

FYS3150 - Project 3

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(Dated: 14. oktober 2020)

Abstract om du vil

I. INTRODUCTION

II. THEORY

A. Newton's laws and the motion of planets

The well-known Newtons second law reads

$$\sum \mathbf{F} = m \frac{d^2 \mathbf{r}}{dt^2}, \quad (1)$$

where $\sum \mathbf{F}$ is the sum of forces acting on an object, m is its mass and \mathbf{r} its position.

Newton's law of gravity says that the total gravitational force acting on an object i from all the other objects in a system is

$$F_{G,i} = -Gm_i \sum_{j \neq i} m_j \frac{\mathbf{r}_i - \mathbf{r}_j}{|\mathbf{r}_i - \mathbf{r}_j|^3} \quad (2)$$

where \mathbf{r}_i are the positions and m_i are the masses of the objects. G is the gravitational constant.

For the motions of planets, on which the gravitational force is the only force acting, the two equations 1 and 2 combined give us the differential equation that describes the motion of an object:

$$\frac{d^2 \mathbf{r}_i}{dt^2} = \frac{1}{m_i} F_{G,i} = -G \sum_{j \neq i} m_j \frac{\mathbf{r}_i - \mathbf{r}_j}{|\mathbf{r}_i - \mathbf{r}_j|^3} \quad (3)$$

B. Velocity Verlet

Classical problem of Newtonian mechanics often involve solving a set of two coupled first order differential equations, namely

$$\begin{aligned} \frac{dx}{dt} &= v \\ \frac{dv}{dt} &= a \end{aligned}$$

where x is position, v is velocity and a is acceleration. Doing a Taylor expansion around a point in time $t = t_0$ gives

$$\begin{aligned} x(t) &= x(t_0) + (t - t_0) \frac{dx}{dt}(t_0) + \frac{1}{2}(t - t_0)^2 \frac{d^2x}{dt^2}(t_0) + O(h^3) \\ &= x(t_0) + (t - t_0)v(t_0) + \frac{1}{2}(t - t_0)^2 a + O(h^3) \end{aligned}$$

An expression can then be found for position at a time $t = t_0 \pm h$:

$$x(t \pm h) = x(t_0) \pm hv(t_0) + \frac{1}{2}h^2a \pm O(h^3) \quad (4)$$

Discretizing equation 4 and letting $x_i = x(t)$, $x_{i+1} = x(t + h)$, $v_i = v(t)$ and $v_{i+1} = v(t + h)$, we get

$$\begin{aligned} x_{i+1} &\approx x_i + hv_i + \frac{h^2}{2}a_i \\ v_{i+1} &\approx v_i + ha_i + \frac{h^2}{2}\dot{a}_i, \end{aligned}$$

where \dot{a}_i is the derivative of the acceleration with respect to time. Similarly to the case of Forward Euler, doing a Taylor expansion of a gives after some manipulation

$$\dot{a}_i \approx \frac{a_{i+1} - a_i}{h} \quad (5)$$

Inserting this into the expressions we have for x_{i+1} and v_{i+1} gives us the equations that describe the Velocity Verlet method:

$$\begin{aligned} x_{i+1} &\approx x_i + hv_i + \frac{h^2}{2}a_i \\ v_{i+1} &\approx v_i + \frac{h}{2}(a_{i+1} + a_i) \end{aligned}$$

From equation 4 we see that the mathematical error in this approximation goes like $O(h^3)$.

C. Forward Euler

From equation 4 we see that by including only the two first terms in the Taylor expansion of $x(t)$ and $v(t)$ we get

$$\begin{aligned} x_{i+1} &\approx x_i + hv_i \\ v_{i+1} &\approx v_i + ha_i \end{aligned}$$

This is referred to as the Forward Euler method of solving differential equations. The mathematical error in this approximation goes like $O(h^2)$.

III. METHODS

A. Discretization of the equations of motion

Equation 3 is the equation of motion that describes the system of planets and stars in motion. We want to

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solve this equation numerically for all the planets in our solar system. In order to do this we must discretize the equations. We can rewrite the equation as a set of first

order differential equations using the velocity:

$$\begin{aligned}\frac{d\mathbf{r}}{dt} &= \mathbf{v} \\ \frac{d\mathbf{v}}{dt} &= -G \sum_{j \neq i} m_j \frac{\mathbf{r}_i - \mathbf{r}_j}{|\mathbf{r}_i - \mathbf{r}_j|^3}\end{aligned}$$

The Verlet is not a self-starting algorithm. Mainly applicable when we are not interested in the velocity Velocity Verlet is energy conserving, Forward Euler is not

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 - [3] Griffiths, D. J., Schroeter, D. F., 2018, *Introduction to Quantum Mechanics, Third edition*, s. 75
 - [4] Griffiths, D. J., Schroeter, D. F., 2018, *Introduction to Quantum Mechanics, Third edition*, s. 133