



# INDIAN INSTITUTE OF TECHNOLOGY MADRAS

A

Roll No.

M M I 6 B 0 2 3

Name :

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Total No. of Pages

Quiz I



Quiz II/ Mid-Sem



End-Semester



Make-up



Date : Feb 7, 2020

Semester & Degree :

Course No.

Part :

Question No.	1	2	3	4	5	6	7	8	9	10
Marks	10	11	8	4	10					
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										33/50

Answer on both sides of the paper including the space below

$$\textcircled{1} (a) \quad \psi_2 \times \left[ -\frac{\hbar^2}{2m} \frac{d^2 \psi_1}{dx^2} + V \psi_1 \right] = E \psi_1 \quad - (1)$$

$$\psi_1 \times \left[ -\frac{\hbar^2}{2m} \frac{d^2 \psi_2}{dx^2} + V \psi_2 \right] = E \psi_2 \quad - (2)$$

Subtracting the 2 eqns,

$$-\frac{\hbar^2}{2m} \left[ \psi_2 \frac{d^2 \psi_1}{dx^2} - \psi_1 \frac{d^2 \psi_2}{dx^2} \right] = E \psi_1 \psi_2 - E \psi_1 \psi_2 = 0 \quad - (3)$$

$$\text{Also, } \frac{d}{dx} \left[ \psi_2 \frac{d\psi_1}{dx} - \psi_1 \frac{d\psi_2}{dx} \right]$$

$$= \psi_2 \frac{d^2 \psi_1}{dx^2} + \frac{d\psi_2}{dx} \frac{d\psi_1}{dx} - \frac{d\psi_1}{dx} \frac{d\psi_2}{dx} - \psi_1 \frac{d^2 \psi_2}{dx^2}$$

$$\Rightarrow -\frac{\hbar^2}{2m} \frac{d}{dx} \left[ \psi_2 \frac{d\psi_1}{dx} - \psi_1 \frac{d\psi_2}{dx} \right] = 0$$

$$\rightarrow \frac{d}{dn} \left[ \psi_1 \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_1}{dx} \right] = 0$$

$$\rightarrow \psi_1 \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_1}{dx} = W \quad (\text{const})$$

(b) If the wave  $f^{ns}$  are normalizable,

$$\psi \rightarrow 0 \quad \text{as } x \rightarrow \infty$$

$$\Rightarrow \left( \psi_1 \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_1}{dx} \right) \rightarrow (0 - 0 = 0) \quad \text{as } x \rightarrow \infty$$

But since  $W$  is a constant, it must hold for all  $x$

$$\rightarrow W = \psi_1 \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_1}{dx} = 0$$

(c)

$$\psi_1 \frac{d\psi_2}{dx} = \psi_2 \frac{d\psi_1}{dx}$$

$$\Rightarrow \frac{1}{\psi_1} \frac{d}{dx} [\psi_1] = \frac{1}{\psi_2} \frac{d}{dx} [\psi_2]$$

$$\rightarrow \int \frac{d\psi_1}{\psi_1} = \int \frac{d\psi_2}{\psi_2} + C$$

$$\rightarrow \ln \psi_1 = \ln \psi_2 + C$$

$$\Rightarrow \ln \left( \frac{\psi_1}{\psi_2} \right) = C \rightarrow \psi_1 = e^C \psi_2 = k \psi_2$$

where  $k$  is a constant

$$\boxed{\psi_1 = k\psi_2}$$

Since  $\psi_1, \psi_2$  are linearly dep, they aren't really unique and therefore the solutions are not distinct

③

a)  $A^+ = \begin{pmatrix} a & 0 & 0 \\ 0 & -a & 0 \\ 0 & a & -a \end{pmatrix} = A$  ✓

$B^+ = \begin{pmatrix} b & 0 & 0 \\ 0 & 0 & ib \\ 0 & -ib & 0 \end{pmatrix}^T = \begin{pmatrix} b & 0 & 0 \\ 0 & 0 & -ib \\ 0 & ib & 0 \end{pmatrix} = B$  ✓

Both  $\hat{A}, \hat{B}$  are Hermitian

→ They do represent observables

b) ~~No, they cannot be measured simultaneously~~

$AB = \begin{pmatrix} a & 0 & 0 \\ 0 & -a & 0 \\ 0 & 0 & -a \end{pmatrix} \begin{pmatrix} b & 0 & 0 \\ 0 & 0 & -ib \\ 0 & ib & 0 \end{pmatrix}$

$$= \begin{pmatrix} ab & 0 & 0 \\ 0 & 0 & iab \\ 0 & -iab & 0 \end{pmatrix}$$

$$BA = \begin{pmatrix} b & 0 & 0 \\ 0 & 0 & -ib \\ 0 & ib & 0 \end{pmatrix} \begin{pmatrix} a & 0 & 0 \\ 0 & -a & 0 \\ 0 & 0 & -a \end{pmatrix}$$

$$= \begin{pmatrix} ab & 0 & 0 \\ 0 & 0 & iab \\ 0 & -iab & 0 \end{pmatrix}$$

$$AB = AB$$

$$\Rightarrow [A, B] = 0 \quad \text{null matrix}$$

Since commutator is zero, the observables can be measured simultaneously

(c)

$$\det(A - \lambda I) = (a - \lambda)(-a - \lambda)(-a - \lambda) = 0$$

$$\Rightarrow \lambda = -a, -a, a \quad \text{are the eigen values}$$

$$\det \begin{pmatrix} b - \lambda & 0 & 0 \\ 0 & -\lambda & -ib \\ 0 & ib & -\lambda \end{pmatrix} = (b - \lambda)(\lambda^2 + i^2 b^2) = 0$$

$$\lambda = b, -b, b$$

$$a = 1, \quad b = -1$$

$$\begin{pmatrix} -a \\ -a \\ a \end{pmatrix} \quad \begin{pmatrix} b \\ b \\ -b \end{pmatrix}$$

$\Rightarrow$  the eigen values are degenerate



(d)

$$\begin{pmatrix} a & 0 & 0 \\ 0 & -a & 0 \\ 0 & 0 & -a \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \\ \gamma \end{pmatrix} = a \begin{pmatrix} \alpha \\ \beta \\ \gamma \end{pmatrix}$$

$$\Rightarrow \begin{aligned} a\alpha &= a\alpha \\ -a\beta &= a\beta \\ -a\gamma &= a\gamma \end{aligned} \Rightarrow \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

another vector  $u = \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$  why?

For  $B$ , it is the same

2

$$v_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad v_2 = \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$$

Gram Schmidt

$$v_2 = x_2 - \frac{\langle x_2, v_1 \rangle}{\langle v_1, v_1 \rangle} v_1$$

Something  
not right!

$$\begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix} - 0 \Rightarrow \tilde{v}_2 = \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}$$

$$\left\{ \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix}, \begin{pmatrix} 0 \\ 1/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} \right\}$$

Speak to  
me, if  
needed!

(2)

$$\hat{T}(a) = \exp\left(\frac{-i\hat{p} \cdot a}{\hbar}\right)$$

(a)

$$[\hat{x}_i, \hat{T}(a)] f$$

$$= x_i e^{-\frac{i(p \cdot a)}{\hbar}} f - \left( e^{-\frac{i\hbar \frac{d}{dx_i} (x_i) \cdot a}{\hbar}} f \right)$$

$$= x_i e^{-\frac{i(p \cdot a)}{\hbar}} f - i \left( -i\hbar \frac{d}{dx_i} \right) e^{-\frac{i(p \cdot a)}{\hbar}} f$$

$$= x_i e^{-\frac{i(p \cdot a)}{\hbar}} f - \hbar \frac{d}{dx_i} e^{-\frac{i(p \cdot a)}{\hbar}} f$$

[ ~~$x_i$~~ ,

$$\hat{p} = (p_1, p_2, p_3)$$

$$x = (x_1, x_2, x_3)$$

$$\left[ x_i, e^{-\frac{i\hat{p} \cdot a}{\hbar}} \right]$$

$$= x_i e^{-\frac{i\hat{p} \cdot a}{\hbar}} - \frac{i\hbar \frac{d}{dx_i} (x_i)}{1} e^{-\frac{i\hat{p} \cdot a}{\hbar}}$$

$$[\hat{x}_i, \hat{T}(a)] = x_i e^{-\frac{\hat{p} \cdot a}{\hbar}} e^{-\frac{i(-i)\hbar}{\hbar}}$$

$$= x_i e^{-\frac{\hat{p} \cdot a}{\hbar}} e^{-\left[\frac{\hbar}{\hbar} \frac{d}{dx} + \frac{\hbar}{\hbar} \frac{d}{dy} + \frac{\hbar}{\hbar} \frac{d}{dz}\right] x_i}$$

$$= x_i e^{-\frac{\hat{p} \cdot a}{\hbar}} e^{-1}$$

(b)

$$\langle \bar{\psi} | \hat{n} | \bar{\psi} \rangle$$

$$= c \langle \bar{\psi} | n_T - T n | \bar{\psi} \rangle$$

$$= c \langle \bar{\psi} | n | \bar{\psi} \rangle - c \langle \bar{\psi} | n | \bar{\psi} \rangle$$

$$= c (\langle \bar{\psi} | \hat{n} | \bar{\psi} \rangle - \langle \bar{\psi} | n | \bar{\psi} \rangle)$$

$$\langle \bar{\psi} | \hat{n} | \bar{\psi} \rangle \quad ?$$

$$\begin{aligned}
 \textcircled{4} \text{ a) } \langle \hat{x} \rangle &= \langle \alpha | x | \alpha \rangle = \sqrt{\frac{\hbar}{2m\omega}} \langle \alpha | a_+ + a_- | \alpha \rangle \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \left( \langle \alpha | a_- | \alpha \rangle + \langle \alpha | a_+ | \alpha \rangle \right) \\
 &= \sqrt{\frac{\hbar}{2m\omega}} \left( \alpha + \alpha^* \right) \checkmark
 \end{aligned}$$

$$\begin{aligned}
 \langle \hat{p}_x \rangle &= \langle \alpha | p | \alpha \rangle = i \sqrt{\frac{\hbar m \omega}{2}} \langle \alpha | a_+ - a_- | \alpha \rangle \\
 &= i \sqrt{\frac{\hbar m \omega}{2}} \left( \langle \alpha | a_- | \alpha \rangle - \langle \alpha | a_+ | \alpha \rangle \right) \\
 &= -i \sqrt{\frac{\hbar m \omega}{2}} \left( \alpha - \alpha^* \right) \checkmark
 \end{aligned}$$

(b)

$$\hat{a} | 0 \rangle = 0$$

↳ ground state

$$\left( \frac{x}{\sqrt{\frac{\hbar}{2m\omega}}} + \frac{ip}{\sqrt{2m\omega\hbar}} \right) | 0 \rangle = 0$$

$$\Rightarrow (m\omega x + ip) | 0 \rangle = 0$$

$$\begin{aligned}
 \sigma_1(x) &= \langle x | 0 \rangle = \langle x | \left( m\omega x + \hbar \frac{d}{dx} \right) | 0 \rangle = 0 \\
 \Rightarrow m\omega x \sigma_1(x) + \hbar \frac{d}{dx} \sigma_1(x) &= 0
 \end{aligned}$$



$$\psi_0(x) = A \exp\left(-\frac{m\omega x^2}{2\hbar}\right)$$

$$\psi_\alpha(x) = A_\alpha H_\alpha\left(x \sqrt{\frac{m\omega}{2\hbar}}\right) e^{-\frac{m\omega x^2}{2\hbar}}$$

$$\psi_\alpha(x) = \langle x | \hat{a}^\dagger | \alpha \rangle$$

$$\psi_\alpha(n) = \langle n | \alpha \rangle$$

$$= \frac{\langle n | \hat{a}^\dagger | \alpha \rangle}{\alpha^n}$$

$$\textcircled{1} \quad c_n = \langle \psi_n | \alpha \rangle = \frac{1}{\sqrt{n!}} \langle (a^\dagger)^n \psi_0 | \alpha \rangle$$

$$= \frac{1}{\sqrt{n!}} \alpha^n \langle 0 | \alpha \rangle = \frac{\alpha^n}{n!} c_0$$

$$1 = \sum_{n=0}^{\infty} |c_n|^2 = |c_0|^2 \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!}$$

$$\Rightarrow c_0 = e^{-|\alpha|^2/2}$$

$$\Rightarrow \psi_\alpha(n) = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle$$

$$\hat{x} \sqrt{\frac{2m\omega}{\hbar}} = \alpha + \alpha^\dagger, \quad \hat{p} \sqrt{\frac{2}{\hbar m\omega}} = \alpha - \alpha^\dagger$$

$$\Rightarrow \psi_\alpha(x) = \sum_{n=0}^{\infty} \frac{\left( \hat{x} \sqrt{\frac{2m\omega}{\hbar}} + \hat{p} \sqrt{\frac{2}{\hbar m\omega}} \right)^n}{2^n \sqrt{n!}} |n\rangle$$

(5)

$$\Delta A^2 \Delta B^2 \geq \frac{1}{4} \left| \langle [\hat{A}, \hat{B}] \rangle \right|^2$$

(a)

$$\Delta S_x^2 \Delta S_y^2 \geq \frac{1}{4} \left| \langle [\hat{S}_x, \hat{S}_y] \rangle \right|^2$$

$$|+\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |-\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$|\chi\rangle = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} a \\ b \end{pmatrix}$$

$$|a|^2 + |b|^2 = 1$$

So without loss of generality,

$$a = \cos \frac{\theta}{2}, \quad b = \sin \frac{\theta}{2} e^{i\phi}$$

$$\rightarrow \chi = \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$\langle S_x \rangle = \langle \chi | S_x | \chi \rangle$$

$$= \begin{pmatrix} \cos \frac{\theta}{2} & \sin \frac{\theta}{2} \end{pmatrix} \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$= \begin{pmatrix} \cos \frac{\theta}{2} & \sin \frac{\theta}{2} \end{pmatrix} \begin{pmatrix} \sin \frac{\theta}{2} \\ \cos \frac{\theta}{2} \end{pmatrix} \frac{\hbar}{2} = \frac{\hbar}{2} \sin \theta$$

$$S_x^2 = \frac{\hbar^2}{4} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\langle S_x^2 \rangle = \frac{\hbar^2}{4} \text{ for spin } \frac{1}{2}$$

Since  $\langle S_x^2 \rangle = \langle S_y^2 \rangle = \langle S_z^2 \rangle$ ,

$$\langle S_x^2 \rangle + \langle S_y^2 \rangle + \langle S_z^2 \rangle = \frac{3\hbar^2}{4}$$

$$\langle S_y \rangle = \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & \sin \frac{\theta}{2} \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos \frac{\theta}{2} \\ \sin \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & \sin \frac{\theta}{2} \\ i & 0 \end{pmatrix} \begin{pmatrix} -i \sin \frac{\theta}{2} \\ i \cos \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \left( -i \sin \frac{\theta}{2} \cos \frac{\theta}{2} + i \sin \frac{\theta}{2} \cos \frac{\theta}{2} \right) = 0$$

Also  $\langle S_y^2 \rangle = \frac{\hbar^2}{4}$

$$\Rightarrow \Delta S_x^2 = \langle S_x^2 \rangle - \langle S_x \rangle^2$$

$$= \frac{\hbar^2}{4} - \frac{\hbar^2}{4} \sin^2 \theta = \frac{\hbar^2}{4} (1 - \sin^2 \theta)$$

$$\Delta S_y^2 = \langle S_y^2 \rangle - \langle S_y \rangle^2 = \frac{\hbar^2}{4} - 0 = \frac{\hbar^2}{4}$$

$$\rightarrow \Delta S_x^2 \Delta S_y^2 = \frac{\hbar^2}{4} (1 - \sin^2 \theta) \frac{\hbar^2}{4}$$

$$= \frac{\hbar^4}{16} (1 - \sin^2 \theta)$$

$$\langle S_x \rangle = \langle \chi^\dagger | S_x | \chi \rangle$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & e^{-i\phi} \sin \frac{\theta}{2} \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & e^{-i\phi} \sin \frac{\theta}{2} \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{i\phi} \sin \frac{\theta}{2} \\ \cos \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \left( e^{i\phi} \sin \frac{\theta}{2} \cos \frac{\theta}{2} + e^{-i\phi} \sin \frac{\theta}{2} \cos \frac{\theta}{2} \right)$$

$$= \frac{\hbar}{2} \cancel{e^{i\phi}} \left( \sin \frac{\theta}{2} \right) \left( \cos \frac{\theta}{2} \right) \left( 2 \cos \cancel{\phi} \right)$$

$$S_x^2 = \frac{\hbar^2}{4} I \Rightarrow \langle S_x^2 \rangle = \frac{\hbar^2}{4}$$

$$\Delta S_x^2 = \langle S_x^2 \rangle - \langle S_x \rangle^2$$

$$= \frac{\hbar^2}{4} - \frac{\hbar^2}{4} \sin^2 \theta \cos^2 \theta$$

$$= \frac{\hbar^2}{4} (1 - \sin^2 \theta \cos^2 \theta)$$





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Additional Sheet

$$\Delta S_y^2 = \langle S_y^2 \rangle - \langle S_y \rangle^2 = \frac{\hbar^2}{4} - \langle S_y \rangle^2$$

$$\langle S_y \rangle = \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & e^{-i\phi} \sin \frac{\theta}{2} \\ 0 & -i \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \begin{pmatrix} \cos \frac{\theta}{2} & e^{-i\phi} \sin \frac{\theta}{2} \\ 0 & -i \end{pmatrix} \begin{pmatrix} -ie \sin \frac{\theta}{2} \\ i \cos \frac{\theta}{2} \end{pmatrix}$$

$$= \frac{\hbar}{2} \left[ -ie \sin \frac{\theta}{2} \cos \frac{\theta}{2} + ie \sin \frac{\theta}{2} \cos \frac{\theta}{2} \right]$$

$$= \frac{\hbar}{2} i \frac{\sin \theta}{2} \left[ e^{-i\phi} - e^{i\phi} \right]$$

$$= \frac{\hbar}{2} i \frac{\sin \theta}{2} \left( -2i \sin \phi \right)$$

$$= \frac{\hbar}{2} \sin^2 \theta \sin \phi (-i^2) = \frac{\hbar}{2} \sin^2 \theta \sin \phi$$

$$\rightarrow \Delta S_y^2 = \frac{\hbar^2}{4} \left( 1 - \frac{\sin^2 \theta}{\sin^2 \theta} \right)$$

$$\Rightarrow \Delta S_x^2 \Delta S_y^2 = \frac{\hbar^4}{16} \left( 1 - \sin^2 \theta \cos^2 \theta \right) \left( 1 - \sin^2 \theta \right)$$

$$= \frac{\hbar^4}{16} \left( 1 - t(1-t) \right) \left( 1 - t^2 \right) \quad \text{where } t = \sin^2 \theta$$

$$f(t) = (1-t^2)(1-t+t^2) = 1-t+t^2 - t^3 + t^4$$

$$f'(t) = 4t^3 + 3t^2 - 1 = 0$$

Product is max when

$$1 - \sin^2 \theta \cos^2 \theta = 1 - \sin^4 \theta$$

$$\Rightarrow \sin^2 \theta \cos^2 \theta = \sin^4 \theta$$

$$\Rightarrow \tan^2 \theta = 1$$

$$\Rightarrow \theta = 45^\circ$$

$$\Rightarrow |X\rangle = \cos(45/2)$$

$$\Delta S_x^2 \Delta S_y^2 = \frac{\hbar^4}{16} \left( 1 - \sin^2 \theta \cos^2 \theta \right) \left( 1 - \sin^2 \theta \sin^2 \theta \right)$$

Product is max when terms in the bracket are

$$\Rightarrow 1 - \sin^2 \theta \cos^2 \theta = 1 - \sin^2 \theta \sin^2 \theta \quad \text{why?}$$

$$\Rightarrow \tan \theta = 1 \Rightarrow \theta = 45^\circ$$

$$\Rightarrow \left. \Delta S_x^2 \Delta S_y^2 \right|_{\text{max}} = \frac{\hbar^4}{16} \left( 1 - \frac{\sin^2 \theta}{2} \right)^2$$

this is max when  $\theta = 0$

$$\Rightarrow (\theta, \theta) = (0, 45^\circ) = \left(0, \frac{\pi}{4}\right)$$

$$\Rightarrow X = \begin{pmatrix} 1 & \\ e^{i\pi/4} & \sin \theta \end{pmatrix} = \begin{pmatrix} 1 & \\ & 0 \end{pmatrix}$$

for  $X$ ,  $\theta \rightarrow \text{any value}$

$$\Delta S_x^2 = \frac{\hbar^2}{4}, \quad \Delta S_y^2 = \frac{\hbar^2}{4}$$

$$[S_x, S_y] = S_z$$

$$\Rightarrow \langle [S_x, S_y] \rangle = |\langle S_z \rangle|^2 = \frac{\hbar^2}{4}$$

$$\Rightarrow \frac{1}{\hbar} |\langle [\hat{S}_x, \hat{S}_y] \rangle|^2 = \frac{\hbar^2}{\hbar}$$

$$\langle S_z \rangle = \frac{\hbar}{4} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\hbar}{4}$$

$$\Rightarrow |\langle S_z \rangle|^2 = \frac{\hbar^2}{4}$$

$\Rightarrow$  Equality is satisfied.