

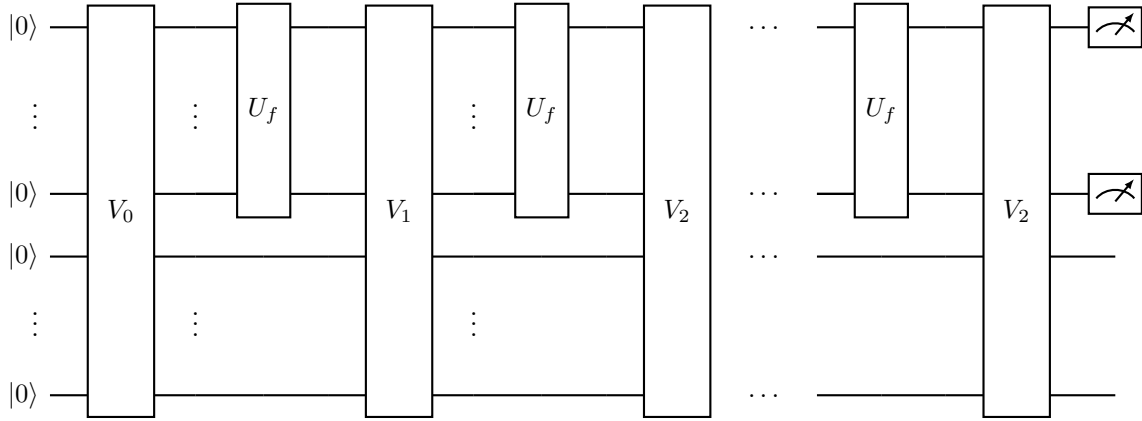
Lecture 10

Unstructured Search and Discrete-time Quantum Walk

- Lower bound of Grover search
- Random walks
- Block encoding

1 Lower Bound of Grover Search

After all, a quantum algorithm for searching looks below: $N = 2^n$



$|\psi_t\rangle :=$ state after V_t , the t -th non-query operation, assuming no marked item.

$|\psi_t^x\rangle :=$ state after V_t , assuming a unique marked item x .

To solve the unstructured search problem, we must have $\| |\psi_T\rangle - |\psi_T^x\rangle \| \geq c$ for a constant c . Since the algorithm must work for any x , we must have

$$\sum_{x=1}^N \| |\psi_T\rangle - |\psi_T^x\rangle \| \geq cN$$

Since $|\psi_0^x\rangle = |\psi_0\rangle$ for all x , $\sum_{x=1}^N \| |\psi_0\rangle - |\psi_0^x\rangle \| = 0$.

$$\begin{aligned}
\| |\psi_t\rangle - |\psi_t^x\rangle \| &= \| V_t (\psi_{t-1}) - V_t U_x |\psi_{t-1}^x\rangle \| && U_x = I - 2|x\rangle\langle x|, \text{ the black box} \\
&= \| |\psi_{t-1}\rangle - U_x |\psi_{t-1}^x\rangle \| && \text{when the marked item is } x. \quad U_x^2 = I \\
&= \| U_x |\psi_{t-1}\rangle - |\psi_{t-1}^x\rangle \| \\
&= \| (U_x |\psi_{t-1}\rangle - |\psi_{t-1}\rangle) + (|\psi_{t-1}\rangle - |\psi_{t-1}^x\rangle) \| && \text{triangle inequality} \\
&\leq \| U_x |\psi_{t-1}\rangle - |\psi_{t-1}\rangle \| + \| |\psi_{t-1}\rangle - |\psi_{t-1}^x\rangle \|.
\end{aligned}$$

$|\psi_t\rangle = \sum_{y=1}^N \alpha_{y,t} |y\rangle |\phi_y\rangle$ collect all the vectors where the first subsystem is y
 $|\phi_y\rangle$ is normalized. $\sum_{y=1}^N |\alpha_{y,t}|^2 = 1$.

$$U_x |\psi_t\rangle = \sum_{y \neq x} \alpha_{y,t} |y\rangle |\phi_y\rangle - \alpha_{x,t} |x\rangle |\phi_x\rangle = |\psi_t\rangle - 2\alpha_{x,t} |x\rangle |\phi_x\rangle.$$

$$\Rightarrow \| U_x (p_t) - |\psi_t\rangle \| = 2 |\alpha_{x,t}|.$$

This further implies

$$\begin{aligned}
\| |\psi_T^x\rangle - |\psi_T\rangle \| &\leq 2 \sum_{j=1}^{T-1} |\alpha_{x,j}|. && (\sum_i a_i^2)(\sum_i b_i^2) \geq (\sum_i a_i b_i)^2 \\
\Rightarrow CN &\leq \sum_{x=1}^N \| |\psi_T\rangle - |\psi_T^x\rangle \| \leq 2 \sum_{x=1}^N \sum_{j=1}^{T-1} |\alpha_{x,j}| = 2 \sum_{j=1}^{T-1} \sum_{x=1}^N |\alpha_{x,j}|. \\
&\leq 2 \sum_{j=1}^{T-1} \sqrt{N} \cdot \sqrt{\sum_{x=1}^N |\alpha_{x,j}|^2} = 2\sqrt{N}(T-1) \\
2\sqrt{N}(T-1) &\geq CN \Rightarrow T = \Omega(\sqrt{N}).
\end{aligned}$$

Corollary 1.1. *Grover search is optimal for unstructured search.*

Remark 1.2. From a high-level picture, OR gives a quadratic quantum speedup. Is this the best we can achieve for total functions?

For a total function $f : \{0,1\}^N \rightarrow \{0,1\}$, denote $D(f), R(f), Q(f)$ to be its classical deterministic, classical randomized, and quantum query complexities.

It was widely conjectured that $R(f) = O(Q(f)^2)$ for a long time.

[Aaronson, Ben-David, and Kothari, STOC 2016](#): $\exists f$ st. $R(f) = \tilde{\Omega}(Q(f)^{2.5})$.

State-of-the-art:

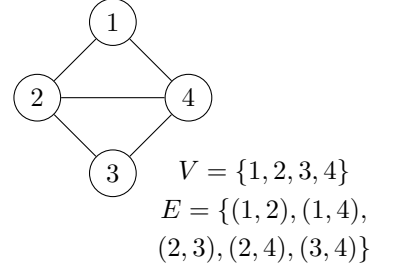
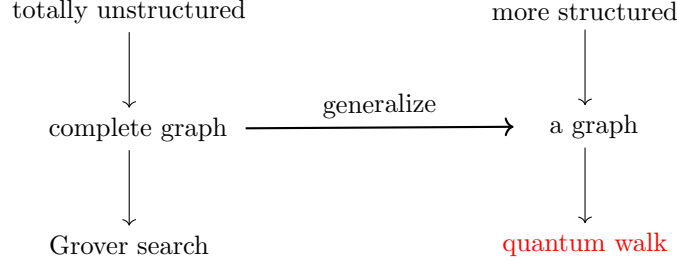
1. \forall total f , $D(f) = O(Q(f)^4)$. [\[Ambainis-Balodis-Belovs-Lee-Santha-Smotrovs, JACM 2017\]](#)
 2. total f , $D(f) = \Omega(Q(f)^4)$. [\[Aaronson-Ben-David-Kothari-Rao-Tal, STOC 2021\]](#)
 3. total f , $R(f) = \Omega(Q(f)^3)$. [\[Bansal-Sinha STOC 2021\]](#)
- (since $D(f) \geq R(f), R(f) = O(Q(f)^4)$. [\[Sherstov-Storozhenko-Wu STOC 2021\]](#))

2 Discrete-time Quantum Walk

Generalize unstructured search? **Break the symmetry in a general sense.**

A common tool to characterize discrete objects: graphs $G = (V, E)$

Grover: basically a problem on complete graph



2.1 Random walks

Classical random walk on graphs. Given a graph $G = (V, E)$, a random walk is given by a transition matrix P , with its entry P_{ij} denoting the probability of the transition from vertex i to vertex j .

Such a P is called a **stochastic matrix**, satisfying: $P_{ij} \geq 0 \quad \forall i, j; \sum_j P_{ij} = 1 \quad \forall i$.

Properties of random walks (Markov Chains):

- **Irreducible:** Any vertex can be reached from any other vertex in a finite number of steps.
- **Aperiodic:** There exists no integer greater than 1 that divides the length of every cycle of the graph.
- **Ergodic:** Both irreducible and aperiodic.

Theorem 2.1 (Perron-Frobenius Theorem). Any ergodic random walk P has a **unique** stationary state π such that $\pi_i > 0 \quad \forall i$, $\sum_i \pi_i = 1$, and $\sum_i \pi_i P_{ij} = \pi_j \quad \forall j$.

In other words, π is the left eigenvector of P with eigenvalue 1.

- **Reversible:** The detailed balance condition $\pi_i P_{ij} = \pi_j P_{ji} \quad \forall i, j$ is satisfied.
- **Discriminant matrix:** A matrix D where $D_{ij} = \sqrt{P_{ij} P_{ji}} \quad \forall i, j$.

D is real symmetric, and hence Hermitian.

Proposition 2.2. If a random walk is ergodic and reversible, then the stationary state $|\pi\rangle = \sum_i \sqrt{\pi_i} |i\rangle$ is an eigenvector of D with eigenvalue 1, i.e. $D|\pi\rangle = |\pi\rangle$. Furthermore, $D = \text{diag}(\sqrt{\lambda}) \cdot P \cdot (\text{diag}(\sqrt{\lambda}))^{-1}$, where

$\text{diag}(\sqrt{\lambda}) = \begin{pmatrix} \sqrt{\lambda_1} & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & \sqrt{\lambda_N} \end{pmatrix}$. Hence, the set of (left) eigenvalues of P and the set of the eigenvalues of D are the same.

Proof. For any i ,

$$\begin{aligned}
 (D|\pi\rangle)_i &= \langle i|D|\pi\rangle & (D = \sum_{i,j} D_{ij} |i\rangle\langle j|) \\
 &= \sum_j D_{ij} \langle j|\pi\rangle \\
 &= \sum_j \sqrt{P_{ij} P_{ji}} \sqrt{\pi_j} & (\text{reversibility: } \pi_i P_{ij} = \pi_j P_{ji}) \\
 &= \sum_j P_{ij} \sqrt{\pi_i} = \sqrt{\pi_i}.
 \end{aligned}$$

Furthermore. for any i, j , $(D)_{ij} = \sqrt{P_{ij}P_{ji}} = \sqrt{\pi_i}P_{ij}(\sqrt{\pi_j})^{-1} = \left(\text{diag}(\sqrt{\lambda}) \cdot P \cdot \text{diag}(\sqrt{\lambda})^{-1}\right)_{ij}$. Therefore, $D = \text{diag}(\sqrt{\lambda}) \cdot P \cdot \text{diag}(\sqrt{\lambda})^{-1}$. ■

2.2 Block encoding

Now, how shall we define quantum walks?

A simplest idea (say for unweighted graph): $|j\rangle \rightarrow |\partial_j\rangle := \frac{1}{\sqrt{\deg(j)}} \sum_{k:(j,k) \in E} |k\rangle$

What's the issue with this? Not a unitary map in general.

$$\langle 2|4\rangle = 0$$

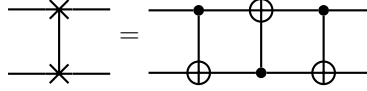
$$\langle \partial_2|\partial_4\rangle = \left(\frac{1}{\sqrt{3}}\langle 1| + \frac{1}{\sqrt{3}}\langle 3| + \frac{1}{\sqrt{3}}\langle 4|\right) \left(\frac{1}{\sqrt{3}}|1\rangle + \frac{1}{\sqrt{3}}|2\rangle + \frac{1}{\sqrt{3}}|3\rangle\right) = \frac{2}{3}$$

Instead, we consider: $\text{Op } |0^n\rangle |j\rangle = \sum_k \sqrt{P_{jk}} |k\rangle |j\rangle$.

The discriminant matrix looks good-it's symmetric. Can we make that?

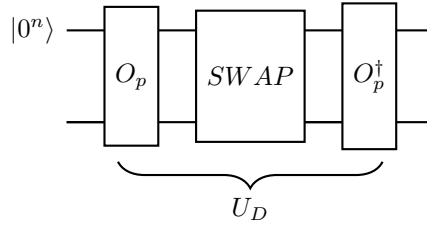
We also introduce the n-qubit swap operator: $\text{SWAP } |i\rangle |j\rangle = |j\rangle |i\rangle \quad \forall i, j \in \{0, 1\}^n$.

(SWAP can be implemented by $\binom{n}{2} = O(n^2)$ 2-qubit swap gates. In total, SWAP can be implemented by $O(n^2)$ CNOTs.)



$$\begin{aligned} \text{SWAP } |i\rangle |j\rangle &= |j\rangle |i\rangle \\ \forall i, j &\in \{0, 1\} \end{aligned}$$

We consider



UD: 1 step quantum walk.

Proposition 2.3. $\forall i, j \in [N]$, $\langle 0^n | \langle i | U_D | 0^n \rangle | j \rangle = D_{ij}$. Intuitively, $U_D = \begin{pmatrix} D & \cdot \\ \cdot & \cdot \end{pmatrix} = |0^n\rangle \langle 0^n| \otimes D + \dots$. Such a form is called a **block encoding** of D .

Proof.

For any j , $|0^n\rangle |j\rangle \xrightarrow{O_p} \sum_k \sqrt{P_{jk}} |k\rangle |j\rangle \xrightarrow{\text{SWAP}} \sum_k \sqrt{P_{jk}} |j\rangle |k\rangle$.

Meanwhile $|0^n\rangle |i\rangle \xrightarrow{O_p} \sum_{k'} \sqrt{P_{ik'}} |k'\rangle |i\rangle$. $\delta_{x,y} = \begin{cases} 1 & x = y \\ 0 & x \neq y \end{cases}$

So the inner product gives

$$\langle 0^n | \langle i | U_D | 0^n \rangle | j \rangle = \sum_{k, k'} \sqrt{P_{ik'}} \sqrt{P_{jk}} \delta_{k', j} \delta_{i, k} = \sqrt{P_{ij}} \sqrt{P_{ji}} = D_{ij}.$$

■

Next: What can be obtain by using the block encoding?

Denote the eigendecomposition of D as $D = \sum_i \lambda_i |v_i\rangle \langle v_i|$.

Far each eigenstate $|v_i\rangle$,

$$U_D |0^n\rangle |v_i\rangle = |0^n\rangle D |v_i\rangle + |\tilde{\perp}_i\rangle = \lambda_i |0^n\rangle |v_i\rangle + |\tilde{\perp}_i\rangle \quad (*)$$

Here $|\tilde{\perp}_i\rangle$ is an unnormalized state satisfying $\Pi |\tilde{\perp}_i\rangle = 0$, where

$$\Pi = |0^n\rangle \langle 0^n| \otimes I$$

(*) should give a normalized state, so we may write $|\tilde{\perp}_i\rangle = \sqrt{1 - \lambda_i^2} |\perp_i\rangle$, where $|\perp_i\rangle$ is a normalized state.

Suppose $\lambda_i \neq \pm 1$. First, notice that $U_D^\dagger = U_D$ ($(ABC)^\dagger = C^\dagger (AB)^\dagger = C^\dagger B^\dagger A^\dagger$)

$$U_D^\dagger = (\text{Op SWAP Op}^\dagger)^\dagger = (\text{Op}^\dagger)^\dagger \text{SWAP}^\dagger \text{Op}^\dagger = \text{Op SWAP Op}^\dagger = U_D$$

Apply U_D to both sides of (*), we have $(U_D^2 = U_D^\dagger U_D = I)$

$$\begin{aligned} |0^n\rangle |v_i\rangle &= \lambda_i U_D |0^n\rangle |v_i\rangle + \sqrt{1 - \lambda_i^2} U_D |\perp_i\rangle \\ &= \lambda_i \left(\lambda_i |0^n\rangle |v_i\rangle + \sqrt{1 - \lambda_i^2} |\perp_i\rangle \right) + \sqrt{1 - \lambda_i^2} U_D |\perp_i\rangle. \\ \Rightarrow (1 - \lambda_i^2) |0^n\rangle |v_i\rangle &= \sqrt{1 - \lambda_i^2} \lambda_i |\perp_i\rangle + \sqrt{1 - \lambda_i^2} U_D |\perp_i\rangle \\ \Rightarrow U_D |\perp_i\rangle &= \sqrt{1 - \lambda_i^2} |0^n\rangle |v_i\rangle - \lambda_i |\perp_i\rangle. \end{aligned}$$