A Brief Introduction to Quantum Graphs

Viktor Qvarfordt

2015-11-30

Table of Contents

- Defining Quantum graphs
 - Metric graphs
 - Differential operators
 - Matching conditions
- Properties of Quantum Graphs
 - First example
 - Scattering: Inverse problems
 - Star graph: Calculating the spectrum

Quantum Graph

Definition

A quantum graph is a triple (Γ, L, MC) , where

- Γ is a metric graph,
- L is a differential expression defined on functions on Γ , and
- MC (matching conditions) is a set of conditions relating the limit values of the functions defined on the edges of Γ .

Metric Graph

Definition

A metric graph is a tuple (E, V) where:

- $E = \{e_n\}_{n=1}^N$ is a set of intervals $(edges) [x_{2n-1}, x_{2n}] \subset \mathbb{R}$.
- $V = \{v_m\}_{m=1}^M$ is a partition of the set $\{x_j\}_{j=1}^{2N}$. Each element in V is called a *vertex*.

Metric Graph

Definition

A metric graph is a tuple (E, V) where:

- $E = \{e_n\}_{n=1}^N$ is a set of intervals $(edges) [x_{2n-1}, x_{2n}] \subset \mathbb{R}$.
- $V = \{v_m\}_{m=1}^M$ is a partition of the set $\{x_j\}_{j=1}^{2N}$. Each element in V is called a *vertex*.

Example

$$\begin{array}{ccc}
 & & & & & & & & & & \\
\hline
v_1 & & & & & & & & \\
V_1 & & & & & & & \\
V_2 & & & & & & \\
E & = \{e_1, e_2\} &= \{[x_1, x_2], [x_3, x_4]\} \\
V & = \{v_1, v_2, v_3\} &= \{\{x_1\}, \{x_2, x_3\}, \{x_4\}\} \\
\end{array}$$

Functions on Metric Graphs

Definition

Metric graph Γ , define the Hilbert space $L_2(\Gamma)$ by

$$L_2(\Gamma) = \bigoplus_{e \in F} L_2(e), \qquad \langle f, g \rangle_{L_2(\Gamma)}^2 = \int_{\Gamma} f(x) \overline{g}(x) \, dx.$$

Differential Expressions

Electric potential q(x) and magnetic potential a(x). Magnetic Schrödinger Operator

$$L_{q,a} = \left(i\frac{d}{dx} + a(x)\right)^2 + q(x).$$

Differential Expressions

Electric potential q(x) and magnetic potential a(x). Magnetic Schrödinger Operator

$$L_{q,a} = \left(i\frac{d}{dx} + a(x)\right)^2 + q(x).$$

Laplace Operator

$$L = L_{0,0} = -\frac{d^2}{dx^2}.$$

Differential Expressions

Electric potential q(x) and magnetic potential a(x). Magnetic Schrödinger Operator

$$L_{q,a} = \left(i\frac{d}{dx} + a(x)\right)^2 + q(x).$$

Laplace Operator

$$L = L_{0,0} = -\frac{d^2}{dx^2}.$$

The equation

$$Lu = \lambda u$$
.

Differential Operators

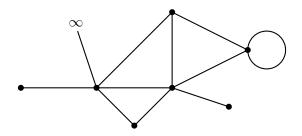
We are insterested in studying self-adjoint differential operators. This is a rather technical discussion, essentially one requires symmetry,

$$\langle Lu, v \rangle = \langle u, Lv \rangle.$$

This implies that the energy of the system is preserved, since a symmetric operator has real eigenvalues. Furthermore the corresponding eigenfunctions form an orthonormal basis.

Questions Arise

So far everything makes sense for single-line graphs. What about general graphs?



How to "connect" the equations at the vertices? Matching conditions.

Matching Conditions

Let us consider under which conditions a differential operator associated with the differential expression $L = -\frac{d^2}{dx^2}$ is symmetric.

$$\begin{split} \langle Lu, v \rangle_{L_{2}(\Gamma)} &= \sum_{e_{n} \in E} \int_{e_{n}} (-u'') \overline{v} \, dx \\ &= \sum_{e_{n} \in E} \left(\int_{e_{n}} u' \overline{v}' dx - \left[u' \overline{v} \right]_{x_{2n-1}}^{x_{2n}} \right) \\ &= \sum_{e_{n} \in E} \left(\int_{e_{n}} u(-\overline{v}'') dx + \left[u \overline{v}' - u' \overline{v} \right]_{x_{2n-1}}^{x_{2n}} \right) \\ &= \langle u, Lu \rangle_{L_{2}(\Gamma)} + \sum_{e_{n} \in E} \left[u \overline{v}' - u' \overline{v} \right]_{x_{2n-1}}^{x_{2n}} \\ &= \langle u, Lu \rangle_{L_{2}(\Gamma)} + \sum_{m=1}^{M} \sum_{x \in V_{m}} \partial u(x) \overline{v}(x) - u(x) \partial \overline{v}(x) \end{split}$$

Matching Conditions

That is, we require

$$0 = \langle Lu, v \rangle - \langle u, Lv \rangle = \sum_{m=1}^{M} \sum_{x_j \in v_m} \partial u(x_j) \overline{v}(x_j) - u(x_j) \partial \overline{v}(x_j).$$

Matching Conditions

That is, we require

$$0 = \langle Lu, v \rangle - \langle u, Lv \rangle = \sum_{m=1}^{M} \sum_{x_j \in v_m} \partial u(x_j) \overline{v}(x_j) - u(x_j) \partial \overline{v}(x_j).$$

Definition

The standard matching conditions, for a vertex v, are given by

$$\begin{cases} u(x_j) = u(x_i) \text{ for all } x_j, x_i \in v \\ \sum_{x_j \in v} \partial u(x_j) = 0. \end{cases}$$

Properties of Quantum Graphs

- Defining Quantum graphs
 - Metric graphs
 - Differential operators
 - Matching conditions
- Properties of Quantum Graphs
 - First example
 - Scattering: Inverse problems
 - Star graph: Calculating the spectrum

First Example

One-dimensional time-independent Schrödinger equation with no potential

$$-u''(x) = \lambda u(x) \iff u(x) = Ae^{\overrightarrow{ikx}} + Be^{-\overrightarrow{ikx}}, \quad k^2 = \lambda.$$

A and B are determined by the boundary conditions.

First Example

One-dimensional time-independent Schrödinger equation with no potential

$$-u''(x) = \lambda u(x) \iff u(x) = Ae^{\overrightarrow{ikx}} + Be^{-\overrightarrow{ikx}}, \quad k^2 = \lambda.$$

A and B are determined by the boundary conditions.

Example (Particle in a Box)

Infinite potential in the region $x < 0, x > \ell$ is modeled by allowing the particle to exist only on the interval $[0, \ell]$.

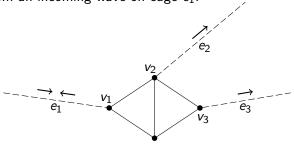
Figure : Finite line graph of length ℓ , parametrized from 0 to ℓ .

The boundary condition $u(0) = u(\ell) = 0$ (Dirichlet condition) gives

$$u(x) = A \sin kx$$
, $\lambda = k^2 = \frac{\pi^2}{\ell^2} n^2$, $n = 1, 2, ...$

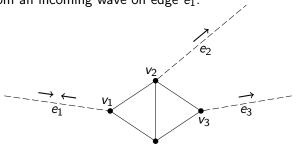
Scattering: Inverse problems

Scattering from an incoming wave on edge e_1 .

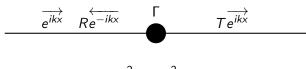


Scattering: Inverse problems

Scattering from an incoming wave on edge e_1 .



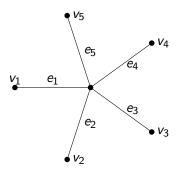
More generally,



$$|R|^2 + |T|^2 = 1$$

Defining the star graph

Consider a star graph with d edges, all having the same length ℓ and standard matching conditions at all vertices.



Viktor Qvarfordt Spectrum 2015-11-30 14 / 18

$$Lu_j = \lambda u_j \iff u_j(x) = A_j \cos kx + B_j \sin kx, \ j = 1, \dots, d, \ \lambda = k^2 > 0$$

Viktor Qvarfordt

$$Lu_j = \lambda u_j \iff u_j(x) = A_j \cos kx + B_j \sin kx, \ j = 1, \dots, d, \ \lambda = k^2 > 0$$

Parametrize every edge so that x=0 at the outer vertices and $x=\ell$ at the inner vertex. Matching conditions at the outer verices require the derivative to vanish,

$$u'_j(0) = kB_j\cos 0 = 0 \implies B_j = 0.$$

Viktor Qvarfordt spectrum 2015-11-30 15 / 18

$$Lu_j = \lambda u_j \iff u_j(x) = A_j \cos kx + B_j \sin kx, \ j = 1, \dots, d, \ \lambda = k^2 > 0$$

Parametrize every edge so that x=0 at the outer vertices and $x=\ell$ at the inner vertex. Matching conditions at the outer verices require the derivative to vanish,

$$u'_j(0) = kB_j\cos 0 = 0 \implies B_j = 0.$$

At the inner vertex, x = 0, the continuity condition require

$$A_j \cos k\ell = A_{j'} \cos k\ell$$
, for all $j \neq j'$.

Assume first $\cos k\ell \neq 0$, then

$$A_j = A$$
, for all j ,

for some constant A.

Viktor Qvarfordt spectrum 2015-11-30 15 / 18

The derivative condition require

$$\sum_{j=1}^{d} \partial u_{j}(\ell) = -dkA\sin(k\ell) = 0 \implies \sin k\ell = 0,$$

since $k \neq 0$ and $A \neq 0$ to avoid the trivial solution. This gives the possible eigenvalues λ_n with corresponding eigenfunctions f_n ;

$$\lambda_n = k_n^2 = \left(\frac{n\pi}{\ell}\right)^2, \quad n = 1, 2, \dots$$
 $f_n(x) = A\cos k_n x.$

Viktor Qvarfordt spectrum 2015-11-30 16 / 18

The derivative condition require

$$\sum_{j=1}^{d} \partial u_{j}(\ell) = -dkA\sin(k\ell) = 0 \implies \sin k\ell = 0,$$

since $k \neq 0$ and $A \neq 0$ to avoid the trivial solution. This gives the possible eigenvalues λ_n with corresponding eigenfunctions f_n ;

$$\lambda_n = k_n^2 = \left(\frac{n\pi}{\ell}\right)^2, \quad n = 1, 2, \dots$$
 $f_n(x) = A\cos k_n x.$

Next, if $\cos k\ell = 0$, the continuity condition for the inner vertex,

$$A_j \cos k\ell = A_{j'} \cos k\ell$$
, for all $j \neq j'$,

is trivially satisfied. Hence the constants A_j on each vertex need not coincide.

Viktor Qvarfordt spectrum 2015-11-30 16 / 18

The condition on the derivatives now reads

$$0 = \sum_{j=1}^d \partial u_j(\ell) = -k \sin(k\ell) \sum_{j=1}^d A_j \implies \sum_{j=1}^d A_j = 0,$$

since we are assuming $k \neq 0$ and $\cos k\ell = 0$. This equation has d-1 linearly independent solutions. This is a degenerate energy level, there are (d-1) eigenstates with the same energy, given by $\cos k\ell = 0$, that is

$$\lambda_n = k_n^2 = \left(\frac{\pi}{2\ell} + \frac{n\pi}{\ell}\right)^2, \quad n = 1, 2, \dots$$

Finally, one easily calculates the eigenfunction for $\lambda=0$, the ground state at zero energy, to be a constant function on the entire graph.

Viktor Qvarfordt spectrum 2015-11-30 17 / 18

To sum up, we can order all eigenvalues in an increasing sequence,

$$\lambda_0 = 0$$

$$\lambda_{dn} = \left(\frac{dn\pi}{\ell}\right)^2, \quad n = 1, 2, \dots$$

$$\lambda_{dn+j} = \left(\frac{\pi}{2\ell} + \frac{dn\pi}{\ell}\right)^2, \quad j = 1, \dots, (d-1), \quad n = 1, 2, \dots$$

Viktor Qvarfordt

To sum up, we can order all eigenvalues in an increasing sequence,

$$\lambda_0 = 0$$

$$\lambda_{dn} = \left(\frac{dn\pi}{\ell}\right)^2, \quad n = 1, 2, \dots$$

$$\lambda_{dn+j} = \left(\frac{\pi}{2\ell} + \frac{dn\pi}{\ell}\right)^2, \quad j = 1, \dots, (d-1), \quad n = 1, 2, \dots$$

Hence the m:th eigenvalue grows as

$$\lambda_m \sim \frac{\pi^2}{(d\ell)^2} m^2$$
 as $m \to \infty$.

Note that $d\ell$ is the total length of the graph. The growth of the eigenvalues is an example of the the Weyl asymptotic law.

Viktor Qvarfordt spectrum 2015-11-30 18 / 18