Painlevé Systems and Eigenvalue Distributions

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Outline

- Painlevé equations
- Hamiltonian formulation
 - Auxiliary Hamiltonians
 - \circ σ -form of the equations
- Relationship to eigenvalue distributions
 - Numerical results
- Discussion

Painlevé Equations (Historical Context)

• Late 19th Century: Fuchs/Poincaré classify first order differential equations:

$$P(y', y, t) = 0$$

Reduced equations to Weierstrauss or Ricatti equations:

$$\left(\frac{dy}{dt}\right)^2 = 4y^3 - g_2 - g_3$$
 $\frac{dy}{dt} = a(t)y^2 + b(t)y + c(t)$

Painlevé wanted to do something similar with 2nd order ODEs

Painlevé Equations (Historical Context)

Painlevé attempted to classify equations for y" rational in y, y', and t:

$$y'' = R(y, y', t)$$

Reduced to either first order equations, linear equations, or 6 equations:

$$PI y'' = 6y^2 + t$$

PII $y'' = 2y^3 + ty + \alpha$,

PIII $y'' = \frac{1}{u} (y')^2 - \frac{1}{t} y' + \gamma y^3 + \frac{1}{t} (\alpha y^2 + \beta) + \frac{\delta}{u}$,

PIV $y'' = \frac{1}{2u} (y')^2 + \frac{3}{2} y^3 + 4ty^2 + 2(t^2 - \alpha) y + \frac{\beta}{u}$,

 $PV y'' = \left(\frac{1}{2u} + \frac{1}{u-1}\right) (y')^2 - \frac{1}{t}y' + \frac{(y-1)^2}{t^2} \left(\alpha y + \frac{\beta}{u}\right) + \frac{\gamma y}{t} + \frac{\delta y(y+1)}{u-1},$

PVI $y'' = \frac{1}{2} \left(\frac{1}{y} + \frac{1}{y-1} + \frac{1}{y-t} \right) (y')^2 - \left(\frac{1}{t} + \frac{1}{t-1} + \frac{1}{y-t} \right) y'$

 $+\frac{y(y-1)(y-t)}{t^2(t-1)^2}\left(\alpha+\frac{\beta t}{u^2}+\frac{\gamma(t-1)}{(u-1)^2}+\frac{\delta t(t-1)}{(u-t)^2}\right).$

σ-form of the equations

• (Auxiliary) Hamiltonians satisfy highly nonlinear differential equations

$$\sigma PII \quad (\sigma_{II}'')^2 + 4\sigma_{II}' \left((\sigma_{II}')^2 - t\sigma_{II}' + \sigma_{II} \right) - a^2 = 0,$$

$$\sigma PIII' \quad (t\sigma''_{III'})^2 - v_1 v_2 (\sigma'_{III'})^2 + \sigma'_{III'} (4\sigma'_{III'} - 1) (\sigma_{III'} - t\sigma'_{III'}) - \frac{1}{4^3} (v_1 - v_2)^2 = 0,$$

$$\sigma \text{PIV} \quad (\sigma_{IV}'')^2 - 4 (t\sigma_{IV}' - \sigma_{IV})^2 + 4\sigma_{IV}' (\sigma_{IV}' + 2\alpha_1) (\sigma_{IV}' - 2\alpha_2) = 0,$$

$$\sigma \text{PV} \quad (t\sigma_{V}'')^2 - \left(\sigma_V - t\sigma_V' + 2 (\sigma_V')^2 + (\nu_0 + \nu_1 + \nu_2 + \nu_3) \sigma_V'\right)^2$$

$$+4 \left(\nu_{0}+\sigma'_{V}\right) \left(\nu_{1}+\sigma'_{V}\right) \left(\nu_{2}+\sigma'_{V}\right) \left(\nu_{3}+\sigma'_{V}\right)=0,$$

$$\sigma \text{PVI} \quad \sigma'_{VI} \left(t(1-t)\sigma''_{VI}\right)^{2}+\left(\sigma'_{VI} \left(2\sigma_{VI}-(2t-1)\sigma'_{VI}\right)+v_{1}v_{2}v_{3}v_{4}\right)^{2}=\prod_{k=1}^{4} \left(\sigma'_{VI}+v_{k}^{2}\right)$$

Hamiltonian Formulation

 Present a Hamiltonian H such that eliminating the momentum from the Hamilton equations of motion gives the desired Painlevé equation

$$H = H(p, q, t; \vec{v})$$
 $q' = \frac{\partial H}{\partial p} \qquad p' = -\frac{\partial H}{\partial q}$

$$H_{II} = \frac{1}{2}p^{2} - 2q^{3} - tq,$$

$$H_{III} = -\frac{1}{2}\left(2q^{2} - p + t\right)p - \frac{v_{1} - v_{2}}{2}q,$$

$$tH_{III'} = q^{2}p^{2} - \left(q^{2} + v_{1}q - t\right)p + \frac{1}{2}\left(v_{1} + v_{2}\right)q,$$

$$H_{IV} = (2p - q - 2t)pq - 2\left(v_{1} - v_{2}\right)p + \left(v_{3} - v_{2}\right)q,$$

$$tH_{V} = q(q - 1)^{2}p^{2} - \left\{\left(v_{1} - v_{2}\right)\left(q - 1\right)^{2} - 2\left(v_{1} + v_{2}\right)q(q - 1) + tq\right\}p + \left(v_{3} - v_{2}\right)\left(v_{4} - v_{2}\right)\left(q - 1\right),$$

$$t(t - 1)H_{VI} = q(q - 1)\left(q - t\right)p^{2} - \left(\left(v_{3} + v_{4}\right)\left(q - 1\right)\left(q - t\right) + \left(v_{3} - v_{4}\right)q(q - t) - \left(v_{1} + v_{2}\right)q(q - 1)\right)p + \left(v_{3} - v_{1}\right)\left(v_{3} - v_{2}\right)\left(q - t\right),$$

Relationships between Painlevé equation parameters and Hamiltonian parameters

PII
$$v_1 + v_2 = 0$$
, $\alpha = v_1 - \frac{1}{2}$,
PIII' $\alpha = -4v_2$, $\beta = 4(v_1 + 1)$, $\gamma = 4$, $\delta = -4$,
PIV $v_1 + v_2 + v_3 = 0$, $\alpha = 1 + 2v_3 - v_1 - v_2$, $\beta = -2\alpha_1^2$,
PV $v_1 + v_2 + v_3 + v_4 = 0$, $\alpha = \frac{1}{2}(v_3 - v_4)^2$, $\beta = -\frac{1}{2}(v_1 - v_2)^2$, $\gamma = 2v_1 + 2v_2 - 1$, $\delta = -\frac{1}{2}$,
PVI $\alpha = \frac{1}{2}(v_1 - v_2)^2$, $\beta = -\frac{1}{2}(v_3 + v_4)^2$, $\gamma = \frac{1}{2}(v_3 - v_4)^2$, $\delta = \frac{1}{2}(1 - (1 - v_1 - v_2)^2)$.

Relations with Auxiliary Hamiltonians and σ functions

$$\begin{split} h_{III}(t) &= H_{III}, \\ h_{III'}(t) &= t H_{III'} + \frac{1}{4}v_1^2 - \frac{1}{2}t, \\ h_{IV}(t) &= H_{IV} - 2v_2t, \\ h_{V}(t) &= t H_{V} + (v_3 - v_2) \left(v_4 - v_2\right) - v_2t - 2v_2^2, \\ h_{VI}(t) &= t (t-1) H_{VI} + e_2 \left[-v_1, -v_2, v_3\right] t, -\frac{1}{2} e_2 \left[-v_1, -v_2, v_3, v_4\right], \quad e_p \left[a_1, \dots, a_s\right] \coloneqq \sum_{1 < j_1 < \dots < j_p < s} a_{j_1} a_{j_2} \cdots a_{j_p} \\ \sigma_{II}(t) &= -2^{1/3} h_{II} \left(-2^{1/3}t\right) \Big|_{\left(v_1, v_2\right) = (a, -a)}, \\ \sigma_{III'}(t) &= -h_{III'}(t/4) + \frac{t}{8} + \frac{v_1 v_2}{4}, \\ \sigma_{IV}(t) &= \left(h_{IV}(t) + 2v_2 t\right) \Big|_{\left(1 + v_3 - v_1, v_1 - v_2, v_2 - v_3\right) = \left(\alpha_0, \alpha_1, \alpha_2\right)}, \\ \sigma_{V}(t) &= h_{V}(t) + v_2 t + 2v_2^2, \quad \nu_{j-1} = v_j - v_2 \quad (j = 1, \dots, 4), \\ \sigma_{VI}(t) &= h_{VI}(t) \end{split}$$

Relationships to Eigenvalue Distributions

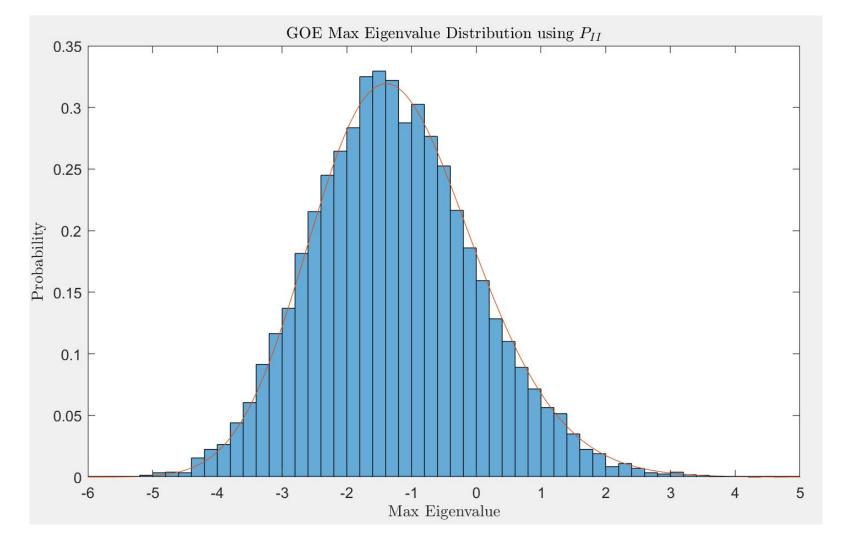
Painlevé II: q(x)

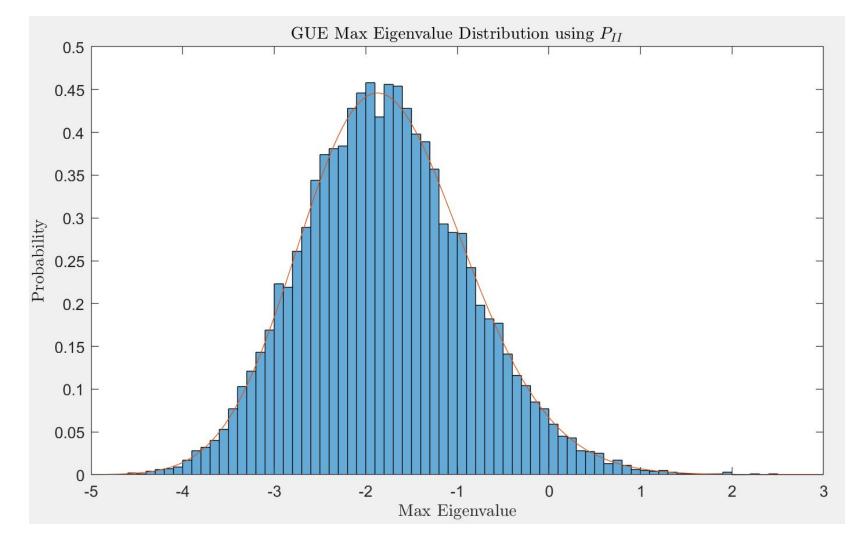
Using boundary condition that $q(x) \sim Ai(x)$ as $x \to \infty$

- Key part of finding Tracy-Widom PDF (for GUEs)
- Can also be used to find limiting distribution of the maximum GOE and GSE eigenvalues

Painlevé III

Hard-edge scaling limit for Laguerre and Jacobi ensembles





Relationships to Eigenvalue Distributions

Painlevé IV (σ-form)

$$\sigma \propto x^{2N-2} \exp(-x^2)$$

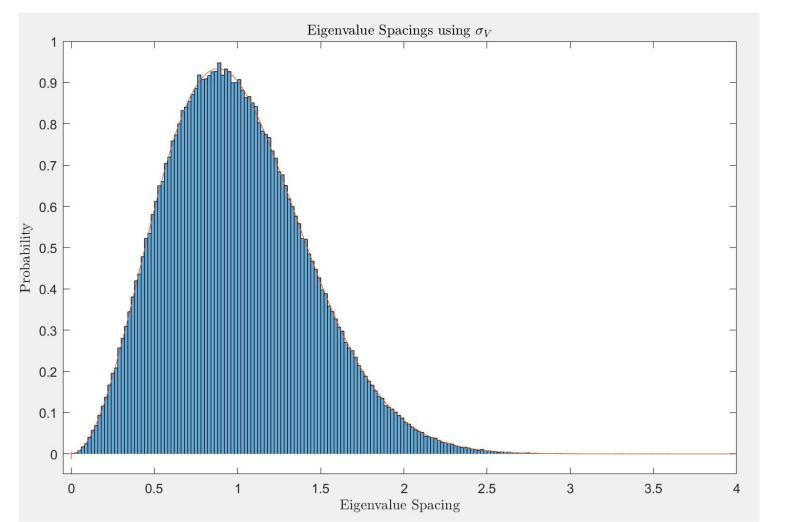
Largest eigenvalue for GUE (finite N)

Painlevé V (σ-form)

- Bulk scaling limit for GUE
- Smallest eigenvalue for Laguerre ensemble

Painlevé VI (σ-form)

Density for eigenvalues in Jacobi ensemble



Relationships to Eigenvalue Distributions

Painlevé VI (σ-form)

Density for eigenvalues in Jacobi-weighted random matrices

Discussion

We now more fully understand issues with stability

- Asymptotic boundary conditions that vanish cause sensitivity to different choices for the boundary
- Highly nonlinear equations (especially the σ -forms) can be very sensitive to these