

PHY 4210-01 Senior Lab  
Lab N4: Rutherford Scattering

Sarah Arends  
Jacquelyne Miksanek  
Ryan Wojtyla

Instructor: Dr. Marcus Hohlmann

March 14, 2019

**Abstract**

# Contents

<b>1</b>	<b>Objective of the Experiment</b>	<b>3</b>
<b>2</b>	<b>Theory of the Experiment</b>	<b>3</b>
<b>3</b>	<b>Equipment Utilized</b>	<b>6</b>
<b>4</b>	<b>Procedure</b>	<b>6</b>
4.1	Procedural Modifications . . . . .	6
<b>5</b>	<b>Data Analysis</b>	<b>6</b>
5.1	Data Analysis I: Gold . . . . .	6
5.2	Data Analysis II: Aluminum . . . . .	6
<b>6</b>	<b>Results</b>	<b>7</b>
6.1	Results I: Gold . . . . .	7
6.2	Results II: Aluminum . . . . .	7
<b>7</b>	<b>Conclusion</b>	<b>7</b>
<b>8</b>	<b>Appendices</b>	<b>7</b>
8.1	Appendix A: Data . . . . .	7
8.2	Appendix B: Source Code . . . . .	7

# 1 Objective of the Experiment

## 2 Theory of the Experiment

When an alpha particle with impact parameter  $b$  approaches a nucleus, it is scattered at an angle  $\theta$ . If the impact parameter is given an infinitesimal range of  $[b, b + db]$ , the resulting scattering angle then has a range of  $[\theta - d\theta, \theta]$ ; the impact parameter and scattering angle are inversely proportional.

Because the alpha particle can be incident within a defined range at any angle relative to the nucleus, a ring of possible incident locations is created in front of the nucleus. This ring is illustrated in Figure 1.

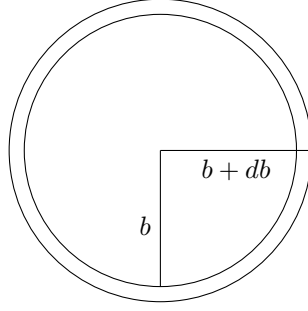


Figure 1: The ring whose area represents the possible region in which alpha particles may be incident on a target nucleus.

The area of this ring is found as the area of any ring is found:

$$\begin{aligned} A &= \pi(b + db)^2 - \pi b^2 \\ &= \pi(b^2 + 2bdb + db^2) - \pi b^2 \\ &= \pi b^2 + 2\pi bdb + \pi db^2 - \pi b^2 \\ &= 2\pi bdb + \pi db^2 \end{aligned}$$

Since  $db$  is infinitesimally small, it can be approximated to be zero. Therefore, the area of the incident ring,  $\Delta\sigma$ , is

$$\Delta\sigma = 2\pi bdb \tag{1}$$

Since the impact parameter  $b$  is directly proportional to the size of the cross section and the scattering angle  $\theta$  is inversely proportional to the impact parameter, the size of the cross section decreases as the scattering angle increases. Therefore, the cross section experiences a negative rate of change as  $\theta$  increases. Hence,

$$\Delta\sigma(\theta) = -d\sigma(\theta) \quad (2)$$

The circumference of a circle is equal to  $2\pi r$ , where  $r$  is the radius of the circle. In the experiment, the radius of the ring onto which the alpha particle is projected after it is scattered is  $R \sin(\theta)$ , where  $\theta$  is the scattering angle and  $R$ , described in Figure 2, is the distance between the point at which the alpha particle was scattered and the edge of the ring.

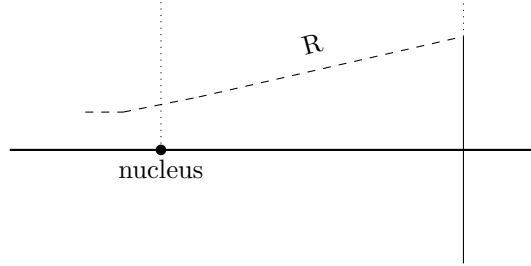


Figure 2: The path of the alpha particle as it is scattered by a nucleus.  $R$  is the path length between the nucleus and the ring of the solid angle.

The outer circumference of this ring is  $2\pi R \sin(\theta)$ . Since the alpha particle is incident within a range whose minimum is  $\theta - d\theta$ , however, the ring has a thickness of  $Rd\theta$ . Since the thickness of the ring is infinitesimal, the ring's area can be approximated to be that of a rectangle. Therefore, the area of the ring is  $A = 2\pi R \sin(\theta) Rd\theta$ .

The solid angle of the scattered alpha particles at an angle  $\theta$  is:

$$\begin{aligned} \Delta\Omega &= \frac{A}{R^2} \\ &= \frac{(2\pi R \sin(\theta) Rd\theta)}{R^2} \end{aligned}$$

$$d\Omega = 2\pi \sin(\theta) \quad (3)$$

An expression for the differential cross section  $\frac{d\sigma}{d\Omega}(\theta)$  can be found by multiplying Equation 2 by Equation 2 divided by itself.

$$\begin{aligned} \Delta\sigma &= -\frac{d\sigma}{d\Omega}(\theta) d\Omega \\ &= -\frac{d\sigma}{d\Omega}(\theta) 2\pi \sin(\theta) d\theta \end{aligned}$$

from Equation 2:

$$-\frac{d\sigma}{d\Omega}(\theta)2\pi\sin(\theta)d\theta = 2\pi bdb$$

$$\frac{d\sigma}{d\Omega}(\theta) = -\frac{b}{\sin(\theta)} \frac{db}{d\theta} \quad (4)$$

Since it is known that  $b = \frac{ZZ'e^2}{2E} \cot\left(\frac{\theta}{2}\right)$ , it can be inserted into Equation 2. Furthermore, since  $b$  is a function of  $\theta$ ,  $\frac{db}{d\theta}$  can also be found:

$$\begin{aligned} b &= \frac{ZZ'e^2}{2E} \cot\left(\frac{\theta}{2}\right) \\ \frac{db}{d\theta} &= \frac{ZZ'e^2}{2E} \left(-\frac{1}{2} \csc^2\left(\frac{\theta}{2}\right)\right) \\ \frac{db}{d\theta} &= -\frac{ZZ'e^2}{4E} \csc^2\left(\frac{\theta}{2}\right) \end{aligned}$$

Now that  $b$  and  $db$  have been found, the full expression for the differential cross section  $\frac{d\sigma}{d\Omega}$  can be determined:

$$\begin{aligned} \frac{d\sigma}{d\Omega}(\theta) &= -\frac{b}{\sin(\theta)} \frac{db}{d\theta} \\ &= -\left(\frac{ZZ'e^2}{2E} \cot\left(\frac{\theta}{2}\right)\right) \frac{1}{\sin(\theta)} \left(-\frac{ZZ'e^2}{4E} \csc^2\left(\frac{\theta}{2}\right)\right) \\ &= 2 \left(\frac{ZZ'e^2}{4E}\right)^2 \frac{\cos\left(\frac{\theta}{2}\right)}{\sin\left(\frac{\theta}{2}\right)} \frac{1}{\sin(\theta)} \frac{1}{\sin^2\left(\frac{\theta}{2}\right)} \\ &= 2 \left(\frac{ZZ'e^2}{4E}\right)^2 \cos\left(\frac{\theta}{2}\right) \frac{1}{(2\sin\left(\frac{\theta}{2}\right)\cos\left(\frac{\theta}{2}\right))} \frac{1}{\sin^3\left(\frac{\theta}{2}\right)} \end{aligned}$$

$$\frac{d\sigma}{d\Omega}(\theta) = \left(\frac{ZZ'e^2}{4E}\right)^2 \frac{1}{\sin^4\left(\frac{\theta}{2}\right)} \quad (5)$$

Equation 2 is the equation for calculating the theoretical differential cross-section.

The differential cross-section can also be found experimentally using a found alpha particle scattering rate at a particular angle  $\theta$ . First, this relation must be constructed. The collimated beam of alpha particles begins its journey with an incident rate of  $\frac{dN_0}{dt}$ . This beam is then incident on a thin foil with an atomic density of  $n = \frac{\rho N_A d}{A}$ , where  $\rho$  is the density of the foil material,  $d$  is the thickness of the foil, and  $A$  is the atomic number of the foil material. By being incident on the foil, the alpha particles are exposed to a differential cross-section at the

particular angle of  $\frac{d\sigma}{d\Omega}(\theta)$ . The alpha particles are scattered by the nuclei across a solid angle  $\Delta\Omega = A_{\text{detector}}r^2$ , where  $A_{\text{detector}}$  is the area of the detector and  $r$  is the distance between the foil and detector. Multiplying these factors together results in the scattering rate of the alpha particles incident on a particular foil at a particular angle:  $\frac{dN}{dt}(\theta) = \frac{dN_0}{dt}n\frac{d\sigma}{d\Omega}(\theta)\Delta\Omega$ . Since the scattering rate is determined experimentally, the equation can be rearranged for the differential cross-section:

$$\frac{d\sigma}{d\Omega}(\theta) = \frac{\frac{dN}{dt}(\theta)}{\frac{dN_0}{dt}n\Delta\Omega} \quad (6)$$

### 3 Equipment Utilized

- List equipment and specifications

Figure 3: Description of schematic here

### 4 Procedure

#### 4.1 Procedural Modifications

### 5 Data Analysis

#### 5.1 Data Analysis I: Gold

#### 5.2 Data Analysis II: Aluminum

## **6 Results**

### **6.1 Results I: Gold**

### **6.2 Results II: Aluminum**

## **7 Conclusion**

## **8 Appendices**

### **8.1 Appendix A: Data**

### **8.2 Appendix B: Source Code**