

# Resurgence of the Airy function and other exponential integrals

Veronica Fantini and Aaron Fenyes

April 1, 2022

## 1 Introduction

### 1.1 Why does Borel resummation work?

Borel resummation is a way of turning a formal power series

$$\varphi_\bullet = z^\sigma \left( \frac{\varphi_0}{z} + \frac{\varphi_1}{z^2} + \frac{\varphi_2}{z^3} + \frac{\varphi_3}{z^4} + \dots \right),$$

with  $\sigma \in [0, 1)$ , into a function which is asymptotic to  $\varphi_\bullet$  as  $z \rightarrow \infty$ . Different functions can be asymptotic to the same power series, and Borel resummation picks one of them, performing an implicit regularization [[arXiv:1705.03071](#), or maybe [arXiv:1412.6614](#)]. When a function matches the Borel sum of its asymptotic series, we'll say it's *Borel regular*. Several familiar kinds of regularity imply Borel regularity, and shed light on why it occurs.

- **Having a good asymptotic approximation**

Let  $R_N$  be the difference between a function and the partial sum

$$\frac{\varphi_0}{z} + \frac{\varphi_1}{z^2} + \frac{\varphi_2}{z^3} + \dots + \frac{\varphi_{N-2}}{z^{N-1}}$$

of its asymptotic series. Watson showed a century ago that the function is Borel regular whenever there's a constant  $c \in (0, \infty)$  with

$$|R_N| \leq \frac{c^{N+1} N!}{|z|^N}$$

over all orders  $N$  and all  $z$  in a wide enough wedge around infinity.

- **Satisfying a singular differential equation**

- Think about conditions where this works.
- Maybe the correct place is the setting of Ecalle's formal integral. See §5.2.2.1 of Delabaere's *Divergent Series, Summability and Resurgence III*.
- Say there's a unique solution (up to scaling) that shrinks as you go right; everything else blows up exponentially. Then this is the only solution that can be expressed as a Laplace transform.

- If the Borel-transformed equation has a subexponential solution  $\hat{f}$  which is “shifted holomorphic” (we called this having a “fractional power singularity” in **airy-resurgence**), then  $\mathcal{L}\hat{f}$  satisfies the original equation, because there are no boundary terms.
- Draw diagram showing formal vs. holomorphic solutions in time vs. frequency domains.

• **Being a thimble integral**

Let  $X$  be a translation surface—a Riemann surface carrying a holomorphic 1-form  $\nu$ . Suppose  $X$  is of *meromorphic type*, meaning that we got it by puncturing a compact Riemann surface  $\bar{X}$  at finitely many points, and  $\nu$  has a pole at each puncture. A *translation coordinate* on  $X$  is a local coordinate whose derivative is  $\nu$ .

Take another meromorphic-type translation surface  $B$  and a holomorphic Morse<sup>1</sup> map  $f: \bar{X} \rightarrow \bar{B}$  that sends punctures to punctures. Suppose every singularity of  $B$  is a critical value of  $f$ . [Typical usage of “Borel plane” seems ambiguous, so maybe we can use “Borel plane” for  $B$  and “Borel cover” for the Riemann surface of the Borel-transformed series. How to handle the OrrSommerfeld functions (DLMF §9.13)? We know  $f = 4u^3 - 3u$  is the pullback of a translation coordinate, but we also need a puncture at  $f(0)\dots$ ] For each critical point  $p$ , let  $\Gamma_p$  be the ray going rightward from  $f(p)$ , and let  $\zeta_p$  be the translation coordinate around  $\Gamma_p$  which vanishes at  $f(p)$ . These are well-defined as long as  $\Gamma_p$  misses the critical values of  $f$ . The preimage  $f^{-1}(\Gamma_p)$  is a bunch of disjoint curves, as long as  $\Gamma_p$  misses the other critical values of  $f$ . The *Lefschetz thimble*  $\Lambda_p$  is the component of  $f^{-1}(\Gamma_p)$  that goes through  $p$ , oriented so that shifting it to its left would make its projection run clockwise around  $\Gamma_p$ . The *thimble integral*

$$I_p = \int_{\Lambda_p} e^{-zf^*\zeta_p} \nu$$

is a holomorphic function on the right half-plane parametrized by  $z$ , and it turns out [we hope] to be Borel regular.

[Talk about exponential integrals and their decomposition into thimble integrals.]

In higher-dimensional complex manifolds, integrals over Lefschetz thimbles are still Borel regular [“Exponential integrals, Lefschetz thimbles and linear resurgence”][“Exponential Integral” lectures?]. This fact plays an important technical role in quantum mechanics, where infinite-dimensional exponential integrals are supposed to give the expectation values of observable quantities. Physicists often use Borel summation and related techniques to assign values to these integrals [Costin & Kruskal, “On optimal truncation...”].

Choose a path  $\gamma: \mathbb{R} \rightarrow X$  whose projection  $f \circ \gamma$  starts out going leftward out of a puncture, ends up going rightward into a puncture, and never touches a critical value of  $f$ . Choose a translation coordinate  $\zeta$  on  $B$  and continue it along  $f \circ \gamma$ , noting that

---

<sup>1</sup>This condition means that the critical points of  $f$  are isolated (the compactness of  $\bar{X}$  guarantees this) and the 2-jet of  $f$  is non-zero at every critical point.

it may become multi-valued if  $f \circ \gamma$  intersects itself. This data defines the *exponential integral*

$$I = \int_{\gamma} e^{-zf^*\zeta} \nu,$$

a holomorphic function on the right half-plane parametrized by  $z$ . It turns out [**we hope**] that we can get  $I$  by summing  $e^{-\alpha_p z} I_p$  over various critical points—as long as none of the  $\Gamma_p$  run into each other. [**We get jumps at phases where the  $\Gamma_p$  do hit each other.**] The constants  $\alpha_p$  are values of  $\zeta$ , continued to the critical points along certain paths.

- Each resummation method for asymptotic series makes some implicit assumption that allows us to reconstruct a holomorphic function from its asymptotic behaviour.
- The resummation method works correctly for functions which satisfy that assumption.
- For the modified Bessel function  $K_{1/3}$ , Borel resummation works because the asymptotic series encodes a second-order differential equation.
  - Different aspects of this example appear in various places (Mariño, Kawai–Takei, Sauzin). We give a detailed, unified treatment.
- We can generalize this argument to all  $K_{1/n}$  and their limit  $K_0$ .
- We can also generalize to all third-order exponential integrals.
  - Most of them are equivalent to the  $K_{1/3}$  integral, but there’s also an interesting degeneration.

## 1.2 Fractional derivative formula

- Theorem ?? says that for a certain class of exponential integrals

$$I(z) = \int_{\Gamma} e^{-zf} \nu,$$

the inverse Laplace [**better to say Borel?**] transform is the  $\frac{3}{2}$  derivative of  $d\zeta/df$ , where  $f^*d\zeta = \nu$  [**check**].

- the asymptotic expansion of  $I(z)$  is a resurgent function.
- Is it always a *simple* resurgent function?
  - **Maxim believes it is in general, and indeed in our examples we get simple resurgent functions. But how to prove it in general?**

### 1.3 Stokes phenomenon

- For Bessel functions, we can see explicitly how solutions jump when the Laplace transform angle crosses a critical value.
- The jump comes from the branch cut difference identity for hypergeometric functions.
- Possible interpretation of the Stokes factors as intersection numbers in Morse–Novikov theory [ask Maxim]

## 2 The Laplace and Borel transforms

### 2.1 The Laplace transform

- Action on differential equations.
  - Can we find a way to prove this when the differential operator spits out a function that's not integrable around zero?
- Global picture?

### 2.2 The Borel transform

- Action on differential equations.
  - No inhomogeneous terms! How is this consistent with the Laplace transform's action? Is there always an inhomogeneous solution with subexponential asymptotics?

## 3 Third-order exponential integrals

- Reduce to

$$I(z) = \int \exp[-z(u^3 + pu + q)] du$$

using change of coordinate.

- When  $p \neq 0$ , can reduce further to

$$I(z) = p^{1/2} e^{-qz} K_{1/3}(p^{3/2}z).$$

- As  $p$  goes to zero,  $I(z)$  degenerates to

$$\left(\frac{1}{2}\right)^{2/3} e^{-qz} \Gamma\left(\frac{1}{3}\right) z^{-1/3} = \left(\frac{1}{2}\right)^{2/3} e^{-qz} \mathcal{L}_{\zeta,0}(\zeta^{-2/3}) = \left(\frac{1}{2}\right)^{2/3} \mathcal{L}_{\zeta-q,q}(\zeta^{-2/3}).$$

#### Veronica's proposal

#### Title: Borel regularity and Resurgence of Exponential Integrals

1. what are exponential integrals? Geometrically they represent a Poincaré pairing (as explained by Kontsevich in **IHES lectures**). In physics ?? has to be done

2. what does it mean being Borel regular? and when does it happen? has to be done
  - Watson condition
  - linear, homogeneous ODE with regular singularity at 0 and irregular singularity at infinity [P.Ramis ?, Loday-Richaud ?, big idea in **airy-resurgence**]
  - our theorem [**draft2**]
3. Resurgence of exponential integrals and Stokes phenomena
  - Thimbles integrals [Kontsevich]: geometric computation of Stokes constants has to be done
  - ODE and fractional derivative formula [**draft2**]: hypergeometric functions has to be done
4. The Airy example
  - $I(z)$  is a solution of a linear ODE. We explicitly find its Borel transform, knowing the nature of singularities and the asymptotic behaviour of a basis of solution for the ODE [**airy-resurgence**]
  - Stokes constants using fractional derivative formula and Borel transform computation [**draft2**]
  - Comparison with the literature (Marino, Sauzin, Kontsevich slides) has to be done
5. Bessel  $1/n$  (follows from Airy) we compute Borel transform and Stokes constants has to be done
6. Bessel 0 (it is different because we have infinite cover) we compute Stokes constants [**draft2**]
7. Bessel  $\mu$  (follows from Bessel 0) we compute Stokes constants [**modified Bessel**]