# V64

# Interferometry

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Durchführung: DATUM Abgabe: DATUM

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## 1 Objective

In this experiment a Sagnac-interferometer is used in combination with a HeNe-laser is used to determine the refractive indices of glass and air. To achieve this the Interferometer is first adjusted and the visibility measured to guarantee a high quality in the interference pattern.

# 2 Theory

intererometry uses the effect of interference between lightwaves to make measurements for example regarding properties of certain materials like in this case the refractive indices of glass and air. To achieve interference several factors have to be taken into consideration, which will be discussed in the following chapter.

#### 2.1 Coherence

When trying to observe interference between lightwaves it is important to consider, that interference occurs only with coherent light. This means, the different lightwaves have to have a constant phase relationship with each other to make interference possible. This property of light is called coherence and is distinguished between spatial and temporal coherence. Real lightsources emitt light in a certain spectrum and therefore constantly shift their phase, which means that a constant phase relation can only be achieved within a certain timeframe. Perfect monochromatic light theoretically is perfectly temporally coherent but can not be achieved in reality. Spatial coherence is achieved when wavefronts do not shift their phase difference due to spatial differences between the propagating wavefronts.

Coherence is a phenomenon that exists on a certain scale, meaning that light can not only be incoherent or perfectly coherent. This is quantified by the Degree of coherence:

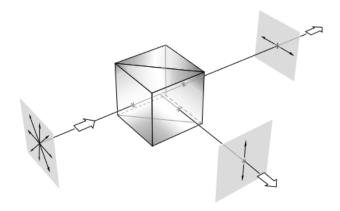
$$\tilde{\gamma}_{12}(\tau) = \frac{\langle \tilde{E}_1(t+\tau)\tilde{E}_2^*(t)\rangle_T}{\sqrt{\langle |\tilde{E}_1|^2\rangle\langle |\tilde{E}_2|^2\rangle}} \tag{1}$$

with the electric field at two points  $E_1$ ,  $E_2$  and the relative offset  $\tau$ . Here  $\langle ... \rangle_T$  denotes the time average over an Intervall T. The degree of coherence takes up values between 1 and 0 with  $\gamma_{12}=1$  indicating perfect coherence.

#### 2.2 Polarisation

Another important factor in interference of transversal waves like lightwaves is the polarisation of the wave. This describes the orientation of the plane in which the electric filed oscillates. The most important distinction in terms of polarisation is between linear and circluar/elliptical polarisation where the electric field vector is not polarised in a single plane but rotates. Linearly polarised light can be described by the polarization angle and generated from unpolarized light by using a polarisation filter, which only allows

the component with the corresponding polarization angle to pass. Linearly polarized lightbeams can only produce interference if they have the same polarization angle. A



**Abbildung 1:** Depiction of a Polarising beam-splitter cube splitting unpolarised light into the s-polarised and p-polarised components. [**Hecht 2018**]

polarizing beam-splitter cube (PBSC) can be used to split light with a certain polarization into two perpendicular components. The transmitted beam is then polarized parallel (p-polarization) and the reflected beam perpendicular (s-polarization) to the incident plane. The effect of a PBSC on unpolarized light is depicted in Figure 1.

#### 2.3 Visibility

A metric to describe the quality of an interference pattern is the so called visibility:

$$V = \frac{I_{max} - I_{min}}{I_{max} + I_{min}} \tag{2}$$

which quantifies the contrast between the minimum Intensity  $I_{min}$  and maximum Intensity  $I_{max}$ . The intensity is proportional to the square of the electric field strength  $I \propto E^2$ . More precisely the intensity of two interfering lightbeams can be described as:

$$I \propto \langle |E_1 \cos(wt) + E_2 \cos(wt + \delta)|^2 \rangle$$

Assuming a phase difference of  $\delta_{max}=n2\pi$  for complete constructive interference at the maximum and  $\delta_{min}=(2n+1)\pi$  with  $n\epsilon N^0$  for destructive interference at the minimum and consequently using  $\cos(wt+\delta_{max})=\cos(wt)$  and  $\cos(wt+\delta_{min})=-\cos(wt)$  the corresponding intensities can be calculated:

$$\begin{split} I_{max/min} &\propto \left\langle \cos^2(wt)(E_1 \pm E_2)^2 \right\rangle \\ &\iff I_{max/min} \propto E_1^2 \pm 2E_1E_2 + E_2^2 \end{split}$$

The strengths of the electric fields  $E_1$  and  $E_2$  is determined by the polarisation angle  $\Phi$  of the initial light beam with the electric field strength  $E_0$  before beeing split by the beam-splitter cube:

$$\begin{pmatrix} E_1 \\ E_2 \end{pmatrix} = E_0 \begin{pmatrix} \cos(\varPhi) \\ \sin(\varPhi) \end{pmatrix}$$

Inserting this into the previous equation yields:

$$I_{max/min} \propto E_0^2 (1 \pm 2\sin(\Phi)\cos(\Phi)) \tag{3}$$

When inserting the results for the Intensities into the visibility it can be expressed depending only on the polarization Angle:

$$V(\Phi) = 2\sin(\Phi)\cos(\Phi) \tag{4}$$

#### 2.4 Refractive indicex of glass

Using a glass holder holding two glass plates placed in the interferometer the refractive index of the glass can be determined based on the number of observed maxima or minima:

$$M = \frac{\Delta \phi}{2\pi}$$

with the phase shift  $\Delta \phi$  depending on the rotation angle of the glass plate  $\Theta$  and the refractive index of the glass  $n_{glass}$ :

$$\Delta \phi = \frac{2\pi}{\lambda_{vac}} D\left(\frac{n_{glass} - 1}{2n_{glass}} \Theta^2 + O(\Theta^4)\right)$$
 (5)

With the vacuum-wavelength of light  $\lambda_{vac}$  and the thickness of the glass plate D. Because the plates in the glass holder are already rotated by an angle of  $\pm 10^{\circ}$  the formula for the phase shift has to be slightly changed and results in a linear relationship between the number of Maxima and the rotation angle:

$$M = \frac{D}{\lambda_{vac}} \frac{n_{glass} - 1}{2n_{glass}} \left( (\Theta + \frac{\pi}{18})^2 - (\Theta - \frac{\pi}{18})^2 \right) = \frac{D}{\lambda_{vac}} \frac{n_{glass} - 1}{n_{glass}} 2\Theta \frac{\pi}{18} \tag{6}$$

$$\iff n = \frac{1}{1 - \frac{M\lambda_{vac}}{D2\Theta\frac{\pi}{12}}} \tag{7}$$

#### 2.5 Refractive index of air

The refractive index of air can be determined analogously by placing a gas cell of Length L in the beam. This again results in a phase shift of

$$\Delta \Phi = \frac{2\pi}{\lambda_{vac}} (n_{air} - 1) L$$

corresponding to a refractive index of:

$$n_{air} = \frac{M\lambda_{vac}}{L} + 1 \tag{8}$$

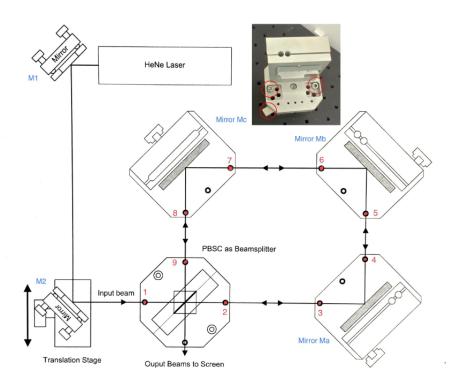
Alternatively the Refractive index of air can also be determined using the Lorentz-Law:

 $\frac{n_{air}^2 - 1}{n_{air}^2 + 1} = \frac{Ap}{RT} \tag{9}$ 

which describes the relationship between refractive index, pressure p and temparature T. In this case A denotes the molrefraction value and R the universal gas constant. [sample]

#### 3 Measurement

For this experiment a Sagnac-Interferometer is chosen because of its high resistance against external disturbances. A schematic picture of the interferometer used can be seen



**Abbildung 2:** Schematic picture of the Sagnac-Interferometer used for the purpose of this experiment. [V64]

in Figure 2. The light for the interferometer is generated by a HeNe-Laser emitting linearly polarised light tilted by 45° from the vertical with a wavelength of 632,990 nm.

#### 3.1 Adjustment

To be able to properly use the interferometer for measurements it first has to be adjustet. For this purpose adjustement plates are used to stepwise adjust the interferometer.

First the mirror M1 is adjusted so that the laser beam hits the center of Mirror M2. This mirror can then be positioned in a way, that the beam transmitted by the PBSC passes through the center of the adjustment plates in positions 1 and 2. Afterwards the reflected beam is adjusted to pass the center of adjustment plates in positions 8 and 9 by moving the PBSC itself. Then the mirrors Ma and MC are adjusted using plates in positions 5 and 6. In a last step the mirror Mb can be adjusted by using plates in the positions 7 and 4. With all adjustement steps finished a polarization filter has to be introduced to the interfering beams to synchronize their polarization angle. The resulting picture should then still show no interference fringes if the interferometer is perferctly adjusted. By moving the mirror M2 away from the adjusted central position the two overlapping beams can be separated so that each beam can be manipulated individually.

#### 3.2 Measuring the visibility

To measure the visibility a glass holder with two tilted glass plates at an adjustable angle is introduced to the beam. Each plate has a thickness of  $D=1\,\mathrm{mm}$ . The polarisation angle of the input beam is also additionally controlled by a polarisation filter in front of the PBSC. By utilizing the angle of the glass plates the interference pattern can now be changed to either show a minimum or a maximum. The intensity is measured using the output voltage of a photo-diode. The polarisation of the input beam is then increased in  $10^{\circ}$  steps from  $0^{\circ}$  to  $180^{\circ}$  and for each angle the intensity at the minimum and the maximum are recorded. This measurement is repeated three times.

#### 3.3 Refractive index of glass

After the visibility is measured in the previous step the polarisation angle of the input beam is set to guarantee the maximal visibility. For the measurements regarding the refractive indices the difference voltage method is used with two diodes. The angle of the glass plates is then increased from  $0^{\circ}$  to  $10^{\circ}$  and an automatic counter in combination with an oscilloscope is used to measure the number of Maxima. This procedure is repeated ten times.

#### 3.4 Refractive index of air

Instead of the glass holder a gas cell with a length of  $L=(100,0\pm0,1)\,\mathrm{mm}$  is then introduced to the beam. The cell can be avacuated using a pump and then slowly filled with air again. The number of Maxima is recorded in steps of 50 mbar. This measurement is repeated five times.

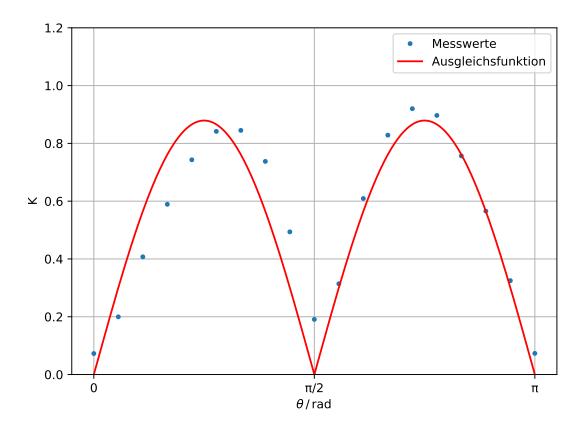
# 4 Analysis

#### 4.1 Kontrast

In order to study the usage of interferrometrie to determine refraction indices, first of all the kontrast of the used Sagnac Interferometer was calculated. The kontrast was calculated with equation (??) and the values of table 1. Table 1 also contains the calculated kontrast values. The maximum contrast K=0.92 was measured at  $\Phi=130^\circ$  Therefore the polarization filter was set to  $\Phi=130^\circ$  for the following measurements. In addition, a fit of the form

$$K = A \cdot |\cos(\Phi)\sin(\Phi)| \tag{10}$$

is performed with the mean values of the measured values. This can be seen for the determined value of  $A=1.76\pm0.07$  in graph 3.



**Abbildung 3:** The measured values of the contrast according to equation (??) and the fit calculation according to equation (??).

**Tabelle 1:** Recorded measured values for contrast measurement, as well as the respective contrast value.

$\Phi/^{\circ}$	$U_{ m min,1}/{ m V}$	$U_{\mathrm{max},1}/\mathrm{V}$	$K_1$	$U_{\mathrm{min},2}/\mathrm{V}$	$U_{\mathrm{max,2}}/\mathrm{V}$	$K_2$	$U_{\mathrm{min,3}}/\mathrm{V}$	$U_{\rm max,3}/{\rm V}$	$K_3$
0	1.73	1.98	0.06	1.8	2.1	0.07	1.76	2.04	0.07
10	1.21	1.8	0.19	1.16	1.78	0.21	1.24	1.83	0.19
20	0.69	1.65	0.41	0.7	1.63	0.39	0.69	1.66	0.41
30	0.36	1.44	0.60	0.39	1.46	0.57	0.38	1.47	0.58
40	0.21	1.48	0.75	0.23	1.47	0.72	0.21	1.46	0.74
50	0.14	1.7	0.84	0.15	1.64	0.83	0.14	1.66	0.84
60	0.17	1.98	0.84	0.16	1.93	0.84	0.16	1.92	0.84
70	0.37	2.35	0.72	0.32	2.15	0.74	0.33	2.25	0.74
80	0.85	2.36	0.47	0.7	2.22	0.52	0.79	2.31	0.49
90	1.64	2.36	0.18	1.48	2.28	0.21	1.62	2.33	0.17
100	1.43	2.93	0.34	1.46	2.68	0.29	1.49	2.79	0.30
110	1.08	4.13	0.58	0.92	3.78	0.60	0.91	4.06	0.63
120	0.52	5.06	0.81	0.49	5.06	0.82	0.44	5.36	0.84
130	0.24	5.53	0.91	0.25	5.56	0.91	0.21	5.76	0.92
140	0.34	5.92	0.89	0.33	5.91	0.89	0.3	5.99	0.90
150	0.73	5.26	0.75	0.77	5.39	0.75	0.73	5.45	0.76
160	1.16	4.13	0.56	1.22	4.25	0.55	1.16	4.38	0.58
170	1.63	3.18	0.32	1.67	3.19	0.31	1.62	3.29	0.34
180	1.86	2.08	0.05	1.85	2.18	0.08	1.79	2.11	0.08

### 4.2 Refraction index of glas

To determine the refractive index of glass, the number of intensity maxima M was recorded. The refractive index was determined using the equation (??). The thickness of the plates is D=1 mm, the wavelength of the laser  $\lambda_0=632.990$  nm and  $\Omega_0=10^\circ$  as the two plates are each  $\Omega_0$  inclined. The measured values, as well as the refractive index determined in each case, can be found in Table 2. On average, the refractive index determined for glass is

$$n_{\rm Glas} = 1.64 \pm 0.13.$$

**Tabelle 2:** AMeasured values to determine the refractive index of glass and the determined refractive index.

Durchgang	M	$n_{ m Glas}$
1	38	1.652392
2	38	1.652392
3	38	1.652392
4	37	1.624503
5	38	1.652392
6	38	1.652392
7	37	1.624503
8	38	1.652392
9	38	1.652392
10	37	1.624503

#### 4.3 Refraction index of air

To determine the refractive index of air, the measured values were recorded as described in chapter ?? and determined using equation ??. The length of the gas chamber is  $L=(100.0\pm0.1)$  mm and the temperature  $T=20.6~^{\circ}\mathrm{C}$ . The recorded values as well as the calculated reflection indices are shown in Table 3.

**Tabelle 3:** Measured values recorded to determine the refractive index of air next to the refractive index calculated according to equation (??). Here,  $M_i$  denotes the number of interference minima or maxima that have passed up to that point, where i indicates the passage.

p/mbar	$M_1$	$n_1$	$M_2$	$n_2$	$M_3$	$n_3$	$M_4$	$n_4$
50	2	1,00001266	2	1,00001266	2	1,00001266	3	1,000 018 99
100	4	1,00002532	4	1,00002532	4	1,00002532	5	1,00003165
150	7	1,00004431	6	1,00003798	6	1,00003798	7	1,00004431
200	9	1,00005697	8	1,00005064	8	1,00005064	9	1,00005697
250	11	1,00006963	10	1,00006330	10	1,00006330	11	1,00006963
300	13	1,00008229	12	1,00007596	12	1,00007596	13	1,00008229
350	15	1,00009495	15	1,00009495	15	1,00009495	15	1,00009495
400	17	1,00010761	17	$1,\!00010761$	17	1,00010761	18	1,00011394
450	20	1,00012660	19	1,00012027	19	1,00012027	20	1,00012660
500	22	1,00013926	21	1,00013293	21	1,00013293	22	1,00013926
550	24	1,00015192	23	1,00014559	23	1,00014559	24	1,00015192
600	26	1,00016458	25	$1,\!00015825$	25	1,00015825	26	1,00016458
650	28	1,00017724	27	1,00017091	27	1,00017091	28	1,00017724
700	30	1,00018990	30	1,00018990	29	1,00018357	30	1,00018990
750	32	1,00020256	31	1,00019623	32	1,00020256	33	1,00020889
800	35	1,00022155	34	1,00021522	34	1,00021522	35	1,00022155
850	37	1,00023421	36	$1,\!00022788$	36	1,00022788	37	1,00023421
900	39	1,00024687	38	$1,\!00024054$	38	1,00024054	39	$1,\!00024687$
950	41	$1,\!00025953$	40	$1,\!00025320$	40	$1,\!00025320$	41	$1,\!00025953$

#### 4.4 Lorentz-Lorenz law

Since the refractive index also dependents on temperature and pressure according to the Lorentz-Lorenz law, an fit calculation is carried out according to equation ?? . The fit has the form

$$n = \frac{a}{TR} \cdot p + b \tag{11}$$

The temperate is T=20.6 °C and R describes the universal gas constant. This results in the in table 4 shown values for a and b.

**Tabelle 4:** The results of the fit for the variables a and b for each run.

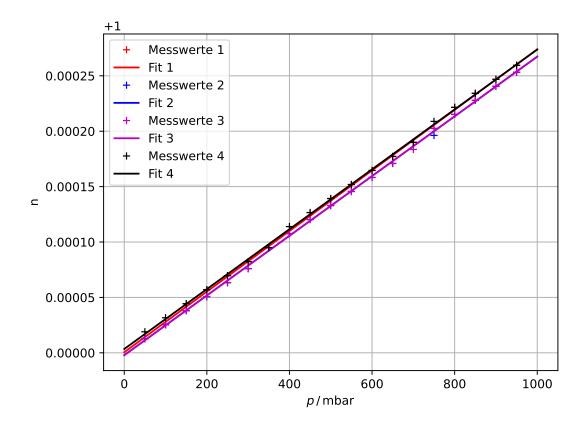
Messung	$a/\left(10^{-2}\mathrm{m}^3/\mathrm{mol}\right)$	b
1	$0.00066776 \pm 0.00000394$	$1.00000056\ \pm0.00000092$
2	$0.00065800 \pm 0.00000397$	$0.99999789 \pm 0.00000093$
3	$0.00065854 \pm 0.00000364$	$0.99999778 \pm 0.00000085$
4	$0.00066017 \pm 0.00000370$	$1.00000344 \pm 0.00000086$

The fits and thus the calculated refractive indices are shown in Figure 4 . The average values for the variables are:

$$a = 0.0006611 \pm 0.0000019 \frac{m^3}{\text{mol}}$$
$$b = 0.9999999 \pm 0.0000004.$$

According to the Loretz-Lorenz law, this results in the following refractive index in a normal atmosphere with T = 15  $^{\circ}$ C and p = 1013 hPa:

$$n = 1.0002795 \pm 0.0000009 \tag{12}$$



**Abbildung 4:** The calculated refractive indices n for air and the fit.

# 5 Diskussion

The measured values and the results obtained are in line with the expectations. An overview of the determined and theoretical values, as well as the respective deviation, can

be found in table Tabelle 5. The determined contrasts follow the expected distribution and no major deviations are recognizable. As expected, the extremes can also be found at multiples of 45°.

For the refractive index of glass,  $n_{\rm glass}=1.64\pm0.13$  was determined. The theoretical value is  $n_{\rm glass,\ theo}=1.45$  [Brechungsindex]. The determined value therefore has a deviation of  $13,10\,\%$ . One reason for this slightly higher deviation may lie in the way the experiment was carried out. Here, the intensity maxima and minima were only recorded with one diode instead of two.

The theoretical value for air is  $n_{\rm air, theo} = 1.000292$  [Brechungsindex]. Averaged over all measurement series, the refractive index of air in a standard atmosphere is  $n_{\rm air} = 1.0002795 \pm 0.0000009$ , which represents a deviation of  $\ll 1\%$  from the theoretical value.

**Tabelle 5:** The refractive indices determined for glass and air compared to the respective theoretical values.

	$n_{ m Glas}$	$n_{ m Luft}$
Theorie	$1,\!45$	1,000292
Versuch	$1.64 \pm 0.13$	$1.0002795 \pm 0.0000009$
Abweichung	13.10~%	0.0012~%