$$

Either relation, together with the assumptions of constant angular momentum $L = \mu r^2 \frac{d\theta}{dt}$, can be solved to derive Kepler's laws.

(exercise for student - can use standard, Lagrangian or Hamiltonian methods)

Solution:

$$r = \frac{h^2}{GM} \frac{1}{1 + e \cos \theta}$$

$$h = r^2 \frac{d\theta}{dt} = \sqrt{GM_1 p}$$

semi latus rectum

is specific angular momentum.

Circular orbits

$e = 0$ - super special.

$$r = \frac{h^2}{GM} = \frac{r^2 v_{\text{orb}}^2}{GM} \Rightarrow v_{\text{orb}} = \sqrt{\frac{GM}{r}}.$$

Elliptical orbits

$0 < e < 1$

If a is the semi-major axis,

$$2a = \frac{p}{1-e} + \frac{p}{1+e} \quad (\text{apoapsis} + \text{periapsis})$$

$$\Rightarrow p = \frac{h^2}{GM} = a(1-e^2).$$

With this substitution, $r = \frac{a(1-e^2)}{1+e \cos \theta},$

$$p^2 = \frac{4\pi^2}{GM} a^3, \quad v_{\text{orb}} = \sqrt{GM \left(\frac{2}{r} - \frac{1}{a} \right)}.$$

Parabolic orbits

$e = 1$ - super special

Consider energy @ periapsis ($\theta = 0$).

$$\text{First, } r_p = \frac{h^2}{2GM} \Rightarrow r_p = \frac{2GM}{v_p^2}.$$

Then, $E_p = \frac{1}{2} \mu v^2 - \frac{GM\mu}{r_p} = 0 !$

Must be 0 everywhere! Then,

$$v = \sqrt{\frac{2GM}{r}} \quad (\text{escape velocity}).$$

Hyperbolic orbits $e > 1$

Excess velocity @ $r = \infty$ $v_\infty = \sqrt{\frac{GM}{a}}$.

Properties @ periaipse (velocity, distance) can then be derived from orbit equations and conservation of energy (exercise for student).