# Resurgence of the Airy function and other exponential integrals

Veronica Fantini and Aaron Fenyes

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## 1 Introduction

#### 1.1 Why does Borel resummation work?

Borel resummation is a way of turning a formal power series

$$\varphi_{\bullet} = z^{\sigma} \left( \frac{\varphi_0}{z} + \frac{\varphi_1}{z^2} + \frac{\varphi_2}{z^3} + \frac{\varphi_3}{z^4} + \ldots \right),$$

with  $\sigma \in [0,1)$ , into a function which is asymptotic to  $\varphi_{\bullet}$  as  $z \to \infty$ . Different functions can be asymptotic to the same power series, and Borel resummation picks one of them, performing an implicit regularization [arXiv:1705.03071, or maybe arXiv:1412.6614]. When a function matches the Borel sum of its asymptotic series, we'll say it's *Borel regular*. Several familiar kinds of regularity imply Borel regularity, and shed light on why it occurs.

#### Having a good asymptotic approximation

Let  $R_N$  be the difference between a function and the partial sum

$$\frac{\varphi_0}{z} + \frac{\varphi_1}{z^2} + \frac{\varphi_2}{z^3} + \ldots + \frac{\varphi_{N-2}}{z^{N-1}}$$

of its asymptotic series. Watson showed a century ago that the function is Borel regular whenever there's a constant  $c \in (0, \infty)$  with

$$|R_N| \le \frac{c^{N+1}N!}{|z|^N}$$

over all orders N and all z in a wide enough wedge around infinity.

#### • Satisfying a singular differential equation

- Think about conditions where this works.
- Maybe the correct place is the setting of Ecalle's formal integral. See §5.2.2.1 of Delabaere's Divergent Series, Summability and Resurgence III.
- Say there's a unique solution (up to scaling) that shrinks as you go right; everything else blows up exponentially. Then this is the only solution that can be expressed as a Laplace transform.

- If the Borel-transformed equation has a subexponential solution  $\hat{f}$  which is "shifted holomorphic" (we called this having a "fractional power singularity" in airy-resugence), then  $\mathcal{L}\hat{f}$  satisfies the original equation, because there are no boundary terms.
- Draw diagram showing formal vs. holomorphic solutions in time vs. frequency domains.

#### • Being a thimble integral

Let X be a translation surface—a Riemann surface carrying a holomorphic 1-form  $\nu$ . Suppose X is of *meromorphic type*, meaning that we got it by puncturing a compact Riemann surface  $\bar{X}$  at finitely many points, and  $\nu$  has a pole at each puncture. A translation coordinate on X is a local coordinate whose derivative is  $\nu$ .

Take another meromorphic-type translation surface B and a holomorphic map  $f \colon \bar{X} \to \bar{B}$  that sends punctures to punctures. Suppose the critical points of f are all simple. For each critical point p, let  $\Gamma_p$  be the ray going rightward from f(p), and let  $\zeta_p$  be the translation coordinate around  $\Gamma_p$  which vanishes at f(p). These are well-defined as long as  $\Gamma_p$  misses the singularities of B. The preimage  $f^{-1}(\Gamma_p)$  is a bunch of disjoint curves, as long as  $\Gamma_p$  misses the other critical values of f. [Is B is the Borel plane? Can we assume all its singularities are critical values of f?] The Lefschetz thimble  $\Lambda_p$  is the component of  $f^{-1}(\Gamma_p)$  that goes through p, oriented so that shifting it to its left would make its projection run clockwise around  $\Gamma_p$ . The thimble integral

$$I_p = \int_{\Lambda_p} e^{-zf^*\zeta_p} \nu$$

is a holomorphic function on the right half-plane parameterized by z, and it turns out [we hope] to be Borel regular.

[Talk about exponential integrals and their decomposition into thimble integrals.]

In higher-dimensional complex manifolds, integrals over Lefschetz thimbles are still Borel regular ["Exponential integrals, Lefschetz thimbles and linear resurgence"] ["Exponential Integral" lectures?]. This fact plays an important technical role in quantum mechanics, where infinite-dimensional exponential integrals are supposed to give the expectation values of observable quantities. Physicists often use Borel summation and related techniques to assign values to these integrals [Costin & Kruskal, "On optimal truncation..."].

Choose a path  $\gamma \colon \mathbb{R} \to X$  whose projection  $f \circ \gamma$  starts out going leftward out of a puncture, ends up going rightward into a puncture, and never touches a critical value of f. [How do we get jumps at the critical values?] Choose a translation coordinate  $\zeta$  on B and continue it along  $f \circ \gamma$ , noting that it may become multi-valued if  $f \circ \gamma$  intersects itself. This data defines the *exponential integral* 

$$I = \int_{\gamma} e^{-zf^*\zeta} \nu,$$

a holomorphic function on the right half-plane parameterized by z. It turns out [we hope] that we can get I by summing  $e^{-\alpha_p z}I_p$  over various critical points. The constants  $\alpha_p$  are values of  $\zeta$ , continued to the critical points along certain paths.

- Each resummation method for asymptotic series makes some implicit assumption that allows us to reconstruct a holomorphic function from its asymptotic behavior.
- The resummation method works correctly for functions which satisfy that assumption.
- For the modified Bessel function  $K_{1/3}$ , Borel resummation works because the asymptotic series encodes a second-order differential equation.
  - Different aspects of this example appear in various places (Mariño, Kawai-Takei, Sauzin). We give a detailed, unified treatment.
- We can generalize this argument to all  $K_{1/n}$  and their limit  $K_0$ .
- We can also generalize to all third-order exponential integrals.
  - Most of them are equivalent to the  $K_{1/3}$  integral, but there's also an interesting degeneration.

## 1.2 Fractional derivative formula

• Theorem ?? says that for a certain class of exponential integrals

$$I(z) = \int_{\Gamma} e^{-zf} \ \nu,$$

the inverse Laplace [better to say Borel?] transform is the  $\frac{3}{2}$  derivative of  $d\zeta/df$ , where  $f^*d\zeta = \nu$  [check].

- the asymptotic expansion of I(z) is a resurgent function.
- Is it always a *simple* resurgent function?
  - Maxim belives it is in general, and indeed in our examples we get simple resurgent functions. But how to prove it in general?

#### 1.3 Stokes phenomenon

- For Bessel functions, we can see explicitly how solutions jump when the Laplace transform angle crosses a critical value.
- The jump comes from the branch cut difference identity for hypergeometric functions.
- Possible interpretation of the Stokes factors as intersections numbers in Morse–Novikov theory [ask Maxim]

# 2 The Laplace and Borel transforms

## 2.1 The Laplace transform

- Action on differential equations.
  - Can we find a way to prove this when the differential operator spits out a function that's not integrable around zero?
- Global picture?

## 2.2 The Borel transform

- Action on differential equations.
  - No inhomogeneous terms! How is this consistent with the Laplace transform's action? Is there always an inhomogeneous solution with subexponential asymptotics?

# 3 Third-order exponential integrals

• Reduce to

$$I(z) = \int \exp\left[-z(u^3 + pu + q)\right] du$$

using change of coordinate.

• When  $p \neq 0$ , can reduce further to

$$I(z) = p^{1/2}e^{-qz}K_{1/3}(p^{3/2}z).$$

• As p goes to zero, I(z) degenerates to

$$\left(\frac{1}{2}\right)^{2/3} e^{-qz} \Gamma\left(\frac{1}{3}\right) z^{-1/3} = \left(\frac{1}{2}\right)^{2/3} e^{-qz} \mathcal{L}_{\zeta,0}(\zeta^{-2/3}) = \left(\frac{1}{2}\right)^{2/3} \mathcal{L}_{\zeta_{-q},q}(\zeta^{-2/3}).$$