Discretizing and Numerically Solving the Wave Equation: documentation for WaveEquationOptimized.py, PhotonGravityWavePerpendicularCase.py, and PhotonGravityWave4D.py

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For a function f(x) whose discretization is represented as a sequence f_n , we define its numerical derivative with respect to x, $nD_x(f_n)$, as the average of the forward difference and backward difference.

$$nD_x(f_n) = \frac{1}{2} \left(\frac{f_n - f_{n-1}}{\Delta x} + \frac{f_{n+1} - f_n}{\Delta x} \right) = \frac{f_{n+1} - f_{n-1}}{2\Delta x}$$
(1)

For differentiable functions, this approximates the true derivative of the function to an arbitrarily high accuracy as the step size $\Delta x \to 0$. As explained in DiscretizedLaplaceEquation.pdf, this logically produces the one-dimensional numerical second derivative $nD_x^2(f_n)$ as well as the numerical Laplacian $nD_{x,y}^2(f_{m,n})$ for our evenly spaced grid $(\Delta x = \Delta y \equiv \Delta)$.

$$nD_x^2(f_n) = nD(nD(f_n)) = \frac{f_{n+1} - 2f_n + f_{n-1}}{\Delta^2}$$
 (2)

$$nD_{x,y}^{2}(f_{m,n}) = \frac{f_{m,n+1} + f_{m,n-1} + f_{m+1,n} + f_{m-1,n} - 4f_{m,n}}{\Delta^{2}}$$
(3)

The Wave Equation for displacement f,

$$\frac{\partial^2 f}{\partial t^2} = c^2 \nabla^2 f,\tag{4}$$

can be numerically approximated using our discretizations as

$$nD_t^2(f_{m,n,t}) = c^2 nD_{x,y}^2(f_{m,n,t}).$$
(5)

Substituting Eq. 2 and Eq. 3 into Eq. 5 gives us

$$\frac{f_{m,n,t+1} - 2f_{m,n,t} + f_{m,n,t-1}}{(\Delta t)^2} = \frac{c^2}{\Delta^2} (f_{m,n+1,t} + f_{m,n-1,t} + f_{m+1,n,t} + f_{m-1,n,t} - 4f_{m,n,t})$$
(6)

Rearranging and subtracting one time step from every term gives us the desired iterative formula that solves the wave equation for given boundary conditions

$$f_{m,n,t} = K(f_{m,n+1,t-1} + f_{m,n-1,t-1} + f_{m+1,n,t-1} + f_{m-1,n,t-1} - 4f_{m,n,t-1}) + 2f_{m,n,t-1} - f_{m,n,t-2}$$
 (7)

where

$$K = c^2 \frac{(\Delta t)^2}{\Delta^2} \tag{8}$$