

# Worksheet 16

Parth Bhargava · A0310667E

**PC3261**  
Classical Mechanics II

November 21, 2025

## Problem 7.4

Obtain the expression for the inertial force term in the d'Alembert's principle.

### Solution

Starting with the position vector as a function of generalized coordinates:

$$\mathbf{r}_\alpha = \mathbf{r}_\alpha(\{q_k(t)\}, t)$$

Taking the total time derivative:

$$\begin{aligned} \frac{d}{dt} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) &= \sum_{i=1}^M \frac{\partial}{\partial q_i} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) \dot{q}_i + \frac{\partial}{\partial t} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) \\ &= \sum_{i=1}^M \frac{\partial}{\partial q_k} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_i} \right) \dot{q}_i + \frac{\partial}{\partial q_k} \left( \frac{\partial \mathbf{r}_\alpha}{\partial t} \right) \\ &= \sum_{i=1}^M \frac{\partial}{\partial q_k} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_i} \dot{q}_i \right) + \frac{\partial}{\partial q_k} \left( \frac{\partial \mathbf{r}_\alpha}{\partial t} \right) \\ &= \frac{\partial}{\partial q_k} \left[ \sum_{i=1}^M \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_i} \dot{q}_i \right) + \frac{\partial \mathbf{r}_\alpha}{\partial t} \right] \\ &= \frac{\partial \dot{\mathbf{r}}_\alpha}{\partial q_k} \end{aligned}$$

For the virtual displacement:

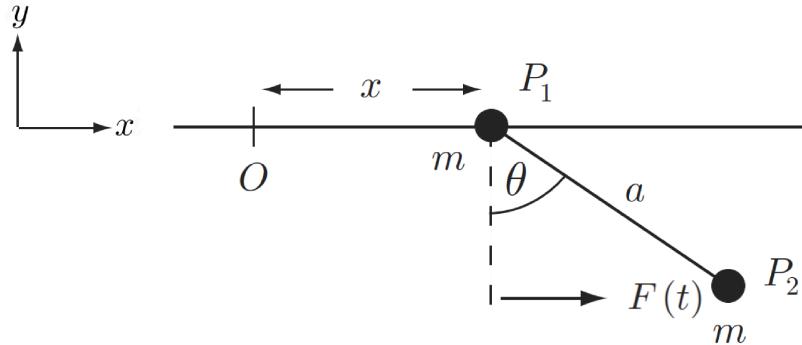
$$\mathbf{r}_\alpha = \mathbf{r}_\alpha(\{q_k(t)\}, t) \Rightarrow \delta \mathbf{r}_\alpha = \sum_{k=1}^M \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \delta q_k$$

The inertial force term in d'Alembert's principle:

$$\begin{aligned} - \sum_{\alpha=1}^N m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \delta \mathbf{r}_\alpha &= - \sum_{\alpha=1}^N \sum_{k=1}^M m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \delta q_k \\ &= - \sum_{k=1}^M \sum_{\alpha=1}^N \left[ \frac{d}{dt} \left( m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) - m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \frac{d}{dt} \left( \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) \right] \delta q_k \\ &= - \sum_{k=1}^M \sum_{\alpha=1}^N \left[ \frac{d}{dt} \left( m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \frac{\partial \mathbf{r}_\alpha}{\partial q_k} \right) - m_\alpha \ddot{\mathbf{r}}_\alpha \cdot \frac{\partial \dot{\mathbf{r}}_\alpha}{\partial q_k} \right] \delta q_k \\ &= - \sum_{k=1}^M \sum_{\alpha=1}^N \left\{ \frac{d}{dt} \left[ \frac{\partial}{\partial \dot{q}_k} \left( \frac{1}{2} m_\alpha \dot{\mathbf{r}}_\alpha \cdot \dot{\mathbf{r}}_\alpha \right) \right] - \frac{\partial}{\partial q_k} \left( \frac{1}{2} m_\alpha \dot{\mathbf{r}}_\alpha \cdot \dot{\mathbf{r}}_\alpha \right) \right\} \delta q_k \end{aligned}$$

### Problem 7.5

Two identical particles,  $P_1$  and  $P_2$ , with mass  $m$  are connected by a light rigid rod of length  $a$ .  $P_1$  is constrained to move along a fixed horizontal frictionless rail and the system moves in the vertical plane through the rail. An external force  $F(t)\hat{e}_x$  is acted on  $P_2$ .



**Generalized coordinates:**  $(q_1, q_2) \equiv (x, \theta)$

Use Lagrange's equation to obtain equations of motions for  $x(t)$  and  $\theta(t)$ .

#### Solution

**Position vectors:**

$$\begin{cases} \mathbf{r}_1(t) = x(t)\hat{e}_x \\ \mathbf{r}_2(t) = [x(t) + a \sin \theta(t)]\hat{e}_x - a \cos \theta(t)\hat{e}_y \end{cases}$$

**Velocity vectors:**

$$\begin{cases} \dot{\mathbf{r}}_1(t) = \dot{x}(t)\hat{e}_x \\ \dot{\mathbf{r}}_2(t) = [\dot{x}(t) + a\dot{\theta}(t) \cos \theta(t)]\hat{e}_x + a\dot{\theta}(t) \sin \theta(t)\hat{e}_y \end{cases}$$

**Kinetic energy:**

$$\begin{aligned} T &\equiv T(x, \theta, \dot{x}, \dot{\theta}, t) = \frac{m}{2}\dot{\mathbf{r}}_1(t) \cdot \dot{\mathbf{r}}_1(t) + \frac{m}{2}\dot{\mathbf{r}}_2(t) \cdot \dot{\mathbf{r}}_2(t) \\ &= \frac{1}{2}m\dot{x}^2(t) + ma^2\dot{\theta}^2(t) + m\dot{x}(t)\dot{\theta}(t) \cos \theta(t) \end{aligned}$$

**Applied forces and generalized forces:**

$$\begin{cases} \mathbf{F}_1(t) = -mg\hat{e}_y \\ \mathbf{F}_2(t) = -mg\hat{e}_y + F(t)\hat{e}_x \\ \mathcal{Q}_x(t) = \mathbf{F}_1(t) \cdot \frac{\partial \mathbf{r}_1}{\partial x} + \mathbf{F}_2(t) \cdot \frac{\partial \mathbf{r}_2}{\partial x} = F(t) \\ \mathcal{Q}_\theta(t) = \mathbf{F}_1(t) \cdot \frac{\partial \mathbf{r}_1}{\partial \theta} + \mathbf{F}_2(t) \cdot \frac{\partial \mathbf{r}_2}{\partial \theta} = a \cos \theta(t)F(t) - mga \sin \theta(t) \end{cases}$$

**Lagrange's equation for  $x(t)$ :**

$$\begin{cases} \frac{\partial T}{\partial x} = 0 \\ \frac{\partial T}{\partial \dot{x}} = 2m\dot{x}(t) + ma\dot{\theta}(t) \cos \theta(t) \end{cases}$$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{x}} \right) - \frac{\partial T}{\partial x} = Q_x(t)$$

$$\Rightarrow \frac{d}{dt} [2m\dot{x}(t) + ma\dot{\theta}(t) \cos \theta(t)] = F(t)$$

$$2m\ddot{x}(t) + ma \cos \theta(t) \ddot{\theta}(t) - ma \sin \theta(t) \dot{\theta}^2(t) = F(t)$$

Lagrange's equation for  $\theta(t)$ :

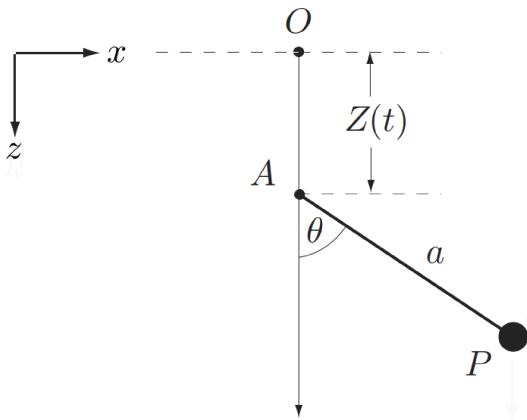
$$\begin{cases} \frac{\partial T}{\partial \dot{\theta}} = -ma \sin \theta(t) \dot{x}(t) \dot{\theta}(t) \\ \frac{\partial T}{\partial \dot{\theta}} = ma^2 \dot{\theta}(t) + ma \cos \theta(t) \dot{x}(t) \end{cases}$$

$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{\theta}} \right) - \frac{\partial T}{\partial \theta} = Q_\theta(t)$$

$$\frac{d}{dt} [ma^2 \dot{\theta}(t) + ma \cos \theta(t) \dot{x}(t)] + ma \sin \theta(t) \dot{x}(t) \dot{\theta}(t) = a \cos \theta(t) F(t) - mga \sin \theta(t)$$

$$ma^2 \ddot{\theta}(t) + ma \cos \theta(t) \ddot{x}(t) = a \cos \theta(t) F(t) - mga \sin \theta(t)$$

### Problem 7.6



A simple pendulum in which the pivot is made to move vertically so that its distance from the fixed origin at time  $t$  is  $Z(t) = Z_0 \cos \Omega t$ . The string is a light rigid rod of length  $a$  that cannot go slack.

**Generalized coordinate:**  $q_1 \equiv \theta$

Use Lagrange's equation to obtain equations of motion for  $\theta(t)$ .

### Solution

**Position vector:**

$$\mathbf{r}(t) = a \sin \theta(t) \hat{e}_x + [Z(t) + a \cos \theta(t)] \hat{e}_z$$

**Velocity vector:**

$$\dot{\mathbf{r}}(t) = a \dot{\theta}(t) \cos \theta(t) \hat{e}_x + [\dot{Z}(t) - a \dot{\theta}(t) \sin \theta(t)] \hat{e}_z$$

**Kinetic energy:**

$$\begin{aligned} T &\equiv T(\theta, \dot{\theta}, t) = \frac{m}{2} \dot{\mathbf{r}}(t) \cdot \dot{\mathbf{r}}(t) \\ &= \frac{m}{2} [a^2 \dot{\theta}^2(t) + \dot{Z}^2(t) - 2a \dot{Z}(t) \dot{\theta}(t) \sin \theta(t)] \end{aligned}$$

**Partial derivatives:**

$$\frac{\partial T}{\partial \theta} = -ma\dot{Z}(t)\dot{\theta}(t) \cos \theta(t), \quad \frac{\partial T}{\partial \dot{\theta}} = ma^2\dot{\theta}(t) - ma\dot{Z}(t) \sin \theta(t)$$

**Applied force and generalized force:**

$$\mathbf{F}(t) = mg\hat{e}_z \Rightarrow \mathcal{Q}_\theta(t) = \mathbf{F}(t) \cdot \frac{\partial \mathbf{r}}{\partial \theta} = -mga \sin \theta(t)$$

**Lagrange's equation:**

$$\begin{aligned} \frac{d}{dt} \left( \frac{\partial T}{\partial \dot{\theta}} \right) - \frac{\partial T}{\partial \theta} &= \mathcal{Q}_\theta(t) \\ \frac{d}{dt} [ma^2\dot{\theta}(t) - ma\dot{Z}(t) \sin \theta(t)] + ma\dot{Z}(t)\dot{\theta}(t) \cos \theta(t) &= -mga \sin \theta(t) \\ \Rightarrow \ddot{\theta}(t) - \frac{1}{a}\ddot{Z}(t) \sin \theta(t) &= -\frac{g}{a} \sin \theta(t) \end{aligned}$$

Substituting  $Z(t) = Z_0 \cos \Omega t$  gives  $\ddot{Z}(t) = -\Omega^2 Z_0 \cos(\Omega t)$ :

$$\ddot{\theta}(t) + \frac{\Omega^2 Z_0}{a} \cos(\Omega t) \sin \theta(t) + \frac{g}{a} \sin \theta(t) = 0$$