

# EE101: Sinusoidal steady state analysis

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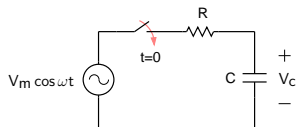


**M. B. Patil**

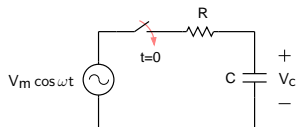
[mbpatil@ee.iitb.ac.in](mailto:mbpatil@ee.iitb.ac.in)

Department of Electrical Engineering  
Indian Institute of Technology Bombay

# Sinusoidal steady state

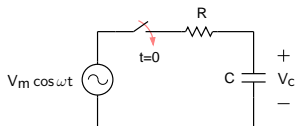


## Sinusoidal steady state



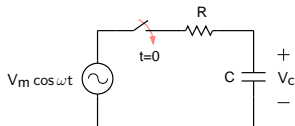
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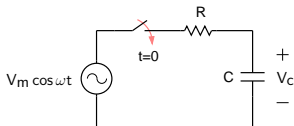


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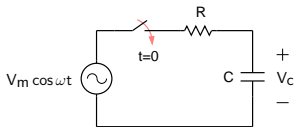
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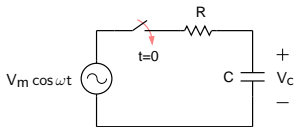
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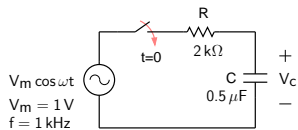
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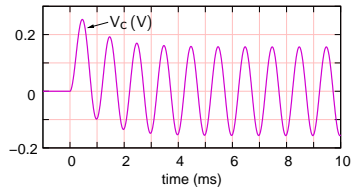
$C_1$  and  $C_2$  can be found by equating the coefficients of  $\sin \omega t$  and  $\cos \omega t$  on the left and right sides.



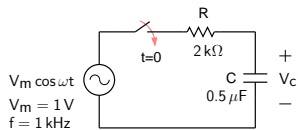
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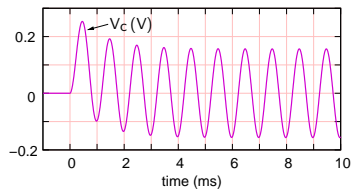
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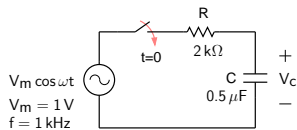


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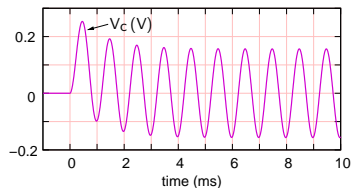


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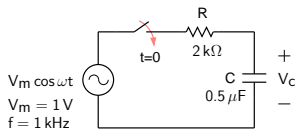


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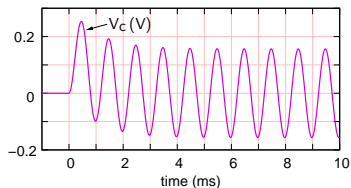


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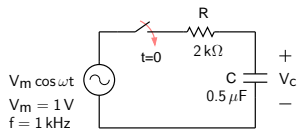


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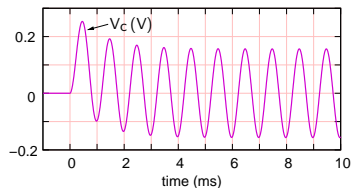


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- \* This is known as the “sinusoidal steady state” response since all quantities (currents and voltages) in the circuit are sinusoidal in nature.
- \* Any circuit containing resistors, capacitors, inductors, sinusoidal voltage and current sources (of the same frequency), dependent (linear) sources behaves in a similar manner, viz., each current and voltage in the circuit becomes purely sinusoidal as  $t \rightarrow \infty$ .

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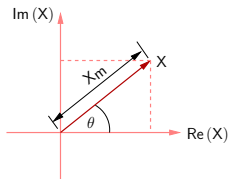
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- \* Use of phasors substantially simplifies analysis of circuits in the sinusoidal steady state.
- \* Note that a phasor can be written in the polar form or rectangular form,  
 $\mathbf{X} = X_m \angle \theta = X_m \exp(j\theta) = X_m \cos \theta + j X_m \sin \theta$ .

The term  $\omega t$  is always *implicit*.



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$i_3(t) = \sqrt{2} \cos(\omega t + 45^\circ) \text{ A}$	$I_3 = 1 + j1 \text{ A}$ $= \sqrt{2} \angle 45^\circ \text{ A}$

# Addition of phasors

Consider addition of two sinusoidal quantities:

$$\begin{aligned}v(t) &= v_1(t) + v_2(t) \\&= V_{m1} \cos(\omega t + \theta_1) + V_{m2} \cos(\omega t + \theta_2)\end{aligned}$$

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Now consider addition of the phasors corresponding to  $v_1(t)$  and  $v_2(t)$ .

$$\begin{aligned}\mathbf{V} &= \mathbf{V}_1 + \mathbf{V}_2 \\&= V_{m1}e^{j\theta_1} + V_{m2}e^{j\theta_2}\end{aligned}$$

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In the time domain,  $\mathbf{V}$  corresponds to  $\tilde{v}(t)$ , with

$$\tilde{v}(t) = \text{Re} [\mathbf{V}e^{j\omega t}]$$

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$$\begin{aligned}\mathbf{V} &= \mathbf{V}_1 + \mathbf{V}_2 \\&= V_{m1}e^{j\theta_1} + V_{m2}e^{j\theta_2}\end{aligned}$$

In the time domain,  $\mathbf{V}$  corresponds to  $\tilde{v}(t)$ , with

$$\begin{aligned}\tilde{v}(t) &= \text{Re} [\mathbf{V}e^{j\omega t}] \\&= \text{Re} [(V_{m1}e^{j\theta_1} + V_{m2}e^{j\theta_2}) e^{j\omega t}]\end{aligned}$$

# Addition of phasors

Consider addition of two sinusoidal quantities:

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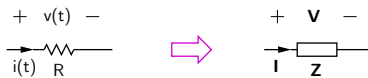
$$\begin{aligned}\tilde{v}(t) &= \text{Re} [\mathbf{V}e^{j\omega t}] \\&= \text{Re} [(V_{m1}e^{j\theta_1} + V_{m2}e^{j\theta_2}) e^{j\omega t}] \\&= \text{Re} [V_{m1}e^{j(\omega t + \theta_1)} + V_{m2}e^{j(\omega t + \theta_2)}] \\&= V_{m1} \cos(\omega t + \theta_1) + V_{m2} \cos(\omega t + \theta_2)\end{aligned}$$

which is the same as  $v(t)$ .

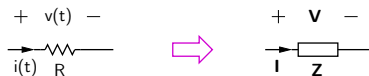
- \* Addition of sinusoidal quantities in the time domain can be replaced by addition of the corresponding phasors in the sinusoidal steady state.

- \* Addition of sinusoidal quantities in the time domain can be replaced by addition of the corresponding phasors in the sinusoidal steady state.
- \* The KCL and KVL equations,  
 $\sum i_k(t) = 0$  at a node, and  
 $\sum v_k(t) = 0$  in a loop,  
amount to addition of sinusoidal quantities and can therefore be replaced by the corresponding phasor equations,  
 $\sum \mathbf{I}_k = \mathbf{0}$  at a node, and  
 $\sum \mathbf{V}_k = \mathbf{0}$  in a loop.

# Impedance of a resistor

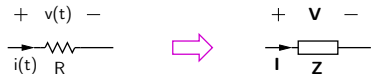


## Impedance of a resistor



Let  $i(t) = I_m \cos(\omega t + \theta)$ .

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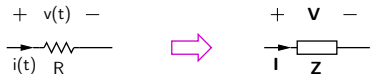
# Impedance of a resistor



Let  $i(t) = I_m \cos(\omega t + \theta)$ .

$$\begin{aligned} v(t) &= R i(t) \\ &= R I_m \cos(\omega t + \theta) \\ &\equiv V_m \cos(\omega t + \theta), \end{aligned}$$

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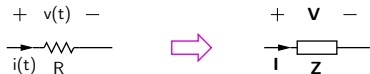
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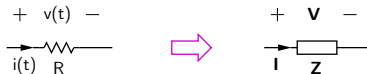
$$\equiv V_m \cos(\omega t + \theta),$$

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$$\text{i.e., } \operatorname{Re} [V_m e^{j\theta} e^{j\omega t}] = R \times \operatorname{Re} [I_m e^{j\theta} e^{j\omega t}],$$

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corresponding to the phasor relationship,

$$\mathbf{V} = R \mathbf{I}.$$

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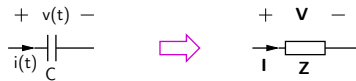
corresponding to the phasor relationship,

$$\mathbf{V} = R \mathbf{I}.$$

Thus, the *impedance* of a resistor, defined as,  $\mathbf{Z} = \mathbf{V}/\mathbf{I}$ , is

$$\mathbf{Z} = R + j0$$

# Impedance of a capacitor

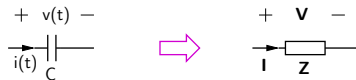


# Impedance of a capacitor



Let  $v(t) = V_m \cos(\omega t + \theta)$ .

# Impedance of a capacitor

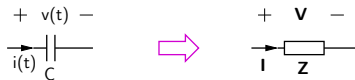


Let  $v(t) = V_m \cos(\omega t + \theta)$ .

$$i(t) = C \frac{dv}{dt} = -C \omega V_m \sin(\omega t + \theta).$$



# Impedance of a capacitor



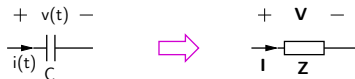
Let  $v(t) = V_m \cos(\omega t + \theta)$ .

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Using the identity,  $\cos(\phi + \pi/2) = -\sin \phi$ , we get

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In terms of phasors,  $\mathbf{V} = V_m \angle \theta$ ,  $\mathbf{I} = \omega C V_m \angle (\theta + \pi/2)$ .

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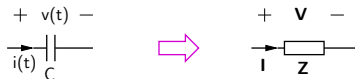
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In terms of phasors,  $\mathbf{V} = V_m \angle \theta$ ,  $\mathbf{I} = \omega C V_m \angle (\theta + \pi/2)$ .

$\mathbf{I}$  can be rewritten as,

$$\mathbf{I} = \omega C V_m e^{j(\theta + \pi/2)} = \omega C V_m e^{j\theta} e^{j\pi/2} = j\omega C (V_m e^{j\theta}) = j\omega C \mathbf{V}$$

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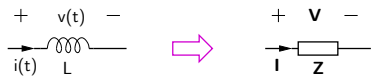
$\mathbf{I}$  can be rewritten as,

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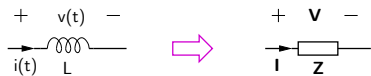
Thus, the *impedance* of a capacitor,  $\mathbf{Z} = \mathbf{V}/\mathbf{I}$ , is  $\boxed{\mathbf{Z} = 1/(j\omega C)}$ ,

and the *admittance* of a capacitor,  $\mathbf{Y} = \mathbf{I}/\mathbf{V}$ , is  $\boxed{\mathbf{Y} = j\omega C}$ .

## Impedance of an inductor

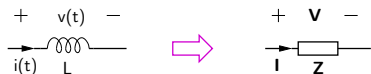


## Impedance of an inductor



Let  $i(t) = I_m \cos(\omega t + \theta)$ .

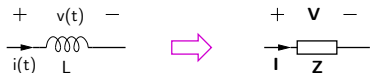
## Impedance of an inductor



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$$v(t) = L \frac{di}{dt} = -L \omega I_m \sin(\omega t + \theta).$$

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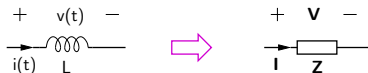
$$v(t) = L \frac{di}{dt} = -L \omega I_m \sin(\omega t + \theta).$$

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## Impedance of an inductor



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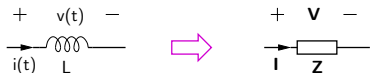
$$v(t) = L \frac{di}{dt} = -L\omega I_m \sin(\omega t + \theta).$$

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In terms of phasors,  $I = I_m \angle \theta$ ,  $V = \omega L I_m \angle (\theta + \pi/2)$ .

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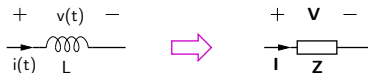
$$v(t) = L\omega I_m \cos(\omega t + \theta + \pi/2).$$

In terms of phasors,  $\mathbf{I} = I_m \angle \theta$ ,  $\mathbf{V} = \omega L I_m \angle (\theta + \pi/2)$ .

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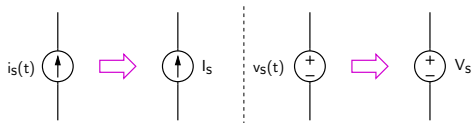
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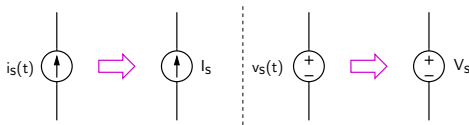
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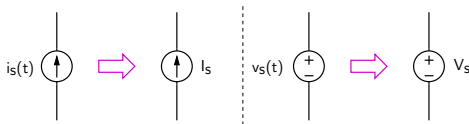
Thus, the *impedance* of an inductor,  $\mathbf{Z} = \mathbf{V}/\mathbf{I}$ , is  $\boxed{\mathbf{Z} = j\omega L}$ ,

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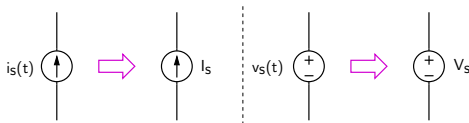




- \* An independent sinusoidal current source,  $i_s(t) = I_m \cos(\omega t + \theta)$ , can be represented by the phasor  $I_m \angle \theta$  (i.e., a *constant* complex number).



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- \* Dependent (linear) sources can be treated in the sinusoidal steady state in the same manner as a resistor, i.e., by the corresponding phasor relationship. For example, for a CCVS, we have,  $v(t) = r i_c(t)$  in the time domain.  
 $\mathbf{V} = r \mathbf{I}_c$  in the frequency domain.

# Use of phasors in circuit analysis



- \* The time-domain KCL and KVL equations  $\sum i_k(t) = 0$  and  $\sum v_k(t) = 0$  can be written as  $\sum \mathbf{I}_k = \mathbf{0}$  and  $\sum \mathbf{V}_k = \mathbf{0}$  in the frequency domain.

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- \* Resistors, capacitors, and inductors can be described by  $\mathbf{V} = \mathbf{Z} \mathbf{I}$  in the frequency domain, which is similar to  $V = R I$  in DC conditions (except that we are dealing with complex numbers in the frequency domain).

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- \* An independent sinusoidal source in the frequency domain behaves like a DC source, e.g.,  $\mathbf{V}_s = \text{constant}$  (a complex number).
- \* For dependent sources, the time-domain relationships such as  $i(t) = \beta i_c(t)$  translate to  $\mathbf{I} = \beta \mathbf{I}_c$  in the frequency domain.

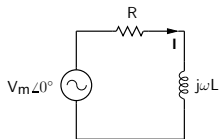
## Use of phasors in circuit analysis

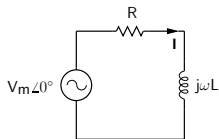
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- \* For dependent sources, the time-domain relationships such as  $i(t) = \beta i_c(t)$  translate to  $\mathbf{I} = \beta \mathbf{I}_c$  in the frequency domain.
- \* Circuit analysis in the sinusoidal steady state using phasors is therefore very similar to DC circuits with independent and dependent sources, and resistors.

## Use of phasors in circuit analysis

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- \* Resistors, capacitors, and inductors can be described by  $\mathbf{V} = \mathbf{Z} \mathbf{I}$  in the frequency domain, which is similar to  $V = R I$  in DC conditions (except that we are dealing with complex numbers in the frequency domain).
- \* An independent sinusoidal source in the frequency domain behaves like a DC source, e.g.,  $\mathbf{V}_s = \text{constant}$  (a complex number).
- \* For dependent sources, the time-domain relationships such as  $i(t) = \beta i_c(t)$  translate to  $\mathbf{I} = \beta \mathbf{I}_c$  in the frequency domain.
- \* Circuit analysis in the sinusoidal steady state using phasors is therefore very similar to DC circuits with independent and dependent sources, and resistors.
- \* Series/parallel formulas for resistors, nodal analysis, mesh analysis, Thevenin's and Norton's theorems can be directly applied to circuits in the sinusoidal steady state.

## *RL* circuit

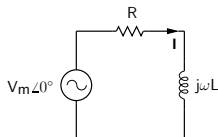




$$I = \frac{V_m \angle 0}{R + j\omega L} \equiv I_m \angle (-\theta),$$

$$\text{where } I_m = \frac{V_m}{\sqrt{R^2 + \omega^2 L^2}}, \text{ and } \theta = \tan^{-1}(\omega L/R).$$

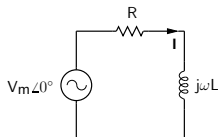




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In the time domain,  $i(t) = I_m \cos(\omega t - \theta)$ , which *lags* the source voltage since the peak (or zero) of  $i(t)$  occurs  $t = \theta/\omega$  seconds after that of the source voltage.



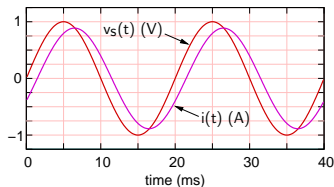
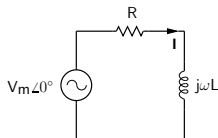
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For  $R = 1 \Omega$ ,  $L = 1.6 \text{ mH}$ ,  $f = 50 \text{ Hz}$ ,  $\theta = 26.6^\circ$ ,  $t_{\text{lag}} = 1.48 \text{ ms}$ .

(SEQUEL file: ee101\_r1\_ac\_1.sqproj)



$$R = 1 \, \Omega$$

$$L = 1.6 \, \text{mH}$$

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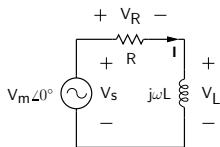
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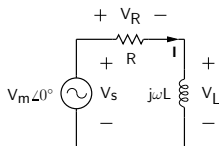
## RL circuit



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## RL circuit

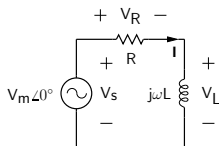


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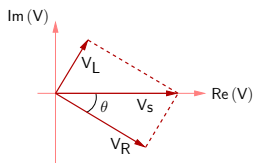
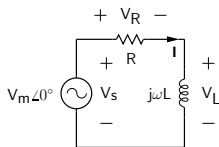
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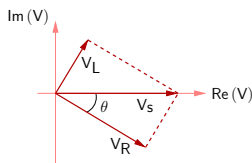
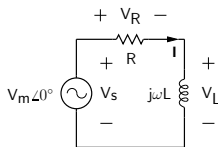
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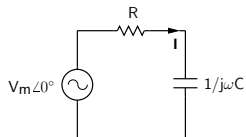
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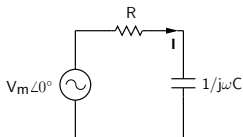
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## RC circuit

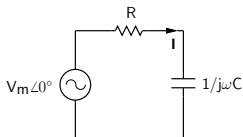




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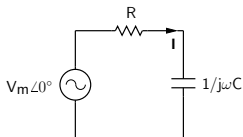
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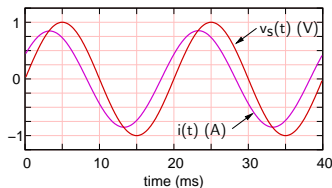
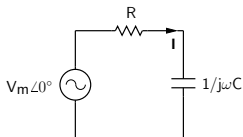
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For  $R = 1\ \Omega$ ,  $L = 5.3\ \text{mF}$ ,  $f = 50\ \text{Hz}$ ,  $\theta = 31^\circ$ ,  $t_{\text{lead}} = 1.72\ \text{ms}$ .

(SEQUEL file: ee101\_rc\_ac\_1.sqproj)

## RC circuit



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$$C = 5.3 \text{ mF}$$

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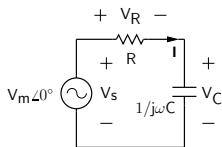
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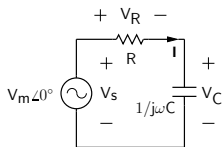
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## RC circuit

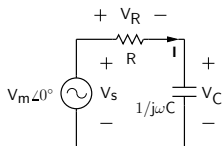


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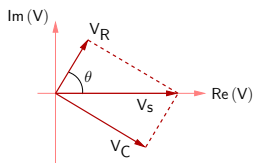
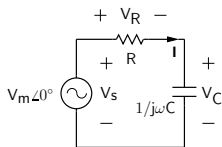
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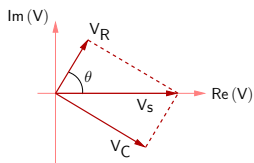
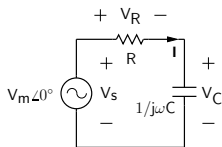
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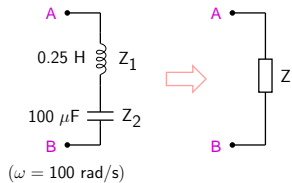
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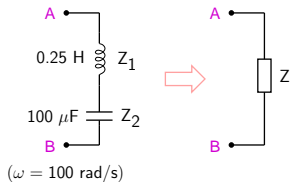
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## Series/parallel connections



## Series/parallel connections

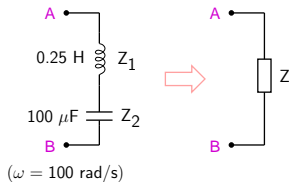


$$Z_1 = j \times 100 \times 0.25 = j 25 \Omega$$

$$Z_2 = -j / (100 \times 100 \times 10^{-6}) = -j 100 \Omega$$

$$Z = Z_1 + Z_2 = -j 75 \Omega$$

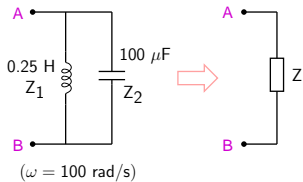
## Series/parallel connections



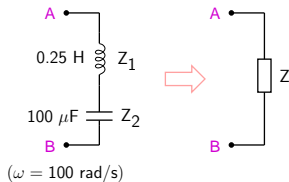
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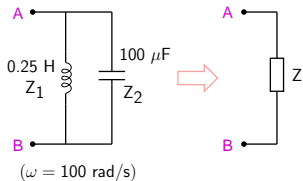
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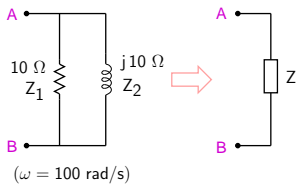
$$Z = Z_1 + Z_2 = -j 75 \Omega$$



$$\begin{aligned} Z &= \frac{Z_1 Z_2}{Z_1 + Z_2} \\ &= \frac{(j 25) \times (-j 100)}{j 25 - j 100} \\ &= \frac{25 \times 100}{-j 75} \\ &= j 33.3 \Omega \end{aligned}$$

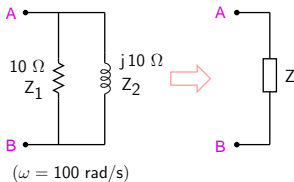
# Impedance example

Obtain  $Z$  in polar form.



## Impedance example

Obtain  $Z$  in polar form.



Method 1:

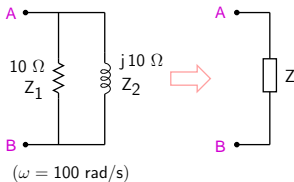
$$\begin{aligned} Z &= \frac{10 \times j10}{10 + j10} = \frac{j10}{1 + j} \\ &= \frac{j10}{1 + j} \times \frac{1 - j}{1 - j} \\ &= \frac{10 + j10}{2} = 5 + j5\ \Omega \end{aligned}$$

Convert to polar form  $\rightarrow Z = 7.07 \angle 45^\circ\ \Omega$



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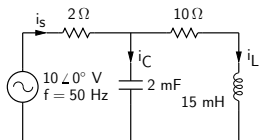
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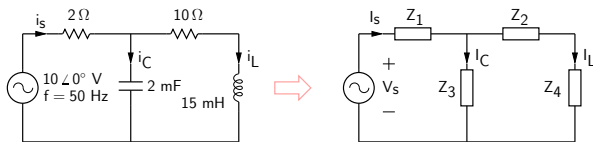
Method 2:

$$\begin{aligned} Z &= \frac{10 \times j10}{10 + j10} = \frac{100 \angle \pi/2}{10\sqrt{2} \angle \pi/4} \\ &= 5\sqrt{2} \angle (\pi/2 - \pi/4) = 7.07 \angle 45^\circ\ \Omega \end{aligned}$$

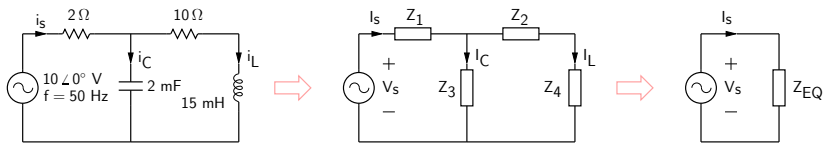
## Circuit example



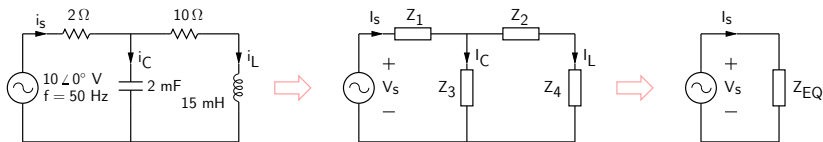
## Circuit example



## Circuit example

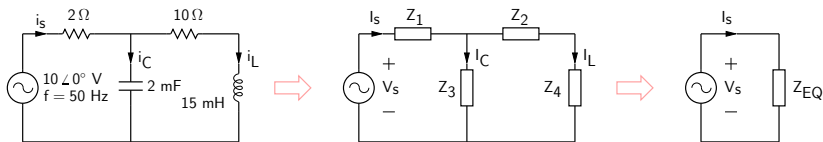


## Circuit example



$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j1.6 \Omega$$

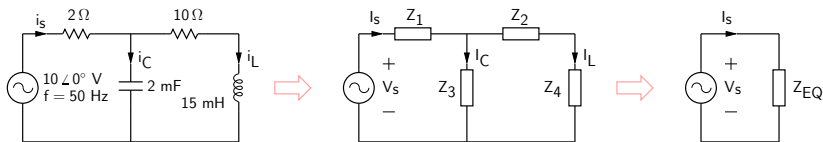
## Circuit example



$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j 1.6 \Omega$$

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## Circuit example

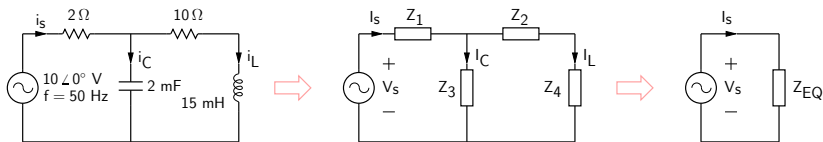


$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j 1.6 \Omega$$

$$Z_4 = 2\pi \times 50 \times 15 \times 10^{-3} = j 4.7 \Omega$$

$$Z_{EQ} = Z_1 + Z_3 \parallel (Z_2 + Z_4)$$

## Circuit example



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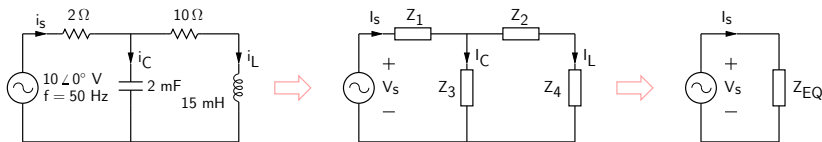
$$Z_4 = 2\pi \times 50 \times 15 \times 10^{-3} = j4.7 \Omega$$

$$Z_{EQ} = Z_1 + Z_3 \parallel (Z_2 + Z_4)$$

$$= 2 + (-j1.6) \parallel (10 + j4.7) = 2 + \frac{(-j1.6) \times (10 + j4.7)}{-j1.6 + 10 + j4.7}$$



## Circuit example



$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j1.6 \Omega$$

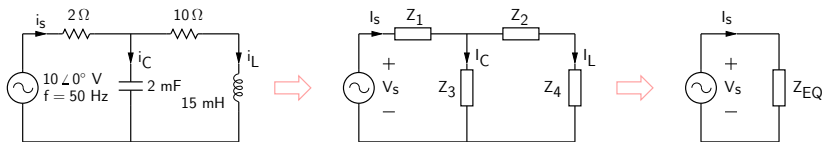
$$Z_4 = 2\pi \times 50 \times 15 \times 10^{-3} = j4.7 \Omega$$

$$Z_{EQ} = Z_1 + Z_3 \parallel (Z_2 + Z_4)$$

$$= 2 + (-j1.6) \parallel (10 + j4.7) = 2 + \frac{(-j1.6) \times (10 + j4.7)}{-j1.6 + 10 + j4.7}$$

$$= 2 + \frac{1.6 \angle (-90^\circ) \times 11.05 \angle (25.2^\circ)}{10.47 \angle (17.2^\circ)} = 2 + \frac{17.7 \angle (-64.8^\circ)}{10.47 \angle (17.2^\circ)}$$

## Circuit example



$$\mathbf{Z}_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j1.6 \, \Omega$$

$$\mathbf{Z}_4 = 2\pi \times 50 \times 15 \times 10^{-3} = j4.7 \, \Omega$$

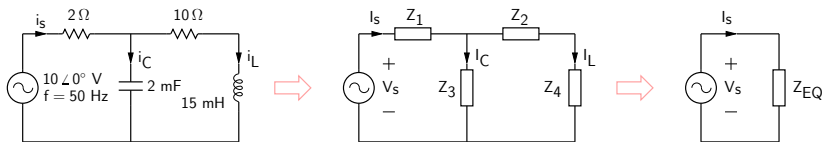
$$\mathbf{Z}_{EQ} = \mathbf{Z}_1 + \mathbf{Z}_3 \parallel (\mathbf{Z}_2 + \mathbf{Z}_4)$$

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$$= 2 + 1.69 \angle (-82^\circ) = 2 + (0.235 - j1.67)$$

## Circuit example



$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j1.6 \Omega$$

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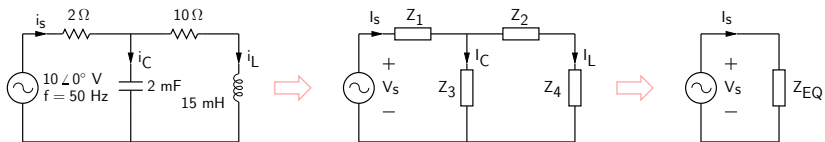
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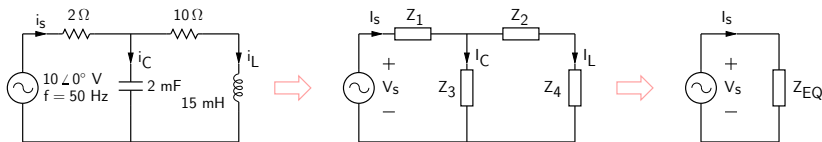
$$= 2 + 1.69 \angle (-82^\circ) = 2 + (0.235 - j1.67)$$

$$= 2.235 - j1.67 = 2.79 \angle (-36.8^\circ) \Omega$$

## Circuit example (continued)

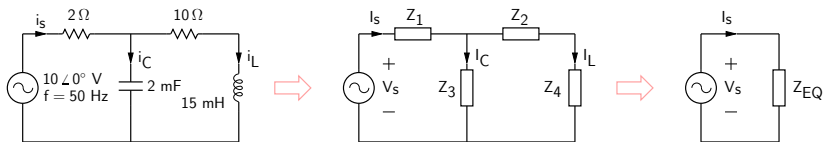


## Circuit example (continued)



$$I_s = \frac{V_s}{Z_{EQ}} = \frac{10 \angle (0^\circ)}{2.79 \angle (-36.8^\circ)} = 3.58 \angle (36.8^\circ) \text{ A}$$

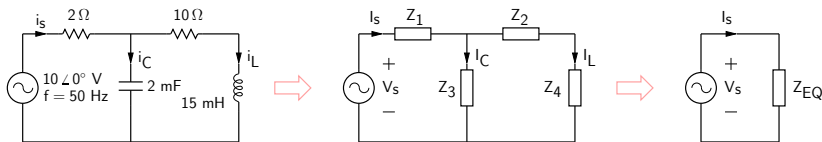
## Circuit example (continued)



$$I_s = \frac{V_s}{Z_{EQ}} = \frac{10 \angle (0^\circ)}{2.79 \angle (-36.8^\circ)} = 3.58 \angle (36.8^\circ) \text{ A}$$

$$I_C = \frac{(Z_2 + Z_4)}{Z_3 + (Z_2 + Z_4)} \times I_s = 3.79 \angle (44.6^\circ) \text{ A}$$

## Circuit example (continued)

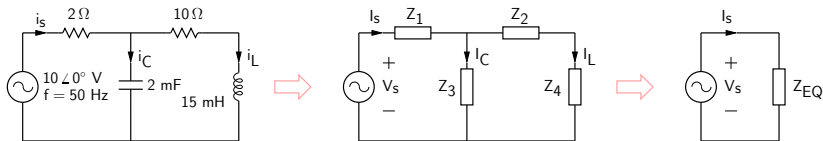


$$I_s = \frac{V_s}{Z_{EQ}} = \frac{10 \angle (0^\circ)}{2.79 \angle (-36.8^\circ)} = 3.58 \angle (36.8^\circ) \text{ A}$$

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$$I_L = \frac{Z_3}{Z_3 + (Z_2 + Z_4)} \times I_s = 0.546 \angle (-70.6^\circ) \text{ A}$$

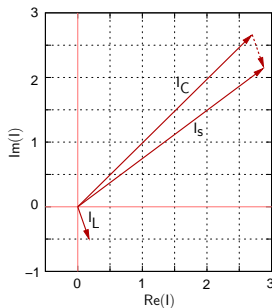
## Circuit example (continued)



$$I_s = \frac{V_s}{Z_{EQ}} = \frac{10 \angle (0^\circ)}{2.79 \angle (-36.8^\circ)} = 3.58 \angle (36.8^\circ) \text{ A}$$

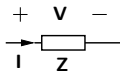
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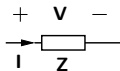




## Sinusoidal steady state: power computation



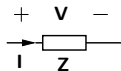
## Sinusoidal steady state: power computation



Let  $v(t) = V_m \cos(\omega t + \theta)$ , i.e.,  $\mathbf{V} = V_m \angle \theta$ ,

$i(t) = I_m \cos(\omega t + \phi)$ , i.e.,  $\mathbf{I} = I_m \angle \phi$ .

## Sinusoidal steady state: power computation



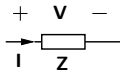
Let  $v(t) = V_m \cos(\omega t + \theta)$ , i.e.,  $\mathbf{V} = V_m \angle \theta$ ,

$i(t) = I_m \cos(\omega t + \phi)$ , i.e.,  $\mathbf{I} = I_m \angle \phi$ .

The *instantaneous* power absorbed by  $\mathbf{Z}$  is,

$$\begin{aligned} P(t) &= v(t) i(t) \\ &= V_m I_m \cos(\omega t + \theta) \cos(\omega t + \phi) \\ &= \frac{1}{2} V_m I_m [\cos(2\omega t + \theta + \phi) + \cos(\theta - \phi)] \end{aligned} \tag{1}$$

## Sinusoidal steady state: power computation



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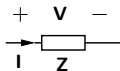
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The *average* power absorbed by  $\mathbf{Z}$  is

$$P = \frac{1}{T} \int_0^T P(t) dt, \text{ where } T = 2\pi/\omega.$$

## Sinusoidal steady state: power computation



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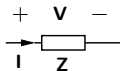
$$\begin{aligned} P(t) &= v(t) i(t) \\ &= V_m I_m \cos(\omega t + \theta) \cos(\omega t + \phi) \\ &= \frac{1}{2} V_m I_m [\cos(2\omega t + \theta + \phi) + \cos(\theta - \phi)] \end{aligned} \quad (1)$$

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The first term in Eq. (1) has an average value of zero and does not contribute to  $P$ .  
Therefore,

## Sinusoidal steady state: power computation



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The *instantaneous* power absorbed by  $\mathbf{Z}$  is,

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The *average* power absorbed by  $\mathbf{Z}$  is

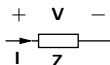
$$P = \frac{1}{T} \int_0^T P(t) dt, \text{ where } T = 2\pi/\omega.$$

The first term in Eq. (1) has an average value of zero and does not contribute to  $P$ . Therefore,

$P = \frac{1}{2} V_m I_m \cos(\theta - \phi)$

 gives the average power absorbed by  $\mathbf{Z}$ .

## Average power for $R$ , $L$ , $C$

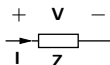


General formula:

$$V = V_m \angle \theta, I = I_m \angle \phi$$

$$P = \frac{1}{2} V_m I_m \cos(\theta - \phi)$$

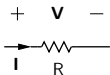
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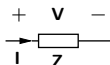
$$V = RI$$

$$\text{For } I = I_m \angle \alpha, V = RI_m \angle \alpha,$$

$$P = \frac{1}{2} (RI_m) I_m \cos(\alpha - \alpha) = \frac{1}{2} I_m^2 R = \frac{1}{2} V_m^2 / R$$



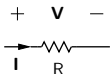
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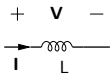
$$P = \frac{1}{2} V_m I_m \cos(\theta - \phi)$$



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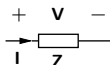


$$V = j\omega LI$$

$$\text{For } I = I_m \angle \alpha, V = \omega L I_m \angle (\alpha + \pi/2),$$

$$P = \frac{1}{2} V_m I_m \cos[(\alpha + \pi/2) - \alpha] = 0$$

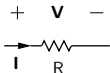
## Average power for $R$ , $L$ , $C$



General formula:

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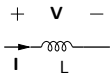
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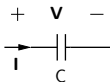
$$P = \frac{1}{2} (RI_m) I_m \cos(\alpha - \alpha) = \frac{1}{2} I_m^2 R = \frac{1}{2} V_m^2 / R$$



$$V = j\omega L I$$

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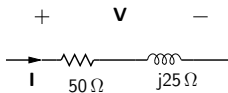


$$I = j\omega C V$$

$$\text{For } V = V_m \angle \alpha, I = \omega C V_m \angle (\alpha + \pi/2),$$

$$P = \frac{1}{2} V_m I_m \cos[\alpha - (\alpha + \pi/2)] = 0$$

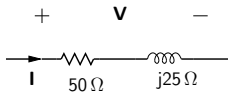
## Average power: example



Given:  $I = 2 \angle 45^\circ\text{ A}$

Find the average power absorbed.

## Average power: example



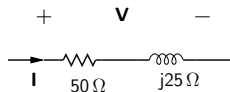
Given:  $I = 2 \angle 45^\circ\text{ A}$

Find the average power absorbed.

Method 1:

$$\begin{aligned} \mathbf{V} &= (50 + j25) \times 2 \angle 45^\circ \\ &= 55.9 \angle 26.6^\circ \times 2 \angle 45^\circ \\ &= 111.8 \angle (45^\circ + 26.6^\circ) \end{aligned}$$

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Given:  $I = 2 \angle 45^\circ\text{ A}$

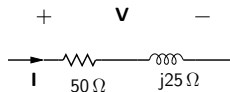
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$$P = \frac{1}{2} \times 111.8 \times 2 \times \cos(26.6^\circ) = 100\text{ W}.$$

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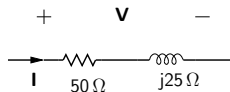
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Method 2:

No average power is absorbed by the inductor.

$\Rightarrow P = P_R$  (average power absorbed by  $R$ )

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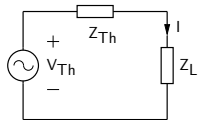
Method 2:

No average power is absorbed by the inductor.

$\Rightarrow P = P_R$  (average power absorbed by  $R$ )

$$\begin{aligned} &= \frac{1}{2} I_m^2 R = \frac{1}{2} \times 2^2 \times 50 \\ &= 100\text{ W}. \end{aligned}$$

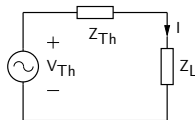
# Maximum power transfer





# Maximum power transfer

Let  $\mathbf{Z}_L = R_L + jX_L$ ,  $\mathbf{Z}_{Th} = R_{Th} + jX_{Th}$ , and  $\mathbf{I} = I_m \angle \phi$ .

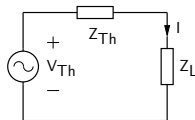


# Maximum power transfer

Let  $\mathbf{Z}_L = R_L + jX_L$ ,  $\mathbf{Z}_{Th} = R_{Th} + jX_{Th}$ , and  $\mathbf{I} = I_m \angle \phi$ .

The power absorbed by  $\mathbf{Z}_L$  is,

$$\begin{aligned} P &= \frac{1}{2} I_m^2 R_L \\ &= \frac{1}{2} \left| \frac{\mathbf{V}_{Th}}{\mathbf{Z}_{Th} + \mathbf{Z}_L} \right|^2 R_L \\ &= \frac{1}{2} \frac{|\mathbf{V}_{Th}|^2}{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2} R_L. \end{aligned}$$



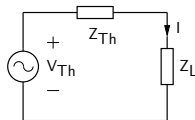
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For  $P$  to be maximum,  $(X_{Th} + X_L)$  must be zero.  $\Rightarrow X_L = -X_{Th}$ .



# Maximum power transfer

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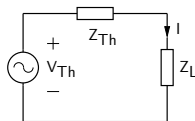
$$\begin{aligned} P &= \frac{1}{2} I_m^2 R_L \\ &= \frac{1}{2} \left| \frac{\mathbf{V}_{Th}}{\mathbf{Z}_{Th} + \mathbf{Z}_L} \right|^2 R_L \\ &= \frac{1}{2} \frac{|\mathbf{V}_{Th}|^2}{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2} R_L. \end{aligned}$$

For  $P$  to be maximum,  $(X_{Th} + X_L)$  must be zero.  $\Rightarrow X_L = -X_{Th}$ .

With  $X_L = -X_{Th}$ , we have,

$$P = \frac{1}{2} \frac{|\mathbf{V}_{Th}|^2}{(R_{Th} + R_L)^2} R_L,$$

which is maximum for  $R_L = R_{Th}$ .



# Maximum power transfer

Let  $\mathbf{Z}_L = R_L + jX_L$ ,  $\mathbf{Z}_{Th} = R_{Th} + jX_{Th}$ , and  $\mathbf{I} = I_m \angle \phi$ .

The power absorbed by  $\mathbf{Z}_L$  is,

$$\begin{aligned} P &= \frac{1}{2} I_m^2 R_L \\ &= \frac{1}{2} \left| \frac{\mathbf{V}_{Th}}{\mathbf{Z}_{Th} + \mathbf{Z}_L} \right|^2 R_L \\ &= \frac{1}{2} \frac{|\mathbf{V}_{Th}|^2}{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2} R_L. \end{aligned}$$

For  $P$  to be maximum,  $(X_{Th} + X_L)$  must be zero.  $\Rightarrow X_L = -X_{Th}$ .

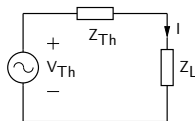
With  $X_L = -X_{Th}$ , we have,

$$P = \frac{1}{2} \frac{|\mathbf{V}_{Th}|^2}{(R_{Th} + R_L)^2} R_L,$$

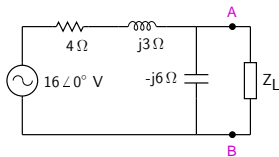
which is maximum for  $R_L = R_{Th}$ .

Therefore, for maximum power transfer to the load  $\mathbf{Z}_L$ , we need,

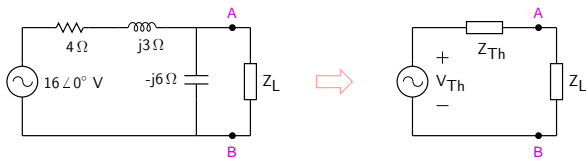
$$R_L = R_{Th}, X_L = -X_{Th}, \text{ i.e., } \boxed{\mathbf{Z}_L = \mathbf{Z}_{Th}^*}.$$



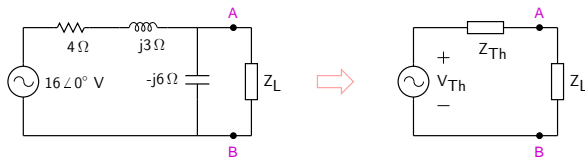
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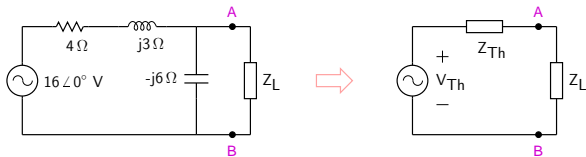
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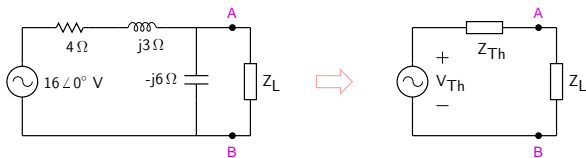
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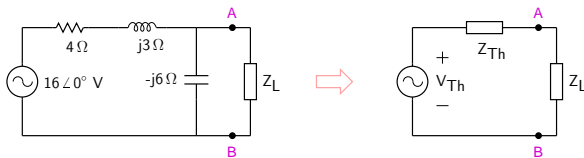


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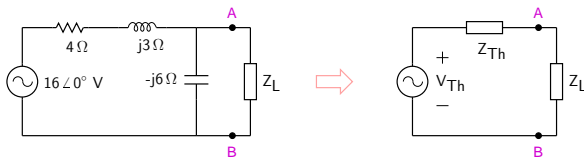
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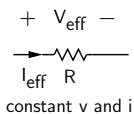
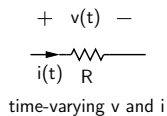
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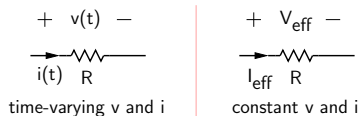
$$P = \frac{1}{2} I_m^2 R_L = \frac{1}{2} \left( \frac{19.2}{2R_L} \right)^2 \times R_L = \frac{1}{2} \frac{(19.2)^2}{4R_L} = 8 \text{ W}.$$

## Effective (rms) values of voltage/current



Consider a periodic current  $i(t)$  passing through  $R$ .

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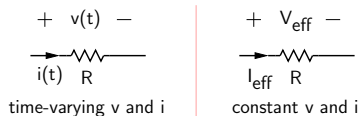
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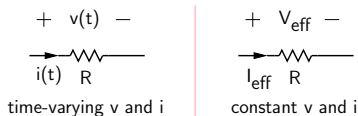
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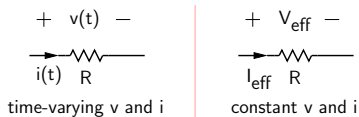
$I_{\text{eff}}$ , the *effective* value of  $i(t)$ , is defined such that  $P_1 = P_2$ , i.e.,

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$$I_{\text{eff}} = \sqrt{\frac{1}{T} \int_0^T [i(t)]^2 dt}.$$



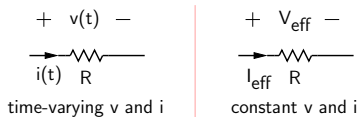
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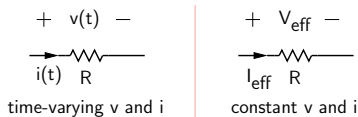
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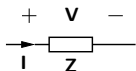
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Similarly,  $V_{\text{eff}} = V_m / \sqrt{2}$ .

# Apparent power, real power, and power factor



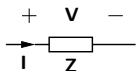
$$V = V_m \angle \theta$$

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The average ("real") power absorbed by  $Z$  is,

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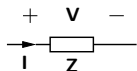
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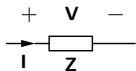
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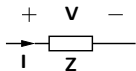
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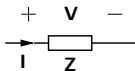
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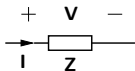


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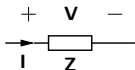
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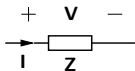


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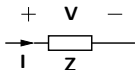
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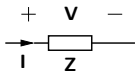
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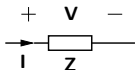
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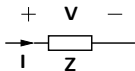
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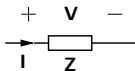
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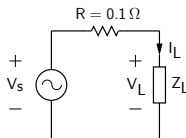
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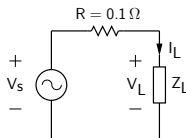


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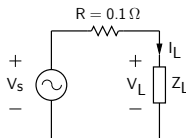
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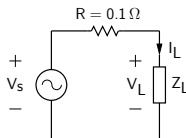
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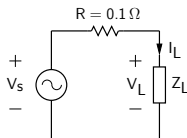
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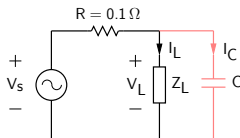
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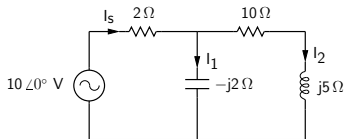
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The effective power factor of an inductive load can be improved by connecting a suitable capacitance in parallel.

## Power computation: home work



- \* Find  $I_1$ ,  $I_2$ ,  $I_s$ .
- \* Compute the average power absorbed by each element.
- \* Verify power balance.

(SEQUEL file: ee101\_phasors\_2.sqproj)