

Figure 3.53: The same as in Fig. 3.52 but for the double differential absorption rates $d^2\Gamma_{pnn}/dx dy$.

The right panel of the Fig. 3.60 shows a cut-off dependence of the predictions obtained with the SMS chiral potential at $N^4\text{LO}^+$. In this case maximum point is interesting as well. We see that predictions with $\Lambda = 500$ MeV and 550 MeV are quite close to each other: the relative difference between them at $E_p = 0.92$ MeV is only 1.5 %. In turn the spread between $\Lambda = 400$ MeV, 450 MeV and 500 MeV is 40 % (at the same point). This cut-off dependence is hidden looking at the colormaps Fig. 3.52, but from this perspective it is clearly presented. *Anyways, the highest Λ , the highest biggest $d\Gamma_{pnn}/dE_p$ is seen.*

Similarly in Fig. 3.61 $d\Gamma_{pnn}/dE_n$ is presented. We observe similar trends which are also shown up at the extremum point which is around $E_n = 66.9$ MeV now. The difference between $N^2\text{LO}$ and $N^4\text{LO}^+$ predictions at this point is 29.0 %. The relative difference between all the predictions except for $N^2\text{LO}$ is 7.5 %. The cut-off predictions are also very similar for $\Lambda = 500$ MeV and 550 MeV (the spread is 1.6 %) while all the rest predictions are quite distinguished - the spread is 39.1 %. *Coming to the next figures, Fig. 3.62 and Fig. 3.63, which show $d\Gamma_{pnn}/dE_p$ dependence of the absorption rate on the Dalitz coordinates $r = \sqrt{x^2 + y^2}$ and $\phi = \arctan \frac{y}{x}$. The similar trend is preserved, namely chiral order figures show that $N^2\text{LO}$ predictions outstand from all other predictions, and noticeable cut-off dependence is observed.*

In general we can conclude that predictions are converged starting from the $N^3\text{LO}$ chiral order as most of the demonstrated results show that the difference between $N^3\text{LO}$, $N^4\text{LO}$ and $N^4\text{LO}^+$ is negligible. At the same time we observe a cut-off dependence where predictions obtained with $\Lambda = 500$ MeV and 550 MeV are very similar, but the spread with all the rest values is there. This nature of the cut-off dependence is also reflected in the total absorption rate, presented in Fig. 3.49.

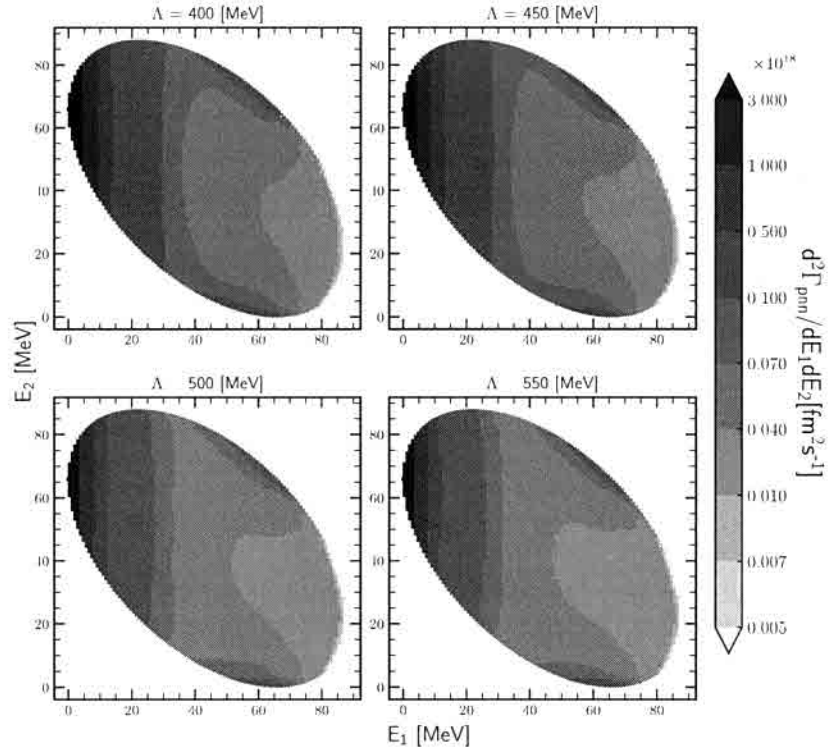


Figure 3.54: Intensity plots for the double differential absorption rates $d^2\Gamma_{pnn}/dE_1dE_2$ for the $\pi^- + {}^3\text{He} \rightarrow p + n + n$ process, obtained using the SMS potential at N^4LO^+ with plane wave part only (without rescattering). All other contributions are the same as in Fig. 3.52: $1\text{NC} + 2\text{N}$ and $2\text{NF} + 3\text{NF}$. Each panel present predictions obtained with different values of the cut-off parameter Λ : from 400 MeV (upper left) to 550 MeV (lower right). Nucleon 1 is a proton.

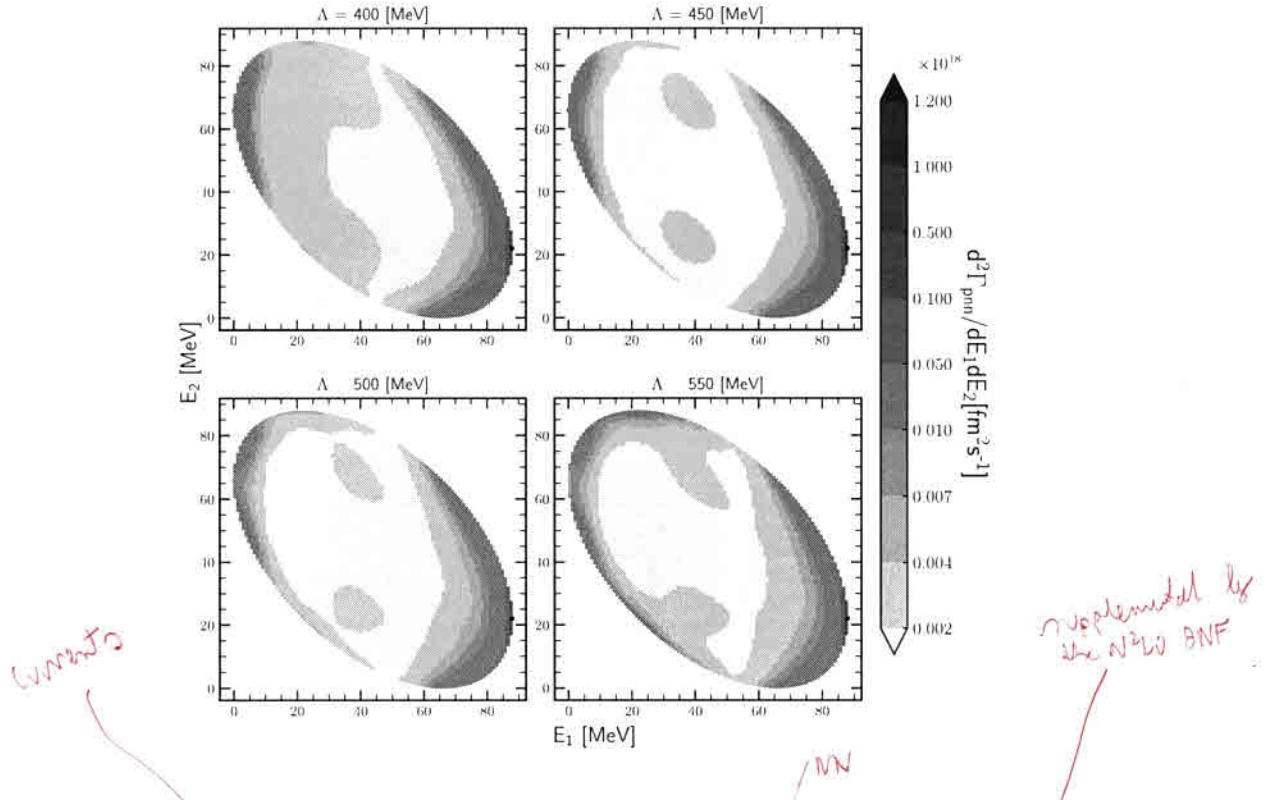


Figure 3.56: Intensity plots for the double differential absorption rates $d^2\Gamma_{pnn}/dE_1 dE_2$ for the $\pi^- + {}^3\text{He} \rightarrow p + n + n$ process, obtained using the SMS potential at $N^4\text{LO}^+$ with 1NC only (without 2N). All other contributions are the same as in Fig. 3.52: PWIAS+RESC and 2NF+3NF. Each panel presents predictions obtained with different values of the cut-off parameter Λ : from 400 MeV (upper left) to 550 MeV (lower right). Nucleon 1 is a proton.

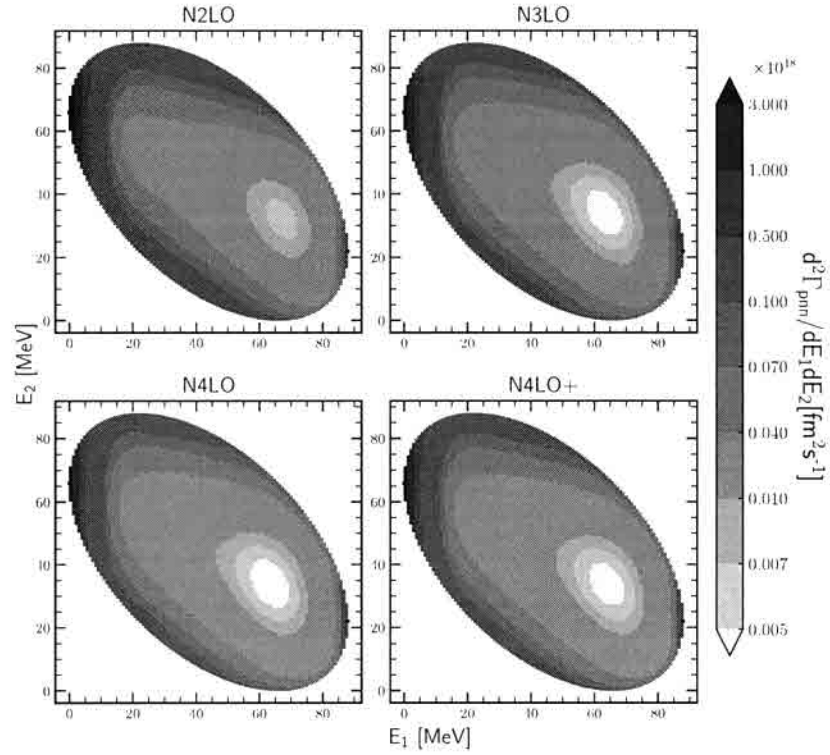


Figure 3.58: Intensity plots for the double differential absorption rates $d^2\Gamma_{pnn}/dE_1dE_2$ for the $\pi^- + {}^3\text{He} \rightarrow p + n + n$ process, obtained using the SMS potential at ~~N⁴LO⁺~~ with all contributions possible: plane wave + rescattering, 1NC + 2N, 2NF + 3NF. Each panel presents predictions obtained with different chiral orders of the SMS potential: from N²LO (upper left) to N⁴LO⁺ (lower right) and with $\Lambda = 450$ MeV. Nucleon 1 is a proton.

SP/F was taken at N²LO in each case.

NN

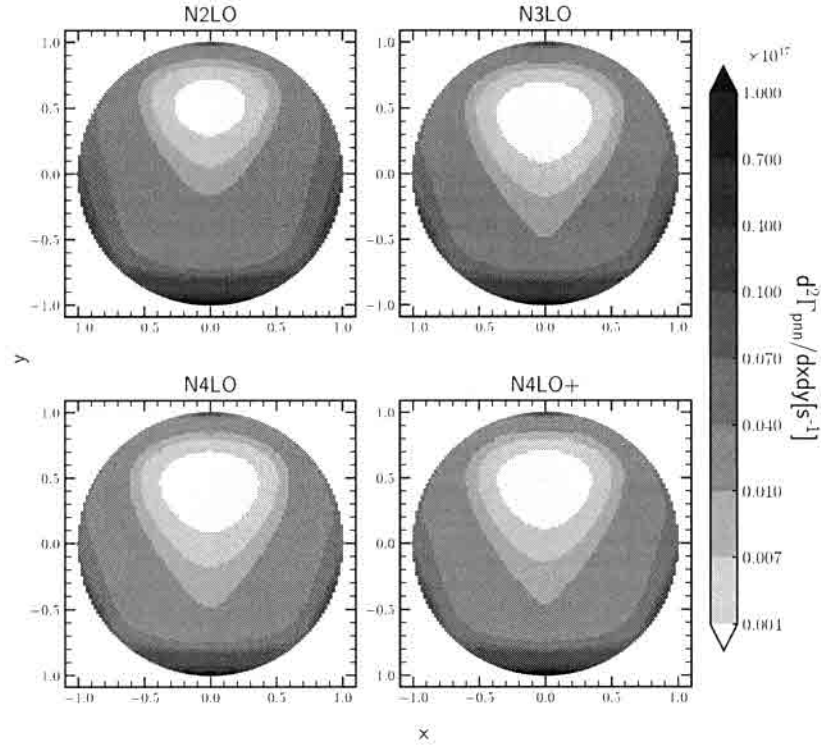


Figure 3.59: The same as in Fig. 3.58 but for the double differential absorption rates $d^2\Gamma_{pnn}/dxdy$.

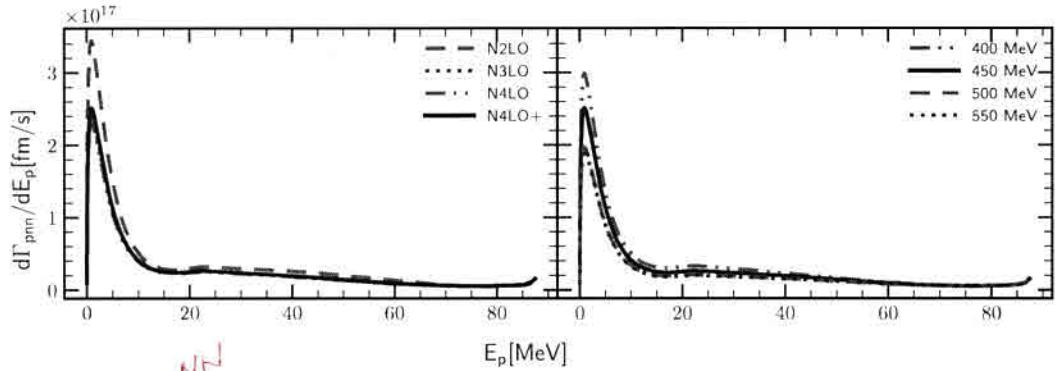


Figure 3.60: Differential absorption rate $d\Gamma_{pnn}/dE_p$ as a function of the proton energy E_p for the $\pi^- + {}^3\text{He} \rightarrow p + n + n$ process. Left panel shows results obtained with N²LO (green dashed line), N³LO (blue dotted line), N⁴LO (red dashed-double-dotted line) and N⁴LO⁺ (black solid line) chiral orders, and with $\Lambda = 450$ MeV. The right panel includes results obtained with the N⁴LO⁺ SMS potential with different values of the Λ : 400 MeV (red dashed-double-dotted line), Λ : 450 MeV (black solid line), Λ : 500 MeV (green dashed line) and Λ : 550 MeV (blue dotted line). All predictions were obtained with "FULL-(1NC+2N)-(2NF+3NF)" setup.

3.4.3 Pion absorption in ^3H

In this subsection I will show results of calculations for the $\pi^- + ^3\text{H} \rightarrow n + n + n$ process. In this case we have only a three-body breakup as no two-body configuration can be composed out of three neutrons.

The total absorption rate $\pi^- + ^3\text{H} \rightarrow n + n + n$ is shown in Fig. 3.64 as a function of the chiral order of the SMS potential while each curve represents different cut-off values used to obtain the prediction. The most advanced configuration was used in this case, namely Plane wave plus rescattering part, both single- and two-nucleon currents and two-nucleon plus three-nucleon forces interaction.

Similarly to the process with ^3He , we see that with each subsequent chiral order predictions become closer to each other, so cut-off dependence gets weaker. Also, the prediction with $\Lambda = 550 \text{ MeV}$ at N^3LO is strangely above the prediction with $\Lambda = 500 \text{ MeV}$. We can also notice, that at N^4LO predictions with cut-off values 500 MeV and 550 MeV are much closer to each other than to other values.

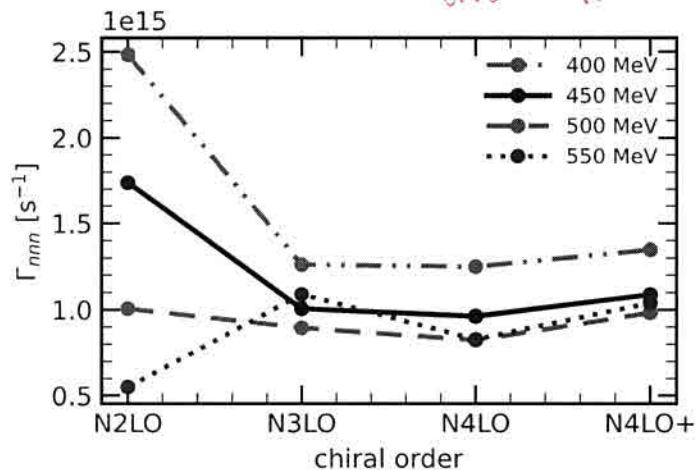


Figure 3.64: Absorption rate for $\pi^- + ^3\text{H} \rightarrow n + n + n$ reaction as a function of the chiral order with different values of the cut-off parameter Λ . Predictions were obtained with 3NF at N^2LO .

1. add remaining discussion of Figs 3.65-3.71

2. Both ~~in~~ in discussion for ^3He and ^3H I miss

details which configuration is at specific (x, y)

For example, in ^{our} paper page 16: $QF_1(n)$ is located at $(x, y) = (0, 1)$. Please add such information to thesis

→ ~~go~~ take them from ₈₁ paper. That ~~make~~ will make paper more quantitative. next page

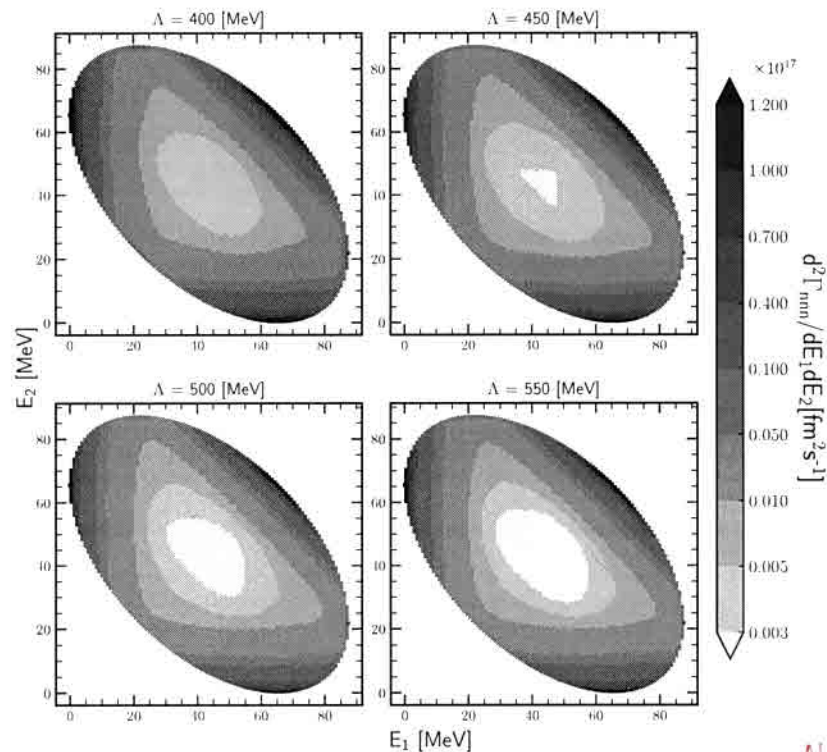


Figure 3.65: Intensity plots for the double differential absorption rates $d^2\Gamma_{nnn}/dE_1dE_2$ for the $\pi^- + {}^3\text{H} \rightarrow n + n + n$ process, obtained using the SMS potential at N^4LO^+ with all contributions possible: plane wave + rescattering, 1NC + 2N, 2NF+3NF. Each panel presents predictions obtained with different values of the cut-off parameter Λ : from 400 MeV (upper left) to 550 MeV (lower right).

3. At some point comparison with data (see Eg 4.00-4.16 from paper) is needed. You can also mention comparison with Guttor et al (Tab V from the paper) + discussion. Also for ${}^2\text{H}$ or info that there is no data.
4. Give number and description for total rates (like Eg 4.00-4.02, 4.03, tab. 1 from paper)

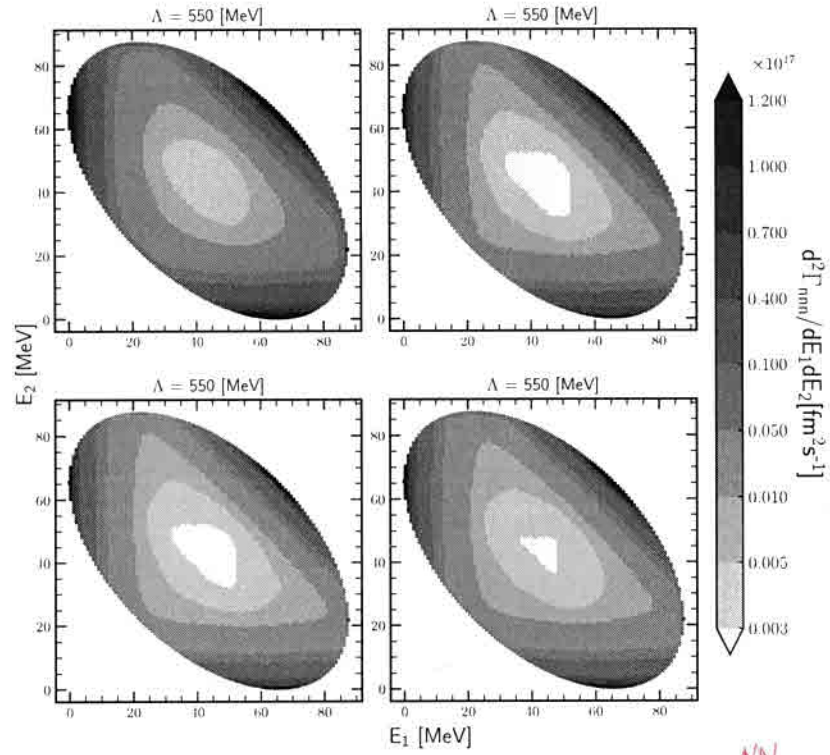


Figure 3.67: Intensity plots for the double differential absorption rates $d^2\Gamma_{nnn}/dE_1 dE_2$ for the $\pi^- + {}^3\text{H} \rightarrow n + n + n$ process, obtained using the SMS potential at N^4LO^+ with all ~~contributions~~ possible: plane wave – rescattering, 1NC – 2N, 2NF+3NF. Each panel present predictions obtained with different chiral orders of the SMS potential: from N^2LO (upper left) to N^4LO^+ (lower right) and with $\Lambda = 450$ MeV.

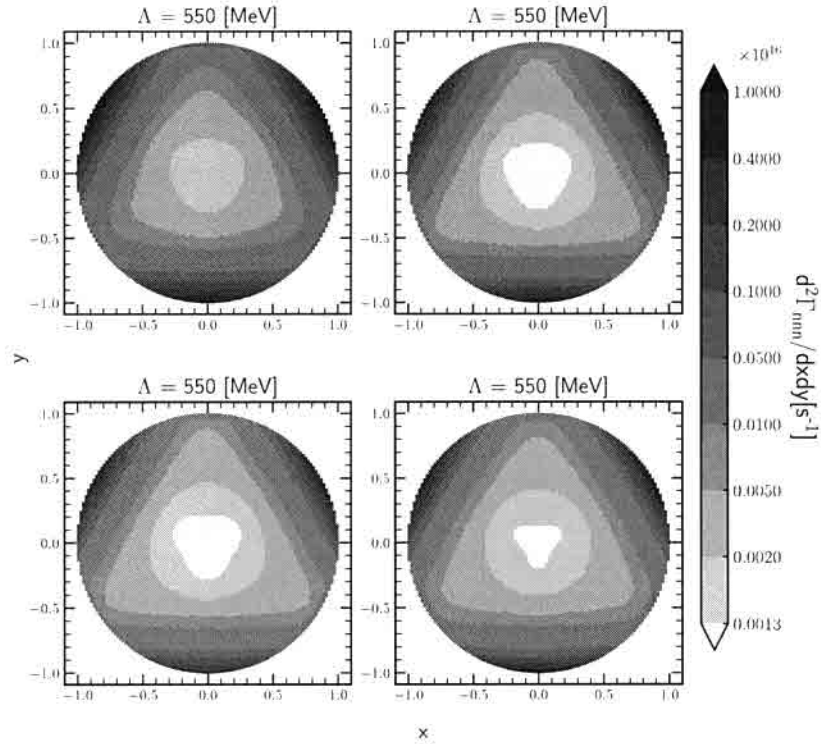


Figure 3.68: The same as in Fig. 3.67 but for the double differential absorption rates $d^2\Gamma_{nnn}/dx dy$.

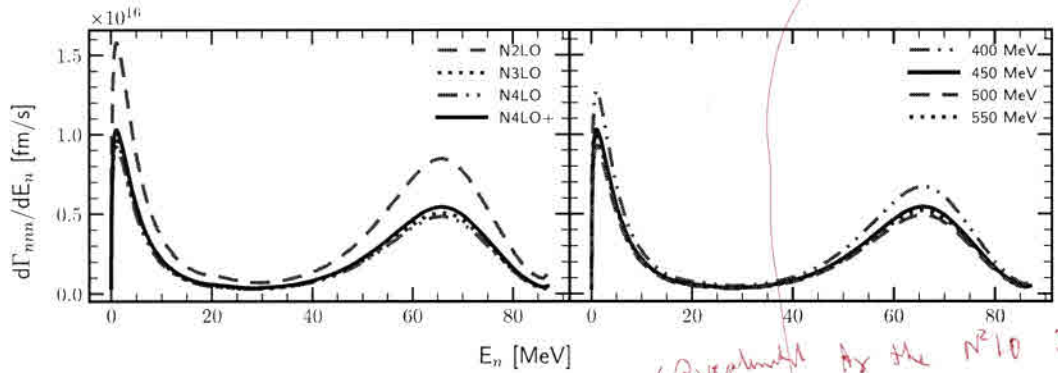


Figure 3.69: Differential absorption rate $d\Gamma_{nnn}/dE_n$ as a function of the neutron energy E_n for the $\pi^- + 3H \rightarrow n + n + n$ process. Left panel shows results obtained with N^2LO (green dashed line), N^3LO (blue dotted line), N^4LO (red dashed-double-dotted line) and N^4LO^+ (black solid line) chiral orders, and with $\Lambda = 450$ MeV. The right panel includes results obtained with the N^4LO^+ SMS potential with different values of the Λ : 400 MeV (red dashed-double-dotted line), $\Lambda = 450$ MeV (black solid line), $\Lambda = 500$ MeV (green dashed line) and $\Lambda = 550$ MeV (blue dotted line). All predictions were obtained with "FULL-(1NC+2N)-(2NF+3NF)" setup.

and only for pion absorption it was supplemented by the π -nucleon current operator
 For the ~~electromagnetic~~ photodisintegration π -like currents via the Siegert theorem

CHAPTER 4

SUMMARY

augmented by the consistently ~~renormalized~~ ~~renormalized~~ ~~renormalized~~ three-nucleon force at N^2 LO.

In this thesis I investigated ${}^2\text{H}$, ${}^3\text{H}$ and ${}^3\text{He}$ photodisintegrations as well as pion absorption by the same nuclei. In order to analyze these reactions and calculate predictions for observables I used a chiral model of interaction namely the most advanced SMS chiral potential up to $N^4\text{LO}^+$ order. Results prepared with AV18 potential have been shown as a reference point. ~~The current operator used is restricted to one single nucleon current~~

The first goal of this work was to investigate if the SMS potential is converged with respect to the chiral order. It would then be a hint whether the development of higher orders is required. In most of the results we observed a very converged predictions which was confirmed by a small difference between the last two chiral orders: $N^4\text{LO}$ and $N^4\text{LO}^+$. This difference in most of regarded situations did not exceed a few percent. Another evidence is a width of the truncation band for $N^4\text{LO}^+$ predictions. For the deuteron photodisintegration process, observables have a maximum width always below 1% for the regarded photon energies ($E_\gamma = 30$ MeV and $E_\gamma = 100$ MeV). The same trend is presented also for other regarded reactions as well. This hints us towards a conclusion that further chiral orders would not improve the predictions much and current potential is converged already. This is also confirmed by the AV18 predictions which are always very close to $N^4\text{LO}^+$ (see e.g. Fig. 3.8, Fig. 3.26 etc.).

The other interesting point in the investigated chiral potential is its dependence on the value of the cutoff parameter Λ , the four values of which (400, 450, 500 and 550 MeV) were investigated. We have observed that the relative spread of predictions with respect to the cutoff value is higher for the larger energies. For example, the spread of the differential cross section for ${}^3\text{He}$ photodisintegration at $E_\gamma = 30$ MeV at the characteristic point (maximum) is around 3%, while at $E_\gamma = 100$ MeV it is three times larger - around 9% (see Fig. 3.28, Fig. 3.29 and discussions). Nevertheless, usually for higher energies the difference between the predictions obtained with different values of Λ is smaller than the difference with experimental data (when it is available) which is evidently visible in Fig. 3.5(b).

We have also checked contributions from the various model components to the predicted values. Namely, predictions obtained with plane wave part only (first term in the Eq. (2.31)), as well as predictions with and without 2N current contributions (obtained via the Siegert approach). For example, in the Fig. 3.4b we see predictions for the deuteron photodisintegration obtained without rescattering part (pink line), without 2N contributions (green line) and full predictions (blue line). Evidently the rescattering

① Both processes were studied in momentum space. ~~and~~

The standard Lippmann-Schwinger equation is solved to get the t -operator and consequently $2N$ scattering state. For three-nucleon processes the formulation of Faddeev equations have been applied. I solved corresponding equations for Faddeev components both for the bound ~~and~~ and scattering $3N$ states. In that way I include all initial state as well as final state interactions. I'm also able to test ~~the~~ importance of FSI by restricting computations to only plane wave approximation.

② ~~if the observable~~ ... to investigate ^{currently available predictions ~~are~~ based on} a quality of ~~value~~ the SMS interaction if applied to the studied here processes. Such information is necessary ~~in~~ due to expected two- and more-nucleon currents at ~~the~~ higher orders of chiral expansion, consistent with the SMS ~~is~~ potential. ~~Various aspects can be~~ Various features of the model can be studied in that context. Firstly, I ~~would be~~ investigated if the predictions based on the SMS interaction are converged ...

The used formulation allows me to study not only the total cross section or capture rates but also various differential cross sections and polarization observables. That in turn allows me to conclude on sensitivity of various observables on studied effects and to pick up observables which after measurement could deliver valuable information.

CHAPTER 4. SUMMARY

processes

part contribution is relatively small for the $E_\gamma = 30$ MeV, but is increasing with larger energies. The analysis of the relative difference at the specific θ_p value does not show this trend (the differences are 10 %, 7 % and 4 % for $E_\gamma = 30, 100$ and 140 MeV respectively at $\theta_p = 80^\circ$), but we see that at the lower energy curves are qualitatively very similar and the difference is mainly around the maximum point (which is close to the analyzed point). In contrast, for the larger energies, curves differ qualitatively and analyzed point is rather depicting the region where the difference is minimal. The difference between the full predictions and the ones, where only 1NC was used is much higher: even for the lowest energy it is around 50 % and for the larger two it grows up to ~ 78 % (at the same angle value). Clearly, both rescattering part and 2NC contributions are very important and bring significant contribution. The latter one is taken into account not completely, since there is no proper 2NC current available yet, but even using the Siegert approach proves that this part is very important in described process. Similar trend is visible at other observables (see e.g. Fig. 3.10) and other processes (see Figs. 3.52, 3.54, 3.56 to compare contributions for the pion absorption at ^3He).

Giving mentioned above results, we can conclude that current version of the chiral potential is good: it rather does not require any additional development in a sense of chiral orders or regularization, but 2NC is expected to bring a large improvements to the predictions.

Our Full model nicely describes data up to $E_\gamma = 70$ MeV.

higher

of very high quality

Contrary, the should be completely derived as it

significant contribution and thus improvements in our understanding of el-mag process

That complex pattern shows that reveals interplay between interaction and the current operator, and is one more recommendation after deriving and derivation and application consistent fully consistent model, i.e. with consistent $2N$ forces, $3N$ forces and one-, two- and three-body currents. Such model should be must be applied only within reliable scheme to compute observables. My work is culminated in the preparation of such scheme, from analytical and numerical aspects sides, and now we are ready for study more sophisticated Hamiltonians.

(in points)

In summary you ~~are~~ should give few ~~few~~
detailed conclusions like

Among the studied processes and observables:

- 1) The most sensitivity to FSI are for ... (observables, energies)
- 2) _____ " _____ $2NC$
- 3) _____ " _____ $3NF$
- 4) ~~the~~ total cross section vs differential cross section
- 5) cross sections vs polarization observables
- 6) Is it better to do measurement on 2H , 3He or 3H ?
- 7) stress that for $\pi + ^3H$ there are three neutrons
in final state what is very unusual but extremely
interesting situation
- 8) for pion absorption write if you suggest measurements
in FSI configuration, QFS or others?

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