

# Worksheet 23

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Classical Mechanics II

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## Problem 9.7

Evaluate  $\{r, \mathbf{n} \cdot \mathbf{L}\}_{q,p}$  where  $r = x\hat{e}_x + y\hat{e}_y + z\hat{e}_z$  and  $\mathbf{n} = n_x\hat{e}_x + n_y\hat{e}_y + n_z\hat{e}_z$  is a constant vector.

**Angular momentum:**  $\mathbf{L} = \mathbf{r} \times \mathbf{p}$ , with components  $L_k = \sum_{i,j} \varepsilon_{ijk} x_i p_j$

### Solution

**Express the Poisson bracket:**

We want to calculate:

$$\{x_i, \mathbf{n} \cdot \mathbf{L}\}_{q,p}$$

for each component  $x_i$  of  $\mathbf{r}$ .

**Expand  $\mathbf{n} \cdot \mathbf{L}$ :**

$$\mathbf{n} \cdot \mathbf{L} = \sum_j n_j L_j = \sum_j n_j \sum_{r,s} \varepsilon_{jrs} x_r p_s$$

**Calculate Poisson bracket:**

$$\begin{aligned} \{x_i, \mathbf{n} \cdot \mathbf{L}\}_{q,p} &= \left\{ x_i, \sum_j \sum_{r,s} \varepsilon_{jrs} n_j x_r p_s \right\}_{q,p} \\ &= \sum_j \sum_{r,s} \varepsilon_{jrs} n_j \{x_i, x_r p_s\}_{q,p} \end{aligned}$$

**Use Leibniz rule:**

$$\{x_i, x_r p_s\}_{q,p} = x_r \{x_i, p_s\}_{q,p} + \{x_i, x_r\}_{q,p} p_s$$

**Apply fundamental Poisson brackets:**

$$\begin{cases} \{x_i, p_s\}_{q,p} = \delta_{is} \\ \{x_i, x_r\}_{q,p} = 0 \end{cases}$$

Therefore:

$$\{x_i, x_r p_s\}_{q,p} = x_r \delta_{is}$$

**Substitute back:**

$$\begin{aligned} \{x_i, \mathbf{n} \cdot \mathbf{L}\}_{q,p} &= \sum_j \sum_{r,s} \varepsilon_{jrs} n_j x_r \delta_{is} \\ &= \sum_j \sum_r \varepsilon_{jri} n_j x_r \end{aligned}$$

**Recognize as cross product:**

The sum  $\sum_{j,r} \varepsilon_{jri} n_j x_r$  is the  $i$ -th component of  $\mathbf{n} \times \mathbf{r}$ .

To see this, recall that the  $i$ -th component of  $\mathbf{n} \times \mathbf{r}$  is:

$$(\mathbf{n} \times \mathbf{r})_i = \sum_{j,r} \varepsilon_{ijr} n_j x_r$$

We have  $\sum_{j,r} \varepsilon_{jri} n_j x_r$ . Using the property  $\varepsilon_{jri} = \varepsilon_{ijr}$  (cyclic permutation):

$$\sum_{j,r} \varepsilon_{jri} n_j x_r = \sum_{j,r} \varepsilon_{ijr} n_j x_r = (\mathbf{n} \times \mathbf{r})_i$$

**Final answer:**

$$\{\mathbf{r}, \mathbf{n} \cdot \mathbf{L}\}_{\mathbf{q}, \mathbf{p}} = \mathbf{n} \times \mathbf{r}$$

**Physical interpretation:**

This result shows that the angular momentum  $\mathbf{L}$  generates infinitesimal rotations in phase space. Specifically, the Poisson bracket with  $\mathbf{n} \cdot \mathbf{L}$  gives the infinitesimal rotation of the position vector about the axis  $\mathbf{n}$ .

More formally, if we consider the canonical transformation generated by  $\mathbf{n} \cdot \mathbf{L}$ :

$$\delta \mathbf{r} = \varepsilon \{\mathbf{r}, \mathbf{n} \cdot \mathbf{L}\}_{\mathbf{q}, \mathbf{p}} = \varepsilon (\mathbf{n} \times \mathbf{r})$$

This is precisely an infinitesimal rotation by angle  $\varepsilon$  about the axis  $\mathbf{n}$ , confirming that angular momentum is the generator of rotations in classical mechanics.

### Problem 9.8

A projectile with mass  $m$  is moving on the vertical  $xy$ -plane in a uniform gravitational field.

**Hamiltonian:**

$$H(x, y, p_x, p_y, t) = \frac{p_x^2 + p_y^2}{2m} + mgy$$

**Given functions:**

$$\begin{cases} F_1 \equiv y - \frac{p_y t}{m} - \frac{1}{2} g t^2 \\ F_2 \equiv x - \frac{p_x t}{m} \end{cases}$$

Show that  $F_1$  and  $F_2$  are constants of motion and find three other constants of motion.

### Solution

**Hamilton's equations:**

From the Hamiltonian, we have:

$$\begin{cases} \dot{x} = \frac{\partial H}{\partial p_x} = \frac{p_x}{m}, & \dot{p}_x = -\frac{\partial H}{\partial x} = 0 \\ \dot{y} = \frac{\partial H}{\partial p_y} = \frac{p_y}{m}, & \dot{p}_y = -\frac{\partial H}{\partial y} = -mg \end{cases}$$

**Check if  $F_1$  is a constant of motion:**

A function  $F(\{q_i, p_i\}, t)$  is a constant of motion if:

$$\frac{dF}{dt} = \{F, H\}_{q,p} + \frac{\partial F}{\partial t} = 0$$

For  $F_1 = y - \frac{p_y t}{m} - \frac{gt^2}{2}$ :

Calculate the partial derivative:

$$\frac{\partial F_1}{\partial t} = -\frac{p_y}{m} - gt$$

Calculate the Poisson bracket:

$$\begin{aligned} \{F_1, H\}_{q,p} &= \frac{\partial F_1}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_1}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_1}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_1}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 0 \cdot \frac{p_x}{m} - 0 \cdot 0 + 1 \cdot \frac{p_y}{m} - \left(-\frac{t}{m}\right) \cdot (-mg) \\ &= \frac{p_y}{m} - \frac{t \cdot mg}{m} \\ &= \frac{p_y}{m} - gt \end{aligned}$$

Therefore:

$$\frac{dF_1}{dt} = \{F_1, H\}_{q,p} + \frac{\partial F_1}{\partial t} = \left(\frac{p_y}{m} - gt\right) + \left(-\frac{p_y}{m} - gt\right) = 0$$

$F_1$  is a constant of motion.

**Check if  $F_2$  is a constant of motion:**

For  $F_2 = x - \frac{p_x t}{m}$ :

$$\begin{aligned}\frac{\partial F_2}{\partial t} &= -\frac{p_x}{m} \\ \{F_2, H\}_{q,p} &= \frac{\partial F_2}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_2}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_2}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_2}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 1 \cdot \frac{p_x}{m} - \left(-\frac{t}{m}\right) \cdot 0 + 0 \cdot \frac{p_y}{m} - 0 \cdot (-mg) \\ &= \frac{p_x}{m}\end{aligned}$$

Therefore:

$$\frac{dF_2}{dt} = \{F_2, H\}_{q,p} + \frac{\partial F_2}{\partial t} = \frac{p_x}{m} - \frac{p_x}{m} = 0$$

$F_2$  is a constant of motion.

### Find three other constants of motion:

Constant 1: The Hamiltonian itself

$$F_3 \equiv H = \frac{p_x^2 + p_y^2}{2m} + mgy$$

Check:

$$\frac{\partial H}{\partial t} = 0 \quad (\text{no explicit time dependence})$$

$$\{H, H\}_{q,p} = 0 \quad (\text{Poisson bracket of any function with itself vanishes})$$

Therefore:

$$\frac{dH}{dt} = \{H, H\}_{q,p} + \frac{\partial H}{\partial t} = 0 + 0 = 0$$

$F_3 = H$  (total energy) is a constant of motion.

Constant 2: Horizontal momentum

$$F_4 \equiv p_x$$

From Hamilton's equations, we already know:

$$\dot{p}_x = -\frac{\partial H}{\partial x} = 0$$

$F_4 = p_x$  (horizontal momentum) is a constant of motion.

Constant 3: Modified vertical momentum

$$F_5 \equiv p_y + mgt$$

The total time derivative is:

$$\frac{dF_5}{dt} = \{F_5, H\}_{q,p} + \frac{\partial F_5}{\partial t}$$

Calculate the partial derivative with respect to time:

$$\frac{\partial F_5}{\partial t} = mg$$

Calculate the Poisson bracket:

$$\begin{aligned}\{F_5, H\}_{q,p} &= \frac{\partial F_5}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_5}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_5}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_5}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 0 \cdot \frac{p_x}{m} - 0 \cdot 0 + 0 \cdot \frac{p_y}{m} - 1 \cdot (mg) \\ &= -mg\end{aligned}$$

Therefore:

$$\frac{dF_5}{dt} = \{F_5, H\}_{q,p} + \frac{\partial F_5}{\partial t} = -mg + mg = 0$$

$F_5 = p_y + mgt$  is a constant of motion.

**Summary of five constants of motion:**

$$\left\{ \begin{array}{ll} F_1 = y - \frac{p_y t}{m} - \frac{gt^2}{2} & \text{(vertical position – free fall)} \\ F_2 = x - \frac{p_x t}{m} & \text{(horizontal position – uniform motion)} \\ F_3 = H = \frac{p_x^2 + p_y^2}{2m} + mgy & \text{(total energy)} \\ F_4 = p_x & \text{(horizontal momentum)} \\ F_5 = p_y + mgt & \text{(modified vertical momentum)} \end{array} \right.$$

**Physical interpretation:** These constants encode the initial conditions and conserved quantities of projectile motion.  $F_1$  and  $F_2$  relate to the initial position,  $F_4$  and  $F_5$  to the initial momentum, and  $F_3$  to the total energy.