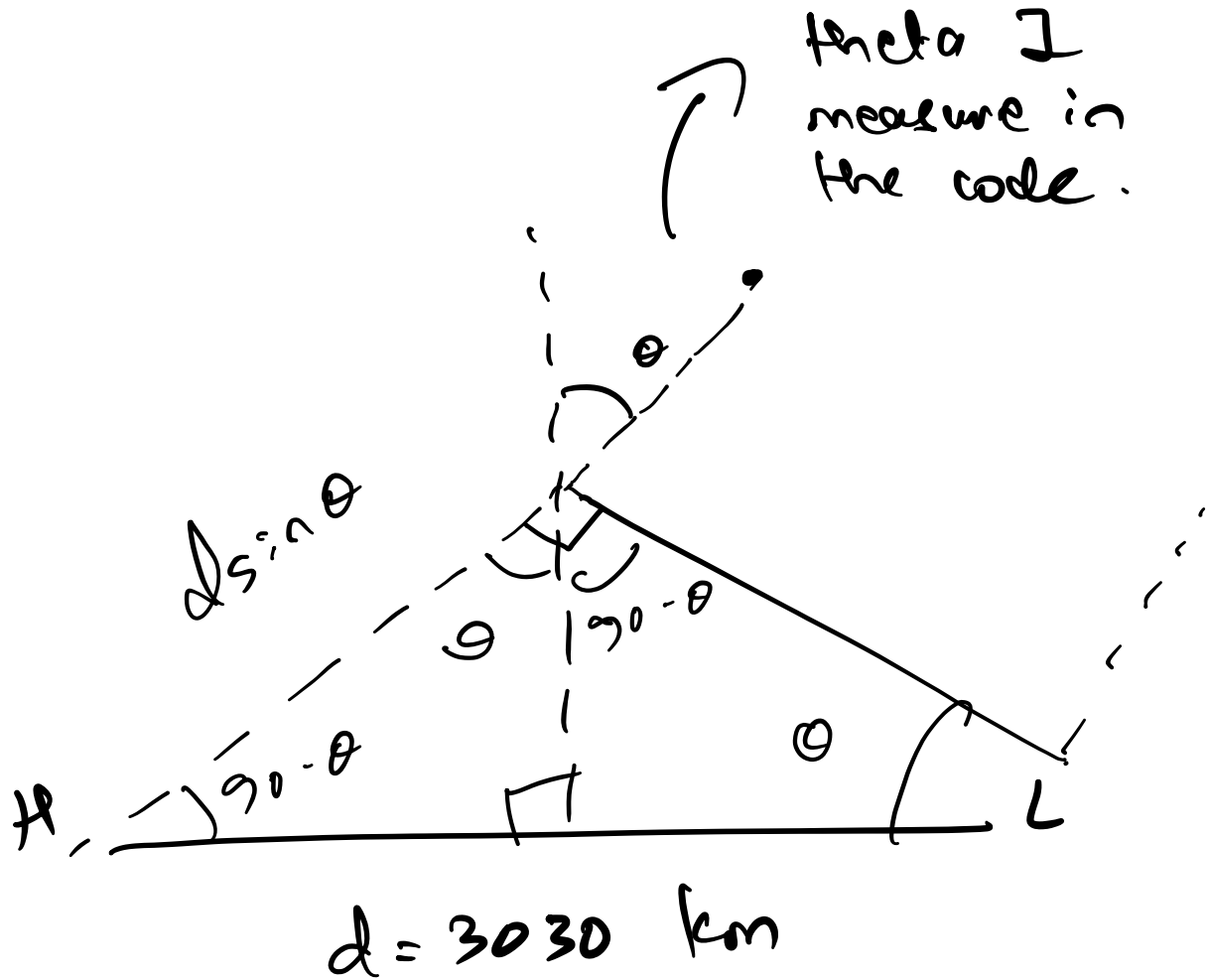


θ / f



$$\frac{d \sin \theta}{\Delta t} = c$$

$$\theta = \sin^{-1} \left(\frac{c \Delta t}{d} \right)$$

Δt is the difference in the

time of arrival of the gravitational wave at Hanford & Livingston detectors.

$$\theta = \sin^{-1} \left(\frac{c \Delta t}{d} \right)$$

$$\sin \theta = \frac{c \Delta t}{d}$$

I get the uncertainty for the TOA of the wave at each detector. Then the uncertainty in Δt is obtained as \Rightarrow

$$\delta \Delta t = \sqrt{\delta \text{TOA}_H^2 + \delta \text{TOA}_L^2}$$

(subtraction rule for uncertainty)

Then I get \Rightarrow

$$\sin(\theta + \delta\theta) = \frac{c(\Delta t + \delta\Delta t)}{d}$$

$$\theta + \delta\theta = \sin^{-1}\left(\frac{c(\Delta t + \delta\Delta t)}{d}\right)$$

$$\delta\theta = \sin^{-1}\left(\frac{c(\Delta t + \delta\Delta t)}{d}\right) - \theta$$

Q2 Since the beam pattern is gaussian,

$$A_{\text{eff}} = \int_0^R e^{-r^2/2\sigma^2} 2\pi r dr$$

Taking $r^2/2\sigma^2 = x$,

$$\frac{2r dr}{2\sigma^2} = dx$$

$$r dr = \sigma^2 dx$$

$$A_{\text{eff}} = \int_0^{R^2/2\sigma^2} e^{-x} 2\pi\sigma^2 dx$$

$$= -2\pi\sigma^2 \left[e^{-x} \right]_0^{R^2/2\sigma^2}$$

$$= 2\pi\sigma^2 \left[1 - e^{-R^2/2\sigma^2} \right]$$

$$\frac{A_{\text{eff}}}{A} = \frac{2\pi\sigma^2 \left[1 - e^{-R^2/2\sigma^2} \right]}{\pi R^2}$$

$$= \frac{2\sigma^2}{R^2} \left[1 - e^{-R^2/2\sigma^2} \right]$$

Fraction of the beam which stays on the dish (since the beam gaussian extends to ∞) =

$$dA_{\text{dish}} = \int_0^R e^{-r^2/2\sigma^2} 2\pi r dr$$

$$\frac{\int_0^R e^{-r^2/2\sigma^2} 2\pi r dr}{\int_0^\infty e^{-r^2/2\sigma^2} 2\pi r dr}$$

$$= \frac{1 - e^{-R^2/2\sigma^2}}{1}$$

$$= \frac{2\pi\sigma^2 (1 - e^{-R^2/2\sigma^2})}{2\pi\sigma^2}$$

$$= (1 - e^{-R^2/2\sigma^2})$$

2) Signal strength at the feed:

$$\frac{A_{eff}}{A} \times dA_{dish}$$

$$= \frac{2\sigma^2}{R^2} (1 - e^{-R^2/2\sigma^2})^2$$

(For $R=1$)

$$= 2\sigma^2 (1 - e^{-1/2\sigma^2})^2$$