

Resurgence of the Airy function

Aaron Fenyes

August 17, 2022

1 The Laplace transform

1.1 Analytic version

1.1.1 Regularity and decay properties

Let ζ be the standard coordinate on \mathbb{C} , and let $z: T^*\mathbb{C} \rightarrow \mathbb{C}$ be the fiberwise-linear map $d\zeta \mapsto 1$ (see “The Geometry of the Laplace Transform” in **draft2**). The Laplace transform in ζ turns a function φ on \mathbb{C} into a function $\mathcal{L}_\zeta \varphi$ on $T^*\mathbb{C}$, defined on the cotangent space of $\zeta = \alpha$ by the integral

$$\mathcal{L}_\zeta \varphi|_{\zeta=\alpha} = \int_{\Gamma_{\zeta,\alpha}} e^{-z\zeta} \varphi d\zeta,$$

where $\Gamma_{\zeta,\alpha}$ is the rightward ray starting at $\zeta = \alpha$ (compare **draft2**). We’ll use the shorthand $\mathcal{L}_{\zeta,\alpha} f := \mathcal{L}_\zeta f|_{\zeta=\alpha}$ throughout this document [but maybe we should get rid of it].

Let’s say a function f is in $O_{\zeta,\alpha\leftarrow}(g)$ or $O_{\zeta,\alpha\rightarrow}(g)$, respectively, if $|f| \lesssim g$ on some neighborhood of the starting point or infinite end of $\Gamma_{\zeta,\alpha}$. A function is *subexponential* along $\Gamma_{\zeta,\alpha}$ if it’s in $O_{\zeta,\alpha\rightarrow}(e^{c\zeta})$ for all $c > 0$. Let $\mathcal{E}_{\zeta,\alpha}$ be the space of functions which are subexponential on $\Gamma_{\zeta,\alpha}$, integrable at the starting point, and locally integrable throughout. If f is in $\mathcal{E}_{\zeta,\alpha}$, then $\mathcal{L}_{\zeta,\alpha} f$ is well-defined and holomorphic for $\operatorname{Re}(z) > 0$ on the part of $T^*\mathbb{C}$ that lies over $\Gamma_{\zeta,\alpha}$ [1, §5.6].

The asymptotics of f at the starting point of $\Gamma_{\zeta,\alpha}$ control the asymptotics of $\mathcal{L}_{\zeta,\alpha} f$ at the infinite end of $\Gamma_{z,0}$. Once we see how this works for $\alpha = 0$, Section 1.2 will do the rest. Let $F = \mathcal{L}_{\zeta,0} f$. Equation 1.8 of ?? shows¹ that

$$f \in O_{\zeta,0\leftarrow}(1) \implies F \in O_{z,0\rightarrow}\left(\frac{1}{z}\right).$$

More generally, for $\tau > -1$ [prove or cite],

$$f \in O_{\zeta,0\leftarrow}(\zeta^\tau) \implies F \in O_{z,0\rightarrow}\left(\frac{1}{z^{1+\tau}}\right).$$

Exact power law asymptotics relate similarly [prove or cite]:

$$f \sim \zeta^\tau \text{ at the start of } \Gamma_{\zeta,0} \implies F \sim \frac{\Gamma(1+\tau)}{z^{1+\tau}} \text{ at the end of } \Gamma_{z,0}.$$

¹The argument cited still works in our generality. For holomorphic f , one can also use Equation 1.5 of *Borel-Laplace Transform and Asymptotic Theory* (Sternin & Shatalov).

[The big- O asymptotics dictionary is interesting, but we might not need it. Consider dropping.]

1.1.2 Action on integral operators

When $\varphi \in \mathcal{E}_\zeta$, you can use differentiation under the integral to show that [2, Theorem 1.34]

$$\mathcal{L}_\zeta(\zeta^n \varphi) = \left(-\frac{\partial}{\partial z}\right)^n \mathcal{L}_\zeta \varphi. \quad (1)$$

for all integers $n \geq 0$. You can also use a 2d integration argument, akin to the one in [2, Theorem 2.39], to show that $\partial_{\zeta,a}^{-\lambda} \varphi \in \mathcal{E}_\zeta$ and

$$\mathcal{L}_{\zeta,a} \partial_{\zeta,a}^{-\lambda} \varphi = z^{-\lambda} \mathcal{L}_{\zeta,a} \varphi$$

for all $\lambda \in (0, \infty)$.

1.2 Change of translation chart

Define a new coordinate ζ_α on \mathbb{C} so that $\zeta = \alpha + \zeta_\alpha$. From the calculation

$$\begin{aligned} \mathcal{L}_{\zeta,a} \varphi &= \int_{\Gamma_{\zeta,\alpha}} e^{-z\zeta} \varphi d\zeta \\ &= \int_{\Gamma_{\zeta_\alpha,0}} e^{-z(\alpha+\zeta_\alpha)} \varphi d\zeta_\alpha \\ &= e^{-\alpha z} \int_{\Gamma_{\zeta_\alpha,0}} e^{-z\zeta_\alpha} \varphi d\zeta_\alpha \\ &= e^{-\alpha z} \mathcal{L}_{\zeta_\alpha,0} \varphi, \end{aligned}$$

we learn that

$$\mathcal{L}_{\zeta_\alpha,0} \varphi = e^{\alpha z} \mathcal{L}_{\zeta,a} \varphi.$$

1.3 Rescaling of translation structure

Let's rescale the translation structure of \mathbb{C} , expanding displacements by a factor of $\mu \in (0, \infty)$. The coordinate $\xi = \mu\zeta$ is a chart for the new translation structure. The corresponding frequency coordinate $x: T^*\mathbb{C} \rightarrow \mathbb{C}$ is given by $d\xi \mapsto 1$, so $x = \mu^{-1}z$. From the calculation

$$\begin{aligned} \mathcal{L}_{\xi,0} \varphi &= \int_{\Gamma_{\xi,0}} e^{-x\xi} \varphi d\xi \\ &= \int_{\Gamma_{\zeta,0}} e^{-z\zeta} \varphi \mu d\zeta \\ &= \mu \mathcal{L}_{\zeta,0} \varphi \end{aligned}$$

we learn that

$$\mathcal{L}_{\xi,0} \varphi = \mu \mathcal{L}_{\zeta,0} \varphi.$$

Note that $\mathcal{L}_{\xi,0}$ is defined in the new translation structure on \mathbb{C} , while $\mathcal{L}_{\zeta,0}$ is defined in the old translation structure. We can still compare them, because they both turn complex-valued functions on \mathbb{C} into holomorphic functions on $T^*\mathbb{C}$.

2 Integro-differential equations

2.1 Existence of solutions

2.1.1 Algebraic integral operators

Take a simply connected open set $\Omega \subset \mathbb{C}$ that touches but doesn't contain $\zeta = 0$. Let $\mathcal{H}L^\infty(\Omega)$ be the space of bounded holomorphic functions on Ω with the supremum norm $\|\cdot\|_\infty$. For any $\sigma \in \mathbb{R}$, multiplying by $\zeta^{-\sigma}$ maps $\mathcal{H}L^\infty(\Omega)$ isomorphically onto another space of holomorphic functions on Ω . We'll call this space $\mathcal{H}L^{\infty,\sigma}(\Omega)$ and give it the norm $\|f\|_{\infty,\sigma} = \|\zeta^\sigma f\|_\infty$, so that

$$\begin{aligned} \mathcal{H}L^\infty(\Omega) &\rightarrow \mathcal{H}L^{\infty,\sigma}(\Omega) \\ \phi &\mapsto \zeta^{-\sigma} \phi \end{aligned}$$

is an isometry. More generally,

$$\begin{aligned} \mathcal{H}L^{\infty,\rho}(\Omega) &\rightarrow \mathcal{H}L^{\infty,\rho+\delta}(\Omega) \\ f &\mapsto \zeta^{-\delta} f \end{aligned}$$

is an isometry for all $\rho \in \mathbb{R}$ and $\delta \in [0, \infty)$. This reduces to the previous statement when $\rho = 0$. For each $\delta \in [0, \infty)$, the functions in $\mathcal{H}L^{\infty,\rho}(\Omega)$ belong to $\mathcal{H}L^{\infty,\rho+\delta}(\Omega)$ too, and the inclusion map $\mathcal{H}L^{\infty,\rho}(\Omega) \hookrightarrow \mathcal{H}L^{\infty,\rho+\delta}(\Omega)$ has norm $\|\zeta^\delta\|_\infty$. Conceptually, $\|\zeta^\delta\|_\infty$ measures of the size of Ω , so let's write it as M^δ with $M = \|\zeta\|_\infty$.

Since $\mathcal{H}L^\infty(\Omega)$ is a Banach algebra, the function space $\mathcal{H}L^{\infty,\infty}(\Omega) := \bigcup_{\sigma \in \mathbb{R}} \mathcal{H}L^{\infty,\sigma}(\Omega)$ is a graded algebra, with a different norm on each grade. For each $\rho, \delta \in \mathbb{R}$, multiplication by a function $m \in \mathcal{H}L^{\infty,\delta}(\Omega)$ gives a map $\mathcal{H}L^{\infty,\rho}(\Omega) \rightarrow \mathcal{H}L^{\infty,\rho+\delta}(\Omega)$ with norm $\|m\|_{\infty,\delta}$.

We'll study integral operators $\mathcal{G}: \mathcal{H}L^{\infty,\rho}(\Omega) \rightarrow \mathcal{H}L^{\infty,\sigma}(\Omega)$ of the form

$$[\mathcal{G}f](a) = \int_{\zeta=0}^a g(a, \cdot) f d\zeta,$$

where the kernel g is an algebraic function over \mathbb{C}^2 which can be singular on Δ , the diagonal.² To avoid ambiguity, we fix a branch of g to use at the start of the integration path. The domain of g is a covering of \mathbb{C}^2 which can be branched over Δ . Continuing g around Δ changes its phase by a root of unity, leaving its absolute value the same **[check]**. That makes $|g|$ a well-defined function on $\mathbb{C}^2 \setminus \Delta$, which we can use to bound $\|\mathcal{G}\|$.

For each $a \in \mathbb{C}$, the expression $|g(a, \cdot) d\zeta|$ defines a *density* on $\Omega \setminus \{a\}$ —a norm on the tangent bundle which is compatible with the conformal structure. The square of a density is a Riemannian metric. Let $\ell_{g,\Omega}^{\sigma,\rho}(a)$ be the distance from $\zeta = 0$ to a with respect to the density $|\zeta(a)^\sigma g(a, \cdot) \zeta^{-\rho} d\zeta|$ on $\Omega \setminus \{a\}$. The bound

$$\begin{aligned} |[\zeta^\sigma \mathcal{G}f](a)| &\leq \left| \zeta(a)^\sigma \int_{\zeta=0}^a g(a, \cdot) f d\zeta \right| \\ &\leq \int_{\zeta=0}^a |\zeta^\sigma f| |\zeta(a)^\sigma g(a, \cdot) \zeta^{-\rho} d\zeta| \\ &\leq \|f\|_{\infty,\sigma} \int_{\zeta=0}^a |\zeta(a)^\sigma g(a, \cdot) \zeta^{-\rho} d\zeta| \end{aligned}$$

²Thanks to Alex Takeda for suggesting this.

holds for any integration path. Taking the infimum over all paths, we see that

$$|[\mathcal{G}f](a)| \leq \ell_{g,\Omega}^{\sigma,\rho}(a) \|f\|_\infty.$$

so $\|\mathcal{G}\| \leq \sup_{a \in \Omega} \ell_{g,\Omega}^{\sigma,\rho}(a)$. Crucially, we can always make $\|\mathcal{G}\|$ a contraction by restricting Ω .

2.1.2 The example of fractional integrals

Setting $g(a, a') = (\zeta(a) - \zeta(a'))^{-\lambda-1}$ with $\lambda \in (-\infty, 0)$, we get the fractional integral $\partial_{\zeta \text{ from } 0}^\lambda$. The shortest path from $\zeta = 0$ to a with respect to $|\zeta(a)^{\rho+\lambda} g(a, \cdot) \zeta^{-\rho} d\zeta|$ is the same as the shortest path with respect to $|d\zeta|$ [check]. It follows that

$$\begin{aligned} \ell_{g,\Omega}^{\sigma,\rho}(a) &= \int_0^{|\zeta(a)|} |\zeta(a)|^{\rho+\lambda} (|\zeta(a)| - r)^{-\lambda-1} r^{-\rho} dr \\ &= |\zeta(a)|^{\rho+\lambda} \int_0^1 (|\zeta(a)| - |\zeta(a)|t)^{-\lambda-1} (|\zeta(a)|t)^{-\rho} |\zeta(a)| dt \\ &= |\zeta(a)|^{\rho+\lambda-\lambda-1-\rho+1} \int_0^1 (1-t)^{-\lambda-1} t^{-\rho} dt \\ &= \int_0^1 (1-t)^{-\lambda-1} t^{-\rho} dt \\ &= B(-\lambda, 1-\rho). \end{aligned}$$

The beta function B can be written more explicitly as

$$B(-\lambda, 1-\rho) = \frac{\Gamma(-\lambda) \Gamma(1-\rho)}{\Gamma(1-\lambda-\rho)}.$$

Now we can see that for each $\lambda \in (-\infty, 0)$ and $\rho \in \mathbb{R}$, the fractional integral $\partial_{\zeta \text{ from } 0}^\lambda$ maps $\mathcal{H}L^{\infty,\rho}(\Omega)$ into $\mathcal{H}L^{\infty,\rho+\lambda}(\Omega)$, with norm $\|\partial_{\zeta \text{ from } 0}^\lambda\| \leq B(-\lambda, 1-\rho)$.

2.1.3 Fractional integral equations near a regular singular point

Angeliki: Maybe Kato–Rellich perturbation theory can give existence immediately. It might not give uniqueness, though.

Consider an integral operator \mathcal{J} of the form

$$p + \partial_{\zeta \text{ from } 0}^{-1} \circ q + \sum_{\lambda \in \Lambda} \partial_{\zeta \text{ from } 0}^\lambda \circ r_\lambda,$$

where:

- p is a function in $\mathcal{H}L^{\infty,-1}(\Omega)$ that extends holomorphically over $\zeta = 0$, and its derivative at $\zeta = 0$ is non-zero.
- q is a function in $\mathcal{H}L^\infty(\Omega)$ that extends holomorphically over $\zeta = 0$.
- r_λ are functions in $\mathcal{H}L^\infty(\Omega)$.
- Λ is a countable subset of $(-\infty, -1)$ whose supremum is less than -1 .

Our demand that p and q have convergent power series at $\zeta = 0$ can probably be relaxed; having convergent Novikov series, for example, should be enough. We could also probably replace $\partial_{\zeta \text{ from } 0}^{-1}$ with $\partial_{\zeta \text{ from } 0}^{-1+\delta} \circ \zeta^\delta$ for some $\delta \in [0, 1)$, or adjust the $\partial_{\zeta \text{ from } 0}^\lambda$ similarly.

We want to solve the equation $\mathcal{J}f = 0$. Let's look for a solution of the form $f = \zeta^{\tau-1} + \tilde{f}$ with $\tau \in (0, \infty)$ and $\tilde{f} \in \mathcal{H}L^{\infty, 1-\tau-\epsilon}(\Omega)$ for some $\epsilon \in (0, 1]$. When ϵ is small enough that $\Lambda \subset (-\infty, -1-\epsilon]$, we'll see that we can always find such a solution, as long as we're willing to shrink Ω . In fact, there's exactly one such solution. [\[Add convergence conditions for \$\partial^\lambda\$ terms.\]](#)

Let p'_0 and q_0 be the values of $\frac{\partial}{\partial \zeta} p$ and q , respectively, at $\zeta = 0$. We're assuming that p and q extend holomorphically over $\zeta = 0$, and the additional assumption that $p \in \mathcal{H}L^{\infty, -1}(\Omega)$ implies that p has a first-order zero at $\zeta = 0$.

Since p and q extend holomorphically over $\zeta = 0$, and p vanishes at $\zeta = 0$, we can write

$$\begin{aligned} p &= p'_0 \zeta + \tilde{p} \\ q &= q_0 + \tilde{q} \end{aligned}$$

with $\tilde{p} \in \mathcal{H}L^{\infty, -2}(\Omega)$ and $\tilde{q} \in \mathcal{H}L^{\infty, -1}(\Omega)$. Then we have

$$\mathcal{J} = p'_0 \zeta + q_0 \partial_{\zeta \text{ from } 0}^{-1} + \tilde{\mathcal{J}}$$

with

$$\tilde{\mathcal{J}} = \tilde{p} + \partial_{\zeta \text{ from } 0}^{-1} \circ \tilde{q} + \sum_{\lambda \in \Lambda} \partial_{\zeta \text{ from } 0}^\lambda \circ r_\lambda$$

For any $\tau \in (0, \infty)$,

$$\mathcal{J} \zeta^{\tau-1} = (p'_0 + q_0/\tau) \zeta^\tau + \tilde{\mathcal{J}} \zeta^{\tau-1}.$$

Setting $\tau = -q_0/p'_0$ makes the first term vanish, leaving

$$\mathcal{J} \zeta^{\tau-1} = \tilde{\mathcal{J}} \zeta^{\tau-1}.$$

Then the equation $\mathcal{J}f = 0$ becomes

$$\begin{aligned} 0 &= \tilde{\mathcal{J}} \zeta^{\tau-1} + \mathcal{J} \tilde{f} \\ 0 &= \tilde{\mathcal{J}} \zeta^{\tau-1} + \left[p'_0 \zeta + q_0 \partial_{\zeta \text{ from } 0}^{-1} + \tilde{\mathcal{J}} \right] \tilde{f} \\ -p'_0 \zeta \tilde{f} &= \tilde{\mathcal{J}} \zeta^{\tau-1} + \left[q_0 \partial_{\zeta \text{ from } 0}^{-1} + \tilde{\mathcal{J}} \right] \tilde{f} \\ \tilde{f} &= \left[-\frac{1}{p'_0} \zeta^{-1} \circ \tilde{\mathcal{J}} \right] \zeta^{\tau-1} + \left[\tau \zeta^{-1} \circ \partial_{\zeta \text{ from } 0}^{-1} - \frac{1}{p'_0} \zeta^{-1} \circ \tilde{\mathcal{J}} \right] \tilde{f}. \end{aligned} \tag{2}$$

$$\begin{array}{c}
\mathcal{HL}^{\infty,\rho}(\Omega) \xrightarrow[\|\tilde{p}\|_{\infty,-2}]{\tilde{p}} \mathcal{HL}^{\infty,\rho-2}(\Omega) \begin{array}{l} \nearrow^{\|\zeta\|_{\infty}} \mathcal{HL}^{\infty,\rho-1}(\Omega) \\ \searrow_{\|\zeta\|_{\infty}^{1-\epsilon}} \mathcal{HL}^{\infty,\rho-1-\epsilon}(\Omega) \end{array} \\
\\
\mathcal{HL}^{\infty,\rho}(\Omega) \xrightarrow[\|\tilde{q}\|_{\infty,-1}]{\tilde{q}} \mathcal{HL}^{\infty,\rho-1}(\Omega) \xrightarrow[B(1,2-\rho)=\frac{1}{2-\rho}]{\partial^{-1}} \mathcal{HL}^{\infty,\rho-2}(\Omega) \begin{array}{l} \nearrow^{\|\zeta\|_{\infty}} \mathcal{HL}^{\infty,\rho-1}(\Omega) \\ \searrow_{\|\zeta\|_{\infty}^2} \mathcal{HL}^{\infty,\rho}(\Omega) \end{array} \\
\\
\mathcal{HL}^{\infty,\rho}(\Omega) \xrightarrow[\|r_{\lambda}\|_{\infty}]{r_{\lambda}} \mathcal{HL}^{\infty,\rho}(\Omega) \xrightarrow[B(-\lambda,1-\rho)]{\partial^{\lambda}} \mathcal{HL}^{\infty,\rho+\lambda}(\Omega) \begin{array}{l} \nearrow^{\|\zeta\|_{\infty}^{-1-\lambda}} \mathcal{HL}^{\infty,\rho-1}(\Omega) \\ \searrow_{\|\zeta\|_{\infty}^{-1-\epsilon-\lambda}} \mathcal{HL}^{\infty,\rho-1-\epsilon}(\Omega) \end{array} \\
\\
\mathcal{HL}^{\infty,\rho}(\Omega) \xrightarrow[B(1,1-\rho)=\frac{1}{1-\rho}]{\partial^{-1}} \mathcal{HL}^{\infty,\rho-1}(\Omega) \xrightarrow[\|\zeta^{-1}\|_{\infty,1}=1]{\zeta^{-1}} \mathcal{HL}^{\infty,\rho}(\Omega)
\end{array}$$

From [our previous discussion], we can work out that

$$\begin{aligned}
& \tilde{\mathcal{J}}: \mathcal{HL}^{\infty,1-\tau}(\Omega) \rightarrow \mathcal{HL}^{\infty,-\tau-\epsilon}(\Omega) \text{ with} \\
& \|\tilde{\mathcal{J}}\| \leq \left(\|\tilde{p}\|_{\infty,-2} + \frac{1}{1+\tau} \|\tilde{q}\|_{\infty,-1} \right) M^{1-\epsilon} + \sum_{\lambda \in \Lambda} B(-\lambda, \tau) \|r_{\lambda}\|_{\infty} M^{-1-\epsilon-\lambda}
\end{aligned}$$

and

$$\begin{aligned}
& \tilde{\mathcal{J}}: \mathcal{HL}^{\infty,1-\tau-\epsilon}(\Omega) \rightarrow \mathcal{HL}^{\infty,-\tau-\epsilon}(\Omega) \text{ with} \\
& \|\tilde{\mathcal{J}}\| \leq \left(\|\tilde{p}\|_{\infty,-2} + \frac{1}{1+\tau+\epsilon} \|\tilde{q}\|_{\infty,-1} \right) M + \sum_{\lambda \in \Lambda} B(-\lambda, \tau + \epsilon) \|r_{\lambda}\|_{\infty} M^{-1-\lambda}.
\end{aligned}$$

Since $\|\zeta^{-1}\|_{\infty,1} = 1$, it follows that

$$\begin{aligned}
& \zeta^{-1} \circ \tilde{\mathcal{J}}: \mathcal{HL}^{\infty,1-\tau}(\Omega) \rightarrow \mathcal{HL}^{\infty,1-\tau-\epsilon}(\Omega) \text{ with} \\
& \|\zeta^{-1} \circ \tilde{\mathcal{J}}\| \leq \left(\|\tilde{p}\|_{\infty,-2} + \frac{1}{1+\tau} \|\tilde{q}\|_{\infty,-1} \right) M^{1-\epsilon} + \sum_{\lambda \in \Lambda} B(-\lambda, \tau) \|r_{\lambda}\|_{\infty} M^{-1-\epsilon-\lambda} \quad (3)
\end{aligned}$$

and

$$\begin{aligned}
& \zeta^{-1} \circ \tilde{\mathcal{J}} \circ \mathcal{HL}^{\infty,1-\tau-\epsilon}(\Omega) \text{ with} \\
& \|\zeta^{-1} \circ \tilde{\mathcal{J}}\| \leq \left(\|\tilde{p}\|_{\infty,-2} + \frac{1}{1+\tau+\epsilon} \|\tilde{q}\|_{\infty,-1} \right) M + \sum_{\lambda \in \Lambda} B(-\lambda, \tau + \epsilon) \|r_{\lambda}\|_{\infty} M^{-1-\lambda}. \quad (4)
\end{aligned}$$

We can also see that

$$\begin{aligned}
& \tau \zeta^{-1} \circ \partial_{\zeta \text{ from } 0}^{-1} \circ \mathcal{HL}^{\infty,1-\tau-\epsilon}(\Omega) \text{ with} \\
& \|\tau \zeta^{-1} \circ \partial_{\zeta \text{ from } 0}^{-1}\| = \frac{\tau}{\tau+\epsilon} < 1. \quad (5)
\end{aligned}$$

Now, let's return to equation 2, which tells us that $f = \zeta^{\tau-1} + \tilde{f}$ satisfies $\mathcal{J}f = 0$ when [and only when?] \tilde{f} is a fixed point of the affine map $\mathcal{A}(\cdot) + b$, where

$$\begin{aligned}\mathcal{A} &= \tau \zeta^{-1} \circ \partial_{\zeta}^{-1} \text{from } 0 - \frac{1}{p'_0} \zeta^{-1} \circ \tilde{\mathcal{J}} \\ b &= \left[-\frac{1}{p'_0} \zeta^{-1} \circ \tilde{\mathcal{J}} \right] \zeta^{\tau-1}\end{aligned}$$

Choosing $\epsilon \in (0, 1]$ so that $\Lambda \subset (-\infty, -1 - \epsilon]$ has given us the domain and codomain statements in bounds 3 and 4, which tell us that $\mathcal{A}(\cdot) + b$ sends $\mathcal{HL}^{\infty, 1-\tau-\epsilon}(\Omega)$ into itself. We'll show that when Ω is small enough, $\mathcal{A}(\cdot) + b$ contracts $\mathcal{HL}^{\infty, 1-\tau-\epsilon}(\Omega)$, and thus—by the contraction mapping theorem—has a unique fixed point.

An affine map is a contraction if and only if its linear part is a contraction. We know from bound 5 that $\tau \zeta^{-1} \circ \partial_{\zeta}^{-1} \text{from } 0$ contracts $\mathcal{HL}^{\infty, 1-\tau-\epsilon}(\Omega)$. Since the supremum of Λ is less than -1 , all the powers of $M = \|\zeta\|_{\infty}$ in bound 4 are positive. Thus, by shrinking Ω , we can make the norm of $\zeta^{-1} \circ \tilde{\mathcal{J}}$ on $\mathcal{HL}^{\infty, 1-\tau-\epsilon}(\Omega)$ as small as we want—small enough to make \mathcal{A} a contraction.

2.2 Order shifting

Consider holomorphic functions on a simply connected open set that touches but doesn't contain $\zeta = 0$. A function is *regular* at $\zeta = 0$ if it extends holomorphically over that point. We'll say a function is *slight* at $\zeta = 0$ if it can be written as [rephrase in terms of $\mathcal{HL}^{\infty, \rho}$]

$$\zeta^{\alpha_1} f_1 + \dots + \zeta^{\alpha_r} f_r + g \tag{6}$$

where f_1, \dots, f_r are regular, $\alpha_1, \dots, \alpha_r \in \mathbb{R} \setminus \mathbb{Z}$, and g, g', g'', \dots go to zero at $\zeta = 0$.

Locally integrable slight functions play a special role in Laplace transform methods for linear differential equations. This is because differential equations in the frequency domain arise most naturally from integral equations in the spatial domain, but we'd like to work with differential equations in the spatial domain too. In the space of locally integrable slight functions, differential and integral equations enjoy their simplest equivalence.

Proposition 1. *When ψ is slight and locally integrable at $\zeta = 0$,*

$$\left[\sum_{k=0}^n \left(\frac{\partial}{\partial \zeta} \right)^k \circ h_k + \sum_{k=1}^m \partial_{\zeta, 0}^{-k} \circ h_{-k} \right] \psi = 0,$$

if and only if

$$\left[\sum_{k=1}^n \left(\frac{\partial}{\partial \zeta} \right)^{k-1} \circ h_k + \sum_{k=0}^m \partial_{\zeta, 0}^{-k-1} \circ h_{-k} \right] \psi = 0,$$

where h_n, \dots, h_{-m} are regular functions.

To prove this result, we'll need a little background.

Each slight function has a unique *normal form*: an expression of the form (6) where $\alpha_1, \dots, \alpha_r$ are distinct modulo \mathbb{Z} and f_1, \dots, f_r are non-vanishing at $\zeta = 0$. The *order* of a slight function is the smallest power α_k in its normal form. If there are no $\zeta^{\alpha} f$ terms, we say the order is ∞ . A slight function vanishes at $\zeta = 0$ if and only if its order is positive, and it's locally integrable at $\zeta = 0$ if and only if its order is greater than -1 .

Multiplication by a regular function and differentiation with respect to ζ both preserve the space of slight functions. Integration from $\zeta = 0$ preserves the space of locally integrable slight functions. Thus, in general, integro-differential operators with regular coefficients send locally integrable slight functions to slight functions. Each basic operation has a simple effect on a slight function's order. Multiplication by $\zeta^n f$, where f is regular and non-vanishing at $\zeta = 0$, adds n to the order. Differentiation and integration add -1 and 1 , respectively.

Proof of Proposition 1. The reverse implication holds without any special condition on ψ , because $\frac{\partial}{\partial \zeta} \partial_{\zeta,0}^{-1}$ acts as the identity on all differentiable functions.

To prove the forward implication, rewrite the first equation in the statement as

$$\frac{\partial}{\partial \zeta} \left[\sum_{k=1}^n \left(\frac{\partial}{\partial \zeta} \right)^{k-1} \circ h_k \right] \psi = - \left[h_0 + \sum_{k=1}^m \partial_{\zeta,0}^{-k} \circ h_{-k} \right] \psi. \quad (7)$$

The function

$$\phi = \left[\sum_{k=1}^n \left(\frac{\partial}{\partial \zeta} \right)^{k-1} \circ h_k \right] \psi$$

is slight. Looking at the right-hand side of equation 7, we can see that ϕ' is slight and locally integrable, so its order is greater than -1 . Hence, ϕ has positive order, which means it vanishes at $\zeta = 0$.

Integrating both sides of equation 7, we get

$$\partial_{\zeta,0}^{-1} \frac{\partial}{\partial \zeta} \phi = - \left[\partial_{\zeta,0}^{-1} \circ h_0 + \sum_{k=1}^m \partial_{\zeta,0}^{-k-1} \circ h_{-k} \right] \psi.$$

Since $\partial_{\zeta,0}^{-1} \frac{\partial}{\partial \zeta}$ acts as the identity on functions that vanish at $\zeta = 0$, this simplifies to

$$\phi = - \left[\sum_{k=0}^m \partial_{\zeta,0}^{-k-1} \circ h_{-k} \right] \psi,$$

which rearranges to the second equation in the statement. \square

3 The Airy equation

3.1 Basics

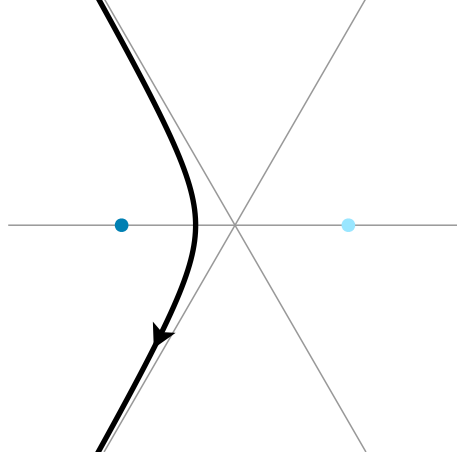
The Airy equation is

$$\left[\left(\frac{\partial}{\partial y} \right)^2 - y \right] \psi = 0. \quad (8)$$

One solution is given by the Airy function,

$$\text{Ai}(y) = \frac{i}{2\pi} \int_{\Gamma} \exp \left(-\frac{1}{3}t^3 + yt \right) dt,$$

where Γ is a path that comes from ∞ at 120° and goes to ∞ at -120° .



The contour Γ in the u plane.

With the substitution $t = 2uy^{1/2}$, we can rewrite the Airy integral as

$$\text{Ai}(y) = y^{1/2} \frac{i}{\pi} \int_{y^{-1/2}\Gamma} \exp \left[-\frac{2}{3}y^{3/2} (4u^3 - 3u) \right] du.$$

We've rescaled the contour by a factor of two, but it still approaches ∞ in the desired way. Note that $4u^3 - 3u$ is the third Chebyshev polynomial.

3.2 Rewriting as a modified Bessel equation

We can distill the most interesting part of the Airy function by writing

$$\text{Ai}(y) = \frac{1}{\pi\sqrt{3}} y^{1/2} K\left(\frac{2}{3}y^{3/2}\right),$$

where

$$K(z) = i\sqrt{3} \int_{z^{-1/3}\Gamma} \exp[-z(4u^3 - 3u)] du. \quad (9)$$

Saying that Ai satisfies the Airy equation is equivalent to saying that K satisfies the modified Bessel equation

$$\left[z^2 \left(\frac{\partial}{\partial z} \right)^2 + z \frac{\partial}{\partial z} - \left[\left(\frac{1}{3} \right)^2 + z^2 \right] \right] \varphi = 0. \quad (10)$$

In fact, K is the modified Bessel function $K_{1/3}$ [3, equation 9.6.1].

The method we'll demonstrate in Section 3.4 works for any differential equation

$$\left[P\left(\frac{\partial}{\partial z}\right) + z^{-1}Q\left(\frac{\partial}{\partial z}\right) + z^{-2}R(z^{-1}) \right] \varphi = 0,$$

where P is a polynomial, Q is a polynomial of one degree lower, and R is an entire function. Let's put equation 10 in that form:

$$\left[\left[\left(\frac{\partial}{\partial z} \right)^2 - 1 \right] + z^{-1} \frac{\partial}{\partial z} - \left(\frac{1}{3} \right)^2 z^{-2} \right] \varphi = 0. \quad (11)$$

3.3 Asymptotic analysis

Equation 10 has a regular singularity at $z = 0$ and an irregular singularity at $z = \infty$. From the general theory of such equations, we know that the space of trans-series solutions has a basis of trans-monomials

$$\{e^{-\alpha z} z^{-\tau_\alpha} \tilde{W}_\alpha \mid \alpha^2 - 1 = 0\}$$

where the \tilde{W}_α are formal power series in z^{-1} with no constant term. From equations 10.40.2 and 10.17.1 of [3], we learn that $K \sim \left(\frac{\pi}{2}\right)^{1/2} e^{-z} z^{1/2} \tilde{W}_1$, with

$$\tilde{W}_1 = z^{-1} - \frac{(\frac{1}{6})_1(\frac{5}{6})_1}{2^1 \cdot 1!} z^{-2} + \frac{(\frac{1}{6})_2(\frac{5}{6})_2}{2^2 \cdot 2!} z^{-3} - \frac{(\frac{1}{6})_3(\frac{5}{6})_3}{2^3 \cdot 3!} z^{-4} + \dots \quad (12)$$

The holomorphic analysis in Section 3.4 will give us holomorphic solutions

$$\{e^{-\alpha z} z^{-\tau_\alpha} W_\alpha \mid \alpha^2 - 1 = 0\},$$

which seem analogous to the trans-monomials above. Borel summation makes the analogy precise. We'll see in Section ?? that each $z^{\tau_\alpha} W_\alpha$ is proportional to the Borel sum of $z^{\tau_\alpha} \tilde{W}_\alpha$.

3.4 Going to the spatial domain

3.4.1 The big idea

We're going to look for functions v_α whose Laplace transforms $\mathcal{L}_{\zeta,\alpha} v_\alpha$ satisfy equation 11. We'll succeed when $\alpha^2 - 1 = 0$, and we'll see that K is a scalar multiple of $\mathcal{L}_{\zeta,1} v_1$.

We can see from Section 1.1.2 that $\mathcal{L}_{\zeta,\alpha} v$ satisfies the differential equation 11 if and only if v satisfies the integral equation

$$\left[[\zeta^2 - 1] - \partial_{\zeta,\alpha}^{-1} \circ \zeta - \left(\frac{1}{3}\right)^2 \partial_{\zeta,\alpha}^{-2} \right] v = 0. \quad (13)$$

It's tempting to differentiate both sides of this equation until we get

$$\left[\left(\frac{\partial}{\partial \zeta}\right)^2 \circ [\zeta^2 - 1] - \frac{\partial}{\partial \zeta} \circ \zeta - \left(\frac{1}{3}\right)^2 \right] v = 0, \quad (14)$$

which is easier to solve. Unfortunately, a solution of equation 14 won't satisfy equation 13 in general. However, as we learned in Section 2.2, a solution of equation 14 *will* satisfy equation 13 if it's slight and locally integrable at $\zeta = \alpha$.

This is great news, because equation 14 has a regular singularity at each root of $\zeta^2 - 1$, and the Frobenius method often gives a slight solution at each regular singular point. We can see the regular singularities by moving the derivatives to the right:

$$\left[(\zeta^2 - 1) \left(\frac{\partial}{\partial \zeta}\right)^2 + 3\zeta \frac{\partial}{\partial \zeta} + \left[1 - \left(\frac{1}{3}\right)^2\right] \right] v = 0.$$

In Sections 3.4.2–3.4.3, we'll see this approach succeed. For each root α , we'll find a solution v_α of equation 14 which is slight and locally integrable at $\zeta = \alpha$. We know the function $\mathcal{L}_{\zeta,\alpha} v_\alpha$ will satisfy equation 11, and we can even find its asymptotics from the order τ_α of v_α . We learned in Section 1.2 that

$$\mathcal{L}_{\zeta,\alpha} v_\alpha = e^{-\alpha z} V_\alpha$$

where $V_\alpha = \mathcal{L}_{\zeta_\alpha, 0} v_\alpha$ and $\zeta = \alpha + \zeta_\alpha$. We can see from Section 1.1.1 that V_α is asymptotic to a scalar multiple of $z^{-1-\tau_\alpha}$ at $z = \infty$, so the further decomposition

$$\mathcal{L}_{\zeta, \alpha} v_\alpha = e^{-\alpha z} z^{-\tau_\alpha} W_\alpha,$$

makes W_α is asymptotic to a scalar multiple of z^{-1} at $z = \infty$.

$$\left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (\zeta - 1)(\zeta + 1) - \frac{\partial}{\partial \zeta} \circ \zeta - \left(\frac{1}{3} \right)^2 \right] v = 0$$

$$\left[\left[2 + 2(2\zeta) \frac{\partial}{\partial \zeta} + (\zeta^2 - 1) \left(\frac{\partial}{\partial \zeta} \right)^2 \right] - \left[1 + \zeta \frac{\partial}{\partial \zeta} \right] - \left(\frac{1}{3} \right)^2 \right] v = 0$$

$$\left[(\zeta^2 - 1) \left(\frac{\partial}{\partial \zeta} \right)^2 + 3\zeta \frac{\partial}{\partial \zeta} + \left[1 - \left(\frac{1}{3} \right)^2 \right] \right] v = 0$$

$$\left[(\zeta - 1)(\zeta + 1) \left(\frac{\partial}{\partial \zeta} \right)^2 + 3\zeta \frac{\partial}{\partial \zeta} + \left[1 - \left(\frac{1}{3} \right)^2 \right] \right] v = 0$$

3.4.2 Focus on $\zeta = 1$

Let's find a solution of equation 14 which is slight and locally integrable at $\zeta = 1$. Define a new coordinate ζ_1 on \mathbb{C} so that $\zeta = 1 + \zeta_1$. In this coordinate, equation 14 looks like

$$\left[\zeta_1(2 + \zeta_1) \left(\frac{\partial}{\partial \zeta_1} \right)^2 + 3(1 + \zeta_1) \frac{\partial}{\partial \zeta_1} + \left[1 - \left(\frac{1}{3} \right)^2 \right] \right] v = 0. \quad (15)$$

With another change of coordinate, given by $\zeta_1 = -2\xi_1$, we can rewrite equation 14 as the hypergeometric equation

$$\left[\xi_1(1 - \xi_1) \left(\frac{\partial}{\partial \xi_1} \right)^2 + 3\left(\frac{1}{2} - \xi_1\right) \frac{\partial}{\partial \xi_1} - \left[1 - \left(\frac{1}{3} \right)^2 \right] \right] v = 0. \quad (16)$$

Looking through the twenty-four expressions for Kummer's six solutions, we find one [3, formula 15.10.12] which is manifestly slight and locally integrable at $\xi_1 = 0$:

$$\begin{aligned} v_1 &= \xi_1^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; \xi_1\right) \\ &= -i\sqrt{2} \zeta_1^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; -\frac{1}{2}\zeta_1\right) \end{aligned}$$

From the argument in Section 3.4.1, we know that $\mathcal{L}_{\zeta, 1} v_1$ satisfies equation 10, and can be written as $e^{-z} V_1$, where $V_1 = \mathcal{L}_{\zeta_1, 0} v_1$. Since v_1 has order $-1/2$, the decomposition $V_1 = z^{1/2} W_1$ makes W_1 asymptotic to a scalar multiple of z^{-1} at $z = \infty$.

3.4.3 Focus on $\zeta = -1$

Let's find a solution of equation 14 which is slight and locally integrable at $\zeta = -1$. In the rescaled coordinate from Section 3.4.2, this is the point $\xi_1 = 1$. Looking again through Kummer's table of solutions, we find another expression [3, formula 15.10.14] which is manifestly slight and locally integrable at $\xi_1 = 1$:

$$\begin{aligned} v_{-1} &= (1 - \xi_1)^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; 1 - \xi_1\right) \\ &= \sqrt{2} \zeta_{-1}^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; \frac{1}{2}\zeta_{-1}\right) \end{aligned}$$

where ζ_{-1} is the coordinate with $\zeta = -1 + \zeta_{-1}$. From the argument in Section 3.4.1, we know that $\mathcal{L}_{\zeta_{-1}} v_{-1}$ satisfies equation 10, and can be written as $e^z V_{-1}$, where $V_{-1} = \mathcal{L}_{\zeta_{-1}, 0} v_{-1}$. Since v_{-1} , like our other solution, has order $-1/2$, the same decomposition $V_{-1} = z^{1/2} W_{-1}$ makes W_{-1} asymptotic to a scalar multiple of z^{-1} at $z = \infty$.

In this example, v_1 and v_{-1} happen to be related by a symmetry: the Möbius transformation that pulls ζ back to $-\zeta$. Kummer's solutions typically come from six different hypergeometric equations, which are related by the Möbius transformations that permute their singularities. In our case, though, exchanging 1 with -1 keeps equation 14 the same.

If we'd followed the routine from Section 3.4.2, rewriting equation 14 in the coordinate ζ_{-1} and then rewriting it again in a more recognizable form, we would've arrived at the hypergeometric equation

$$\left[\xi_{-1}(1 - \xi_{-1}) \left(\frac{\partial}{\partial \xi_{-1}} \right)^2 + 3 \left(\frac{1}{2} - \xi_{-1} \right) \frac{\partial}{\partial \xi_{-1}} - \left[1 - \left(\frac{1}{3} \right)^2 \right] \right] v = 0,$$

where $\zeta_{-1} = 2\xi_{-1}$. This is the same as what we'd get by substituting ξ_{-1} for ξ_1 in Section 3.4.2! In other words, the holomorphic map that pulls ξ_{-1} back to ξ_1

4 The Airy-Lucas equation

4.1 Basics

The Airy-Lucas equation is

$$\left[\left(\frac{\partial}{\partial y} \right)^2 - (m-1)y^{-1} \frac{\partial}{\partial y} - y^{n-2} \right] \psi = 0 \quad (17)$$

with $n \in \{3, 4, 5, \dots\}$ and $m \in \{1, 2, \dots, r-1\}$. A few solutions are given by the Airy-Lucas functions [Charbonnier et al., equation 3.6]

$$\widehat{\text{Ai}}_{n,m-1}^{(k)}(y) = \begin{Bmatrix} 1 & j \text{ even} \\ i & j \text{ odd} \end{Bmatrix} \frac{y^{m/2}}{\pi} \int_{\Lambda^{(j)}} \exp \left[\frac{2}{n} y^{n/2} T_n(u) \right] U_{m-1}(u) du,$$

where $\Lambda^{(k)}$ is the Lefschetz thimble through $u = \cos \left(\frac{k}{n} \pi \right)$.

4.2 Rewriting as a modified Bessel (?) equation

We can distill the most interesting parts of the Airy-Lucas function by writing

$$\widehat{\text{Ai}}_{n,m-1}^{(k)}(y) = \text{const.} y^{m/2} K \left(\frac{2}{n} y^{n/2} \right),$$

where

$$K(z) = \text{const.} \int_{z^{-1/n} \Lambda^{(k)}} \exp [-z T_n(u)] U_{m-1}(u) du. \quad (18)$$

Saying that $\widehat{\text{Ai}}_{n,m-1}^{(k)}$ satisfies the Airy-Lucas equation is equivalent to saying that K satisfies the modified Bessel (?) equation

$$\left[z^2 \left(\frac{\partial}{\partial z} \right)^2 + z \frac{\partial}{\partial z} - \left[\left(\frac{m}{n} (?) \right)^2 + z^2 \right] \right] \varphi = 0. \quad (19)$$

In fact, as we'll see in Section [?](#), K is the modified Bessel function $K_{m/n}$.

Like we did in equation 11, we can rewrite the modified Bessel equation above as

$$\left[\left(\frac{\partial}{\partial z} \right)^2 - 1 \right] + z^{-1} \frac{\partial}{\partial z} - \left(\frac{m}{n} \right)^2 z^{-2} \varphi = 0. \quad (20)$$

5 The higher Airy equation

5.1 Basics

The higher Airy equation is

$$\left[\left(-\frac{\partial}{\partial y} \right)^{n-1} - y \right] \psi = 0 \quad (21)$$

with $n \in \{3, 4, 5, \dots\}$. A few solutions are given by the hyper-Airy functions [\[Charbonnier et al., equation 3.6\]](#)

With

$$z = (-1)^{n-1} \frac{n-1}{n} y^{n/(n-1)} \quad w = (-1)^n (n-1) y^{1/(n-1)} u,$$

we have

$$\begin{aligned} \widetilde{\text{Ai}}_n^{(k)}(y) &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[\frac{1}{n} w^n - y w\right] dw \\ &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[\frac{1}{n} w (w^{n-1} - ny)\right] dw \\ &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[\frac{1}{n} w ((n-1)^{n-1} y u^{n-1} - ny)\right] dw \\ &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[\frac{1}{n} y w ((n-1)^{n-1} u^{n-1} - n)\right] dw \\ &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[(-1)^n \frac{n-1}{n} y^{n/(n-1)} u ((n-1)^{n-1} u^{n-1} - n)\right] dw \\ &= \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} \int_{\Lambda^{(j)}} \exp\left[-z ((n-1)^{n-1} u^n - nu)\right] dw \\ &= (-1)^n (n-1) \frac{\exp\left(\pi i k \frac{n-2}{n-1}\right)}{2\pi i} y^{1/(n-1)} \int_{\Lambda^{(j)}} \exp\left[-z ((n-1)^{n-1} u^n - nu)\right] du \\ &= (-1)^n (n-1) \frac{\exp\left(\pi i k \left(1 - \frac{1}{n-1}\right)\right)}{2\pi i} y^{1/(n-1)} \int_{\Lambda^{(j)}} \exp\left[-z ((n-1)^{n-1} u^n - nu)\right] du \\ &= (-1)^{n+k} (n-1) \frac{\exp\left(-\pi i \frac{k}{n-1}\right)}{2\pi i} y^{1/(n-1)} \int_{\Lambda^{(j)}} \exp\left[-z ((n-1)^{n-1} u^n - nu)\right] du \end{aligned}$$

$$\widetilde{\text{Ai}}_n^{(k)}(y) = (-1)^{n+k} (n-1) \frac{\exp\left(-\pi i \frac{k}{n-1}\right)}{2\pi i} y^{1/(n-1)} \int_{\Lambda^{(j)}} \exp\left[-z ((n-1)^{n-1} u^n - nu)\right] du,$$

where $\Lambda^{(k)}$ is the Lefschetz thimble through $u = \cos\left(\frac{k}{n}\pi\right)$.

5.2 Rewriting as a (???) equation

We can distill the most interesting parts of the hyper-Airy function by writing

$$\widetilde{\text{Ai}}_n^{(k)}(y) = \text{const.} y^{1/(n-1)} K\left((-1)^n \frac{n-1}{n} y^{n/(n-1)}\right),$$

where

$$K(z) = \text{const.} \int_{\text{const.}(n)z^{-1/n}\Lambda^{(k)}} \exp\left[-z\left((n-1)^{n-1}u^n - nu\right)\right] du. \quad (22)$$

Saying that $\widetilde{\text{Ai}}_n^{(k)}$ satisfies the higher Airy equation is equivalent to saying that K satisfies an equation of the form

$$\left[\left(-\frac{\partial}{\partial z}\right)^{n-1} - 1\right] - c_n^{(1)}z^{-1}\left(-\frac{\partial}{\partial z}\right)^{n-2} - c_n^{(2)}z^{-2}\left(-\frac{\partial}{\partial z}\right)^{n-3} - \dots - c_n^{(n-1)}z^{-(n-1)}\Big] \varphi = 0. \quad (23)$$

The sub-leading coefficients are the triangular numbers

$$c_n^{(1)} = \frac{n(n-1)}{2}.$$

The later coefficients can be written as³

$$c_n^{(k)} = \frac{b_n^{(k)}}{n^k} (n+1)^{\underline{k+2}}$$

in terms of the polynomials

$$\begin{aligned} b_n^{(2)} &= \frac{1}{24} \\ b_n^{(3)} &= \frac{1}{48}n \\ b_n^{(4)} &= \frac{73}{5760}n^2 + \frac{1}{1152}n - \frac{1}{2880} \\ b_n^{(5)} &= \frac{11}{1280}n^3 + \frac{768}{61}n^2 - \frac{1920}{181}n \\ b_n^{(6)} &= \frac{3625}{580608}n^4 + \frac{61}{41472}n^3 - \frac{181}{322560}n^2 - \frac{1}{41472}n + \frac{1}{181440}. \end{aligned}$$

Searching for 580608 in the OEIS turns up the leading coefficients $\beta^{(k)}$ of these polynomials, which are listed as A249276 and A249277. They're defined by the identity [Yang, "Approximations for Constant e and Their Applications"]

$$\frac{1}{e} \left(\frac{n}{n-1}\right)^{n-1} = 1 - \frac{1/2}{n} - \frac{\beta^{(2)}}{n^2} - \frac{\beta^{(3)}}{n^3} - \frac{\beta^{(4)}}{n^4} - \frac{\beta^{(5)}}{n^5} - \dots,$$

which tells us that

$$\frac{1}{e} \left(\frac{n}{n-1}\right)^{n-1} = 1 - \frac{1/2}{n} - \frac{b_n^{(2)}}{n^2} - \frac{b_n^{(3)}}{n^4} - \left[\frac{b_n^{(4)}}{n^6} + o\left(\frac{1}{n^5}\right)\right] - \left[\frac{b_n^{(5)}}{n^8} + o\left(\frac{1}{n^6}\right)\right] - \dots$$

The last coefficient can be written as

$$c_n^{(n-1)} = \left(\frac{n-1}{n}\right)^{n-1} \left(\frac{1}{n-1}\right)^{\underline{n-1}},$$

³Many thanks to Peter Taylor for noticing this [<https://mathoverflow.net/q/422337/1096>].

giving

$$b_n^{(n-1)} = (n-1)^{n-1} \left(\frac{1}{n-1} \right)^{\frac{n-1}{2}} / (n+1)^{\frac{n+1}{2}}.$$

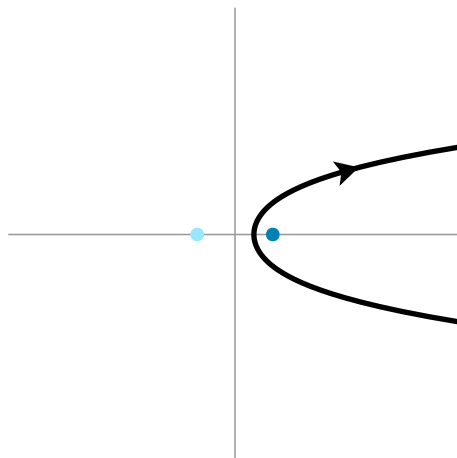
6 Sketches

6.1 Contour argument for the Airy function

We can recast integral 9 into the ζ plane by setting $\zeta = 4u^3 - 3u$. Projecting $z^{-1/3}\Gamma$ to a contour γ_z in the ζ plane and choosing the branch of u that lifts γ_z back to $z^{-1/3}\Gamma$, we have

$$K = \frac{i}{\sqrt{3}} \int_{\gamma_z} e^{-z\zeta} \frac{d\zeta}{4u^2 - 1}. \quad (24)$$

For $z \in (0, \infty)$, the contour γ_z runs clockwise around $[1, \infty)$, as shown below. Let's assume $z \in (0, \infty)$ for the rest of the section. [Our conclusions should probably hold whenever $\text{Re}(z) > 0$.]



The contour γ_1 in the ζ plane.

It happens⁴ that for our desired branch of u ,

$$\frac{1}{4u^2 - 1} = -F\left(\frac{1}{3}, \frac{2}{3}; \frac{1}{2}; \zeta^2\right),$$

so we can rewrite integral 24 as

$$K = \frac{1}{i\sqrt{3}} \int_{\gamma_z} e^{-z\zeta} F\left(\frac{1}{3}, \frac{2}{3}; \frac{1}{2}; \zeta^2\right) d\zeta.$$

This gives us an alternate route to the conclusion of Section 3.4, which we'll follow below.

⁴Veronica: This comes from [3, equation 15.4.14].

In addition to the solutions g_1 and f_0 from Section ??, equation ?? has the solutions

$$\begin{aligned} g_0 &= F\left(\frac{2}{3}, \frac{4}{3}; \frac{3}{2}; 1 - \xi\right) \\ f_1 &= (1 - \xi)^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; 1 - \xi\right), \end{aligned}$$

given by formulas 15.10.13 and 15.10.14 from [3].

The quadratic transformation identity 15.8.27 from [3] shows [verified numerically] that⁵

$$F\left(\frac{1}{3}, \frac{2}{3}; \frac{1}{2}; \zeta^2\right) = \frac{1}{3}(g_1 + g_0),$$

so we have

$$K = \frac{1}{i3\sqrt{3}} \int_{\gamma_z} e^{-z\zeta} (g_1 + g_0) d\zeta.$$

The solution g_1 is holomorphic on $\zeta \in [1, \infty)$, so it integrates to zero. The solution g_0 , in contrast, is non-meromorphic at $\zeta = 1$. Along the branch cut $\zeta \in [1, \infty)$, its above-minus-below difference is $-\frac{3\sqrt{3}}{2} f_0$, as given⁶ by equation 15.2.3 from [3]. Hence,

$$\begin{aligned} K &= \frac{i}{2} \int_1^\infty e^{-z\zeta} f_0 d\zeta \\ e^z K &= \frac{i}{2} \int_1^\infty e^{-z(\zeta-1)} f_0 d\zeta \\ K_1 &= \frac{i}{2} \mathcal{L}_{\zeta_1} f_0, \end{aligned}$$

just as we found in Section ??.

6.2 Another solution

Section 6.1 associates the solution K of equation 10 with the solution g_0 of equation ??, which contributes the pole at $\zeta = 1$ of

$$\frac{du}{d\zeta} = \frac{1}{4u^2 - 1} = \frac{1}{3}(g_1 + g_0).$$

The solution g_1 , which contributes the pole at $\zeta = -1$, is associated with another solution of equation 10.

To express this other solution as a Laplace transform, following the method of Section ??, we would use the solution

$$f_1 = (1 - \xi)^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; 1 - \xi\right)$$

of equation ??, given by formula 15.10.14 from [3]. This is the only solution, up to scale, which has a fractional power singularity at $\zeta = -1$.

⁵Note that $2\Gamma(\frac{1}{2})\Gamma(\frac{3}{2}) = 2\Gamma(\frac{1}{2})\frac{1}{2}\Gamma(\frac{1}{2}) = \pi$ and $[\Gamma(\frac{5}{6})\Gamma(\frac{7}{6})]^{-1} = [\Gamma(\frac{5}{6})\frac{1}{6}\Gamma(\frac{1}{6})]^{-1} = \frac{6\sin(\frac{1}{6}\pi)}{\pi} = \frac{3}{\pi}$.

⁶Note that $\Gamma(\frac{3}{2})\Gamma(\frac{1}{2})^{-1} = \frac{1}{2}$ and $[\Gamma(\frac{2}{3})\Gamma(\frac{4}{3})]^{-1} = [\Gamma(\frac{2}{3})\frac{1}{3}\Gamma(\frac{1}{3})]^{-1} = \frac{3\sin(\frac{1}{3}\pi)}{\pi} = \frac{3\sqrt{3}}{2\pi}$.

In summary, the contour integration method of solving equation 10 is associated with the basis

$$\begin{aligned} g_1 &= F\left(\frac{2}{3}, \frac{4}{3}; \frac{3}{2}; \xi\right) \\ g_0 &= F\left(\frac{2}{3}, \frac{4}{3}; \frac{3}{2}; 1 - \xi\right) \end{aligned}$$

of solutions for equation ??, given by formulas 15.10.11 and 15.10.13 from [3]. These solutions contribute the poles at $\xi = 1$ and $\xi = 0$, respectively, of a generic solution.

The Laplace transformation method of solving equation 10, on the other hand, is associated with the basis

$$\begin{aligned} f_1 &= (1 - \xi)^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; 1 - \xi\right) \\ f_0 &= \xi^{-1/2} F\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; \xi\right) \end{aligned}$$

given by formulas 15.10.14 and 15.10.12 from [3]. These solutions, up to scale, are the only ones with fractional power singularities.

Identities 15.10.18, and 15.10.22 from [3] give the change of basis

$$\begin{aligned} f_1 &= \frac{1}{\sqrt{3}} g_1 + \frac{1}{2} f_0 \\ f_0 &= \frac{1}{\sqrt{3}} g_0 + \frac{1}{2} f_1. \end{aligned}$$

Summing these identities, we see that

$$g_1 + g_0 = \frac{\sqrt{3}}{2} (f_1 + f_0),$$

giving the alternate decomposition

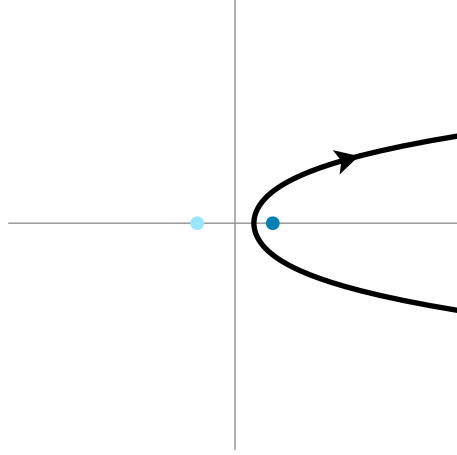
$$\frac{du}{d\zeta} = \frac{1}{2\sqrt{3}} (f_1 + f_0).$$

6.3 Contour argument for Airy-Lucas functions

Generalizing from Section 6.1, we can recast integral 18 into the ζ plane by setting $-\zeta = T_n(u)$, which implies that $-d\zeta = nU_{n-1}(u) du$. Projecting as before, we get [sign flipped] [see identity in `cyl-resurgence.tex`]

$$\begin{aligned} K_{m/n}(z) &= \frac{n}{2 \sinh\left(\frac{m}{n} i\pi\right)} \int_{z^{-1/n} \Lambda^{(k)}} \exp[zT_n(u)] U_{m-1}(u) du \\ &= -\frac{1}{2 \sinh\left(\frac{m}{n} i\pi\right)} \int_{\gamma_z} e^{-z\zeta} \frac{U_{m-1}(u)}{U_{n-1}(u)} d\zeta \\ &= -\frac{1}{2 \frac{n}{m} \sinh\left(\frac{m}{n} i\pi\right)} \int_{\gamma_z} e^{-z\zeta} F\left(1 - \frac{m}{n}, 1 + \frac{m}{n}, \frac{3}{2}, \frac{1}{2} \pm \frac{1}{2}\zeta\right) d\zeta. \end{aligned}$$

where the sign must be chosen so that $\pm\zeta$ stays in the left half-plane over the whole integration path (?). For $z \in (0, \infty)$, the contour γ_z runs counterclockwise around $[1, \infty)$, as shown below, so we have to choose the negative sign above (?). Let's assume $z \in (0, \infty)$ for the rest of the section. [Our conclusions should probably hold whenever $\text{Re}(z) > 0$.]



The contour γ_1 [reversed] in the ζ plane.

The integrand is non-meromorphic at $\zeta = 1$. Along the branch cut $\zeta \in [1, \infty)$, its above-minus-below difference is

$$\begin{aligned} & - (2\pi i) \frac{n}{2\pi m} \sin(\frac{m}{n}\pi) (\pm \frac{1}{2}\zeta - \frac{1}{2})^{-1/2} F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2} \mp \frac{1}{2}\zeta) \\ & = - \frac{n}{m} \sinh(\frac{m}{n} i\pi) (\pm \frac{1}{2}\zeta - \frac{1}{2})^{-1/2} F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2} \mp \frac{1}{2}\zeta), \end{aligned}$$

as given⁷ by equation 15.2.3 from [3]. Hence, $K_{m/n}$ turns out to be the Laplace transform along $(1, \infty)$ of

$$\frac{1}{2} (-\frac{1}{2}\zeta - \frac{1}{2})^{-1/2} F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2} + \frac{1}{2}\zeta).$$

Guessing the branch of the square root for consistency with the numerically checked result in Section 6.1, we get

$$\frac{i}{2} (\frac{1}{2} + \frac{1}{2}\zeta)^{-1/2} F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2} + \frac{1}{2}\zeta).$$

In other words, $K_{m/n} = \frac{i}{2} \mathcal{L}_{\zeta, -1} v_{-1}$ with

$$v_{-1} = \sqrt{2} \zeta_{-1}^{-1/2} F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2} \zeta_{-1}),$$

where $\zeta = -1 + \zeta_{-1}$.

6.4 Lifting to a countable cover

Formula 18 expresses the modified Bessel function $K_{m/n}$ as an exponential integral on a finite cover of \mathbb{C} . Lifting to a countable cover reveals this formula as a special case of a general integral formula for modified Bessel functions.

Setting $u = \cosh(t/n)$ and recalling that

$$\begin{aligned} \cosh(n\tau) &:= T_n(\cosh(\tau)) \\ \sinh(m\tau) &:= U_{m-1}(\cosh(\tau)) \sinh(\tau), \end{aligned}$$

⁷Note that $\Gamma(\frac{3}{2})\Gamma(\frac{1}{2})^{-1} = \frac{1}{2}$ and $[\Gamma(1 - \frac{m}{n})\Gamma(1 + \frac{m}{n})]^{-1} = [\Gamma(1 - \frac{m}{n}) \frac{m}{n} \Gamma(\frac{m}{n})]^{-1} = \frac{n}{m\pi} \sin(\frac{m}{n}\pi)$.

we can rewrite formula 18 as [switching to the conventional sign for the projection map, so $\Lambda^{(3)}$ now comes from ∞ at -60° and goes to ∞ at 60°]

$$\begin{aligned}
K_{m/n}(z) &= \frac{n}{2 \sinh\left(\frac{m}{n} i\pi\right)} \int_{z^{-1/n} \Lambda^{(k)}} \exp[z T_n(u)] U_{m-1}(u) du \\
&= \frac{n}{2 \sinh\left(\frac{m}{n} i\pi\right)} \int_{\Omega} \exp[z \cosh(t)] U_{m-1}(\cosh(t/n)) \sinh(t/n) d(t/n) \\
&= \frac{1}{2 \sinh\left(\frac{m}{n} i\pi\right)} \int_{\Omega} \exp[z \cosh(t)] \sinh\left(\frac{m}{n} t\right) dt.
\end{aligned} \tag{25}$$

For any $\nu \in \mathbb{C} \setminus \mathbb{Z}$, formulas 10.27.4 and 10.32.12 from [3] tell us that

$$\begin{aligned}
K_\nu(z) &= \frac{\pi}{\sin(\nu\pi)} \cdot \frac{1}{2} [I_{-\nu}(z) - I_\nu(z)] \\
&= \frac{1}{2i \sin(\nu\pi)} \int_{\Omega} \exp[z \cosh(t)] \frac{1}{2} [e^{\nu t} - e^{-\nu t}] dt \\
&= \frac{1}{2 \sinh(\nu i\pi)} \int_{\Omega} \exp[z \cosh(t)] \sinh(\nu t) dt,
\end{aligned} \tag{26}$$

where Ω is a path that comes from ∞ along $-i\pi + (0, \infty)$ and goes to ∞ along $i\pi + (0, \infty)$. The integral converges when z is in the right half-plane. We get formula 25 when we choose a rational parameter $\nu = m/n$.

When ν goes to 0, formula 25 becomes

$$K_\nu(z) = \frac{1}{2\pi i} \int_{\Omega} \exp[z \cosh(t)] t dt.$$

Choosing Ω to be the unit-speed path that runs from ∞ leftward to $-i\pi$, upward to $i\pi$, and rightward back to ∞ , we can rewrite this formula as

$$\begin{aligned}
K_\nu(z) &= \frac{1}{2\pi i} \int_0^\infty \exp[-z \cosh(t)] 2\pi i dt \\
&= \int_0^\infty \exp[-z \cosh(t)] dt \\
&= \int_1^\infty \exp\left[-z \frac{1}{2} \left(s + \frac{1}{s}\right)\right] \frac{ds}{s},
\end{aligned}$$

with $s = e^t$. This is a special case of formula 10.32.9 from [3].

6.5 Link to dihedral triangulations

Here are the local normal forms of Kummer's first and second solutions. In the table, $[\bullet]$ denotes a local holomorphic function of order 0 [3, identities 15.8.2–4], and

$$\begin{aligned}
a' &= a + (1 - c) \\
b' &= b + (1 - c) \\
c' &= 2 - c
\end{aligned}$$

Function	0	1	∞
$F(a, b, c, \xi)$	$[\bullet]$	$[\bullet] + (1 - \xi)^{c-a-b}[\bullet]$	$(1/\xi)^a[\bullet] + (1/\xi)^b[\bullet]$
$\xi^{1-c}F(a', b', c', \xi)$	$\xi^{1-c}[\bullet]$	$[\bullet] + (1 - \xi)^{c-a-b}[\bullet]$	$(1/\xi)^a[\bullet] + (1/\xi)^b[\bullet]$

For example, at $\zeta = \infty$, the function

$$u_{-1} = F(1 - \frac{m}{n}, 1 + \frac{m}{n}, \frac{3}{2}, 1 - \frac{1}{2}\zeta_{-1})$$

is slight, with normal form $(1/\zeta_{-1})^{1-m/n}f_1 + (1/\zeta_{-1})^{1+m/n}f_2$ [3, identity 15.8.3]. At $\zeta = 1$, it's holomorphic with order 0. At $\zeta = -1$, it has the form $h + \zeta_{-1}^{-1/2}f$, where h is holomorphic with order 0.

On the other hand, the function

$$v_{-1} = \sqrt{2}\zeta_{-1}^{-1/2}F(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, \frac{1}{2}\zeta_{-1})$$

is slight at $\zeta = \pm 1$ with order $-\frac{1}{2}$, and slight at ∞ with order $1 - \frac{m}{n}$.

Hence, the ratio u_{-1}/v_{-1} looks like a holomorphic multiple of $\zeta_1^{1/2}$ near $\zeta = 1$, a holomorphic function plus a holomorphic multiple of $\zeta_{-1}^{1/2}$ near $\zeta = -1$, and a holomorphic function plus a holomorphic multiple of $(1/\zeta_{-1})^{2m/n}$ near $\zeta = \infty$. This is consistent with u_{-1}/v_{-1} being a conformal map that sends the upper half-plane to a triangle with interior angle $\pi/2$ at the images of $\zeta = \pm 1$ and interior angle $2\pi \frac{m}{n}$ at the image of ∞ .

Let's try to get the small tetrahedral tiling by setting

$$\begin{aligned} 1 - c &= \frac{1}{3} \\ c - a - b &= \frac{1}{3} \\ b - a &= \frac{1}{2}, \end{aligned}$$

so $(a, b, c) = (-\frac{1}{12}, \frac{5}{12}, \frac{2}{3})$. The general hypergeometric equation is

$$\begin{aligned}
& \left[\xi_1(1 - \xi_1) \left(\frac{\partial}{\partial \xi_1} \right)^2 + [c - (a + b + 1)\xi_1] \frac{\partial}{\partial \xi_1} - ab \right] v = 0 \\
& \left[-\frac{1}{2}\xi_1(1 + \frac{1}{2}\xi_1)4 \left(\frac{\partial}{\partial \xi_1} \right)^2 + [c + \frac{1}{2}(a + b + 1)\xi_1](-2) \frac{\partial}{\partial \xi_1} - ab \right] v = 0 \\
& \left[-\xi_1(2 + \xi_1) \left(\frac{\partial}{\partial \xi_1} \right)^2 - [2c + (a + b + 1)\xi_1] \frac{\partial}{\partial \xi_1} - ab \right] v = 0 \\
& \left[(1 - \zeta)(1 + \zeta) \left(\frac{\partial}{\partial \zeta} \right)^2 - [2c + (a + b + 1)(\zeta - 1)] \frac{\partial}{\partial \zeta} - ab \right] v = 0 \\
& \left[(1 - \zeta^2) \left(\frac{\partial}{\partial \zeta} \right)^2 - [(c - a - b) + (c - 1) + (a + b + 1)\zeta] \frac{\partial}{\partial \zeta} - ab \right] v = 0 \\
& \left[(1 - \zeta^2) \left(\frac{\partial}{\partial \zeta} \right)^2 - [\mu - \lambda + (a + b + 1)\zeta] \frac{\partial}{\partial \zeta} - ab \right] v = 0 \\
& \left[\left[\frac{\partial}{\partial \zeta} \circ (1 - \zeta^2) + 2\zeta \right] \frac{\partial}{\partial \zeta} - (\mu - \lambda) \frac{\partial}{\partial \zeta} - (a + b + 1) \left[\frac{\partial}{\partial \zeta} \circ \zeta - 1 \right] - ab \right] v = 0 \\
& \left[\frac{\partial}{\partial \zeta} \circ \left[\frac{\partial}{\partial \zeta} \circ (1 - \zeta^2) + 2\zeta \right] + 2 \left[\frac{\partial}{\partial \zeta} \circ \zeta - 1 \right] - (\mu - \lambda) \frac{\partial}{\partial \zeta} - (a + b + 1) \left[\frac{\partial}{\partial \zeta} \circ \zeta - 1 \right] - ab \right] v = 0 \\
& \left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (1 - \zeta^2) + 4 \frac{\partial}{\partial \zeta} \circ \zeta - 2 - (\mu - \lambda) \frac{\partial}{\partial \zeta} - (a + b + 1) \frac{\partial}{\partial \zeta} \circ \zeta + [(a + b + 1) - ab] \right] v = 0 \\
& \left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (1 - \zeta^2) - (\mu - \lambda) \frac{\partial}{\partial \zeta} + (3 - a - b) \frac{\partial}{\partial \zeta} \circ \zeta + [-1 + a + b - ab] \right] v = 0 \\
& \left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (1 - \zeta^2) - (\mu - \lambda) \frac{\partial}{\partial \zeta} + (3 - a - b) \frac{\partial}{\partial \zeta} \circ \zeta - (a - 1)(b - 1) \right] v = 0 \\
& \left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (\zeta^2 - 1) + (\mu - \lambda) \frac{\partial}{\partial \zeta} + (a + b - 3) \frac{\partial}{\partial \zeta} \circ \zeta + (a - 1)(b - 1) \right] v = 0,
\end{aligned}$$

For the tetrahedral tiling, we get

$$\left[\left(\frac{\partial}{\partial \zeta} \right)^2 \circ (\zeta^2 - 1) - \frac{8}{3} \frac{\partial}{\partial \zeta} \circ \zeta + \frac{7 \cdot 13}{12^2} \right] v = 0,$$

which comes from the integral equation

$$\left[(\zeta^2 - 1) - \frac{8}{3} \partial_{\zeta, \alpha}^{-1} \circ \zeta + \frac{7 \cdot 13}{12^2} \partial_{\zeta, \alpha}^{-2} \right] v = 0.$$

In the frequency domain, this equation becomes

$$\left[\left[\left(\frac{\partial}{\partial z} \right)^2 - 1 \right] + \frac{8}{3} z^{-1} \frac{\partial}{\partial z} + \frac{7 \cdot 13}{12^2} z^{-2} \right] \mathcal{L}_{\zeta, \alpha} v = 0.$$

Set

$$\mathcal{L}_{\zeta, \alpha} v = \int_{\Lambda} e^{-zf} p \, du.$$

Observe that

$$\left[\left(\frac{\partial}{\partial z} \right)^2 - 1 \right] \int_{\Lambda} e^{-zf} p \, du = \int_{\Lambda} e^{-zf} (f^2 - 1) p \, du$$

and

$$\begin{aligned}
z^{-1} \frac{\partial}{\partial z} \int_{\Lambda} e^{-zf} p \, du &= z^{-1} \int_{\Lambda} e^{-zf} (-fp) \, du \\
&= -z^{-1} \int_{\Lambda} e^{-zf} (-zf') q_1 \, du \\
&= \int_{\Lambda} e^{-zf} f' q_1 \, du
\end{aligned}$$

where $q'_1 = -fp$. and

$$\begin{aligned}
z^{-2} \int_{\Lambda} e^{-zf} p \, du &= -z^{-2} \int_{\Lambda} e^{-zf} (-zf') p_1 \, du \\
&= z^{-1} \int_{\Lambda} e^{-zf} f' p_1 \, du \\
&= -z^{-1} \int_{\Lambda} e^{-zf} (-zf') p_2 \, du \\
&= \int_{\Lambda} e^{-zf} f' p_2 \, du
\end{aligned}$$

with $p'_1 = p$ and $p'_2 = f'p_1$.

$$\left[\left(\frac{\partial}{\partial z} \right)^2 - 1 \right] + \frac{8}{3} z^{-1} \frac{\partial}{\partial z} + \frac{7 \cdot 13}{12^2} z^{-2} \Big] \mathcal{L}_{\zeta, \alpha} v = 0.$$

In DLMF notation, let's try

$$\begin{aligned}
u_1 &= F\left(1 - \frac{m}{n}, 1 + \frac{m}{n}, \frac{3}{2}, \xi_{-1}\right) \\
&= w_1(\xi_{-1}) \\
u_{-1} &= w_3(\xi_{-1}) \\
v_{-1} &= w_2(\xi_{-1}) \\
v_1 &= -i\sqrt{2} \zeta_1^{-1/2} F\left(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, -\frac{1}{2}\zeta_1\right) \\
&= \left(-\frac{1}{2}\zeta_1\right)^{-1/2} F\left(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, -\frac{1}{2}\zeta_1\right) \\
&= \left(1 - \frac{1}{2}\zeta_{-1}\right)^{-1/2} F\left(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, 1 - \frac{1}{2}\zeta_{-1}\right) \\
&= \left(1 - \xi_{-1}\right)^{-1/2} F\left(\frac{1}{2} + \frac{m}{n}, \frac{1}{2} - \frac{m}{n}, \frac{1}{2}, 1 - \xi_{-1}\right) \\
&= w_4(\xi_{-1}),
\end{aligned}$$

where $\xi_{-1} = \frac{1}{2}\zeta_{-1}$.

Let's guess that we're working with the map ϕ/ψ with

$$\begin{aligned}
\phi &= \frac{\Gamma(-\frac{1}{2})\Gamma(\frac{3}{2})}{\Gamma(\frac{1}{2}-\frac{m}{2n})\Gamma(\frac{1}{2}+\frac{m}{2n})} (\mp\zeta) F(1-\frac{m}{2n}, 1+\frac{m}{2n}, \frac{3}{2}; \zeta^2) \\
&= \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2} \pm \frac{1}{2}\zeta) - \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2} \mp \frac{1}{2}\zeta) \\
&= \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, -\frac{1}{2}\zeta_1) - \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2}\zeta_{-1}) \quad [\text{choosing negative sign}] \\
\psi &= \frac{\Gamma(\frac{1}{2})\Gamma(\frac{3}{2})}{\Gamma(1-\frac{m}{2n})\Gamma(1+\frac{m}{2n})} F(\frac{1}{2}-\frac{m}{2n}, \frac{1}{2}+\frac{m}{2n}, \frac{1}{2}; \zeta^2) \\
&= \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2} \pm \frac{1}{2}\zeta) + \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2} \mp \frac{1}{2}\zeta) \\
&= \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, -\frac{1}{2}\zeta_1) + \frac{1}{2}F(1-\frac{m}{n}, 1+\frac{m}{n}, \frac{3}{2}, \frac{1}{2}\zeta_{-1}) \quad [\text{choosing negative sign}]
\end{aligned}$$

6.6 Comparison with Ghate and Venkataramana

$$\begin{aligned}
\mu_{\text{GV}} &= 1 - \left\{ \frac{k_1}{d} \right\} - \left\{ \frac{k_3}{d} \right\} & \lambda &= 1 - c \\
\nu_{\text{GV}} &= 1 - \left\{ \frac{k_2}{d} \right\} - \left\{ \frac{k_3}{d} \right\} & \mu &= c - a - b \\
\lambda_{\text{GV}} &= 1 - \left\{ \frac{k_1}{d} \right\} - \left\{ \frac{k_2}{d} \right\} & \nu &= b - a
\end{aligned}$$

The hyperelliptic curve in Section 4 is

$$y^d = (x - \alpha_1)^{k_1} (x - \alpha_2)^{k_2} (x - \alpha_3)^{k_3}.$$

Dihedral case (Theorem 2):

$$\begin{aligned}
\frac{k_1}{d} &= \frac{p}{2m} \\
\frac{k_2}{d} &= \frac{p}{2m} \\
\frac{k_3}{d} &= \frac{m-p}{2m}
\end{aligned}$$

Note that $\lambda = p'/m$ with $p + p' = m$. The hyperelliptic curve is

$$y^d = (x - \alpha_1)^{k_1} (x - \alpha_2)^{k_2} (x - \alpha_3)^{k_3}.$$

Cosines appear in Section 2.3: $x_j = \exp(2\pi i \frac{k_j}{d})$, and we work in the totally real field $\mathbb{Q}(\cos(\frac{2\pi}{m}))$.

Airy case:

$$\begin{aligned}\frac{k_1}{d} &= \frac{1}{3} \\ \frac{k_2}{d} &= \frac{1}{3} \\ \frac{k_3}{d} &= \frac{1}{6},\end{aligned}$$

so $(k_1, k_2, k_3, d) = (2, 2, 1, 6)$, and the hyperelliptic curve is

$$y^6 = (x - \alpha_1)^2(x - \alpha_2)^2(x - \alpha_3).$$

6.7 Correspondence with Mariño's series

Let $F_1(z)$ be the holomorphic function corresponding to Mariño's formal power series $\varphi_1(z^{-1})$. The formal power series corresponding to F_1 will be written in the variable z .

$$\begin{aligned}\text{Ai}(x) &= \frac{1}{2\sqrt{\pi}} x^{-1/4} e^{-z} \varphi_1\left(\frac{2}{3}z^{-1}\right) \\ &= \frac{1}{2\sqrt{\pi}} x^{-1/4} e^{-z} F_1\left(\frac{3}{2}z\right) \\ \text{Ai}(x) &= \frac{1}{\pi\sqrt{3}} x^{1/2} K\left(\frac{2}{3}x^{3/2}\right)\end{aligned}$$

Putting together,

$$\begin{aligned}\frac{1}{2\sqrt{\pi}} x^{-1/4} e^{-z} F_1\left(\frac{3}{2}z\right) &= \frac{1}{\pi\sqrt{3}} x^{1/2} K\left(\frac{2}{3}x^{3/2}\right) \\ \frac{\sqrt{3\pi}}{2} x^{-3/4} e^{-z} F_1\left(\frac{3}{2}z\right) &= K\left(\frac{2}{3}x^{3/2}\right) \\ \frac{\sqrt{3\pi}}{2} \left(\frac{3}{2}z\right)^{-1/2} e^{-z} F_1\left(\frac{3}{2}z\right) &= K(z) \\ \sqrt{\frac{\pi}{2}} z^{-1/2} e^{-z} F_1\left(\frac{3}{2}z\right) &= K(z) \\ \sqrt{\frac{\pi}{2}} \left[\mathcal{L}^{-1} z^{-1/2}\right] * \left[\mathcal{L}^{-1} F_1\left(\frac{3}{2}z\right)\right](\zeta - 1) &= k(\zeta) \\ \sqrt{\frac{\pi}{2}} \left[\Gamma\left(-\frac{1}{2}\right)^{-1} \zeta^{-1/2}\right] * \frac{2}{3} f_1\left[\frac{2}{3}(\zeta - 1)\right] &= k(\zeta) \\ -\frac{1}{3\sqrt{2}} \left[\zeta^{-1/2}\right] * f_1\left[\frac{2}{3}(\zeta - 1)\right] &= k(\zeta)\end{aligned}$$

Notice that if the hypergeometric differentiation formula holds for fractional derivatives,

$$\left(\frac{\partial}{\partial \xi}\right)^{1/2} F\left(\frac{2}{3}, \frac{4}{3}; \frac{3}{2}; \xi\right) \propto F\left(\frac{7}{6}, \frac{11}{6}; 2; \xi\right)$$

References

- [1] C. Mitschi and D. Sauzin, *Divergent Series, Summability and Resurgence I: Monodromy and Resurgence*. No. 2153 in Lecture Notes in Mathematics. 2016.
- [2] J. L. Schiff, *The Laplace Transform: Theory and Applications*. Springer, 1999.
- [3] “NIST digital library of mathematical functions.” <http://dlmf.nist.gov>, release 1.1.3 of 2021-09-15. F. W. J. Olver, A. B. Olde Daalhuis, D. W. Lozier, B. I. Schneider, R. F. Boisvert, C. W. Clark, B. R. Miller, B. V. Saunders, H. S. Cohl, and M. A. McClain, editors.