

Simulating QCD Jet Production in e^+e^- annihilation

– concepts of Monte Carlo event generators for collider experiments –

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Intro Lecture – part I

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Methods & Physics of Monte-Carlo event generators for the simulation of high-energy scattering events

Particle Physics at High-Energy Colliders

- investigate nature of short-range phenomena by scattering experiments
- bring high-energetic particle beams to collision, e.g. e^+e^- , ep , or pp
- measure & analyse the resulting final states

→ **experimental challenge**: identify & measure all particles emerging from individual beam clashes, reconstruct energies & momenta, specific detector components for γ , e^\pm , μ^\pm , n , p , ..., E_T

→ **theoretical challenge**: confront measurements with hypotheses for the underlying physics

The HEP trinity

Theory

QFT in Lagrangian formulation

of quantum nature, non-deterministic

benchmark Standard Model \mathcal{L}_{SM}

hypothetical New Physics \mathcal{L}_{BSM}

Experiment

complex multi-component detector

ATLAS, CMS, LHCb, ...

reconstruction of individual events

degrading, upgrades, improvements

Simulation: linking theory & experiment

multi-purpose Monte Carlo generators

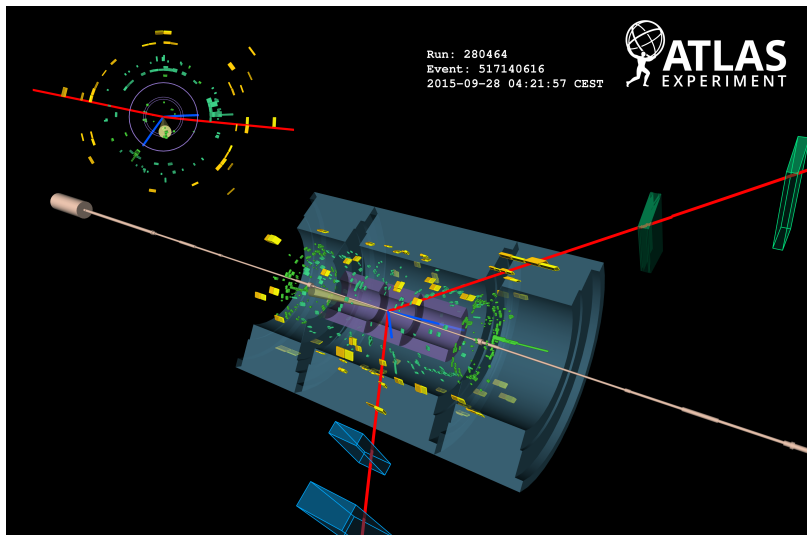
- stochastic simulation of particle-level events
- HERWIG, PYTHIA, SHERPA, ...

detector simulation codes

- detailed models for all detector components
- Geant4 based, Neural Network approaches, ...

provide fully-simulated event samples

Particle detectors & Real collider data



ATLAS event display: $pp \rightarrow H \rightarrow ZZ^* \rightarrow 2e2\mu + X$ event

Monte Carlo Event Generators: the global picture

- perturbative methods

- **Hard interaction**

- exact matrix elements $|\mathcal{M}|^2$
LO, NLO, NNLO – QCD, NLO – EW

- **Radiative corrections**

- parton showers in the initial and final state
resummation of soft-collinear logs: LL, NLL

- non-perturbative models

- **Multiple Interactions**

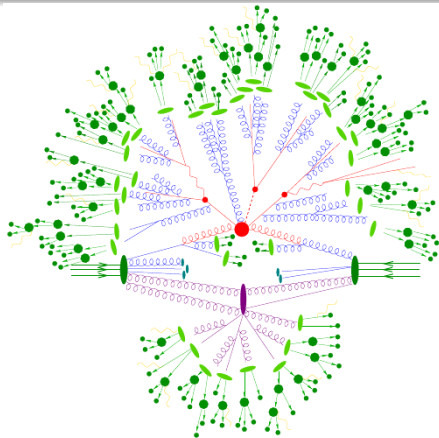
- beyond factorization: modelling

- **Hadronization**

- parton-hadron transition

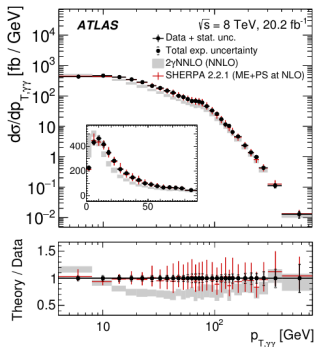
- **Hadron Decays**

- phase space or effective theories

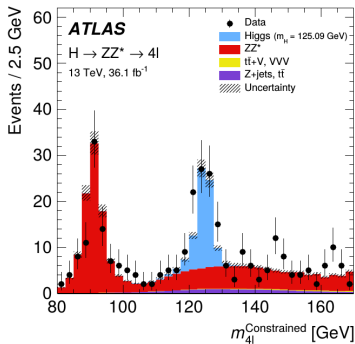


- predict fully exclusive final states
- factorize short & long range physics
- subject to approximations/compromises
- variety of specialized methods & tools

Di-Photon production



Higgs-Boson production

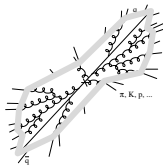
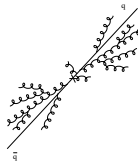
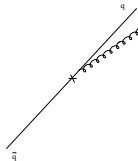
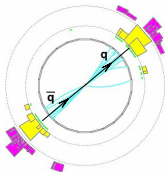


- confronting experiment with Standard Model theory predictions
- precision Standard Model measurements and searches for New Physics

Outline of the project

physics issues related to the modelling of high-energy collisions

- **Part I:** hard process generation, here $e^+e^- \rightarrow q\bar{q}$ at lowest order
 - integrate cross section for $e^+e^- \rightarrow q\bar{q}$
 - generate corresponding scattering events fully differentially
- **Part II:** final-state parton showering, *i.e.* soft- & collinear gluon emissions
 - simulate $q\bar{q}$ -initiated final-state parton cascade
 - analyse QCD jet observables, compare to reference data

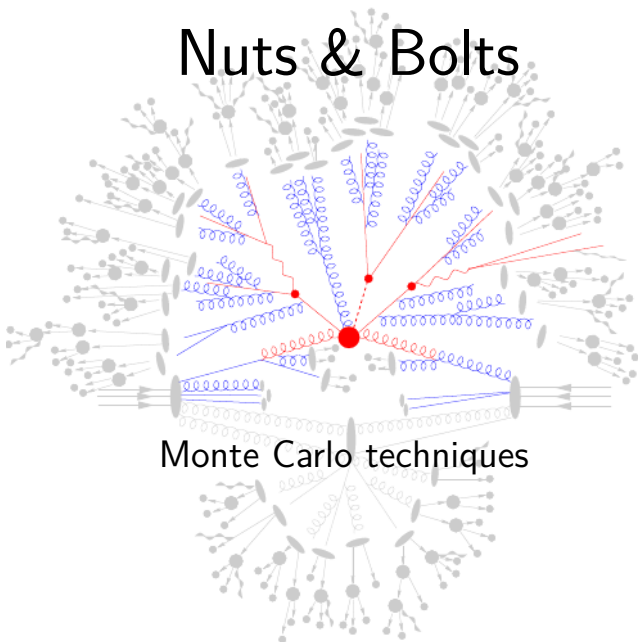


aspects of numerical methods/solutions used to tackle these

- Monte Carlo Importance Sampling
- Markov Chain parton-branching simulation
- sequential jet reconstruction algorithms



Nuts & Bolts



Monte-Carlo techniques

“Spatial” problems: without memory

- What is the volume of a given body?
Pick a point at random, use uniform probability in this area.
- What is the integrated cross section of a given process?
Pick an event at random, according to the differential cross section.

“Temporal” problems: have memory

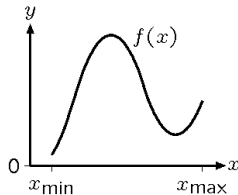
- Financial systems: What is the probability for a stock to have price X at time t , given the price Y at t_0 ?
- Parton shower: What is the probability for a parton to branch at a scale Q , given that it was created at a scale Q_0 ?

In particle physics often combined problems

What is the probability for a parton to branch at Q , with the daughters sharing the mother momentum in some specific way?

Monte-Carlo techniques: “Spatial” problems

Assume function $f(x)$ in range $x_{\min} \leq x \leq x_{\max}$,
where $f(x) \geq 0$ everywhere
(in practise x is multi-dimensional)



Two standard tasks

- 1 Calculate definite integral (approximatively)

$$I = \int_{x_{\min}}^{x_{\max}} f(x') dx' = (x_{\max} - x_{\min}) \langle f(x) \rangle$$

$$\simeq I_N = (x_{\max} - x_{\min}) \frac{1}{N} \sum_{i=1}^N f(x_i) \pm (x_{\max} - x_{\min}) \sqrt{\frac{\langle f^2 \rangle - \langle f \rangle^2}{N}}$$

random points $\in [x_{\min}, x_{\max}]$

↙

↪ error scales with $1/\sqrt{N}$, independent of d

↪ estimate improves by adding single point to the current estimate

- 2 Select x at random according to $f(x)$

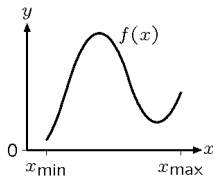
Monte-Carlo techniques: Basics

Select x according to $f(x)$

$$\text{as } P(x) \sim f(x) = \frac{f(x)}{\int_0^{f(x)} dy}$$

equivalent to uniform selection of (x, y) in area
 $x_{\min} \leq x \leq x_{\max}$ & $0 \leq y \leq f(x)$

$$\leadsto \int_{x_{\min}}^x f(x') dx' = \# \int_{x_{\min}}^{x_{\max}} f(x') dx'$$

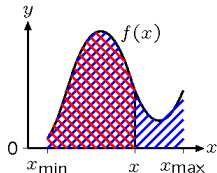


Analytical solution: inverse transform

known primitive function $F(x)$ and inverse F^{-1}

$$F(x) - F(x_{\min}) = \#(F(x_{\max}) - F(x_{\min})) = \#A_{\text{tot}}$$

$$\leadsto x = F^{-1}(F(x_{\min}) + \#A_{\text{tot}})$$



Example: $f(x) = e^{-x}$ for $x \geq 0$

$$F(x) = 1 - e^{-x}$$

$$1 - e^{-x} = \#1 \leadsto x = -\ln(\#)$$

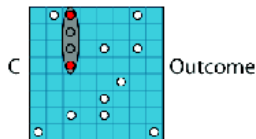
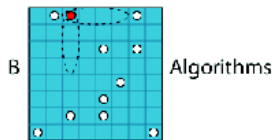
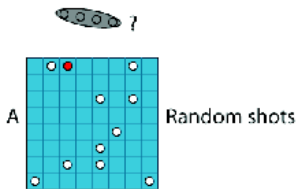
Monte-Carlo techniques: Variance Reduction

Variance reduction techniques: Motivation

- MC error estimate scales as $1/\sqrt{N}$, independent of dimension
- second limiting factor, the intrinsic variance of the integrand, i.e. V_N
 \leadsto devise techniques to reduce variance, thus enhancing MC convergence:
stratified sampling, control variates, importance sampling

Example: Battleships

Use of (cleverly chosen) random points.



Monte-Carlo techniques: Variance Reduction

Importance Sampling

- large variation of $f(x)$ leads to large uncertainty on MC estimate
- MC most efficient when each point x_i has nearly same weight $f(x_i)$
- arrange by choosing large number of points where function is largest
- compensate for overpopulation by reducing function value accordingly
- mathematically, IS corresponds to change of integration variable

$$\int f(x') dx' = \int \frac{f(x')}{g(x')} g(x') dx' = \int \frac{f(x')}{g(x')} dG(x')$$

$$\text{with } g(x) = \frac{dG}{dx} = \frac{\partial^d}{\partial x_1 \dots \partial x_d} G(x)$$

↪ points chosen according to $G(x)$ instead of uniformly

↪ f weighted inversely by $g(x) = dG/dx$

↪ relevant variance is $V(f/g)$, small if $g(x)$ & $f(x)$ similar in shape

Monte-Carlo techniques: Variance Reduction

Importance Sampling cont'd

conditions on the mapping function $g(x)$

- (i) $g(x)$ is probability density: $g(x) \geq 0$ and $\int g(x')dx' = 1$
- (ii) $G(x)$, the integral of $g(x)$, is known analytically
(integrated distribution, increasing monotonically with x)
- (iii) either function $G(x)$ can be inverted (solved) for x analytically or, alternatively, a g -distributed random-number generator is available
- (iv) the ratio $f(x)/g(x)$ is ideally nearly constant
 \leadsto variance $V(f/g)$ will be small compared to $V(f)$

$$E_N = \frac{1}{N} \sum_{i=1}^N \frac{f(x_i)}{g(x_i)}, \quad V_N \left(\frac{f}{g} \right) = \frac{1}{N} \sum_{i=1}^N \left(\frac{f(x_i)}{g(x_i)} \right)^2 - E_N^2$$

Importance Sampling: disadvantages

dangerous if function $g(x)$ approaches zero where f is non-vanishing, then $V(f/g)$ becomes infinite and usual techniques to estimate the variance from sample points may not detect this, if region of $g = 0$ is small

Monte-Carlo techniques: Variance Reduction

Breit-Wigner propagator – particle with mass M and decay width Γ

$$I = \int_{s_{\min}}^{s_{\max}} \frac{ds}{(s - M^2)^2 + M^2\Gamma^2}$$

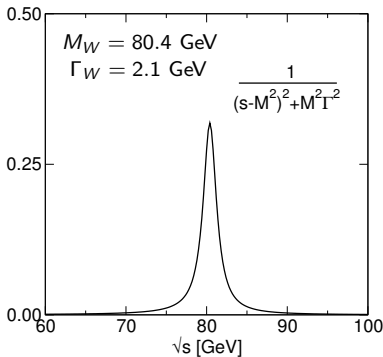
use mapping (variable transform)

$$s = M\Gamma \tan \rho + M^2$$

$$\leadsto ds = M\Gamma \sec^2 \rho d\rho$$

such that

$$\begin{aligned} I &= \int_{\rho_{\min}}^{\rho_{\max}} \frac{M\Gamma \sec^2 \rho d\rho}{M^2\Gamma^2 \tan^2 \rho + M^2\Gamma^2} \\ &= \frac{1}{M\Gamma} \int_{\rho_{\min}}^{\rho_{\max}} d\rho \leadsto \delta I = 0 \end{aligned}$$



\leadsto pick ρ uniformly from

$$\begin{aligned} \rho &= \rho_{\min} + \#(\rho_{\max} - \rho_{\min}) \\ \rho_{\min/\max} &= \arctan \left(\frac{s_{\min/\max} - M^2}{M\Gamma} \right) \end{aligned}$$

Monte-Carlo techniques: Hit-or-Miss method

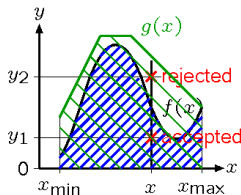
Sampling according to $f(x)$

known $g(x)$ as local overestimate for target $f(x)$:

$g(x) \geq f(x)$ in $x_{\min} \leq x \leq x_{\max}$

and $G(x) = \int_{x_{\min}}^x g(x')dx'$ & $G^{-1}(y)$ are simple

- 1 select $x =$ according to distribution $g(x)$
- 2 select $y = \#g(x)$
- 3 if $y \leq f(x)$ accept point [$N_{\text{acc}}++$, $N_{\text{try}}++$]
if $y > f(x)$ reject point [$N_{\text{fail}}++$, $N_{\text{try}}++$]



Monte-Carlo techniques: Hit-or-Miss method

Example: $f(x) = xe^{-x}$ and $g(x) = Ne^{-x/2}$

ensure that $g(x)$ is local overestimate:

$$\frac{f(x)}{g(x)} = \frac{xe^{-x/2}}{N} \stackrel{!}{\leq} 1 \rightarrow \text{find maximum}$$

$$\frac{d}{dx} \left(\frac{f(x)}{g(x)} \right) = \frac{1}{N} \left(1 - \frac{x}{2} \right) e^{-x/2} \stackrel{!}{=} 0 \leadsto x = 2$$

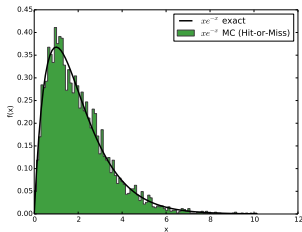
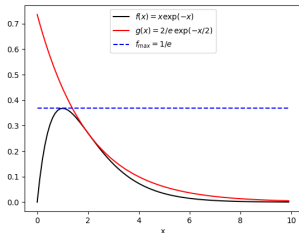
normalize such that $g(2) = f(2) \leadsto N = 2/e$

from $G(x) \sim 1 - e^{-x/2} = \#_1$

$$\leadsto x = -2 \ln(\#_1), \text{ red } y = \#_2 g(x) = \#_2 2e^{-(1+x/2)}$$

\leadsto efficiency IS: $N_{\text{acc}}/N_{\text{try}} = 68.4\%$

\leadsto efficiency $f_{\text{max}} = e^{-1}$: $N_{\text{acc}}/N_{\text{try}} = 13.5\%$



Monte-Carlo techniques: “Temporal” Problems

The radioactive decay problem

- known probability $f(t) \geq 0$ that *something will happen* at time t
[nucleus decays, parton branches, transistor fails]
- *something happens* at t **only** if it didn't happen at $t' < t$

Monte-Carlo techniques: “Temporal” Problems

The radioactive decay problem

- known probability $f(t) \geq 0$ that *something will happen* at time t
[nucleus decays, parton branches, transistor fails]
- *something happens* at t **only** if it didn't happen at $t' < t$

Define: $N(t)$: probability that *nothing* happend until t [$N(0) = 1$]

$P(t) = -dN(t)/dt = f(t)N(t)$: probability for decay at time t

Monte-Carlo techniques: “Temporal” Problems

The radioactive decay problem

- known probability $f(t) \geq 0$ that *something will happen* at time t
[nucleus decays, parton branches, transistor fails]
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Define: $N(t)$: probability that *nothing* happend until t [$N(0) = 1$]

$P(t) = -dN(t)/dt = f(t)N(t)$: probability for decay at time t

Solution:

$$N(t) = \exp \left\{ - \int_0^t f(t') dt' \right\} \quad [f(t) = \lambda \text{ simple radioactive decay}]$$

$$P(t) = f(t) \exp \left\{ - \int_0^t f(t') dt' \right\}$$

- naïve answer $P(t) = f(t)$ modified by exponential suppression
[$P(t) \approx f(t)$ at small t only, otherwise damping]
- in parton-shower picture, $N(t)$ is called the **Sudakov form factor**

Monte-Carlo techniques: “Temporal” Problems

The radioactive decay problem: solution

- assume $f(t)$ has primitive function $F(t)$ with known inverse F^{-1}
- standard solution to find “decay” time t :

$$\int_0^t P(t') dt' = N(0) - N(t) = 1 - \exp \{-(F(t) - F(0))\} = 1 - \#$$

$$\exp \{-(F(t) - F(0))\} = \#$$

$$\leadsto t = F^{-1}(F(0) - \ln(\#))$$

- simple radioactive decay, i.e. $f(t) = \lambda$

$$F(t) = \lambda t = -\ln(\#) \leadsto t = -\ln(\#)/\lambda$$

- if $f(t)$ not sufficiently “nice”, use importance sampling & Hit-or-Miss
- however, simple condition $f(t) \leq g(t) \quad \forall \quad t \geq 0$ is not sufficient

$$\leadsto \hat{P}(t) = g(t) \exp \left\{ - \int_0^t g(t') dt' \right\}$$

\leadsto does not yield proper exponential factor

Monte-Carlo techniques: “Temporal” Problems

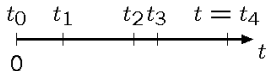
The Veto Algorithm [Seymour 1995]

if $f(t)$ has no simple $F(t)$ or F^{-1} : use “nice” $g(t) \geq f(t)$

- ❶ start with $i = 0$ and $t_0 = 0$
- ❷ increment i and select $t_i = G^{-1}(G(t_{i-1}) - \ln(\#_1))$
[according to $g(t)$, but with constraint $t_i > t_{i-1}$]
- ❸ if $f(t_i)/g(t_i) \leq \#_2$ go back to step 2
otherwise, keep t_i as “decay” time

↪ next step only depends on the very previous one

↪ Markovian process



Monte-Carlo techniques: The Veto Algorithm

Iterative Proof

consider the various ways in which a specific time t can be selected

- probability that first try is accepted, i.e. $t = t_1$, such that no intermediate t values need to be rejected

$$P_0^g(t) = \underbrace{\exp \left\{ - \int_0^t g(t') dt' \right\} g(t)}_{P^g(t)} \times \underbrace{\frac{f(t)}{g(t)}}_{\text{prob. to accept } t}$$

- consider case, where one intermediate step, t_1 , got rejected, only second step, $t = t_2$, is accepted, i.e. $0 \leq t_1 \leq t_2 = t$

$$P_1^g(t) = \underbrace{\int_0^t \exp \left\{ - \int_0^{t_1} g(t') dt' \right\} g(t_1) \left(1 - \frac{f(t_1)}{g(t_1)} \right)}_{\text{prob. to select } t_1 < t} \times \underbrace{\exp \left\{ - \int_{t_1}^t g(t'') dt'' \right\} g(t) \frac{f(t)}{g(t)}}_{\text{prob. to select \& accept } t = t_2 > t_1} dt_1$$

$$\leadsto P_1^g(t) = P_0^g(t) \int_0^t (g(t_1) - f(t_1)) dt_1$$

Monte-Carlo techniques: The Veto Algorithm

Iterative Proof cont'd

- considering arbitrary intermediate steps the probability to accept t yields

$$\begin{aligned}P^g(t) &= \sum_{i=0}^{\infty} P_i^g(t) \\&= P_0^g(t) \sum_{i=0}^{\infty} \frac{1}{i!} \left(\int_0^t (g(t') - f(t')) dt' \right)^i \\&= f(t) \exp \left\{ - \int_0^t g(t') dt' \right\} \exp \left\{ \int_0^t (g(t') - f(t')) dt' \right\} \\&= f(t) \exp \left\{ - \int_0^t f(t') dt' \right\} \\&= P(t) \quad \text{q.e.d.}\end{aligned}$$

Monte-Carlo techniques: The Veto Algorithm

Discussion

- Veto Algorithm standard tool in event generation, e.g. QCD parton showers
- numerical implementation of true Markovian decay/emission process
- process can be stopped at any intermediate scale/time t_{\max}
- usually $f(t)$ function of additional variables x , however, easy to generalise
 - find suitable $g(t, x)$ with $f(t, x) \leq g(t, x)$, the $g(t)$ used in the veto algorithm is then simply given by the integral

$$g(t) = \int g(t, x') dx'$$

- each time a value t_i is selected also the x_i need to be picked, following density $g(t_i, x) dx$
- the point (t_i, x_i) gets accepted with probability $f(t_i, x_i)/g(t_i, x_i)$

Summary

- use Monte Carlo methods to perform integrals and sample distributions
 - need only few points to estimate $\int f(x')dx'$
 - each additional point increases accuracy
- techniques generalise to many dimensions
 - typical LHC phase space $\sim d^3\vec{p} \times 100$'s particles
 - error scales as $1/\sqrt{N}$ vs. $1/N^{2/d}$ or $1/N^{4/d}$ (Trapezoidal or Simpson's Rule)
- suitable for complicated integration regions
 - kinematic cuts or detector cracks
- can sample distributions where exact solutions cannot be found
- use of adaptive techniques to reduce variances
- Veto Algorithm applicable for parton-shower simulations

Questions/Self-Tests

- Implement the above discussed variance reduction example, *i.e.*
 $f(x) = xe^{-x}$ with $g(x) = 2e^{(-1+x/2)}$
- How to sample points $x \in [\epsilon, 1 - \epsilon]$ according to $f(x) = \frac{(1-x)}{x}$?
- Plot the Breit–Wigner distribution for the SM Higgs-boson propagator
($M_H = 125.1$ GeV, $\Gamma_H = 4.2$ MeV)
- Repeat the proof of the Veto Algorithm, see also backup

Literature/Further Reading

- Buckley et al. “General-purpose event generators for LHC physics”
Phys. Rept. **504** (2011), 145-233
arXiv:1101.2599 [hep-ph]
- Salam “Elements of QCD for hadron colliders”
CERN Yellow Rep. School Proc. **5** (2020), 1-56
arXiv:1011.5131 [hep-ph]
- Ellis, Stirling, Webber “QCD and collider physics”
Camb. Monogr. Part. Phys. Nucl. Phys. Cosmol. **8** (1996), 1-435

Backup slides

Monte-Carlo techniques: Variance Reduction

Breit-Wigner propagator – Jacobian

consider
$$\int_{s_{\min}}^{s_{\max}} ds f(s) \rightarrow \int_0^1 d\# \frac{1}{g_{\text{prop}}(s(\#))} f(s(\#))$$

1 Breit-Wigner mapping

$$\int_{s_{\min}}^{s_{\max}} ds = \int_{\rho_{\min}}^{\rho_{\max}} M\Gamma \sec^2 \rho d\rho$$

$$\text{with } \rho(\#) = \rho_{\min} + \#(\rho_{\max} - \rho_{\min})$$

$$\frac{d\rho}{d\#} = \rho_{\max} - \rho_{\min}$$

$$\int_{s_{\min}}^{s_{\max}} ds = \int_0^1 d\# \frac{(s_{\#} - M^2)^2 + M^2\Gamma^2}{M\Gamma} (\rho_{\max} - \rho_{\min})$$

$$\text{where } s_{\#} = s(\rho(\#)) = M\Gamma \tan(\rho(\#)) + M^2$$

Monte-Carlo techniques: The Veto Algorithm

Iterative Proof cont'd

- for $P_2^g(t)$ need to consider two intermediate times, i.e. $0 \leq t_1 \leq t_2 \leq t_3 = t$

$$\begin{aligned} P_2^g(t) &= P_0^g(t) \underbrace{\int_0^t dt_1 (g(t_1) - f(t_1)) \int_{t_1}^t dt_2 (g(t_2) - f(t_2))}_{\text{nested integrals}} \\ &= P_0^g(t) \frac{1}{2} \left(\int_0^t (g(t') - f(t')) dt' \right)^2 \end{aligned}$$

Monte-Carlo techniques: The Veto Algorithm

Iterative Proof cont'd

- for $P_2^g(t)$ need to consider two intermediate times, i.e. $0 \leq t_1 \leq t_2 \leq t_3 = t$

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Proof – nested integrals

to prove the last equality consider inclusion of the region where $t_2 < t_1$

$$\begin{aligned} &\int_0^t dt_1 (g(t_1) - f(t_1)) \int_{t_1}^t dt_2 (g(t_2) - f(t_2)) + \int_0^t dt_2 (g(t_2) - f(t_2)) \int_{t_2}^t dt_1 (g(t_1) - f(t_1)) \\ &= \int_0^t dt' (g(t') - f(t')) \left[2 \int_{t'}^t dt'' (g(t'') - f(t'')) \right] = \left(\int_0^t dt' (g(t') - f(t')) \right)^2 \end{aligned}$$

Monte-Carlo techniques: The Veto Algorithm

Iterative Proof cont'd

- for $P_2^g(t)$ need to consider two intermediate times, i.e. $0 \leq t_1 \leq t_2 \leq t_3 = t$

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Proof – nested integrals

generalisation to n factors

$$\frac{1}{n!} \int_0^t \left(\int_{t''}^t f(t') dt' \right)^n f(t'') dt'' = \frac{1}{(n+1)!} \left(\int_0^t f(t) dt \right)^{n+1}$$