

Mathematical notes for the Lefschetz–thimble plotter

1 Problem class

We consider one-dimensional complex integrals of the form

$$I(\kappa) = \int_{\Gamma} \exp(-\kappa W(z)) dz, \quad (1)$$

where $W: \mathbb{C} \rightarrow \mathbb{C}$ is holomorphic, $\kappa \in \mathbb{C}$ is a (typically large) parameter, and Γ is an integration contour in the complex plane.

The steepest-descent (Picard–Lefschetz) geometry is controlled by:

- the critical points (“saddles”) z_σ of W , defined by $W'(z_\sigma) = 0$;
- the phase of κ (a non-zero argument rotates the constant-phase condition).

In the code in this repository we take $\kappa = 1$ by default. This loses no geometric information when $\kappa \in \mathbb{R}_{>0}$ (it merely rescales the flow time), and it is straightforward to reinsert κ if required.

2 Saddles and non-degeneracy

A saddle is any $z_\sigma \in \mathbb{C}$ satisfying

$$W'(z_\sigma) = 0. \quad (2)$$

For the standard local picture we assume the saddle is non-degenerate,

$$W''(z_\sigma) \neq 0. \quad (3)$$

In particular, near such a saddle one has the quadratic approximation

$$W(z) = W(z_\sigma) + \frac{1}{2}W''(z_\sigma)(z - z_\sigma)^2 + \mathcal{O}((z - z_\sigma)^3). \quad (4)$$

3 Holomorphic gradient flow

Write $z(u) = x(u) + iy(u)$. The holomorphic gradient flow associated with the weight $\exp(-\kappa W)$ is

$$\frac{dz}{du} = \pm \overline{\kappa W'(z)}. \quad (5)$$

The choice of sign corresponds to increasing or decreasing $\text{Re}(\kappa W)$ along the flow.

A one-line computation explains why this flow is natural. Along any solution of (5),

$$\frac{d}{du} (\kappa W(z(u))) = \kappa W'(z(u)) \frac{dz}{du} = \pm \kappa W'(z) \overline{\kappa W'(z)} = \pm |\kappa W'(z)|^2 \in \mathbb{R}_{\geq 0}. \quad (6)$$

Consequently,

$$\frac{d}{du} \text{Im}(\kappa W(z(u))) = 0, \quad \frac{d}{du} \text{Re}(\kappa W(z(u))) = \pm |\kappa W'(z(u))|^2. \quad (7)$$

Thus each flow line is a constant-phase curve ($\text{constant } \text{Im}(\kappa W)$), and $\text{Re}(\kappa W)$ varies monotonically along it.

4 Thimbles and dual thimbles (one variable)

Fix the *upward* sign convention in (5),

$$\frac{dz}{du} = +\overline{\kappa W'(z)}. \quad (8)$$

Then:

- the *thimble* \mathcal{J}_σ is the unstable manifold of (8) at z_σ , i.e. the set of points whose backward flow approaches the saddle:

$$\mathcal{J}_\sigma = \{z(0) : z(u) \text{ solves (8)} \text{ and } \lim_{u \rightarrow -\infty} z(u) = z_\sigma\}; \quad (9)$$

- the *dual thimble* \mathcal{K}_σ is the stable manifold of (8) at z_σ , i.e. points whose forward flow approaches the saddle:

$$\mathcal{K}_\sigma = \{z(0) : z(u) \text{ solves (8)} \text{ and } \lim_{u \rightarrow +\infty} z(u) = z_\sigma\}. \quad (10)$$

In one complex dimension, \mathcal{J}_σ and \mathcal{K}_σ are real one-dimensional curves through z_σ . In the full Picard–Lefschetz story, intersection numbers between Γ and \mathcal{K}_σ determine which thimbles contribute to the original integral.

5 Local direction and the two “arms”

Because the saddle is a fixed point of the flow, one cannot initialise the ODE exactly at z_σ ; the numerical solver would simply remain there. Instead, we start a small distance away, in the correct tangent directions.

Let $\lambda = \kappa W''(z_\sigma)$. Near z_σ we have

$$\kappa W(z) = \kappa W(z_\sigma) + \frac{1}{2}\lambda(z - z_\sigma)^2 + \dots. \quad (11)$$

Imposing the constant-phase condition $\text{Im}(\kappa W(z)) = \text{Im}(\kappa W(z_\sigma))$ at quadratic order gives

$$\text{Im}(\lambda(z - z_\sigma)^2) = 0, \quad \text{i.e.} \quad \lambda(z - z_\sigma)^2 \in \mathbb{R}. \quad (12)$$

Writing $z - z_\sigma = r e^{i\theta}$ and $\lambda = |\lambda|e^{i\phi}$, this becomes $\sin(\phi + 2\theta) = 0$, hence

$$\theta = -\frac{\phi}{2} + \frac{k\pi}{2}, \quad k \in \mathbb{Z}. \quad (13)$$

A convenient unit vector in the thimble direction is therefore

$$v = \exp\left(-\frac{i}{2} \text{Arg}(\lambda)\right). \quad (14)$$

The dual direction is rotated by $\pi/2$. In the implementation this is encoded by an additional phase shift.

Finally, each manifold has two local branches through the saddle, so we take two opposite initial points

$$z_{0,\pm} = z_\sigma \pm \varepsilon v, \quad (15)$$

and integrate the flow from both to draw the full curve.

6 Speed normalisation for high-degree potentials

For high-degree polynomials, the raw flow speed can grow very rapidly because $|W'(z)|$ grows like a high power of $|z|$. For plotting, it is often useful to bound the speed while keeping the same geometric trajectories.

If $f(z)$ is a vector field and $g(z) > 0$ is any positive real function, then the ODEs

$$\frac{dz}{du} = f(z), \quad \frac{dz}{ds} = g(z) f(z) \quad (16)$$

have the *same integral curves* in the complex plane (only the parametrisation changes). This is the mechanism behind the “speed-normalised” option in the code:

$$\frac{dz}{du} = \pm \frac{\overline{W'(z)}}{1 + |\overline{W'(z)}|}. \quad (17)$$

The denominator is positive and bounded away from zero, so the geometry is unchanged, but the numerical solver is far less likely to run away.

7 Mapping to the repository code

With $\kappa = 1$, the code implements the upward flow

$$\frac{dz}{du} = \pm \overline{W'(z)} \quad (18)$$

by evolving $x(u), y(u)$ with

$$x'(u) = \pm \operatorname{Re} \overline{W'(x + iy)}, \quad y'(u) = \pm \operatorname{Im} \overline{W'(x + iy)}. \quad (19)$$

The local direction v is chosen using (14), and the two initial points are $z_{0,\pm} = z_\sigma \pm \varepsilon v$. The plotting window `box` is used as a stop condition (we halt integration once the curve leaves the visible region).

Practical note. For a general complex phase of κ , replace $W'(z)$ by $\kappa W'(z)$ everywhere in the flow and in the definition $\lambda = \kappa W''(z_\sigma)$. This rotates the thimble/dual directions exactly as expected.