Tricomi Problem for the Lavrentiev-Bitsadze Equation with a Semi-Strip in the Elliptic Part

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1 Problem Statement

The Tricomi problem for the Lavrentiev-Bitsadze equation is considered

$$(sgn(y))\frac{\partial^2 u}{\partial x^2}(x,y) + \frac{\partial^2 u}{\partial y^2}(x,y) = 0$$
 (1)

in the domain $D = D^+ \cup D^-$, where $D^+ = \{(x,y): \ 0 < x < \pi, \ 0 < y < +\infty\}$, $D^- = \{(x,y): \ -y < x < y + \pi, \ -\pi/2 < y < 1\}$ in the class of functions $u(x,y) \in C^2(D^+) \cap C^2(D^-) \cap C(\overline{D^+ \cup D^-})$ with boundary conditions

$$u(0,y) = 0, \ u(\pi,y) = 0, \ 0 < y < +\infty,$$
 (2)

$$u(x, -x) = f(x), \ 0 \le x \le \pi/2, \ f(0) = 0,$$
 (3)

$$u(x,y) \rightrightarrows 0, \ y \to +\infty$$
 (4)

and the Frankl matching condition

$$\frac{1}{k}\frac{\partial u}{\partial y}(x,+0) = \frac{\partial u}{\partial y}(x,-0), \ 0 < x < \pi, \tag{5}$$

where $k \in (-\infty, +\infty), k \neq 0$.

2 Main Results

Theorem 1. The solution to the problem (1) - (5) is unique.

Proof. Suppose there exist two solutions $u_1(x,y), u_2(x,y)$ to the problem (1)-(5). Then $u(x,y) = u_1(x,y) - u_2(x,y)$ is a solution to the problem (1)-(5) with the function $f(x) \equiv 0$. In this case, u(x,y) = F(x+y) - F(0).

From this, it follows that the equality $\frac{\partial u}{\partial y} - \frac{\partial u}{\partial x} = 0$ holds for all points x and y in the domain of hyperbolicity. Using the matching condition (5), we get

$$\frac{1}{k}\frac{\partial u}{\partial y} - \frac{\partial u}{\partial x}|_{y=0+0} = 0. \tag{6}$$

As a result, we obtain the problem of finding a harmonic function u(x, y) in the domain D^+ with boundary conditions (2), (4), (6).

By the Zaremba-Giraud principle and equality (6), the extremum cannot be attained on the interval $\{(x,y): 0 < x < \pi, y = 0\}$. On the closed lateral sides and at infinity, the extremum cannot be attained due to conditions (2) and (4). The theorem is proved.

It is known that the general solution in D^- of equation (1) has the form

$$u(x,y) = F(x+y) + f(\frac{x-y}{2}) - F(0).$$
(7)

Differentiate equality (7):

$$\frac{\partial u}{\partial y}(x,y) - \frac{\partial u}{\partial x}(x,y)|_{y=0+0} = -f'\left(\frac{x}{2}\right), \ 0 < x < \pi.$$

Using the matching condition (5), we arrive at the equality

$$\frac{1}{k}\frac{\partial u}{\partial y}(x,0+0) - \frac{\partial u}{\partial x}(x,0+0) = -f'\left(\frac{x}{2}\right), \ 0 < x < \pi.$$

Then we obtain in the domain D^+ the auxiliary problem for the Laplace operator

$$\frac{\partial^2 u}{\partial x^2}(x,y) + \frac{\partial^2 u}{\partial y^2}(x,y) = 0 \tag{8}$$

with boundary conditions

$$u(0,y) = 0, \ u(\pi,y) = 0, \ 0 < y < +\infty,$$
 (9)

$$\frac{1}{k}\frac{\partial u}{\partial y}(x,0+0) - \frac{\partial u}{\partial x}(x,0+0) = -f'\left(\frac{x}{2}\right),\tag{10}$$

$$u(x,y) \rightrightarrows 0, \ y \to +\infty$$
 (11)

Theorem 2. Let |k| < 1, $k \neq 0$, $f(x) \in C[0, \pi/2] \cap C^2(0, \pi/2)$, $f'(x) \in L_2(0, \pi/2)$. Then the solution to the problem (8)-(11) exists and can be represented as a series

$$u(x,y) = \sum_{n=1}^{\infty} A_n e^{-ny} \sin nx,$$
(12)

where the condition (10) is understood in the integral sense

$$\int_{0}^{\pi} \left[\frac{1}{k} \frac{\partial u}{\partial y}(x, y) - \frac{\partial u}{\partial x}(x, y) + f'\left(\frac{x}{2}\right) \right]^{2} dx \to 0, \ y \to 0 + 0,$$

and the coefficients A_n are determined from the equality

$$\sum_{n=1}^{\infty} nA_n \sin\left[nx + \arctan k\right] = \frac{k}{\sqrt{1+k^2}} f'\left(\frac{x}{2}\right)$$
 (13)

Proof.

The system $\{\sin[nx + \arctan k]\}_{n=1}^{\infty}$ forms a Riesz basis in $L_2(0, \pi)$ for $k \in (-1, 1)$ according to the main result of [2]. Therefore, the two-sided Bessel inequality holds

$$C_1 \|f'\|_{L_2(0,\pi)}^2 \le \sum_{n=1}^{\infty} n^2 A_n^2 \le C_2 \|f'\|_{L_2(0,\pi)}^2, \ 0 < C_1 < C_2,$$

where the constants C_1, C_2 do not depend on f'. Therefore, the series $\sum_{n=1}^{\infty} |A_n|$ converges and the series (12) converges uniformly. The function (12) satisfies equation (8) with boundary conditions (9) by construction. Condition (11) is satisfied since $\sum_{n=1}^{\infty} e^{-ny} = \frac{e^{-y}}{1-e^{-y}} = \frac{1}{e^y-1}$. Let's check the fulfillment of condition (10). Let

$$M(x) = \frac{1}{k} \frac{\partial u}{\partial y} - \frac{\partial u}{\partial x} + f'\left(\frac{x}{2}\right)$$

$$M(x) = -\frac{1}{k} \sum_{n=1}^{\infty} n A_n e^{-ny} \sin nx - \sum_{n=1}^{\infty} n A_n e^{-ny} \cos nx + f'\left(\frac{x}{2}\right) =$$
$$= -\sum_{n=1}^{\infty} n A_n e^{-ny} \left[\frac{1}{k} \sin nx + \cos nx\right] + f'\left(\frac{x}{2}\right) =$$

$$\begin{split} &=-\frac{\sqrt{1+k^2}}{k}\sum_{n=1}^{\infty}nA_ne^{-ny}\left[\frac{1}{\sqrt{1+k^2}}\sin nx+\frac{k}{\sqrt{1+k^2}}\cos nx\right]+f'\left(\frac{x}{2}\right)=\\ &=-\frac{\sqrt{1+k^2}}{k}\sum_{n=1}^{\infty}nA_ne^{-ny}\sin\left[nx+\operatorname{arctg}k\right]+f'\left(\frac{x}{2}\right)=\\ &=\frac{\sqrt{1+k^2}}{k}\sum_{n=1}^{\infty}nA_n\left(1-e^{-ny}\right)\sin\left[nx+\operatorname{arctg}k\right].\\ &\operatorname{Let's\ show\ that\ }\lim_{y\to 0+0}I(y)=0. \end{split}$$

$$I(y) = \int_{0}^{\pi} M(x)^{2} dx \le I_{1}(y) + I_{2}(y),$$

$$I_{1}(y) = \frac{2\sqrt{1+k^{2}}}{k} \int_{0}^{\pi} \left[\sum_{n=1}^{m} nA_{n} \sin\left[nx + \operatorname{arctg} k\right] \left(1 - e^{-ny}\right) \right]^{2} dx$$

$$I_{2}(y) = \frac{2\sqrt{1+k^{2}}}{k} \int_{0}^{\pi} \left[\sum_{n=m+1}^{+\infty} nA_{n} \sin\left[nx + \operatorname{arctg} k\right] \left(1 - e^{-ny}\right) \right]^{2} dx$$

Fix an arbitrary positive ε , then

$$I_2(y) \le C_3 \sum_{n=m+1}^{\infty} n^2 A_n^2 (1 - e^{-ny})^2 \le C_3 \sum_{n=m+1}^{\infty} n^2 A_n^2 < \frac{\varepsilon}{2}.$$

This is true if m is sufficiently large, since the series is convergent.

$$I_1(y) \le C_4 \sum_{n=1}^m n^2 A_n^2 (1 - e^{-ny})^2 < \frac{\varepsilon}{2}$$

This is true for $0 < y < \delta$, if δ is sufficiently small. The theorem is proved.

Theorem 3. Let k > 0, then the solution to the problem (8) - (11) is unique

Proof. Let's prove the uniqueness of the solution to this problem. Let u(x,y) be a solution to the homogeneous problem. Introduce the notation $C_{\varepsilon} = (0,\varepsilon), C_R = (0,R), D_R = (\pi,R), D_{\varepsilon} = (\pi,\varepsilon).$ It is the rectangle $C_{\varepsilon}C_RD_RD_{\varepsilon}$. The following relations hold:

$$0 = \iint_{\Pi_{R\varepsilon}} (R - y)(u_{xx} + u_{yy}) dx dy.$$

Note that

$$(R-y)(u_{xx} + u_{yy})u = ((R-y)u_xu)_x + ((R-y)u_yu)_y - (R-y)(u_x^2 + u_y^2) + u_yu =$$

$$= (R-y)(u_{xx}u + u_x^2) + (-u_y + (R-y)u_{yy}u + (R-y)u_y^2) - (R-y)(u_x^2 + u_y^2) + u_yu$$

Substitute this expression into the integral:

$$I = \iint_{\Pi_{R\varepsilon}} ((R - y) u_x u)_x dx dy + \iint_{\Pi_{R\varepsilon}} ((R - y) u_y u)_y dx dy - \iint_{\Pi_{R\varepsilon}} (R - y) (u_x^2 + u_y^2) + \iint_{\Pi_{R\varepsilon}} u_y u dx dy.$$

Simplify these integrals:

$$\iint\limits_{\prod_{R\varepsilon}}\left(\left(R-y\right)u_{x}u\right)_{x}dxdy=\int\limits_{\left[\varepsilon,R\right]}\left[\left(R-y\right)u_{x}u\right]|_{0}^{\pi}dy=\int\limits_{\left[\varepsilon,R\right]}\left[\left(R-y\right)u_{x}(\pi,y)u(\pi,y)-\left(R-y\right)u_{x}(0,y)u(0,y)\right]dy=0$$

since both integrands are zero due to condition (9)

$$\begin{split} \iint\limits_{\prod_{R\varepsilon}} \left(\left(R - y \right) u_y u \right)_y dx dy &= \int\limits_{[0,\pi]} \left[\left(R - y \right) u_y u \right] \big|_{\varepsilon}^R dx = \int\limits_{[0,\pi]} \left[0 - \left(R - \varepsilon \right) u_y (x,\varepsilon) u (x,\varepsilon) \right] dx = - \int\limits_{C_\varepsilon D_\varepsilon} \left(R - \varepsilon \right) u_y u dx \\ \iint\limits_{\prod_{R\varepsilon}} u_y u dx dy &= \iint\limits_{\prod_{R\varepsilon}} \left(\frac{u^2}{2} \right)_y' dx dy = \int\limits_{[0,\pi]} \left[\frac{u^2 (x,R)}{2} - \frac{u^2 (x,\varepsilon)}{2} \right] dx \end{split}$$

In the end, we get

$$I = -\iint_{\prod_{R\varepsilon}} \left(R - y\right) \left(u_x^2 + u_y^2\right) dx dy - \int_{C_\varepsilon D_\varepsilon} \left(R - \varepsilon\right) u_y u dx - \int_{C_\varepsilon D_\varepsilon} \frac{u^2}{2} dx + \int_{C_R D_R} \frac{u^2}{2} dx$$

Add and subtract $\int_{C_{\varepsilon}D_{\varepsilon}}k\left(R-\varepsilon\right)u_{x}udx$, then

$$I = -\iint_{\Pi_{R\varepsilon}} (R - y) \left(u_x^2 + u_y^2 \right) dx dy - \int_{C_{\varepsilon} D_{\varepsilon}} (R - \varepsilon) \left(u_y - k u_x \right) u dx - \int_{C_{\varepsilon} D_{\varepsilon}} (R - \varepsilon) k u_x u dx - \int_{C_{\varepsilon} D_{\varepsilon}} \frac{u^2}{2} dx + \int_{C_R D_R} \frac{u^2}{2} dx.$$

From this, it follows that

$$\iint\limits_{\prod_{R\varepsilon}}\left(R-y\right)\left(u_{x}^{2}+u_{y}^{2}\right)dxdy+\frac{1}{2}\int\limits_{C_{\varepsilon}D_{\varepsilon}}u^{2}dx+k\frac{R-\varepsilon}{2}u^{2}(\pi,\varepsilon)=\int\limits_{C_{\varepsilon}D_{\varepsilon}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}u^{2}dx\leq \frac{1}{2}\int\limits_{C_{\varepsilon}D_{\varepsilon}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C_{R}D_{R}}\left(R-\varepsilon\right)\left(u_{y}-ku_{x}\right)udx+\frac{1}{2}\int\limits_{C$$

$$\leq \{\text{Cauchy-Bunyakovsky Inequality}\} \leq (R - \varepsilon) \left[\int\limits_{C_{\varepsilon}D_{\varepsilon}} \left(u_y - ku_x \right)^2 dx \right]^{\frac{1}{2}} \left[\int\limits_{C_{\varepsilon}D_{\varepsilon}} u^2 dx \right]^{\frac{1}{2}} + \frac{1}{2} \int\limits_{C_RD_R} u^2 dx = M$$

Consider the following inequality: $(2ar - b)^2 \ge 0 \Rightarrow ra^2r^2 - 4abr + b^2 \ge 0 \Rightarrow ab \le ra^2 + b/(4r)$. Let $a = \left[(R - \varepsilon) \int_{C_\varepsilon D_\varepsilon} (u_y - ku_x)^2 dx \right]^{\frac{1}{2}}, b = \left[(R - \varepsilon) \int_{C_\varepsilon D_\varepsilon} u^2 dx \right]^{\frac{1}{2}}, r = R - \varepsilon$, then

$$M \le (R - \varepsilon)^2 \int_{C_{\varepsilon}D_{\varepsilon}} (u_y - ku_x)^2 dx + \frac{1}{4} \int_{C_{\varepsilon}D_{\varepsilon}} u^2 dx + \frac{1}{2} \int_{C_RD_R} u^2 dx.$$

Regroup

$$\iint\limits_{\prod_{R\varepsilon}} \left(R-y\right) \left(u_x^2+u_y^2\right) dx dy + \frac{1}{4} \int\limits_{C_\varepsilon D_\varepsilon} u^2 dx + k \frac{R-\varepsilon}{2} u^2(\pi,\varepsilon) \\ \leq \left(R-\varepsilon\right)^2 \int\limits_{C_\varepsilon D_\varepsilon} \left(u_y - k u_x\right)^2 dx + \frac{1}{2} \int\limits_{C_R D_R} u^2 dx.$$

Due to (10), the equality holds

$$\lim_{\varepsilon \to 0+0} \int_{C_{\varepsilon}D_{\varepsilon}} (u_y - ku_x)^2 dx = 0,$$

from which it follows that

$$\lim_{\varepsilon \to 0+0} \iint_{\prod_{R\varepsilon}} (R-y) \left(u_x^2 + u_y^2 \right) dx dy + \frac{1}{4} \int_0^{\pi} u^2(x,0) dx + k \frac{R}{2} u^2(\pi,0) \le \frac{1}{2} \int_{C_R D_R} u^2 dx.$$

Now let $R \to \infty$, then $\int_{C_R D_R} u^2 dx \to 0$, and in the left part all terms are non-negative, hence $u(x,y) \equiv 0$ in \overline{D} . The theorem is proved.

Theorem 4. Let |k| < 1, $k \neq 0$ and u(x,y) be a solution to the problem (8) - (11), then u_x, u_y can be represented as

$$u_y(x,y) = -\frac{2k}{\pi\sqrt{1+k^2}} \operatorname{Im} \left(\frac{1-e^{iz}}{1+e^{iz}}\right)^{\gamma/\pi} e^{iz} \int_0^{\pi} \frac{1}{\left(\operatorname{tg} t/2\right)^{\gamma/\pi}} \frac{\sin t}{\left(1-e^{i(z+t)}\right) \left(1-e^{i(z-t)}\right)} f'(\frac{t}{2}) dt,$$

$$u_x(x,y) = \frac{2k}{\pi\sqrt{1+k^2}} \operatorname{Re} \left(\frac{1-e^{iz}}{1+e^{iz}}\right)^{\gamma/\pi} e^{iz} \int_0^{\pi} \frac{1}{(\operatorname{tg} t/2)^{\gamma/\pi}} \frac{\sin t}{\left(1-e^{i(z+t)}\right)\left(1-e^{i(z-t)}\right)} f'(\frac{t}{2}) dt,$$

where $\gamma = 2 \operatorname{arctg} k$, z = x + iy.

Proof. Consider equality (12). The system of sines $\{\sin[nx + \operatorname{arctg} k]\}_{n=1}^{\infty}$ forms a basis in $L_2(0, \pi)$ for $k \in (-1, 1)$. Therefore, for the coefficients nA_n , the following representation holds [2]:

$$nA_n = \int_{0}^{\pi} h_n(t)F(t)dt,$$

where

$$F(x) = \frac{k}{\sqrt{1+k^2}} f'(\frac{x}{2}), h_n(t) = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\operatorname{tg} t/2)^{\gamma/\pi}} \sum_{k=1}^{n} \sin kt B_{n-k}, B_l = \sum_{m=0}^{l} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-l}^{m} (-1)^{l-m}, C_l^{n} = \frac{l(l-1)\dots(l-n+1)}{n!}.$$

Let u(x,y) be a solution to the problem (8)-(11), then

$$u(x,y) = \sum_{n=1}^{\infty} A_n e^{-ny} \sin[nx]$$

and accordingly

$$u_y(x,y) = -\sum_{n=1}^{\infty} nA_n e^{-ny} \sin\left[nx\right] =$$

$$\sum_{n=1}^{\infty} \int_{-\pi}^{\pi} \Gamma(x) e^{-ny} \cdot \left[-\frac{1}{2}x\right]$$

$$= -\sum_{n=1}^{\infty} \int_{0}^{h} F(t)h_n(t)e^{-ny}\sin[nx]dt,$$

or

$$u_y(x,y) = -Im \sum_{n=1}^{\infty} \int_{0}^{\pi} F(t)h_n(t)e^{-ny}e^{inx}dt =$$
$$= -Im \sum_{n=1}^{\infty} \int_{0}^{\pi} F(t)h_n(t)e^{inz}dt$$

Swap the order of integration and summation

$$u_y(x,y) = -\operatorname{Im} \int_0^{\pi} F(t) \sum_{n=1}^{\infty} h_n(t) e^{inz} dt.$$

Introduce new notation:

$$I(t,z) = \sum_{n=1}^{\infty} h_n(t)e^{inz}$$

$$I(t,z) = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\operatorname{tg} t/2)^{\gamma/\pi}} \sum_{n=1}^{\infty} \sum_{k=1}^{n} \sin kt B_{n-k} e^{inz} = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\operatorname{tg} t/2)^{\gamma/\pi}} \sum_{k=1}^{\infty} \sin kt \sum_{n=k}^{\infty} e^{inz} B_{n-k} e^{inz} = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\operatorname{tg} t/2)^{\gamma/\pi}} \sum_{k=1}^{\infty} \sin kt \sum_{n=k}^{\infty} e^{inz} B_{n-k} e^{inz}$$

and new index m = n - k

$$I(t,z) = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\lg t/2)^{\gamma/\pi}} \sum_{k=1}^{\infty} \sin kt \sum_{m=0}^{\infty} e^{i(m+k)z} B_m = \frac{2}{\pi} \frac{(2\cos t/2)^{\beta}}{(\lg t/2)^{\gamma/\pi}} \sum_{k=1}^{\infty} e^{ikz} \sin kt \sum_{m=0}^{\infty} e^{imz} B_m,$$

$$\sum_{k=1}^{\infty} e^{ikz} \sin kt = \frac{e^{iz} \sin t}{(1 - e^{i(z+t)}) (1 - e^{i(z-t)})}$$

Consider the second series:

$$\sum_{l=0}^{\infty} e^{ilz} B_l = \sum_{l=0}^{\infty} e^{ilz} \sum_{m=0}^{l} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k = l - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{-\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = |k - m| = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi}^{l-m} C_{\gamma/\pi-\beta}^m (-1)^{l-m} = \sum_{m=0}^{\infty} \sum_{l=m}^{\infty} e^{ilz} C_{\gamma/\pi}^{l-m} C_{\gamma$$

$$\sum_{m=0}^{\infty} \sum_{k=0}^{\infty} e^{i(m+k)z} C_{\gamma/\pi}^k C_{-\gamma/\pi-\beta}^m (-1)^k = \sum_{m=0}^{\infty} e^{imz} C_{-\gamma/\pi-\beta}^m \sum_{k=0}^{\infty} C_{\gamma/\pi}^k (-1)^k e^{ikz} = (1+e^{iz})^{-\gamma/\pi-\beta} (1-e^{iz})^{\gamma/\pi} (1-e^{iz})^{\gamma/$$

since in our case $\beta = 0$, $\gamma = 2 \arctan k$. Finally, we obtain the formula

$$\begin{split} u_y(x,y) &= -\text{Im} \int\limits_0^T F(t)I(t,z)dt = \\ &= -\text{Im} \int\limits_0^\pi \frac{2}{\pi} \frac{(2\cos t/2)^\beta}{(\operatorname{tg} t/2)^{\gamma/\pi}} \frac{e^{iz}\sin t}{\left(1 - e^{i(z+t)}\right)\left(1 - e^{i(z-t)}\right)} (1 + e^{iz})^{-\gamma/\pi - \beta} (1 - e^{iz})^{\gamma/\pi} F(t)dt = \\ &= -\frac{2}{\pi} \text{Im} \left(\frac{1 - e^{iz}}{1 + e^{iz}}\right)^{\gamma/\pi} e^{iz} \int\limits_0^\pi \frac{1}{(\operatorname{tg} t/2)^{\gamma/\pi}} \frac{\sin t}{\left(1 - e^{i(z+t)}\right)\left(1 - e^{i(z-t)}\right)} F(t)dt = \\ &= -\frac{2k}{\pi\sqrt{1 + k^2}} \text{Im} \left(\frac{1 - e^{iz}}{1 + e^{iz}}\right)^{\gamma/\pi} e^{iz} \int\limits_0^\pi \frac{1}{(\operatorname{tg} t/2)^{\gamma/\pi}} \frac{\sin t}{\left(1 - e^{i(z+t)}\right)\left(1 - e^{i(z-t)}\right)} f'(\frac{t}{2})dt \end{split}$$

i.e., the representation:

$$u_y(x,y) = -\frac{2k}{\pi\sqrt{1+k^2}} \operatorname{Im} \left(\frac{1-e^{iz}}{1+e^{iz}}\right)^{\gamma/\pi} e^{iz} \int_0^{\pi} \frac{1}{\left(\operatorname{tg} t/2\right)^{\gamma/\pi}} \frac{\sin t}{\left(1-e^{i(z+t)}\right)\left(1-e^{i(z-t)}\right)} f'(\frac{t}{2}) dt.$$

By similar reasoning, we obtain the representation

$$u_x(x,y) = \frac{2k}{\pi\sqrt{1+k^2}} \operatorname{Re} \left(\frac{1-e^{iz}}{1+e^{iz}}\right)^{\gamma/\pi} e^{iz} \int_0^{\pi} \frac{1}{(\operatorname{tg} t/2)^{\gamma/\pi}} \frac{\sin t}{\left(1-e^{i(z+t)}\right)\left(1-e^{i(z-t)}\right)} f'(\frac{t}{2}) dt.$$

The theorem is proved.

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