

Introduction to Quantum Mechanics

Problem 1.1

For the distribution of ages in the example in Section 1.3.1:

- (a) Compute $\langle j^2 \rangle$ and $\langle j \rangle^2$
- (b) Determine Δj for each j , and use Equation 1.11 to compute the standard deviation
- (c) Use your results in (a) and (b) to check Equation 1.12

Problem 1.2

- (a) Find the standard deviation of the distribution in Example 1.2
- (b) What is the probability that a photograph, selected at random, would show a distance x more than one standard deviation away from the average?

Problem 1.3

Consider the gaussian distribution

$$\rho(x) = Ae^{-\lambda(x-a)^2}$$

where A , a , and λ are positive real constants. (The necessary integrals are inside the back cover.)

- (a) Use Equation 1.16 to determine A
- (b) Find $\langle x \rangle$, $\langle x^2 \rangle$, and σ
- (c) Sketch the graph of $\rho(x)$

Problem 1.4

At time $t = 0$, a particle is presented by the wave function

$$\Psi(x, 0) = \begin{cases} A \frac{x}{a} & 0 \leq x \leq a \\ A \frac{b-x}{b-a} & a \leq x \leq b \\ 0 & \text{otherwise} \end{cases}$$

where A , a , and b are positive constants.

- (a) Normalize Ψ (that is, find A , in terms of a and b)
- (b) Sketch $\Psi(x, 0)$, as a function of x
- (c) Where is the particle most likely to be found at $t = 0$?
- (d) What is the probability of finding the particle to the left of a ? Check your result in the limiting cases $b = a$ and $b = 2a$
- (e) What is the expectation value of x

Solution

- (a) We begin with the normalization condition:

$$\begin{aligned}
1 &= \int_{-\infty}^{\infty} |\Psi(x, 0)|^2 dx \\
&= \int_0^a \left| A \frac{x}{a} \right|^2 dx + \int_a^b \left| A \frac{b-x}{b-a} \right|^2 dx \\
&= \frac{A^2}{a^2} \int_0^a x^2 dx + \frac{A^2}{(b-a)^2} \int_a^b (b-x)^2 dx \\
&= A^2 \left(\frac{a}{3} + \frac{b-a}{3} \right) \\
&= A^2 \frac{b}{3}
\end{aligned}$$

Therefore, $A = \sqrt{\frac{3}{b}}$

- (a) We skip the sketch as requested
- (b) The particle is most likely to be found at $x = a$
- (c) The probability of finding the particle to the left of a is:

$$\begin{aligned}
P &= \int_{-\infty}^a |\Psi(x, 0)|^2 dx \\
&= \int_0^a \left| A \frac{x}{a} \right|^2 dx \\
&= \frac{A^2}{a^2} \int_0^a x^2 dx \\
&= A^2 \frac{a}{3} \\
&= \frac{3}{b} \frac{a}{3} \\
&= \frac{a}{b}
\end{aligned}$$

When $b = a$, $P = 1$ When $b = 2a$, $P = \frac{1}{2}$

- (a) The expectation value of x is:

$$\begin{aligned}
\langle x \rangle &= \int_{-\infty}^{\infty} x |\Psi(x, 0)|^2 dx \\
&= \int_0^a x \left| A \frac{x}{a} \right|^2 dx + \int_a^b x \left| A \frac{b-x}{b-a} \right|^2 dx \\
&= \frac{A^2}{a^2} \int_0^a x^3 dx + \frac{A^2}{(b-a)^2} \int_a^b x(b-x)^2 dx \\
&= A^2 \frac{b(b^3 - 3a^2b + 2a^3)}{12(b-a)^2} \\
&= \frac{3}{b} \frac{b(2a+b)}{12} \\
&= \frac{2a+b}{4}
\end{aligned}$$

Problem 1.5

Consider the wave function

$$\Psi(x, t) = Ae^{-\lambda|x|}e^{-i\omega t},$$

where A , λ , and ω are positive real constants. (We'll see in Chapter 2 for what potential (V) this wave function satisfies the Schrödinger equation.)

- Normalize Ψ .
- Determine the expectation values of x and x^2 .
- Find the standard deviation of x . Sketch the graph of $|\Psi|^2$, as a function of x , and mark the points $\langle x \rangle + \sigma$ and $\langle x \rangle - \sigma$, to illustrate the sense in which σ represents the “spread” in x . What is the probability that the particle would be found outside this range?

Solution

- We begin with the normalization condition:

$$\begin{aligned}
1 &= \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx \\
&= \int_{-\infty}^{\infty} |Ae^{-\lambda|x|}e^{-i\omega t}|^2 dx \\
&= A^2 \int_{-\infty}^{\infty} e^{-2\lambda|x|} dx \\
&= 2A^2 \int_0^{\infty} e^{-2\lambda x} dx \\
&= 2A^2 \frac{1}{2\lambda} \\
&= \frac{A^2}{\lambda}.
\end{aligned}$$

Thus, we have $A = \sqrt{\lambda}$.

(b) The expectation value of x is

$$\begin{aligned}
 \langle x \rangle &= \int_{-\infty}^{\infty} x |\Psi(x, t)|^2 dx \\
 &= \int_{-\infty}^{\infty} x |A e^{-\lambda|x|} e^{-i\omega t}|^2 dx \\
 &= A^2 \int_{-\infty}^{\infty} x e^{-2\lambda|x|} dx \\
 &= 0.
 \end{aligned}$$

The expectation value of x^2 is

$$\begin{aligned}
 \langle x^2 \rangle &= \int_{-\infty}^{\infty} x^2 |\Psi(x, t)|^2 dx \\
 &= \int_{-\infty}^{\infty} x^2 |A e^{-\lambda|x|} e^{-i\omega t}|^2 dx \\
 &= A^2 \int_{-\infty}^{\infty} x^2 e^{-2\lambda|x|} dx \\
 &= 2A^2 \int_0^{\infty} x^2 e^{-2\lambda x} dx \\
 &= 2A^2 \frac{1}{4\lambda^3} \\
 &= \frac{1}{2\lambda^2}.
 \end{aligned}$$

(c) The standard deviation of x is

$$\begin{aligned}
 \sigma &= \sqrt{\langle x^2 \rangle - \langle x \rangle^2} \\
 &= \sqrt{\frac{1}{2\lambda^2}} \\
 &= \frac{1}{\sqrt{2}\lambda}.
 \end{aligned}$$

We skip the sketch as requested.

The probability that the particle would be found outside the range $[-\sigma, \sigma]$ is

$$\begin{aligned}
P &= \int_{-\infty}^{-\sigma} |\Psi(x, t)|^2 dx + \int_{\sigma}^{\infty} |\Psi(x, t)|^2 dx \\
&= 2 \int_{\sigma}^{\infty} |\Psi(x, t)|^2 dx \\
&= 2A^2 \int_{\sigma}^{\infty} e^{-2\lambda x} dx \\
&= 2A^2 \frac{1}{2\lambda} e^{-2\lambda\sigma} \\
&= e^{-2\lambda\sigma} \\
&= e^{-\sqrt{2}} \\
&\approx 0.2431.
\end{aligned}$$

Problem 1.6

Why can't you do integration-by-parts directly on the middle expression in Equation 1.29—pull the time derivative over onto x , note that $\frac{\partial x}{\partial t} = 0$, and conclude that $\frac{d\langle x \rangle}{dt} = 0$?

Solution

$$\begin{aligned}
\frac{\partial}{\partial t} x |\Psi|^2 &= x \frac{\partial}{\partial t} |\Psi|^2 + |\Psi|^2 \frac{\partial x}{\partial t} \\
&= x \frac{\partial}{\partial t} |\Psi|^2
\end{aligned}$$

Problem 1.7

Calculate $\frac{d\langle p \rangle}{dt}$. Answer:

$$\frac{d\langle p \rangle}{dt} = \left\langle -\frac{\partial V}{\partial x} \right\rangle.$$

This is an instance of *Ehrenfest's theorem*, which asserts that expectation values obey the classical laws.

Solution

$$\begin{aligned}
\frac{d\langle p \rangle}{dt} &= -i\hbar \frac{\partial}{\partial t} \int_{-\infty}^{\infty} \Psi^* \frac{\partial \Psi}{\partial x} dx \\
&= -i\hbar \int_{-\infty}^{\infty} \left(\frac{\partial \Psi^*}{\partial t} \frac{\partial \Psi}{\partial x} + \Psi^* \frac{\partial}{\partial x} \frac{\partial \Psi}{\partial t} \right) dx \\
&= -i\hbar \int_{-\infty}^{\infty} \left[\left(\frac{-i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \right) \frac{\partial \Psi}{\partial x} + \Psi^* \frac{\partial}{\partial x} \left(\frac{i\hbar}{2m} \frac{\partial^2 \Psi}{\partial x^2} - \frac{i}{\hbar} V \Psi \right) \right] dx \\
&= \frac{\hbar^2}{2m} \int_{-\infty}^{\infty} \left(\Psi^* \frac{\partial^3 \Psi}{\partial x^3} - \frac{\partial^2 \Psi^*}{\partial x^2} \frac{\partial \Psi}{\partial x} \right) dx + \int_{-\infty}^{\infty} \left[V \Psi^* \frac{\partial \Psi}{\partial x} - \Psi^* \frac{\partial}{\partial x} (V \Psi) \right] dx \\
&= \frac{\hbar^2}{2m} \left[\left(\Psi^* \frac{\partial^2 \Psi}{\partial x^2} \right) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{\partial \Psi^*}{\partial x} \frac{\partial^2 \Psi}{\partial x^2} dx - \left(\frac{\partial \Psi^*}{\partial x} \frac{\partial \Psi}{\partial x} \right) \Big|_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \frac{\partial \Psi^*}{\partial x} \frac{\partial^2 \Psi}{\partial x^2} dx \right] \\
&\quad + \int_{-\infty}^{\infty} \left[V \Psi^* \frac{\partial \Psi}{\partial x} - \Psi^* V \frac{\partial \Psi}{\partial x} - \Psi^* \frac{\partial V}{\partial x} \Psi \right] dx \\
&= \int_{-\infty}^{\infty} -\Psi^* \left[\frac{\partial V}{\partial x} \right] \Psi dx \\
&= \left\langle -\frac{\partial V}{\partial x} \right\rangle.
\end{aligned}$$

Problem 1.8

Suppose you add a constant V_0 to the potential energy (by “constant” I mean independent of x as well as t). In classical mechanics this doesn’t change anything, but what about *quantum mechanics*? Show that the wave function picks up a time-dependent phase factor: $\exp(-i \frac{V_0 t}{\hbar})$. What effect does this have on the expectation value of a dynamical variable?

Solution

Suppose the wave function Ψ satisfies the Schrödinger equation without the constant V_0 :

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V \Psi.$$

Then, for the wave function $\Psi' = \Psi e^{-i \frac{V_0 t}{\hbar}}$, we have

$$\begin{aligned}
i\hbar \frac{\partial \Psi'}{\partial t} &= i\hbar \frac{\partial}{\partial t} \left(\Psi e^{-i \frac{V_0 t}{\hbar}} \right) \\
&= i\hbar \left(\frac{\partial \Psi}{\partial t} e^{-i \frac{V_0 t}{\hbar}} - \frac{i V_0}{\hbar} \Psi e^{-i \frac{V_0 t}{\hbar}} \right) \\
&= \left(-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V \Psi \right) e^{-i \frac{V_0 t}{\hbar}} + V_0 \Psi e^{-i \frac{V_0 t}{\hbar}} \\
&= -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi'}{\partial x^2} + (V + V_0) \Psi',
\end{aligned}$$

which is the Schrödinger equation with the constant V_0 . Thus, the wave function Ψ picks up a time-dependent phase factor $\exp\left(-i \frac{V_0 t}{\hbar}\right)$. The expectation value of a dynamical variable is not affected by the phase factor, since the x -independent phase factor is canceled out when taking the expectation value.