

Wavelet p-Leader Non Gaussian Multiscale Expansions for Heart Rate Variability Analysis in Congestive Heart Failure Patients

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Abstract-Objective: Numerous indices were devised for the statistical characterization of temporal dynamics of heart rate variability (HRV) with the aim to discriminate between healthy subjects and nonhealthy patients. Elaborating on the concepts of (multi)fractal and nonlinear analyses, the present contribution defines and studies formally novel non Gaussian multiscale representations. Methods: A methodological framework for non Gaussian multiscale representations constructed on wavelet p-leaders is developed, relying a priori neither on exact scale-free dynamics nor on predefined forms of departure from Gaussianity. Its versatility in quantifying the strength and nature of departure from Gaussian is analyzed theoretically and numerically. The ability of the representations to discriminate between healthy subjects and congestive heart failure (CHF) patients, and between survivors and nonsurvivor CHF patients, is assessed on a large cohort of 198 subjects. Results: The analysis leads to conclude that i) scale-free and multifractal dynamics are observed, both for healthy subjects and CHF patients, for time scales shorter than 170 s; ii) a circadian evolution of multifractal and non Gaussian properties of HRV is evidenced for healthy subjects, but not for CHF patients; iii) non Gaussian multiscale indices possess high discriminative abilities between survivor and nonsurvivor CHF patients, at specific time scales (\sim 20 s and \sim 85 s). Conclusions: The non Gaussian multiscale representations provide evidence for the existence of short-term cascadetype multifractal mechanisms underlying HRV for both healthy and CHF subjects. A circadian evolution of this mechanism is only evidenced for the healthy group, suggesting an alteration of the sympathetic-parasympathetic

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balance for CHF patients. Significance: Results obtained for a large cohort of subjects suggest that the novel non Gaussian indices might robustly quantify crucial information for clinical risk stratification in CHF patients.

Index Terms—Multiscale, non Gaussian, wavelet leaders, heart rate variability, congestive heart failure.

I. INTRODUCTION

Heart Rate Variability analysis for risk stratification in Congestion Heart Failure: The fluctuations along time of the Heart Rate, often referred to as Heart Rate Variability (HRV), are known to constitute a valuable information regarding the health status of human beings. It is thus often used to distinguish patients suffering from pathologies from healthy subjects (cf. e.g., [1] and the numerous references therein for a review). Despite HRV monitoring being routinely used in clinical practice, refined analysis of the temporal dynamics underlying HRV have been and remain the subject of numerous and sustained academic research efforts. Indeed, HRV is known to be driven by intertwined interactions between physiological, neurological mechanisms and mechanical constraints, which result in temporal dynamics commonly accepted to be highly irregular and complex, thus requiring advanced statistical signal processing methodologies for analysis.

HRV analysis has notably been used for characterization of Congestive Heart Failure (CHF) of interest here (cf. e.g., [1]–[9]). Important recent advances in therapy for CHF patients were reported in medical studies showing a significant decrease in mortality for patients equipped with implantable cardioverter-defibrillators or undergoing resynchronization therapy. Deciding a priori which patients should benefit from such costly treatments therefore constitutes a major clinical stake. However, this requires to devise reliable *a priori* risk stratification strategies, which remains a difficult and challenging task that recently received growing academic interest and research efforts (cf. e.g., [2], [3], [10], [11]).

Related works: HRV time series were analyzed in several different ways, naturally starting with classical static (or marginal or first order) descriptive statistics such as the variance of the interarrival times. However, the seminal contribution [12] evidenced that temporal dynamics were key to HRV analysis. This triggered the massive use of spectral analysis (hence second order statistics) for HRV characterization, with indices such as

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energy in low (LF) and high (HF) frequency bands and spectral slope α [13], [14]. It has since been pointed out repeatedly that achieving classification with sufficient performance and potential use in practice requires to depart from such linear (or first-second order) analyses that implicitly assume Gaussian, stationary short-time correlated models (cf. e.g., [15], [16] and reference therein). Along that line, time-frequency representations and non stationary analysis tools were used, cf. e.g., [17]. Also, several different information theoretic quantities were massively used as complexity measures in HRV analysis, including entropy and entropy rate, usually computed using the sample entropy (sampEN) and approximate entropy (apEN) algorithms, respectively, cf. e.g., [4], [5]. More recently, the concepts of fractal and multifractal were also put forward for modeling long-term dependencies (cf. e.g., [2], [6], [7], [18]–[22]). The recently proposed multifractal and scattering analyses also marry non linear complexity measures with scalefree dynamics [23], [24]. Alternatively, departures from Gaussianity were also probed as diagnostic tools [3], [8], [9], [25].

Goals, contributions and outline: In this context, elaborating on the concepts of fractal and non linear analyses, the overall goal of the present contribution is to construct flexible non Gaussian multiscale representations that benefit from: i) the use of a large continuum of time scales, yet without requiring a priori exact scale-free dynamics and hence permitting potential departure from multifractal models; ii) the possibility for selecting the (range of) statistical orders for which analysis is to be made, hence clearly departing from the linear, or correlation, or second order, analysis and its underlying Gaussian assumption. These non Gaussian multiscale representations are embedded within the wavelet p-leader formalism, recently defined in [26] and shown to achieve robust estimation of the statistical characteristics of real-world data [27]. They are presented theoretically and illustrated on synthetic time series in Section II. The technique to design non Gaussian indices generalizes the intuition developed in a specific case in [3], which however relied on the less versatile detrended fluctuation analysis (DFA) multiscale quantities (cf., [27]-[29] for theoretical studies of DFA and comparisons with p-leaders).

Moreover, the present analysis intends to investigate the potential benefits of such non Gaussian multiscale representations for studying the differences in HRV between healthy subjects and patients suffering from congestive heart failure (CHF), as well as between survivor and non-survivor CHF patients. To that end, these tools are applied to a high quality database, described in Section III, consisting of 24 h HRV data for 108 CHF patients and 90 healthy subjects collected at Fujita Health University Hospital, Japan.

Results are presented and discussed in Section IV. Performance obtained from multifractal and non Gaussian multiscale indices are compared, both for 6-hour day-time period and in terms of evolution along the 24 h. Powers of statistical tests and Kaplan-Meier Survival Curves are used to quantify the significance of differences, with the following main findings: i) multifractal short-term cascade-type temporal dynamics underlying heart rate complexity exist not only for healthy, but also for CHF subjects; ii) for the later group, there is no circadian

evolution of multifractality and hence night time deactivation of sympathetic activity; iii) the combined use of non Gaussian multiscale indices is evidenced to possess high discriminative abilities between Non-Survivor (NS) and Survivor (SV) CHF patients at specific time scales and substantially outperforms multifractal and spectral and entropy based indices.

II. METHODOLOGY

A. Multiscale Quantities

1) Wavelet Coefficients: Let X(t) denote the time series to be analyzed. Given a mother wavelet ψ , i.e., a function that is well localized both in time and frequency and characterized by its number of vanishing moments N_{ψ} , a strictly positive integer defined as $\forall n=0,\ldots,N_{\psi}-1,\int_t^k\psi(t)dt\equiv 0$ and $\int_t^{N_{\psi}}\psi(t)dt\neq 0$, the $(L^1$ -normalized) discrete wavelet transform (DWT) coefficients $d_X(j,k)$ of X are defined by the inner products $d_X(j,k)=\langle \psi_{j,k}|X\rangle$, with $\{\psi_{j,k}(t)=2^{-j}\psi(2^{-j}t-k)\}_{(j,k)\in\mathbb{N}^2}$. We refer the reader to, e.g., [30], for precise definitions and details.

2) Wavelet p-Leaders: The p-leaders are defined as local ℓ^p -norms of the wavelet coefficients computed from the primitive of $X, Y(t) = \int_0^t X(s) ds$:

$$\ell^{(p)}(j,k) \triangleq \left(2^{j} \sum_{\lambda_{j',k'} \subset 3\lambda_{j,k}} |d_{Y}(j',k')|^{p} 2^{-j'}\right)^{1/p} \tag{1}$$

with $\lambda_{j,k} = [k2^j, (k+1)2^j)$ and $3\lambda_{j,k} = \bigcup_{m\{-1,0,1\}} \lambda_{j,k+m}$. The p-leaders, with $p \geq 0$, thus consist of weighted averages of wavelet coefficients existing at all finer scales $2^{j'} \leq 2^j$ and within a local temporal neighborhood $\lambda_{j',k'} \subset 3\lambda_{j,k}$. For p-leaders, estimation benefits from smaller variance, and both positive and negative statistical moments are well defined and numerically stable (unlike, e.g., those of DWT or DFA coefficients), hence permitting the practical assessment of their scale-by-scale distributions, see [26] for precise definitions and details, beyond the scope of this contribution. Following [27], p=1 is used in all analyses reported below. The superscript p is therefore dropped for notational simplification.

B. Scale-Free Dynamics and Multifractal Analysis

Scale invariance in HRV has already been massively studied, mainly using (second order) spectral analysis and 1/f models, cf. e.g., [15], [16]. The p-leaders enable to go beyond such a second order analysis of linear properties and permit to quantify multifractal scale invariance.

1) Moments, Cumulants and Scale Invariance: Let ℓ_s denote a p-leader coefficient $\ell(j,k)$ at scale $s=2^j$. A standard generating function expansion argument [31], [32] yields a relation for the finite moments $\mathbb{E}\ell_s^q < \infty$ of ℓ_s with the cumulants $C_m(s) = \operatorname{Cum}_m \log(\ell_s)$ of order m of $\log(\ell_s)$

$$L_q(s) \triangleq \log(\mathbb{E}\ell_s^q) = \log \mathbb{E}e^{q\log(\ell_s)} = \sum_{m=1}^{\infty} C_m(s) \frac{q^m}{m!}.$$
 (2)

Scale-free temporal dynamics implies a specific form of dependence with respect to scale:

$$L_q(s) = \kappa_q + \zeta(q)\log(s) \tag{3}$$

$$C_m(s) = c_m^0 + c_m \log(s) \tag{4}$$

where (4) follows by substituting (3) in (2) and implies $\zeta(q) = \sum_{m=1}^{\infty} c_m \frac{q^m}{m!}$, cf., [31]. The *scaling exponents* $\zeta(q)$ and *log-cumulants* c_m provide two alternative ways for quantifying the scale invariance properties of X.

2) Scale Invariance and Multifractal Spectrum: The scale-free properties of X(t) are intimately connected with the fluctuations of its point-wise regularity h(t) along t, as measured by the multifractal spectrum D(h): It is theoretically defined as the fractal (Hausdorff) dimension \dim_H of the set of instants t where X(t) has identical regularity h and can in practice be computed from either $\zeta(q)$ or c_m [33], [34]

$$D(h) \triangleq \dim_{H} \{t | h(t) = h\} \leq \inf_{q} (1 + qh - \zeta(q))$$
 (5)

$$\approx 1 + \frac{c_2}{2!} \left(\frac{h - c_1}{c_2}\right)^2 + \frac{-c_3}{3!} \left(\frac{h - c_1}{c_2}\right)^3 + \dots$$
 (6)

When based on p-leaders, the regularity at time t is measured by the p-exponent h(t), a generalization of the usual Hölder exponent, cf., [26], [27]. Expressions (3–6) highlight the strong link between D(h) and the scale invariance properties of X. In particular, c_1 quantifies the average local regularity of X; c_m , $m \geq 2$, relate to the multifractal properties of data, notably c_2 quantifies the degree of fluctuation of local regularity, and c_3 the asymmetry of these fluctuations about c_1 .

Finally, (3–6) permit the robust assessment of $\zeta(q)$, c_m and D(h) for discrete data X: by replacing theoretical moments $\mathbb{E}(\cdot)^q$ or cumulants Cum_m in (3) or (4) with sample moments or cumulants, estimates of $\zeta(q)$, c_m and D(h) can be obtained from linear regression of $L_q(s)$ or $C_m(s)$ for a range of scales $\log_2(s) = j \in [j_1, j_2]$, cf., [32], [34]. The p-leader multiscale and multifractal analysis of a HRV time series is schematically represented in Fig. 1. Other multiscale quantities have been previously proposed, notably also based on (discrete) wavelet transforms [32], [34], [35] (cf., [36] for a review).

3) Cumulants $C_m(s)$ and Departures From Gaussianity: The cumulants $C_m(s)$ underlying the multifractal formalism, cf., (2), provide a description of the distribution of $\log(\ell_s)$: e.g., C_1, C_2, C_3 and C_4 correspond to the mean, variance, skewness and kurtosis, etc. This can be illustrated using three prominent Gaussian and non Gaussian model processes for scale-free dynamics: The seminal fractional Brownian motion (fBm) [37], which is Gaussian, self-similar and satisfies $C_m(s) = const$ for $m \geq 2$; Multifractal random walk (MRW), which is non Gaussian of a special type of departure from Gaussian referred to here as log-normal, cf., [38], [39], specifically, $C_2(s) \neq const$ while $C_m(s) = const$ for $m \ge 3$; Compound Poisson cascade (CPC), which is another multifractal process with a more general form of non Gaussianity than MRW, i.e., $C_m(s) \neq const \ \forall m \geq const \$ 2, cf. [40] (process parameters are chosen to yield identical $C_1(s)$ for all processes and identical $C_2(s)$ for MRW and CPC). Fig. 2 compares $C_m(s)$ for these processes: $C_1(s)$ only captures

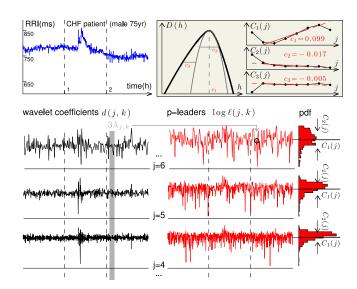


Fig. 1. Illustration of wavelet p-leader based multiscale and multifractal analysis.

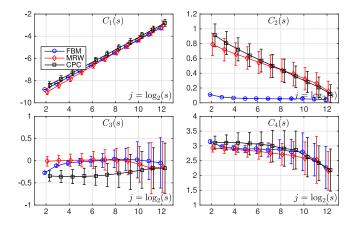


Fig. 2. Cumulants $C_m(s)$ for Gaussian self-similar fBm (blue solid) and non Gaussian multifractal MRW and CPC (red and black solid lines, respectively).

the linear, Gaussian (mean and variance/covariance) properties and thus does not permit to distinguish the three processes; $C_2(s)$ differs between Gaussian fBm and non Gaussian MRW and CPC but collapses on one curve for the latter two processes; $C_3(s)$ is essentially identical for Gaussian fBm and non Gaussian MRW as both processes have symmetrical distributions, yet differs for CPC which is designed to have a non symmetrical distribution. While theoretically different for all three processes, $C_4(s)$ are observed to fall within error bars. Therefore, beyond permitting a valuable descriptions of scale-free and multifractal dynamics, the $C_m(s)$ also relate to departures from Gaussian.

4) Limitations: Scale-free and multifractal analyses have already been massively employed to analyze HRV both in adults [6], [36] and fetuses [23], [34]. Nevertheless, the a priori assumption of exact scale-free dynamics implies a special form of scale dependence as in (3–4), and the sole use of scaling exponents $\zeta(q)$ or c_m can be limiting in applications. In addition, non Gaussian properties over scales s were reported to be use-

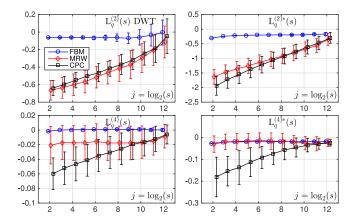


Fig. 3. Non Gaussian indices $L_s^{(2)}(\boldsymbol{q}), L_s^{(4)}(\boldsymbol{q})$ for Gaussian self-similar fBm (blue solid) and non Gaussian multifractal MRW and CPC (red and black solid lines, respectively).

ful for characterizing HRV, cf., [3], [8], [9]. Cumulants $C_m(s)$ quantify non Gaussianity of the $\log(\ell_s)$, yet there is no reason a priori that one single cumulant of a specific order should be relevant for a particular application. Further, $C_m(s)$ of high order m may be difficult to estimate accurately in practice. This motivates the definition of novel, versatile non Gaussian multiscale representations.

C. Non Gaussian Multiscale Representations

1) Definition and Key Properties: Inspired by the generalized moments approach introduced in the turbulence literature [41] and motivated by the non Gaussian index studied in [3], we propose here to define non Gaussian representations as linear combinations of cumulants. Specifically, we define a non Gaussian multiscale expansion of X, with $P \in \mathbb{N}^+$ and $q = (q_1, \ldots, q_{2P}), q_i \neq 0$ and $q_i \neq q_j$, as:

$$\mathbf{L}_{q}^{(2P)}(s) \triangleq \log \left(\prod_{i=1}^{P} \frac{\left(\mathbb{E}\ell_{s}^{q_{2i-1}}\right)^{\frac{1}{q_{2i-1}}}}{\left(\mathbb{E}\ell_{s}^{q_{2i}}\right)^{\frac{1}{q_{2i}}}} \right) \tag{7}$$

Making use of (2), it is straighforward to show that $\mathbf{L}_{q}^{(2P)}(s)$ equals an infinite weighted sum of the cumulants $C_{m}(s)$

$$\mathbf{L}_{q}^{(2P)}(s) = \sum_{i=1}^{P} \frac{L_{s}(q_{2i-1})}{q_{2i-1}} - \frac{L_{s}(q_{2i})}{q_{2i}}$$

$$= \sum_{m=2}^{\infty} C_{m}(s) \frac{\sum_{i=1}^{P} q_{2i-1}^{m-1} - q_{2i}^{m-1}}{m!}.$$
 (8)

Property 0: The evolution across scales s of $\mathbf{L}_{q}^{(2P)}(s)$ is an indicator for the non Gaussian properties of X. Indeed, for Gaussian data, it is straightforward to show that $\mathbf{L}_{q}^{(2P)}(s)$ equals a constant, whose value depends only on the choice of multiresolution quantity (here, p-leaders). Fig. 3 illustrates this for Gaussian fBm (blue lines): $\forall P, \forall q, \mathbf{L}_{q}^{(2P)}(s) = const.$

Property 1: $\mathbf{L}_{\mathbf{q}}^{(2P)}(s)$ does not depend on the value of $C_1(s)$ because all $q_{2i}^{m-1}-q_{2i+1}^{m-1}$ equal zero for m=1. Therefore,

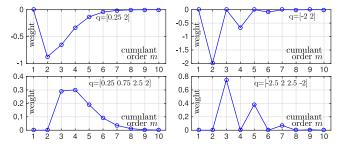


Fig. 4. Weights for the cumulants in $\mathbf{L}^{(2)}(q)$ with q=(0.25,2) (top left) and q=(-2,2) (top right), and $\mathbf{L}^{(4)}(q)$ with q=(0.25,0.75,2.5,2) (bottom left), and q=(-2.5,2,2.5,-2) (bottom right). Choices reported in the right column focus on either even (top) or odd (bottom) cumulants, those in the bottom row exclude C_2 from the analysis.

 $\forall P, \forall q, \mathbf{L}_{q}^{(2P)}(s)$ solely measures the *non linear* properties or departures from Gaussianity for X (cf., Fig. 3).

Property 2: The relative weighting of the cumulants $C_m(s)$ in the expansion $\mathbf{L}_q^{(2P)}(s)$ can be tuned and customized by choosing P and q appropriately. This versatility can be used to quantify, e.g., the *strength* of deviation from Gaussianity, or to finely characterize the *nature* of this departure.

Property 3: $\mathbf{L}_q^{(2P)}(s)$ permits to probe *cumulants of high order*, that would turn difficult to estimate, using only moments $L_q(s) = \log(\mathbb{E}\ell_s^q)$ of low orders q.

2) Versatility: The possible uses of this non Gaussian multiscale expansion are now illustrated on several examples.

Historical example: In [3], a first non Gaussian index was proposed and used for the multiscale assessment of the strength of deviation from Gaussian of HRV. It relied on detrended increments (i.e., quantities reminiscent of DWT coefficients) instead of ℓ_s , and on the specific choice P=1, $\mathbf{q}=(0.25,2)$. Fig. 4 (top left) shows that it mostly probes $C_2(s)$ and $C_3(s)$, while Fig. 3 (1st row, left column) shows that its evolution across scales effectively discriminates Gaussian fBm (blue line) from non Gaussian MRW and CPC (red and black lines, respectively). Yet, it does not permit the discrimination of the latter two non Gaussian processes.

Removing $C_2(s)$: We can refine the non Gaussian multiscale assessment and cancel any cumulant of order $m=M\geq 2$ in (8) by setting $P\geq 2$ and ${\boldsymbol q}$ such that $\sum_{i=1}^P q_{2i-1}^{m-1}-q_{2i}^{m-1}=0$, subject to $q_{2i-1}\neq q_{2j},\ i,j=1,\ldots,P$. The specific choice $P=2,{\boldsymbol q}=(u,\ u+r,\ v,v-r),r\neq\{0,-u,v,v-u,\}$ leads to removal of $C_2(s)$ in $\mathbf{L}_{\boldsymbol q}^{(4)}(s)$, as illustrated in Fig. 4 (bottom left). This enables us to discriminate non Gaussian of \log -normal nature from other types of non Gaussian. Fig. 3 (bottom left) shows that this permits distinguishing the log-normal type MRW (red line) from the more general departure from Gaussian found in CPC (black line).

Removing odd cumulants $C_{2m+1}(s)$: Moreover, the use of p-leader coefficients ℓ_s enables us to use moments of negative order q in the construction of the non Gaussian expansion, leading to additional desirable properties. Specifically, the choice $q_{2i-1} = -q_{2i}$, $q_{2i-1} \neq q_{2j}$, $i, j = 1, \ldots, P$, for $P \geq 1$ (i.e., q = (q, -q) for P = 1) removes all $C_m(s)$ of odd order

m from (8), as illustrated in Fig. 4 (top right). It thus permits to focus mainly on non Gaussian properties of symmetric nature. What is more, Fig. 3 (1st row, right column) shows that the negative-moment-based index leads to smaller (relative to variation with scale) error bars and facilitates distinguishing Gaussian (blue line) from non Gaussian (red and black lines).

Removing even cumulants $C_{2m}(s)$: All cumulants of even order can be cancelled from $\mathbf{L}_{\mathbf{q}}^{(2P)}(s)$ by setting $P \geq 2$ and $q_i = -q_{i+P}, \ i=1,\ldots,P$ subject to $q_{2i-1} \neq q_{2j}, \ i,j=1,\ldots,P$. This choice hence permits to jointly quantify the asymmetry properties; for P=2, this reads $\mathbf{q}=(q_1,\ q_2,\ -q_1,\ -q_2),q_1 \neq q_2$ and is illustrated in Fig. 4 (bottom right). Fig. 3 (bottom right) shows that this reduces relative error bar size and increases the contrast between log-normal type (red line) and general (black line) non Gaussian.

Remark 1: The proposed expressions of form (8) permit to jointly evaluate higher-order cumulants $C_m(s)$ by using moments of smaller order $q_i < m$. As an example, for the weights in Fig. 4 (bottom row), $\mathbf{L}_q^{(4)}$ is essentially a combination of the higher-order cumulants C_3 to C_7 , but is computed with moments of order $|q| \leq 2.5$.

Remark 2: With the assumption of scale invariance for X (i.e., (3) and (4)), (8) is of the form $\mathbf{L}_{q}^{(2P)}(s) = K_{q} + \log(s) \sum_{m=2}^{\infty} c_{m} \frac{\sum_{i=1}^{P} q_{2i-1}^{m-1} - q_{2i}^{m-1}}{m!}$ where we have gathered the terms independent of scale s induced by $\log \kappa_{q_{i}}$ in the constant K_{q} . The weights in this sum are identical to those in (8).

III. MULTISCALE HRV CHARACTERISTICS

During a period of two years (2000 and 2001), a cohort of 61 male and 47 female CHF patients, with age ranging from 21 to 92 (average 66.1 ± 14.8) years was monitored at Fujita Health University Hospital, Japan, out of which 39 (36.1%) died within the follow-up period of 33 ± 17 months (range 1–59 months). Medication status before hospital discharge did not significantly differ between survivor (SV) and non-survivor (NS) patients (cf. [3] for further details).

Before discharge from the Hospital, a 24-hour Holter ECG recording was collected for each single patient. R peaks were extracted from the HRV recordings and thoroughly reviewed for detection error correction and outlier removal. When atrial or ventricular premature complexes occurred, they were handled by median interpolation using two successive beats. Moreover, it has been checked that no sustained tachyarrhythmias were present in the recordings. The RR inter-arrival time values extracted from the HRV recordings are resampled at 4 Hz using cubic splines and analyzed as a time series, denoted as $X \equiv \{x_n, n = 1, \dots, N\}$. An example of a 3 hours block of one raw time series (male, CHF, 75yr, not resampled) and a sketch of its multifractal and multiscale analyses are plotted in Fig. 1. For comparison, we use an age-matched healthy (HM) control group of 90 subjects. The study was approved by the ethics committee of Fujita Health University and conformed to the principles outlined in the Declaration of Helsinki. All patients provided written informed consent.

TABLE I NON GAUSSIAN MULTISCALE INDICES $\mathbf{L}_{m{q}}^{(2P)}(s)$ USED IN THE PRESENT WORK

notation	moments q	cumulants C_m in (8) and interpretation
${\bf L}_{m q}^{(2)}(s)$	(0.25, 2)	$m \ge 2$
		any deviation from Gaussian
$\mathbf{L}_{\boldsymbol{q}}^{(2)*}(s)$		even order $m = 2, 4, \cdots$ mainly symmetric properties of non Gaussian
$\mathbf{L}_{\boldsymbol{q}}^{(4)}(s)$	(0.25, 0.75, 2.5, 2)	$m \ge 3$
		non log-normal nature of non Gaussian
$\mathbf{L}_{\boldsymbol{q}}^{(4)*}(s)$	(-2.5, 2, 2.5, -2)	$m \ge 3$ of odd order: $m = 3, 5, \cdots$
		asymmetry of non Gaussian

IV. RESULTS AND DISCUSSION

A. Estimation Setting and Statistical Testing

The multifractal and multiscale analyses are performed using a Daubechies wavelet with $N_{\psi}=3$ vanishing moments (the use of a different wavelet leads to similar conclusions). Linear regressions for multifractal coefficients c_m are computed for octaves $j_1=4\leq j\leq j_2=9$, corresponding to scales 5.3 s to 170.7 s, cf., Section IV-B1; the multiscale indices $\mathbf{L}_q^{(2P)}(s)$ are defined in Table I. Two different analyses will be performed: i) analysis of the daytime HRV for a 6 hour long block; ii) analysis of the evolution of HRV properties over 24 h for 2 h long blocks with 1 h overlap.

To assess differences between groups, pairwise Wilcoxon rank-sum (WRS) test p-values are computed. Between group differences for NS vs. SV are further evaluated using Kaplan-Meier cumulative survival curves (KMCS) and the Mantel-Haenszel (MH) log rank test.

B. Multifractal Properties of HRV

1) Scale Invariance and Multifractality in Daytime HRV: Fig. 6 (top row) plots the functions $C_{-}(s)$ m = 1, 2, 3, 4 (from

Fig. 6 (top row) plots the functions $C_m(s)$, m=1,2,3,4 (from left to right; averages and 95% error bars) for NS, SV, CHF and HM, for the 6 hour long blocks recorded between 17h00 and 23h00. All groups exhibit scale-free dynamics (i.e., power laws) for scales ranging from about 5 s to 170 s, that is, in scales corresponding to LF and VLF of traditional spectral analysis [1]. While scale-free dynamics have been reported for healthy subjects in [8], [25], [42], our findings suggest that such scale-free dynamics are also observed for CHF patients, cf., Fig. 6.

The manifestation of scale-free dynamics for all groups justifies the calculation and comparison of the multifractal spectra D(h), plotted in Fig. 5. The corresponding average multifractal coefficients c_m , m=1,2,3,4, are given in Table II. Both CHF and HM are characterized by nicely shaped average multifractal spectra, indicating well-behaved concave scaling functions $\zeta(q)$ as predicted by multifractal theory [40], [43] and an excellent fit for multifractal models for all groups, including CHF. Together, these findings permit to postulate that there is a multifractal mechanism underlying heart rate complexity, not only for healthy subjects as reported in [8], [25], [42] but also for CHF subjects. This suggests cascade-type dynamics in HRV data, both for healthy and CHF subjects, at time scales below 3 minutes.

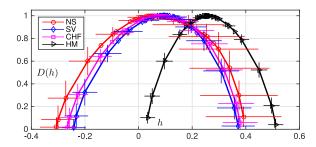


Fig. 5. Multifractal spectra D(h) (averages and 95% error bars) for scales (5.3 s, 170.7 s) for daytime HRV.

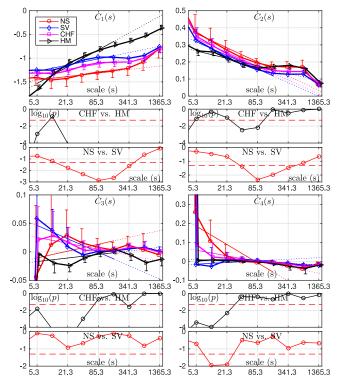


Fig. 6. Cumulants $C_m(s)$ as a function of scale s (in seconds; top subplots; averages and 95% error bars), and WRS p-values for CHF vs. HM and NS vs. SV (bottom subplots) for daytime HRV.

Further, the mode c_1 of the multifractal spectrum, quantifying the average degree of regularity of HRV, is found to be larger for HM than for CHF. For the width of D(h), related to $|c_2|$ and quantifying the range of fluctuations of the regularity of HRV along time, the converse holds: it is smaller for HM than for CHF Together, these two observations indicate that HRV of CHF is characterized by overall more irregular, erratic time evolution, than that of HM, both globally (smaller c_1) and locally (more multifractility, larger $|c_2|$), and corroborates similar findings on the self-similarity exponent H [9, Table 5]. Note that similar observations hold for SV and NS (smaller c_1 and larger $|c_2|$ for NS than SV), but differences are not significant, as discussed next.

2) Discriminative Power of Multifractal Properties of HRV: WRS p-values between the groups HM, CHF and NS, SV for c_m , m = 1, 2, 3, 4, reported in Table II, bottom rows),

TABLE II AVERAGE MULTIFRACTAL COEFFICIENTS c_m FOR SCALES (5.3 s, 170.7 s) (TOP) AND PAIRWISE WRS P-VALUES (BOTTOM)

Estimates	c_1	c_2	c_3	c_4	
HM	0.255	-0.032	0.005	-0.005	
CHF	0.079 -0.058		-0.008	-0.022	
NS	0.053	-0.060	0.009	-0.070	
SV	0.095	-0.057	-0.017	0.006	
WRS p-value	c_1	c_2	c_3	c_4	
CHF vs. HM	< 1e-8	0.276	3.0e-4	0.111	
NS vs. SV	0.156	0.720	0.932	0.068	

TABLE III TOP: SPECTRAL (Col. 2-5, $(V/ms)^2$) AND ENTROPY (Col. 6-7) FEATURES; BOTTOM: PAIRWISE WRS P-VALUES

Estimates	LF	HF	LF/HF	α_{PSD}	sampEN	apEN
HM	0.057	0.028	3.001	2.314	0.267	0.324
CHF	0.064	0.160	0.713	1.362	0.431	0.600
NS	0.064	0.132	0.692	1.476	0.378	0.538
SV	0.064	0.176	0.725	1.300	0.463	0.637
WRS p-value	LF	HF	LF/HF	α_{PSD}	sampEN	apEN
CHF vs. HM	0.435	< 1e-8	< 1e–8	< 1e-8	< 1e-8	< 1e-8
NS vs. SV	0.406	0.173	0.927	0.327	0.062	0.127

assess their independent discriminative power. First, for HM vs. CHF, c_1 and c_3 have extremely high (p < 1e - 8) and high (p < 0.001) significance, respectively, while c_2 and c_4 are not significant. Second, for NS vs. SV, none of the multifractal coefficients are significant (p > 0.05). We conclude that the differences in multifractal dynamics of HRV are discriminative for healthy vs. CHF, but not for NS vs. SV subjects, excluding their use as indicators of mortality risk stratification in CHF.

For comparison, Table III shows average estimates and p-values for the spectral and information theoretic analysis methods described in Section I: energy in the low (LF) and high (HF) frequency band, LF/HF ratio and spectral slope α_{PSD} , sampEN and apEN (computed using standard parameters, embedding dimension m=2 and ball radius $r=0.2\times$ data standard deviation). The results lead to similar conclusions: Differences are significant for HM vs. CHF (p<1e-8, except for LF) but not discriminative for NS vs. SV subjects.

3) Circadian Evolution of Multifractal Properties: To assess the evolution of the multifractal properties of HRV along 24 hours of a day, estimates of c_m , m = 1, 2, 3, 4 (from top to bottom) obtained for two hour long blocks with one hour overlap, for HM, CHF and NS, SV (left and right column) are plotted as a function of the block center in Fig. 7. For the control group HM, there is a clear day-night evolution of multifractal parameters, characterized by a significant decrease for the magnitudes of c_1 and c_2 during sleep. This decrease is caused by the *deactivation* of sympathetic predominance during sleep [44]. Interestingly, there also appears to be a significant increase of the magnitude of c_4 during sleep, which has not been reported before. During sleep, because of the intermittent changes of sleep stages, non Gaussian behavior having heavy tails was observed (see [42]). These intermittent changes of sleep stages may result in the negative c_4 .

For CHF patients, a comparable circadian modulation of the multifractal properties of HRV cannot be observed. At best, a

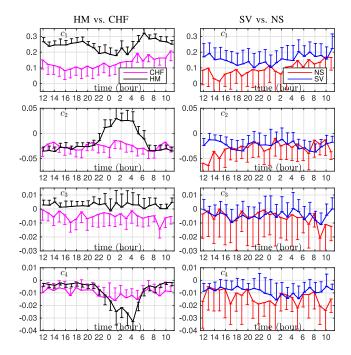


Fig. 7. Estimates for for c_m , m= 1, 2, 3, 4 (top to bottom row; median and 2.91 \mod/\sqrt{N}) obtained for 2 h long windows with 1 h overlap, plotted as a function of window central time for HM & CHF (left column) and NS & SV (right column).

weak difference is found between early morning and evening for the values of c_1 for SV. This suggests that for CHF patients, the sympathetic-parasympathetic balance does not follow the circadian cycle of healthy subjects. In particular, it lacks a daily period with reduced or without sympathetic activity.

C. Multiscale Properties of HRV

1) Cumulants $C_m(s)$: The cumulants $C_m(s)$, plotted in Fig. 6 (top rows), together with WRS p-values for HM vs. CHF and NS vs. SV, as a function of each independent scale s (bottom rows) show that: i) The cumulant $C_1(s)$ strongly differs between HM and CHF and between NS and SV, for a large range of time scales from 20 to 350 seconds. ii) The cumulants $C_2(s)$ and $C_4(s)$ also have significant p-values (yet larger than C_1) for both pairs of groups. iii) There are no significant difference between NS and SV for the third cumulant $C_3(s)$. iv) Interestingly, the scales yielding smallest WRS p-values are different for $\{C_1, C_2\}$ and $\{C_3, C_4\}$ and are identified as 85.3 s and 21.3 s, respectively.

2) Non Gaussian Indices $\mathbf{L}_{q}^{(2P)}(s)$: Fig. 9 plots the wavelet p-leader based multiscale non Gaussian indices $\mathbf{L}_{q}^{(2P)}(s)$. A first important observation is that the non Gaussian indices $\mathbf{L}_{q}^{(2)}(s)$ and $\mathbf{L}_{q}^{(2)*}(s)$ display a significant evolution across scales, confirming that HRV has non Gaussian multiscale statistics. Yet, they tend to a constant value - and hence the distribution of ℓ_s converges to Gaussian - for large scales \geq 170 s. Such a behavior is inherent to and a key characteristic of multifractal cascades, cf. e.g., [31]. It thus further strengthens our conclusions on the

cascade-type multifractal nature of heart rate dynamics for both healthy and CHF subjects, cf., Section IV-B1.

The magnitude of $\mathbf{L}_{q}^{(2)}$ is an indicator for the *degree* of non Gaussianity and is found to be larger for CHF than for HM, and larger for NS than for SV: Thus, CHF is stronger non Gaussian than HM, and NS stronger than SV. These differences are found to be significant for HM vs. CHF at fine scales, and for NS vs. SV in the vicinity of scale 85.3 s. The observation of significantly more non Gaussian distributions for HRV of NS than SV corroborates findings reported in [3].

Indications on the *nature* of the departure from Gaussian for the multiscale distributions of heart rate dynamics are obtained from $\mathbf{L}_{q}^{(4)}(s)$ and $\mathbf{L}_{q}^{(4)*}(s)$. They systematically take on nonzero values and thus suggest that the departures from Gaussian are not of log-normal type. Moreover, these observations are consistent between $\mathbf{L}_{q}^{(4)}(s)$ and $\mathbf{L}_{q}^{(4)*}(s)$. This suggests that the difference in multiscale statistics between NS and SV is to large extent driven by asymmetry because the index $\mathbf{L}_{q}^{(4)*}(s)$ does not depend on cumulants of odd order. These deviations from simple log-normal type non Gaussian are stronger for CHF than for HM, and stronger for NS than for SV. They are significant for CHF vs. HM, and to a lesser extent also for NS vs. SV, at scales in the vicinity of 21.3 s.

To conclude, differences in the *amount* and in the *nature* of the deviations from Gaussian are found significant between groups. Interestingly, the most significant time scales do not coincide for non Gaussian amount (85.3 s for $\mathbf{L}_{q}^{(2)}$ and $\mathbf{L}_{q}^{(2*)}$) and nature (21.3 s for $\mathbf{L}_{q}^{(4)}$ and $\mathbf{L}_{q}^{(4*)}$). This is in agreement with earlier observations on C_m (since $\mathbf{L}_{q}^{(2)}$ mostly depend on C_2 , and $\mathbf{L}_{q}^{(4)}$ only on C_m , $m \geq 3$).

3) Circadian Evolution of Non Gaussian Indices: To assess the evolution of the multiscale non Gaussian properties of HRV along 24 hours of a day, estimates of $\mathbf{L}_{q}^{(2P)}(s)$ obtained for two hour long blocks with one hour overlap, for the scales identified as significant in the paragraph section, for HM, CHF and NS, SV (left and right column) are plotted as a function of the block center in Fig. 8. For the control group HM, there is a clear and distinct day-night evolution for all non Gaussian indices \mathbf{L}_{q} . Specifically, the values for $\mathbf{L}_{q}^{(2)}(85.3 \text{ s})$ and $\mathbf{L}_{q}^{(2)*}(85.3 \text{ s})$ decrease during night time, while $\mathbf{L}_{\boldsymbol{q}}^{(4)}(21.3 \text{ s})$ and in particular $\mathbf{L}_{q}^{(4)*}(21.3 \text{ s})$ increase. This indicates a night time increase in non Gaussianity (i.e., magnitude of $\mathbf{L}_{\boldsymbol{q}}^{(2)}$) that is driven both by purely asymmetric $(\mathbf{L}_{m{q}}^{(4)*})$ and even order $(\mathbf{L}_{m{q}}^{(4)})$ properties and characterized by stronger deviation from log-normal (larger $\mathbf{L}_{q}^{(4)}$). In [42], the non Gaussian behavior was reported to have the heaviest tails at scale $\approx 100 \ \text{s}$ during sleep, and the observed $\mathbf{L}_{q}^{(2)}$ in healthy subjects is consistent with this observation. The non Gaussian distributions with non log-normal-type heavy tails could be a consequence of the intermittent changes of sleep stages, as discussed above.

Such a circadian modulation of the non Gaussian properties of HRV is not present for CHF patients. CHF is characterized by autonomic imbalance including decreased parasympathetic activity and sympathetic hyperactivity. Thus,

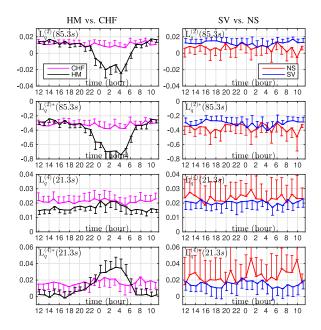


Fig. 8. Estimates for for $\mathbf{L}_{q}^{(2)}(s)$ and $\mathbf{L}_{q}^{(4)}(s)$ (median and 2.91 · mad/ \sqrt{N}) obtained for 2 h long windows with 1 h overlap, plotted as a function of window central time for HM & CHF (left column) and NS & SV (right column).

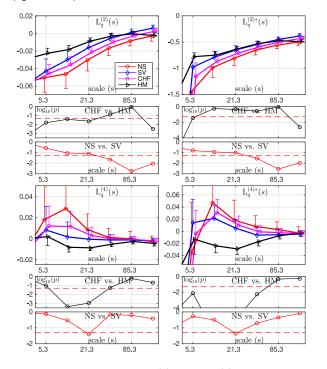


Fig. 9. Non Gaussian indices $\mathbf{L}_{m{q}}^{(2)}(s)$ and $\mathbf{L}_{m{q}}^{(4)}(s)$ as a function of scale s (in seconds; averages and 95% error bars, top rows), and WRS p-values for CHF vs. HM and NS vs. SV (bottom rows) for daytime HRV.

sympathetic-dominant autonomic imbalance may weaken the circadian pattern of the HRV modulation.

4) Discriminative Power of Multiscale Statistics of HRV: We illustrate the discriminative power of the observed differences in multiscale distributions as indicators of mortality risk in CHF. In Fig. 10 (left), KMCS and MH statistics are plotted for $\mathbf{L}_q^{(2)}(85.3~\mathrm{s})$, which yielded most significant WRS p-value

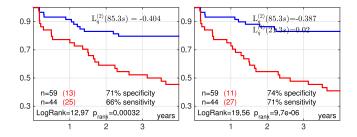


Fig. 10. KMCS for $\mathbf{L}_{q}^{(2)}$ (left) and combined $(\mathbf{L}_{q}^{(2)}, \mathbf{L}_{q}^{(4)})$ (right) with MH log rank test.

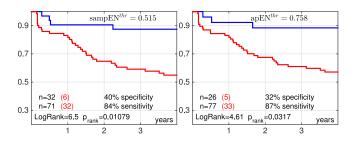


Fig. 11. KMCS for sampEN (left) and apEN (right).

in the previous sections and is similar to the historical example [3] (there, computed for detrended increments). The threshold of the survival curves was chosen such as to maximize the MH log rank. The non Gaussian index $\mathbf{L}_q^{(2)}(85.3)$ yields a p-value of 0.0003 and the MH log rank of 13.0, which is consistent with the results reported in [3], and leads to sensitivity (specificity) of 66% (71%), respectively. Similar results are obtained for $\mathbf{L}_q^{(2)*}(85.3)$, $\mathbf{L}_q^{(4)}(21.3)$ and $\mathbf{L}_q^{(4)*}(21.3)$.

For comparison, KMCS and MH statistics for spectral and entropy based analyses (best results obtained for sampEN and apEN, reported in Fig. 11) are significant but substantially less predictive (p-value > 0.01, MH log rank < 7, specificity $\le 40\%$; results for spectral analysis are worse, with p-value ≥ 0.04 , MH log rank < 4.5, and not reported for space reason).

Fig. 10 (right) further investigates the combined use of non Gaussian expansion indices and plots the KMCS and MH statistics when $\mathbf{L}_{q}^{(2)}(85.3)$ and $\mathbf{L}_{q}^{(4)}(21.3)$ are used together. It indicates that the combined predictive power of non Gaussian indices significantly improves predictive power, leading to pvalue smaller than 10^{-5} , log rank of 19.5, increasing sensitivity by 5% and specificity by 3%, respectively. This suggests that the precise nature of non Gaussian is a discriminative feature for CHF mortality risk stratification and leads us to conclude that non Gaussian expansion indices are useful indicators of mortality risk in CHF.

V. CONCLUSION

This work proposed the construction of novel, versatile non Gaussian multiscale representations for the use of a large continuum of scales without a priori assuming scale-free dynamics assumptions and the joint analysis of customizable ranges of (higher) statistical orders. The representations remain embedded in the recently proposed wavelet *p*-leader multifractal

formalism and inherit from its theoretical grounding and practical robustness. These non Gaussian multiscale representations were used here for investigating the non linear dynamics of HRV for a database of 24 hour long HRV recordings from 90 healthy subjects and 108 CHF patients. It enabled us to evidence the existence of short-term cascade-type multifractal mechanisms underlying HRV for both healthy and CHF subjects for time scales below 3 minutes. Further evidence was provided for the circadian evolution of this mechanism for the healthy group, while the constancy of multifractal parameters suggests a difference in the sympathetic-parasympathetic balance for CHF patients. Moreover, several specific non Gaussian indices were constructed and individually evidenced to possess high discriminative abilities between NS and SV CHF patients for specific time scales (21.3 s and 85.3 s), substantially improving upon spectral, entropy and multifractal based indices. Future work will include studying the use of different combinations of higher order cumulants in the non Gaussian multiscale representation, and the precise assessment of the combined predictive power of non Gaussian indices for mortality risk in CHF patients.

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