

Worksheet 23

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Classical Mechanics II

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Problem 9.7

Evaluate $\{r, n \cdot L\}_{q,p}$ where $r = x\hat{e}_x + y\hat{e}_y + z\hat{e}_z$ and $n = n_x\hat{e}_x + n_y\hat{e}_y + n_z\hat{e}_z$ is a constant vector.

Angular momentum: $L = r \times p$, with components $L_k = \sum_{i,j} \varepsilon_{ijk} x_i p_j$

Solution

Express the Poisson bracket:

We want to calculate:

$$\{x_i, n \cdot L\}_{q,p}$$

for each component x_i of r .

Expand $n \cdot L$:

$$n \cdot L = \sum_j n_j L_j = \sum_j n_j \sum_{r,s} \varepsilon_{jrs} x_r p_s$$

Calculate Poisson bracket:

$$\begin{aligned} \{x_i, n \cdot L\}_{q,p} &= \left\{ x_i, \sum_j n_j \sum_{r,s} \varepsilon_{jrs} x_r p_s \right\}_{q,p} \\ &= \sum_j \sum_{r,s} \varepsilon_{jrs} n_j \{x_i, x_r p_s\}_{q,p} \end{aligned}$$

Use Leibniz rule:

$$\{x_i, x_r p_s\}_{q,p} = x_r \{x_i, p_s\}_{q,p} + \{x_i, x_r\}_{q,p} p_s$$

Apply fundamental Poisson brackets:

$$\begin{cases} \{x_i, p_s\}_{q,p} = \delta_{is} \\ \{x_i, x_r\}_{q,p} = 0 \end{cases}$$

Therefore:

$$\{x_i, x_r p_s\}_{q,p} = x_r \delta_{is}$$

Substitute back:

$$\begin{aligned} \{x_i, n \cdot L\}_{q,p} &= \sum_j \sum_{r,s} \varepsilon_{jrs} n_j x_r \delta_{is} \\ &= \sum_j \sum_r \varepsilon_{jri} n_j x_r \end{aligned}$$

Recognize as cross product:

The sum $\sum_{j,r} \varepsilon_{jri} n_j x_r$ is the i -th component of $n \times r$.

To see this, recall that the i -th component of $\mathbf{n} \times \mathbf{r}$ is:

$$(\mathbf{n} \times \mathbf{r})_i = \sum_{j,r} \varepsilon_{ijr} n_j x_r$$

We have $\sum_{j,r} \varepsilon_{jri} n_j x_r$. Using the property $\varepsilon_{jri} = \varepsilon_{ijr}$ (cyclic permutation):

$$\sum_{j,r} \varepsilon_{jri} n_j x_r = \sum_{j,r} \varepsilon_{ijr} n_j x_r = (\mathbf{n} \times \mathbf{r})_i$$

Final answer:

$$\{\mathbf{r}, \mathbf{n} \cdot \mathbf{L}\}_{q,p} = \mathbf{n} \times \mathbf{r}$$

Physical interpretation:

This result shows that the angular momentum \mathbf{L} generates infinitesimal rotations in phase space. Specifically, the Poisson bracket with $\mathbf{n} \cdot \mathbf{L}$ gives the infinitesimal rotation of the position vector about the axis \mathbf{n} .

More formally, if we consider the canonical transformation generated by $\mathbf{n} \cdot \mathbf{L}$:

$$\delta \mathbf{r} = \varepsilon \{\mathbf{r}, \mathbf{n} \cdot \mathbf{L}\}_{q,p} = \varepsilon (\mathbf{n} \times \mathbf{r})$$

This is precisely an infinitesimal rotation by angle ε about the axis \mathbf{n} , confirming that angular momentum is the generator of rotations in classical mechanics.

Problem 9.8

A projectile with mass m is moving on the vertical xy -plane in a uniform gravitational field.

Hamiltonian:

$$H(x, y, p_x, p_y, t) = \frac{p_x^2 + p_y^2}{2m} + mgy$$

Given functions:

$$\begin{cases} F_1 \equiv y - \frac{p_y t}{m} - \frac{1}{2}gt^2 \\ F_2 \equiv x - \frac{p_x t}{m} \end{cases}$$

Show that F_1 and F_2 are constants of motion and find three other constants of motion.

Solution

Hamilton's equations:

From the Hamiltonian, we have:

$$\begin{cases} \dot{x} = \frac{\partial H}{\partial p_x} = \frac{p_x}{m}, & \dot{p}_x = -\frac{\partial H}{\partial x} = 0 \\ \dot{y} = \frac{\partial H}{\partial p_y} = \frac{p_y}{m}, & \dot{p}_y = -\frac{\partial H}{\partial y} = -mg \end{cases}$$

Check if F_1 is a constant of motion:

A function $F(\{q_i, p_i\}, t)$ is a constant of motion if:

$$\frac{dF}{dt} = \{F, H\}_{q,p} + \frac{\partial F}{\partial t} = 0$$

For $F_1 = y - \frac{p_y t}{m} - \frac{gt^2}{2}$:

Calculate the partial derivative:

$$\frac{\partial F_1}{\partial t} = -\frac{p_y}{m} - gt$$

Calculate the Poisson bracket:

$$\begin{aligned} \{F_1, H\}_{q,p} &= \frac{\partial F_1}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_1}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_1}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_1}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 0 \cdot \frac{p_x}{m} - 0 \cdot 0 + 1 \cdot \frac{p_y}{m} - \left(-\frac{t}{m}\right) \cdot (-mg) \\ &= \frac{p_y}{m} - \frac{t \cdot mg}{m} \\ &= \frac{p_y}{m} - gt \end{aligned}$$

Therefore:

$$\frac{dF_1}{dt} = \{F_1, H\}_{q,p} + \frac{\partial F_1}{\partial t} = \left(\frac{p_y}{m} - gt\right) + \left(-\frac{p_y}{m} - gt\right) = 0$$

F_1 is a constant of motion.

Check if F_2 is a constant of motion:

For $F_2 = x - \frac{p_x t}{m}$:

$$\begin{aligned}\frac{\partial F_2}{\partial t} &= -\frac{p_x}{m} \\ \{F_2, H\}_{q,p} &= \frac{\partial F_2}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_2}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_2}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_2}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 1 \cdot \frac{p_x}{m} - \left(-\frac{t}{m}\right) \cdot 0 + 0 \cdot \frac{p_y}{m} - 0 \cdot (-mg) \\ &= \frac{p_x}{m}\end{aligned}$$

Therefore:

$$\frac{dF_2}{dt} = \{F_2, H\}_{q,p} + \frac{\partial F_2}{\partial t} = \frac{p_x}{m} - \frac{p_x}{m} = 0$$

F_2 is a constant of motion.

Find three other constants of motion:

Constant 1: The Hamiltonian itself

$$F_3 \equiv H = \frac{p_x^2 + p_y^2}{2m} + mgy$$

Check:

$$\frac{\partial H}{\partial t} = 0 \quad (\text{no explicit time dependence})$$

$$\{H, H\}_{q,p} = 0 \quad (\text{Poisson bracket of any function with itself vanishes})$$

Therefore:

$$\frac{dH}{dt} = \{H, H\}_{q,p} + \frac{\partial H}{\partial t} = 0 + 0 = 0$$

$F_3 = H$ (total energy) is a constant of motion.

Constant 2: Horizontal momentum

$$F_4 \equiv p_x$$

From Hamilton's equations, we already know:

$$\dot{p}_x = -\frac{\partial H}{\partial x} = 0$$

$F_4 = p_x$ (horizontal momentum) is a constant of motion.

Constant 3: Modified vertical momentum

$$F_5 \equiv p_y + mgt$$

The total time derivative is:

$$\frac{dF_5}{dt} = \{F_5, H\}_{q,p} + \frac{\partial F_5}{\partial t}$$

Calculate the partial derivative with respect to time:

$$\frac{\partial F_5}{\partial t} = mg$$

Calculate the Poisson bracket:

$$\begin{aligned}\{F_5, H\}_{\mathbf{q}, \mathbf{p}} &= \frac{\partial F_5}{\partial x} \frac{\partial H}{\partial p_x} - \frac{\partial F_5}{\partial p_x} \frac{\partial H}{\partial x} + \frac{\partial F_5}{\partial y} \frac{\partial H}{\partial p_y} - \frac{\partial F_5}{\partial p_y} \frac{\partial H}{\partial y} \\ &= 0 \cdot \frac{p_x}{m} - 0 \cdot 0 + 0 \cdot \frac{p_y}{m} - 1 \cdot (mg) \\ &= -mg\end{aligned}$$

Therefore:

$$\frac{dF_5}{dt} = \{F_5, H\}_{\mathbf{q}, \mathbf{p}} + \frac{\partial F_5}{\partial t} = -mg + mg = 0$$

$$F_5 = p_y + mgt \text{ is a constant of motion.}$$

Summary of five constants of motion:

$$\left\{ \begin{array}{ll} F_1 = y - \frac{p_y t}{m} - \frac{gt^2}{2} & \text{(vertical position – free fall)} \\ F_2 = x - \frac{p_x t}{m} & \text{(horizontal position – uniform motion)} \\ F_3 = H = \frac{p_x^2 + p_y^2}{2m} + mgy & \text{(total energy)} \\ F_4 = p_x & \text{(horizontal momentum)} \\ F_5 = p_y + mgt & \text{(modified vertical momentum)} \end{array} \right.$$

Physical interpretation: These constants encode the initial conditions and conserved quantities of projectile motion. F_1 and F_2 relate to the initial position, F_4 and F_5 to the initial momentum, and F_3 to the total energy.