

FYS4150

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Abstract

Introduction

Method

The implicit schemes are differential equations that can be rewritten as a set of linear equations. The Euler forward, Euler backward and Crank-Nicolson scheme are all implicit and can therefore be written in term of linear equations after scaling:

Euler backwards:

$$u_t \approx \frac{u(x_i, t_j) - u(x_i, t_j - \Delta t)}{\Delta t}$$

$$u_{xx} \approx \frac{u(x_i + \Delta x, t_j) - 2u(x_i, t_j) + u(x_i - \Delta x, t_j)}{\Delta x^2}$$

One can scale the above equation by $\alpha = \Delta t / \Delta x^2$, so the equation only depends on one scaled variable such that:

$$u_{i,j-1} = -\alpha u_{i-1,j} + (1 + 2\alpha)u_{i,j} - \alpha u_{i+1,j}$$

Now the differential equation can be written as a set of linear equations with a matrix A times a vector V_j such that: $AV_j = V_{j-1}$, where A defined from the above differential equations:

$$A = \begin{bmatrix} 1 + 2\alpha & -\alpha & 0 & 0 & \cdots \\ -\alpha & 1 + 2\alpha & -\alpha & 0 & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & 0 & \cdots & -\alpha & 1 + 2\alpha \end{bmatrix}$$

It is now possible to find the previous vector V_{j-1} when we already know what V_j , which we already know due to our initial conditions. A more generalized equation can be written as:

$$A^{-1}(AV_j) = A^{-1}(V_{j-1})$$

and if we keep multiplying by A^{-1} we get the implicit scheme:

$$V_j = A^{-j}V_0$$

A very similar process can be applied to the Euler forward method, but this scheme is explicit:

$$u_t = \frac{u(x_i, t_j + \Delta t) - u(x_i, t_j)}{\Delta t}$$

$$u_{xx} \approx \frac{u(x_i + \Delta x, t_j) - 2u(x_i, t_j) + u(x_i - \Delta x, t_j)}{\Delta x^2}$$

$$u_{i,j-1} = \alpha u_{i-1,j} + (1 - 2\alpha)u_{i,j} + \alpha u_{i+1,j}$$

$$A = \begin{bmatrix} 1 - 2\alpha & \alpha & 0 & 0 & \cdots \\ \alpha & 1 - 2\alpha & \alpha & 0 & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & 0 & \cdots & \alpha & 1 - 2\alpha \end{bmatrix}$$

such that:

$$A^{-1}(AV_j) = A^{-1}(V_{j-1})$$

We generalize again and get:

$$V_j = A^{-j}V_0$$

The Crank-Nicolson scheme is is a combination of both implicit and explicit schemes, namely the Euler forward and Euler backward method.

$$\frac{\theta}{\Delta x^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) + \frac{1-\theta}{\Delta x^2}(u_{i+1,j-1} - 2u_{i,j-1} + u_{i-1,j-1}) = \frac{1}{\Delta t}(u_{i,j} - u_{i,j-1})$$

where θ determines whether the scheme is explicit when $\theta = 0$, or implicit when $\theta = 1$. However, it is when $\theta = 1/2$ where we have the actual Crank-Nicolson scheme which is stable for all Δx and Δt . To derive the Crank-Nicolson scheme we begin with the forward Euler method and Taylor expand $u(x, t + \delta t)$, $u(x + \delta x, t)$, $u(x - \delta x, t)$, $u(x + \delta x, t + \Delta t)$ and $u(x - \delta x, t + \Delta t)$ for $t + \Delta t/2$.

Again we scale the equation with $\alpha = \frac{\Delta t}{\Delta x^2}$ which results in the following equation:

$$-\alpha u_{i-1,j} + (2 + 2\alpha)u_{i,j} - \alpha u_{i+1,j} = \alpha u_{i-1,j-1} + (2 - 2\alpha)u_{i,j-1} + \alpha u_{i+1,j-1}$$

which can be rewritten as

$$(2I + \alpha B)V_j = (2I - \alpha B)V_{j-1}$$

$$V_j = (2I + \alpha B)^{-1}(2I - \alpha B)V_{j-1}$$

where I is the identity matrix and B is given by:

$$B = \begin{bmatrix} 2 & -1 & 0 & \cdots & 0 \\ -1 & 2 & -1 & 0 & \vdots \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ \vdots & 0 & \ddots & \ddots & -1 \\ 0 & \cdots & \cdots & -1 & 2 \end{bmatrix}$$

Results

The truncation errors and stability is calculated in the Taylor expansion and these values are shown in the table below:

Method	Truncation error	Stability for
Euler Forward	$\Delta x^2, \Delta t$	Δx^2 and Δt^2
Euler Backward	$\Delta x^2, \Delta t$	Δx^2 and Δt^2
Crank-Nicolson	$\Delta x^2, \Delta t^2$	$\frac{1}{2}\Delta x^2 \geq \Delta t$

Table 1: Truncation errors and stability for the three methods

Discussion

From Table 1

Concluding remarks

Reference list