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Lecture 9

Unstructured Search

- Grover's algorithm
- Amplitude amplification

1 Recap & Preview

$$\begin{aligned} & \text{Deutsch-Jozsa} \\ & f: \{0,1\}^n \longmapsto \{0,1\}. \\ & \text{constant or balanced.} \end{aligned}$$

$$\begin{cases}
0 \cdots 0 \\
\vdots \\
1-1
\end{cases} 2^n = N$$

$$s: 2^n \text{ binary string.}$$

Simon's problem
$$f: \{0,1\}^n \longmapsto X$$
. $f(x) = f(y)$ iff $x = y$ or $x = y \otimes s$.

Less requirement than Deutsch-Jozsa but still structured. Phase estimation $U |\psi\rangle = e^{i\theta} |\psi\rangle$. Find θ .

A genuinely quantum problem. $O(1/\epsilon)$ queries, w.p. $\geq \frac{8}{\pi^2}$. Going beyond Hadamard. QFT from $Z_2 \otimes \cdots \otimes Z_2$ to Z_{2^n} .

Application: Order finding ⇒ Shor's algorithm.

For all of these, we have structural assumptions:

- Deutsch-Jozsa and Simon's problem: Special f.
- Order finding and Shor's algorithm: Cyclic group \mathbb{Z}_N .

How about we characterize "very general functions"? Start with boolean functions.

Total function: A function $f: \{0,1\}^N \to \{0,1\}$ which has definition on all 2^N inputs. Deutsch-Jozsa: $N = 2^n$, But the definition domain is only $\{s = s_1 \cdots s_{2^n} \mid \# \text{ of } 1 \text{ in } s_i \text{ is } 0, 2^{n-1}, \text{ or } 2^n\}$. If a function is defined on a proper subset of $\{0,1\}^n$. it's called a partial function.

A very typical problem is to compute the OR function (AND is symmetric).

$$f: \{0,1\}^n \to \{0,1\}, N=2^n.$$
 $OR(s_1,\ldots,s_N) = s_1 \cup s_2 \ldots \vee s_N = \begin{cases} 0 & \text{if } s_1,\ldots s_n = 0, \\ 1 & \text{otherwise.} \end{cases}$

Classically: Need $\Theta(n)$ queries to compute the value of OR.

Quantum: $|i, z\rangle \xrightarrow{U_f} |i, z \oplus s_i\rangle, \forall i \in [n], z \in \{0, 1\}$ (*).

Equivalent to having an oracle:

$$f(x) = s_x, \forall x \in [N] \ (f:[N] \to \{0,1\}) \qquad |x,z\rangle \xrightarrow{U_f} |x,f(x) \otimes z\rangle, \quad \forall x \in [N], z \in \{0,1\}.$$

Main focus today: $O(\sqrt{N})$ quantum queries suffice.

2 Grover's algorithm

Phase kick-back: $|x\rangle|-\rangle \xrightarrow{U_f} (-1)^{f(x)}|x\rangle|-\rangle$.

"phase query":
$$|x\rangle \longmapsto (-1)^{f(x)}|x\rangle \quad |\psi\rangle = \frac{1}{\sqrt{N}} \sum_{x=1}^{N} |x\rangle \stackrel{\mathrm{U}}{\longmapsto} \frac{1}{\sqrt{N}} \left(\sum_{x,f(x)=0} |x\rangle - \sum_{x\cdot f(x)=1} |x\rangle \right)$$

How can we find x such that f(x) = 1?

This looks like a reflection.

For simplicity, consider a unique marked item w s.t. f(w) = 1, $f(x) = 0 \ \forall x \neq w$.

$$\begin{array}{l} U|w\rangle = -|w\rangle \\ U|x\rangle = |x\rangle \quad \forall x \neq w. \end{array} \right\} \ U = I - 2|w\rangle\langle w|.$$

$$(I - 2|w\rangle\langle w|)|w\rangle = |w\rangle - 2|w\rangle\langle w|w\rangle = -|w\rangle.$$

$$(I - 2(w)\langle w|)|x\rangle = |x\rangle - 2|w\rangle\langle w|x\rangle = |x\rangle$$

$$(I-2|w\rangle\langle w|)(I-2|w\rangle\langle w|) = I-2|w\rangle\langle w|-2|w\rangle\langle w|+4|w\rangle\langle w|w\rangle\langle w|$$

=I.

We also consider $|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{x=1}^{N} |x\rangle$.

We consider the unitary $V=2|\psi\rangle\langle\psi|-I$. V is independent of the queries, and V can be efficiently implemented with cost $O(\log N)$.

Say $N = 2^n$. Otherwise find the smallest power of 2 larger than N $(2^n < N < 2^{n+1})$, and set $f(N + 1) \dots, f(2^{n+1}) = 0$.

$$\frac{1}{\sqrt{2^n}} \sum_{x=1}^{2^n} |x\rangle = H^{\otimes n} |0\rangle \quad V = H^{\otimes n} R_0 H^{\otimes n}, \text{ where } R_0 = 2 |0^n\rangle \langle 0^n| - I.$$

$$H^{\otimes n} \left(2 |0^n\rangle \langle 0^n| - I \right) H^{\otimes n}$$

 $R_0: |0^n\rangle \to |0^n\rangle, |x\rangle \to -|x\rangle \quad \forall x \in \{0,1\}^n/\{0^n\}.$

$$\begin{array}{c|c} & & & & \\ \hline & & & & \\ \hline -I \\ \hline -I \\ \hline \end{array}$$

As a conclusion, V can be implemented with cost $O(\log N)$. Same as preparing $|\psi\rangle = H^{\otimes n}|0\rangle$.

Grover algorithm:

• - Prepare
$$|\psi\rangle$$

$$U|\psi\rangle = (I - 2|\omega\rangle\langle\omega|)|\psi\rangle = |\psi\rangle - 2\langle\omega|\psi\rangle|\omega\rangle = |\psi\rangle - \frac{2}{\sqrt{N}}|\omega\rangle.$$
• - Repeat $t = \lceil \frac{\pi}{4}\sqrt{n} \rceil$ times
$$U|\omega\rangle = -|\omega\rangle$$
Apply U ;
$$V|\psi\rangle = (2|\psi\rangle\langle\psi| - I)|\psi\rangle = |\psi\rangle$$

$$V|\psi\rangle = (2|\psi\rangle\langle\psi| - I)|\omega\rangle = 2(\langle\psi|\omega\rangle|\psi\rangle - |\omega\rangle) = \frac{2}{\sqrt{N}}|\psi\rangle - |\omega\rangle$$

• - Measure in the computational basis

Therefore, the subspace span $\{|\psi\rangle, |\omega\rangle\}$ is invariant under U and V.

However, $\langle \psi | \omega \rangle \neq 0$. It will read better to consider an orthonormal basis span $\{ |\omega\rangle, |\omega^{\perp}\rangle \}$:

$$|\omega^{\perp}\rangle = \frac{|\psi\rangle - \langle\omega|\psi\rangle |\omega\rangle}{\text{normalization}}, \quad \langle\omega|\omega^{\perp}\rangle = \langle\omega|\psi\rangle - \langle\omega|\psi\rangle |\omega|\omega\rangle = 0.$$

$$U = I - 2|w\rangle \langle w|, \qquad |w\rangle |w^{\perp}\rangle$$

$$U|w\rangle = (I - 2|w\rangle \langle w|)|w\rangle = -1|w\rangle + 0|w^{\perp}\rangle$$

$$U|w^{\perp}\rangle = (I - 2|w\rangle \langle w|)|w^{\perp}\rangle = 0|w\rangle + 1|w^{\perp}\rangle$$

$$U|\psi\rangle = \frac{1}{\sqrt{N}}|\omega\rangle + \sqrt{1 - \frac{1}{N}}|\omega^{\perp}\rangle = \sin\theta |\omega\rangle + \cos\theta |\omega^{\perp}\rangle \quad \left(\theta = \arcsin\frac{1}{\sqrt{N}}\right).$$

$$V = 2|\psi\rangle \langle \psi| - I = 2\left(\frac{\sin\theta}{\cos\theta}\right) (\sin\theta - \cos\theta) - I = \left(\frac{2\sin^2\theta - 1}{2\sin\theta\cos\theta} - \frac{2\sin\theta\cos\theta}{2\cos^2\theta - 1}\right) = \left(\frac{-\cos 2\theta - \sin 2\theta}{\sin 2\theta} - \frac{\sin 2\theta}{\cos 2\theta}\right).$$

$$VU = \left(\frac{-\cos 2\theta - \sin 2\theta}{\sin 2\theta} - \frac{\sin(2(t - 1)\theta)}{\cos(2(t - 1)\theta)} - \frac{\cos(2(t - 1)\theta)}{\sin(2(t - 1)\theta)} - \frac{\cos(2(t - 1)\theta)\sin 2\theta}{\cos(2(t - 1)\theta)\sin 2\theta} - \frac{\cos(2(t - 1)\theta)\cos 2\theta}{\sin(2(t - 1)\theta)\cos 2\theta}.$$

$$= \left(\frac{\cos(2(t - 1)\theta)\cos 2\theta - \sin(2(t - 1)\theta)\sin 2\theta}{-\sin(2(t - 1)\theta)\cos 2\theta - \cos(2(t - 1)\theta)\sin 2\theta} + \frac{\sin(2(t - 1)\theta)\cos 2\theta}{\cos(2(t - 1)\theta)\cos 2\theta}\right)$$

$$= \left(\frac{\cos 2t\theta}{\sin 2t\theta} - \frac{\sin 2t\theta}{\cos 2t\theta}\right).$$

$$(VU)^t|\psi\rangle = \left(\frac{\cos 2t\theta}{\sin 2t\theta} - \frac{\sin 2t\theta}{\cos 2t\theta}\right).$$

$$= \left(\frac{\cos 2t\theta}{-\sin 2t\theta} - \frac{\sin 2t\theta}{\cos 2t\theta}\right).$$

$$(VU)^t|\psi\rangle = \left(\frac{\cos 2t\theta}{-\sin 2t\theta} - \frac{\sin 2t\theta}{\cos 2\theta}\right).$$

$$= \left(\frac{\cos 2t\theta}{-$$

What if there are M marked items?

$$|\omega\rangle \longmapsto \frac{1}{\sqrt{M}} \sum_{x: f(x)=1} |x\rangle \quad \sin\theta = \langle \psi | w \rangle = \langle \psi | \frac{1}{\sqrt{M}} \sum_{x: f(x)=1} |x\rangle = \frac{1}{\sqrt{N}} \cdot \frac{1}{\sqrt{M}} \cdot M = \sqrt{\frac{M}{N}}.$$

 $\Rightarrow O\left(\sqrt{\frac{N}{M}}\right)$ steps suffice for Grover's algorithm.

Remark 1. What if M is unknown? Without knowing M, we may "overshoot".

Research paper: Yoder, Low, Chuang. Fixed-point quantum search with an optimal number of queries. PRL 2014. arxiv: 1409.3305.

Contribution: An algorithm without overshooting, nor knowing M. Cost: $O\left(\sqrt{\frac{N}{M}}\right)$.

Or naively: Guess M using a geometric series: $M = \frac{n}{2} - \frac{\pi}{4} \sqrt{\frac{N}{M}}$. $M \leftarrow M \cdot \frac{3}{4}$. Iteration: $\log N$.

3 Amplitude amplification

In general, we can consider a form of having a unitary U acting on l qubits s.t.

$$U |0^l\rangle = \sqrt{p}|1\rangle |\psi_1\rangle + \sqrt{1-p}|0\rangle |\psi_0\rangle$$

where $|\psi_1\rangle$ and $|\psi_0\rangle$ are normalized (l-1)-qubit quantum states.

We can think of $|1\rangle |\psi_1\rangle$ as the "good state" and $|0\rangle |\psi_0\rangle$ as the "bad state".

Similar to Grover:

- 1. A reflection with respect to the bad state $|0\rangle |\psi_0\rangle$.
- 2. A refection with respect to $U|0^l\rangle$.
- 1: This is basically putting a "-" in front of $|1\rangle |\psi_1\rangle$ and leaving $|0\rangle |\psi_0\rangle$ alone. Solution: Put a Z on the first quit.

$$Z|0\rangle = |0\rangle$$

$$Z|1\rangle = -|1\rangle$$

2: Same to Grover, apply UR_0U^{\dagger} .

Applying 1 and 2 alternatively for t times, we get

$$\sin((2t+1)\theta) |1\rangle |\psi_1\rangle + \cos((2t+1)\theta) |0\rangle |\psi_0\rangle$$
, where $\theta = \arcsin\sqrt{p} \approx \sqrt{p}$.

Taking $(2t+1)\theta \approx \frac{\pi}{2} \Rightarrow t = O\left(\frac{1}{\sqrt{p}}\right)$, we can (approximately) get $|1\rangle |\psi_1\rangle$.

This is known as amplitude amplification.

Remark 2. In fact, there's also a procedure called amplitude estimation, where we can quantitatively output a \tilde{p} s.t. $|\tilde{p}-p| \leq \varepsilon p$, with high probability (say 2/3). using $O(1/\varepsilon)$ queries to U.

On the other hand: Classically, tossing a coin, getting a head with prob. p, need $O(1/\varepsilon^2)$ tries to estimate such a \tilde{p} . Quantum: Amp Est has quadratic speedup.

Amp Amp: Classically, tossing a coin, getting a head with prob. p, need O(1/p) tries to see the first head.

$$\Pr[\text{fail}] = 1 - (1 - p)^m \quad (1 - p)^{\frac{1}{p}} \approx \frac{1}{e}$$

$$\geqslant \frac{2}{3} \quad m \sim \frac{1}{p}$$

Quantum: Amp Amp has quadratic speedup.

Amp Est: Brassard, Hoyer, Mosca, and Tapp. Quantum Amplitude Amplification and Estimation. Contemporary Mathematics, 2002. arxiv: quant-ph/0005055.