

# Worksheet 24

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Classical Mechanics II

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## Problem 9.9

Obtain the canonical transformation generated by  $\Lambda_1(q, Q, t) = qQ$  and the Kamiltonian equations of motion.

**Harmonic oscillator Hamiltonian:**

$$H(q, p) = \frac{1}{2}m\omega^2 q^2 + \frac{p^2}{2m}$$

## Solution

**Type 1 generating function:**

Given  $\Lambda_1(q, Q, t) = qQ$ .

For a Type 1 generating function  $F_1(q, Q, t)$ , the canonical transformation is obtained from:

$$p = \frac{\partial F_1}{\partial q} \quad \text{and} \quad P = -\frac{\partial F_1}{\partial Q}$$

**Derive transformation equations:**

$$\Lambda_1(q, Q, t) = qQ$$

$$p = \frac{\partial \Lambda_1}{\partial q} = Q$$

$$P = -\frac{\partial \Lambda_1}{\partial Q} = -q$$

To find the inverse transformation  $(q, p) \rightarrow (Q, P)$ , we solve:

$$Q = p$$

and

$$P = -q$$

, which gives:

$$Q(q, p) = p$$

and

$$P(q, p) = -q$$

However, the standard convention is to express  $(Q, P)$  in terms of  $(q, p)$ . Using the phase space Lagrangian gauge transformation:

$$P\dot{Q} - K = p\dot{q} - H + \frac{d\Lambda_1}{dt}$$

where:

$$\frac{d\Lambda_1}{dt} = \frac{\partial \Lambda_1}{\partial q}\dot{q} + \frac{\partial \Lambda_1}{\partial Q}\dot{Q} + \frac{\partial \Lambda_1}{\partial t} = Q\dot{q} + q\dot{Q}$$

Substituting:

$$P\dot{Q} - K = p\dot{q} - H + Q\dot{q} + q\dot{Q}$$

Rearranging:

$$P\dot{Q} - K = (p + Q)\dot{q} + q\dot{Q} - H$$

For this to hold for all  $\dot{q}$  and  $\dot{Q}$ , we must have:

$$\begin{cases} p + Q = 0 & \Rightarrow & Q = -p \\ P - q = 0 & \Rightarrow & P = q \end{cases}$$

Therefore:

$$\begin{cases} Q(q, p, t) = -p \\ P(q, p, t) = q \end{cases}$$

**Kamiltonian:**

$$K(Q, P, t) = H(q(Q, P), p(Q, P), t) - \frac{\partial \Lambda_1}{\partial t}$$

Since  $\frac{\partial \Lambda_1}{\partial t} = 0$ ,  $q = P$ , and  $p = -Q$ :

$$K(Q, P, t) = \frac{1}{2}m\omega^2 P^2 + \frac{Q^2}{2m}$$

$$K(Q, P, t) = \frac{1}{2}m\omega^2 P^2 + \frac{Q^2}{2m}$$

**Hamilton's equations in new coordinates:**

$$\begin{cases} \dot{Q} = \frac{\partial K}{\partial P} = m\omega^2 P \\ \dot{P} = -\frac{\partial K}{\partial Q} = -\frac{Q}{m} \end{cases}$$

Taking time derivatives:

$$\ddot{Q} = m\omega^2 \dot{P} = m\omega^2 \left(-\frac{Q}{m}\right) = -\omega^2 Q$$

$$\ddot{P} = -\frac{\dot{Q}}{m} = -\frac{m\omega^2 P}{m} = -\omega^2 P$$

$$\begin{cases} \dot{Q} = m\omega^2 P \\ \dot{P} = -\frac{Q}{m} \end{cases} \Rightarrow \begin{cases} \ddot{Q} = -\omega^2 Q \\ \ddot{P} = -\omega^2 P \end{cases}$$

Both  $Q$  and  $P$  satisfy the harmonic oscillator equation. The transformation exchanges the roles of coordinate and momentum, revealing the symmetry  $H(q, p) = K(P, Q)$  (up to constants).

**Problem 9.10**

Solve for  $q(t)$  and  $p(t)$  via  $Q(t)$  and  $P(t)$  using the canonical transformation:

$$\begin{cases} q(Q, P, t) = \sqrt{2\frac{P}{m\omega}} \sin Q \\ p(Q, P, t) = \sqrt{2m\omega P} \cos Q \end{cases}$$

with Kamiltonian  $K(Q, P, t) = \omega P$ .

**Solution**

**Verify canonicity via Poisson brackets:**

The transformation is canonical if:

$$\{q, p\}_{Q,P} = 1, \quad \{q, q\}_{Q,P} = 0, \quad \{p, p\}_{Q,P} = 0$$

Compute partial derivatives:

$$\begin{cases} \frac{\partial q}{\partial Q} = \sqrt{2\frac{P}{m\omega}} \cos Q \\ \frac{\partial q}{\partial P} = \frac{1}{\sqrt{2m\omega P}} \sin Q \\ \frac{\partial p}{\partial Q} = -\sqrt{2m\omega P} \sin Q \\ \frac{\partial p}{\partial P} = \sqrt{m\frac{\omega}{2P}} \cos Q \end{cases}$$

Check  $\{q, p\}_{Q,P}$ :

$$\begin{aligned} \{q, p\}_{Q,P} &= \frac{\partial q}{\partial Q} \frac{\partial p}{\partial P} - \frac{\partial q}{\partial P} \frac{\partial p}{\partial Q} \\ &= \sqrt{2\frac{P}{m\omega}} \cos Q \cdot \sqrt{m\frac{\omega}{2P}} \cos Q - \frac{1}{\sqrt{2m\omega P}} \sin Q \cdot (-\sqrt{2m\omega P} \sin Q) \\ &= \cos^2 Q + \sin^2 Q = 1 \quad \checkmark \end{aligned}$$

$\{q, q\}_{Q,P} = 0$  and  $\{p, p\}_{Q,P} = 0$  trivially. The transformation is canonical.

**Solve Kamiltonian equations:**

$$K(Q, P, t) = \omega P$$

$$\begin{cases} \dot{Q} = \frac{\partial K}{\partial P} = \omega \\ \dot{P} = -\frac{\partial K}{\partial Q} = 0 \end{cases}$$

Integrating:

$$\begin{cases} Q(t) = \omega t + Q(0) \\ P(t) = P(0) \end{cases}$$

**Transform back to  $(q, p)$ :**

$$\begin{cases} q(t) = \sqrt{2\frac{P(0)}{m\omega}} \sin[\omega t + Q(0)] \\ p(t) = \sqrt{2m\omega P(0)} \cos[\omega t + Q(0)] \end{cases}$$

$$\begin{cases} q(t) = \sqrt{2\frac{P(0)}{m\omega}} \sin[\omega t + Q(0)] \\ p(t) = \sqrt{2m\omega P(0)} \cos[\omega t + Q(0)] \end{cases}$$

**Determine constants from initial conditions:**

At  $t = 0$ :

$$\begin{cases} q(0) = \sqrt{2\frac{P(0)}{m\omega}} \sin Q(0) \\ p(0) = \sqrt{2m\omega P(0)} \cos Q(0) \end{cases}$$

Squaring and adding:

$$q(0)^2 + \frac{p(0)^2}{m^2\omega^2} = \frac{2P(0)}{m\omega} [\sin^2 Q(0) + \cos^2 Q(0)] = \frac{2P(0)}{m\omega}$$

$$P(0) = \frac{m\omega}{2} \left[ q(0)^2 + \frac{p(0)^2}{m^2\omega^2} \right] = \frac{m\omega q(0)^2}{2} + \frac{p(0)^2}{2m\omega}$$

Dividing the equations:

$$\tan Q(0) = \frac{m\omega q(0)}{p(0)} \Rightarrow Q(0) = \arctan\left(\frac{m\omega q(0)}{p(0)}\right)$$

**Compact form:**

Define amplitude  $A = \sqrt{2\frac{P(0)}{m\omega}}$  and phase  $\varphi = Q(0)$ :

$$\begin{cases} q(t) = A \sin(\omega t + \varphi) \\ p(t) = m\omega A \cos(\omega t + \varphi) \end{cases}$$

where:

$$\begin{cases} A = \sqrt{q(0)^2 + \frac{p(0)^2}{m^2\omega^2}} \\ \varphi = \arctan\left(\frac{m\omega q(0)}{p(0)}\right) \end{cases}$$

The canonical transformation converts the problem to one with cyclic coordinate  $Q$ , where  $P = \frac{E}{\omega}$  is conserved (proportional to energy). The coordinate  $Q$  evolves uniformly in time, giving the familiar sinusoidal solution.