

II. Quantum Circuits

- "circuit model": sequence of building blocks that carry out elementary computations, called gates



Single qubit gates

- classical example: NOT $|1\rangle \rightarrow |0\rangle$
- quantum examples: as quantum theory is unitary, quantum gates are represented by

unitary matrices: $U^\dagger U = 11$

$$\text{recall Diac notation: } U = \begin{pmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{pmatrix} = u_{00}|0\rangle\langle 0| + u_{01}|0\rangle\langle 1| + u_{10}|1\rangle\langle 0| + u_{11}|1\rangle\langle 1|$$

$$-\quad \tilde{\sigma}_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

Diac notation

$$\hookrightarrow \tilde{\sigma}_x|0\rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = |1\rangle, \quad \tilde{\sigma}_x|1\rangle = \underbrace{(|0\rangle\langle 1| + |1\rangle\langle 0|)}_{\downarrow} \cdot \underbrace{|1\rangle}_{1} = |0\rangle \underbrace{\langle 1|}_{0} + |1\rangle \underbrace{\langle 0|}_{0} = |0\rangle$$

\Rightarrow bit flip $\hat{=}$ NOT-gate, e.g. $|0\rangle \xrightarrow{\tilde{\sigma}_x} |1\rangle \Rightarrow$ rotation around x-axis by π

$$-\quad \tilde{\sigma}_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = |0\rangle\langle 0| - |1\rangle\langle 1|$$

$$\hookrightarrow \tilde{\sigma}_z|+\rangle = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = |-\rangle, \quad \tilde{\sigma}_z|-\rangle = (|0\rangle\langle 0| - |1\rangle\langle 1|) \cdot \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

\Rightarrow phase flip \Rightarrow rotation around z-axis by π

$$= \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) = |+\rangle$$

$$-\quad \tilde{\sigma}_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = i \cdot \tilde{\sigma}_x \cdot \tilde{\sigma}_z \quad \Rightarrow \text{bit \& phase flip}$$

$\Rightarrow \tilde{\sigma}_x, \tilde{\sigma}_y \& \tilde{\sigma}_z$ are the so-called Pauli matrices and $\tilde{\sigma}_i^2 = 11 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ (does nothing)

\Rightarrow together with identity 11 they form a basis of 2×2 matrices

(\rightarrow any 1-qubit rotation can be written as a linear combination of them)

- Hadamard gate: one of the most important gates for quantum circuits

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{\sqrt{2}} (|0\rangle\langle 0| + |0\rangle\langle 1| + |1\rangle\langle 0| - |1\rangle\langle 1|)$$

$$\hookrightarrow H|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = |+\rangle, \quad H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle\langle 0| + |0\rangle\langle 1| + |1\rangle\langle 0| - |1\rangle\langle 1|) \cdot |1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) = |-\rangle$$

\Rightarrow creates superposition! also $H|+\rangle = |0\rangle, H|-\rangle = |1\rangle \Rightarrow$ used to change between X & Z basis

- similarly, as $S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$ adds 90° to the phase φ : $S \cdot |+\rangle = |+\rangle, S|-\rangle = |-\rangle$

$\Rightarrow S \cdot H$ is applied to change from Z to Y basis

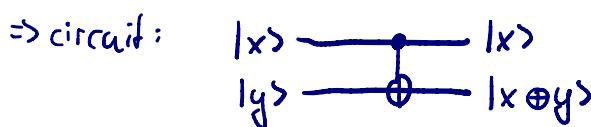
Multipartite quantum states

- we use tensor products to describe multiple states: $|a\rangle \otimes |b\rangle = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \otimes \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} a_1 b_1 \\ a_1 b_2 \\ a_2 b_1 \\ a_2 b_2 \end{pmatrix}$
- example: system A is in state $|1\rangle_A$ and system B is in state $|0\rangle_B$
 \Rightarrow the total (bi-partite) state is $|10\rangle_{AB} = |1\rangle_A \otimes |0\rangle_B = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$
- ↳ remark: states of this form are called **uncorrelated**, but there are also bi-partite states that cannot be written as $|\psi\rangle_A \otimes |\psi\rangle_B$. These states are **correlated** and sometimes even **entangled** (\rightarrow very strong correlation), e.g. $|\Psi\rangle_{AB}^{(0)} = \frac{1}{\sqrt{2}} (|00\rangle_{AB} + |11\rangle_{AB}) = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 1 \end{pmatrix}$
 a so-called **Bell state**, used for teleportation, cryptography, Bell tests, etc.

Two-qubit gates

- classical example: XOR $x = \boxed{\text{XOR}} = x \oplus y \rightarrow \text{irreversible}$ (\rightarrow given the output we cannot recover the input)
 BUT: as quantum theory is unitary, we only consider unitary and therefore **reversible** gates
- quantum example:
 $CNOT = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 11| + |11\rangle\langle 10|$
 $\hookrightarrow CNOT \cdot |00\rangle_{xy} = CNOT \cdot \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = |00\rangle_{xy}, \quad CNOT \cdot |10\rangle_{xy} = |11\rangle_{xy}$
- \Rightarrow

input	output
$x \ y$	$x \ x \oplus y$
$0 \ 0$	$0 \ 0$
$0 \ 1$	$0 \ 1$
$1 \ 0$	$1 \ 1$
$1 \ 1$	$1 \ 0$

 \Rightarrow circuit:

 $\hat{=}$ reversible XOR

\Rightarrow we can show that every function f can be described by a reversible circuit

\Rightarrow quantum circuits can perform all functions that can be calculated classically