

# Airy function: Kawai+Takei vs. Mariño

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Kawai and Takei want to solve

$$\left[ \left( \frac{d}{dx} \right)^2 - \eta^2 x \right] \psi(x, \eta) = 0.$$

They define  $\psi_B(x, y)$  as the inverse Laplace transform of  $\psi(x, \eta)$  with respect to  $\eta$ .

With  $w = x\eta^{2/3}$ , the equation above is equivalent to

$$\left[ \left( \frac{d}{dw} \right)^2 - w \right] \psi(w\eta^{-2/3}, \eta) = 0.$$

Proof: substitute back to get

$$\begin{aligned} \left[ \eta^{-4/3} \left( \frac{d}{dx} \right)^2 - \eta^{2/3} x \right] \psi(x, \eta) &= 0 \\ \left[ \eta^{-4/3} \left( \frac{d}{dx} \right)^2 - \eta^{-4/3} \eta^2 x \right] \psi(x, \eta) &= 0 \\ \eta^{-4/3} \left[ \left( \frac{d}{dx} \right)^2 - \eta^2 x \right] \psi(x, \eta) &= 0. \end{aligned}$$

Hence,  $\psi(w\eta^{-2/3}, \eta) = k(\eta)\text{Ai}(w)$  is a solution for any holomorphic function  $k$ .

## 1 Veronica's change of coordinates

Kawai and Takei study the WKB analysis of the equation

$$\left[ \left( \frac{d}{dx} \right)^2 - \eta^2 x \right] \psi(x, \eta) = 0 \tag{1}$$

as  $\eta \rightarrow \infty$ . They define  $\psi_B(x, y)$  as the inverse Laplace transform of  $\psi(x, \eta)$  with respect to  $\eta$ . In the coordinates  $t = yx^{-3/2}$  they find an explicit formula for  $\psi_B(x, y)$  in terms of

Gauss hypergeometric functions:

$$\begin{aligned}\psi_{+,B}(x,y) &= \frac{1}{x}\phi_+(t) = \frac{\sqrt{3}}{2\sqrt{\pi}} \frac{1}{x} s^{-1/2} {}_2F_1\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; s\right) \\ \psi_{-,B}(x,y) &= \frac{1}{x}\phi_-(t) = \frac{\sqrt{3}}{2\sqrt{\pi}} \frac{1}{x} (1-s)^{-1/2} {}_2F_1\left(\frac{1}{6}, \frac{5}{6}; \frac{1}{2}; 1-s\right)\end{aligned}$$

where  $s = 3t/4 + 1/2$ . The same hypergeometric functions have been computed in Section ?? as the Borel transform of the formal solutions of the Airy equation

$$\left[\left(\frac{d}{dw}\right)^2 - w\right]f(w) = 0. \quad (2)$$

Although the two equations look closely related (they are equivalent by the change of coordinates  $w = x\eta^{2/3}$ ), the Borel transform of  $\psi$  is computed with respect to  $\eta x^{3/2}$  (which is the conjugate variable of  $t$ ) while the Borel transform of  $f(w)$  is computed with respect to  $w$ . So we need to find a different change of coordinates to explain why the Borel transforms of  $\psi(x, \eta)$  and  $f(w)$  are given by the same hypergeometric function.

First of all notice that if  $\eta$  and  $y$  are conjugate variables under Borel transform, meaning

$$\sum_{n \geq 0} a_n \eta^{-n-1} \xrightarrow{\mathcal{B}} \sum_{n \geq 0} \frac{a_n}{n!} y^n$$

then  $t = yx^{-3/2}$  is the conjugate variable of  $q = \eta x^{3/2}$  up to correction by a factor of  $x^{-3/2}$

$$\sum_{n \geq 0} a_n q^{-n-1} = \sum_{n \geq 0} a_n x^{-3/2(n+1)} \eta^{-n-1} \xrightarrow{\mathcal{B}} \sum_{n \geq 0} \frac{a_n x^{-3/2(n+1)}}{n!} y^n = x^{-3/2} \sum_{n \geq 0} \frac{a_n}{n!} t^n.$$

In addition,  $\psi_{B,\pm}(x,y) = \frac{1}{x}\phi_{\pm}(t)$ , therefore we expect that  $\psi(x, \eta) = x^{1/2}\Phi(q)$ . Assume that  $\psi(x, y)$  is a solution of (1), then  $\Phi(q)$  solves

$$\left[\left(\frac{d}{dx}\right)^2 + x^{-1}\frac{d}{dx} - \frac{1}{4}x^{-2} - \eta^2 x\right]\Phi(q) = 0 \quad (3)$$

*Proof.*

$$\begin{aligned}
& \left[ \left( \frac{d}{dx} \right)^2 - \eta^2 x \right] \psi(x, \eta) = 0 \\
& \left[ \left( \frac{d}{dx} \right)^2 - \eta^2 x \right] x^{1/2} \Phi(q) = 0 \\
& \frac{d}{dx} \left[ \frac{1}{2} x^{-1/2} \Phi + x^{1/2} \frac{d}{dx} \Phi \right] - \eta^2 x^{3/2} \Phi = 0 \\
& -\frac{1}{4} x^{-3/2} \Phi + \frac{1}{2} x^{-1/2} \frac{d}{dx} \Phi + \frac{1}{2} x^{-1/2} \frac{d}{dx} \Phi + x^{1/2} \left( \frac{d}{dx} \right)^2 \Phi - \eta^2 x^{3/2} \Phi = 0 \\
& \left[ x^{1/2} \left( \frac{d}{dx} \right)^2 + x^{-1/2} \frac{d}{dx} - \frac{1}{4} x^{-3/2} - \eta^2 x^{3/2} \right] \Phi = 0 \\
& \left[ \left( \frac{d}{dx} \right)^2 + x^{-1} \frac{d}{dx} - \frac{1}{4} x^{-2} - \eta^2 x \right] \Phi = 0
\end{aligned}$$

□

Now rewrite (3) in the coordinates  $q = \eta x^{3/2}$ :

$$\begin{aligned}
& \left[ \left( \frac{d}{dx} \right)^2 + x^{-1} \frac{d}{dx} - \frac{1}{4} x^{-2} - \eta^2 x \right] \Phi = 0 \\
& \left[ \frac{9}{4} \eta^2 x \left( \frac{d}{dq} \right)^2 + \frac{3}{4} \eta x^{-1/2} \frac{d}{dq} + x^{-1} \cdot \frac{3}{2} \eta x^{1/2} \frac{d}{dq} - \frac{1}{4} x^{-2} - \eta^2 x \right] \Phi = 0 \\
& \left[ \eta^2 x \left( \frac{d}{dq} \right)^2 + \frac{1}{3} \eta x^{-1/2} \frac{d}{dq} + \frac{2}{3} \eta x^{-1/2} \frac{d}{dq} - \frac{1}{9} x^{-2} - \frac{4}{9} \eta^2 x \right] \Phi = 0 \\
& \left[ \eta^2 \left( \frac{d}{dq} \right)^2 + \eta x^{-3/2} \frac{d}{dq} - \frac{1}{9} x^{-3} - \frac{4}{9} \eta^2 \right] \Phi = 0 \\
& \left[ \left( \frac{d}{dq} \right)^2 + \eta^{-1} x^{-3/2} \frac{d}{dq} - \frac{1}{9} \eta^{-2} x^{-3} - \frac{4}{9} \right] \Phi = 0 \\
& \left[ \left( \frac{d}{dq} \right)^2 + q^{-1} \frac{d}{dq} - \frac{1}{9} q^{-2} - \frac{4}{9} \right] \Phi = 0
\end{aligned}$$

therefore  $\Phi(q)$  is a solution of the transform Airy equation (see draft2).